

CHARM: 55 YEARS AND COUNTING

J. Rosner (U of Chicago) – MaryK Fest – September 2019

A fourth quark has played a key role in understanding the weak and strong interactions since its proposal in 1964

Today: some early history ([hep-ph/9811359](https://arxiv.org/abs/hep-ph/9811359) for more details) and progress over the subsequent $>$ half-century

“Charm” bestowed by Bjorken and Glashow [PL **11**, 255 (1964)] for a fourth hadron-lepton-symmetric quark

Quartets: Gell-Mann, PL **8**, 214 (1964); Maki-Ohnuki [PTP **32**, 144 (1964)]; Hara [PR **134**, B701 (1964)]; Amati+, Nuovo Cim. **34**, 1732 (1964) & PL **11**, 190

Quartet assignment \Rightarrow suppression of strangeness-changing neutral currents; good behavior of higher-order weak interactions; higher-order effects $\Rightarrow m_{\text{charm}} \lesssim O(2 \text{ GeV})$: Glashow-Iliopoulos-Maiani, PR D **2**, 1285 (1970)

A SIMPLIFIED MOTIVATION

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Weak charge-raising current J_μ^+ : $Q^+ \equiv \int d^3x J_0^+$

Cabibbo-Gell-Mann-Lévy: In basis (u, d, s) :

$$Q^+ = \begin{bmatrix} 0 & \cos \theta & \sin \theta \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}; \quad Q^- = \begin{bmatrix} 0 & 0 & 0 \\ \cos \theta & 0 & 0 \\ \sin \theta & 0 & 0 \end{bmatrix}$$

SU(2) algebra: $Q^3 \equiv (1/2)[Q^+, Q^-]$, $Q^\pm = \pm[Q^3, Q^\pm]$

Q^3 has strangeness-changing neutral currents:

$$Q^3 = \frac{1}{2} \begin{bmatrix} 1 & 0 & 0 \\ 0 & -\cos^2 \theta & -\sin \theta \cos \theta \\ 0 & -\sin \theta \cos \theta & -\sin^2 \theta \end{bmatrix} \begin{array}{l} \text{Undesirable} \\ \text{off-diagonal} \\ \text{terms} \end{array}$$

Can't ignore these in a gauge theory of weak interactions

ADDING CHARM

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New quark c in basis (u, c, d, s) leads to:

$$Q^+ = \begin{bmatrix} 0 & 0 & \cos \theta & \sin \theta \\ 0 & 0 & -\sin \theta & \cos \theta \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}; \quad Q^- = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ \cos \theta & -\sin \theta & 0 & 0 \\ \sin \theta & \cos \theta & 0 & 0 \end{bmatrix}$$

Now $Q^3 = (1/2)[Q^+, Q^-] = (1/2)\text{diag}(1, 1, -1, -1)$ and strangeness-changing neutral currents have been cancelled

Mary K and Ben Lee [PR D **10**, 897 (1974)]: charm in loop diagrams for K - \bar{K} mixing, $K_L \rightarrow \mu^+ \mu^-$, $K_L \rightarrow \gamma \gamma$ within electroweak unification $\Rightarrow m_c \lesssim 2 \text{ GeV}$, $\gg m_u$

Rich spectroscopy implied: mesons $D^0 = c\bar{u}$, $D^+ = c\bar{d}$ [D = “doublet”], $F = c\bar{s}$ [“funny”; renamed D_s by PDG], baryons like $C_0^+ = udc$ [renamed Λ_c by PDG]

GIM PREDICTIONS

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Branching ratio $\mathcal{B}(D^0 = c\bar{u} \rightarrow K^- \pi^+)$ of a few %

Strong production of charm-anticharm pairs

Direct leptons in charm decays

Charm production in neutrino reactions

Neutral flavor-preserving currents

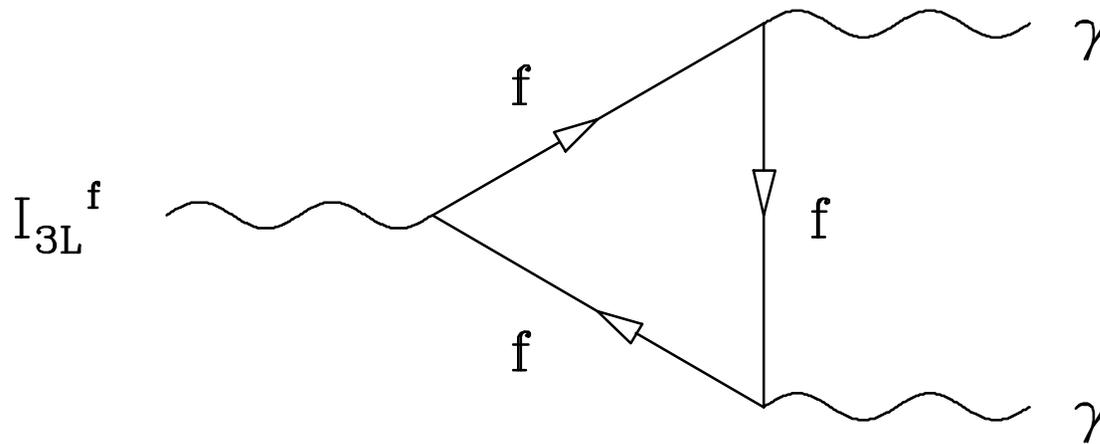
Z^0 observability in e^+e^- direct channel annihilations

Searches advocated by Carlson and Freund; Snow

Glashow at Meson 1974: Charm lifetimes between 10^{-13} and 10^{-12} s, similar semileptonic and hadronic decay rates, strange particles in final state, dileptons in neutrino reactions. Offered to eat his hat if charm was not found by the next Meson Conference (it was)

TRIANGLE ANOMALIES

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Bouchiat-Iliopoulos-Meyer, PL **38B**, 519 (1972); Georgi-Glashow, PR D **6**, 429 (1972); Gross-Jackiw, PR D **6**, 477 (1972): $\sum_f Q_f^2 I_{3f} = 0$ (I_{3f} = weak isospin):

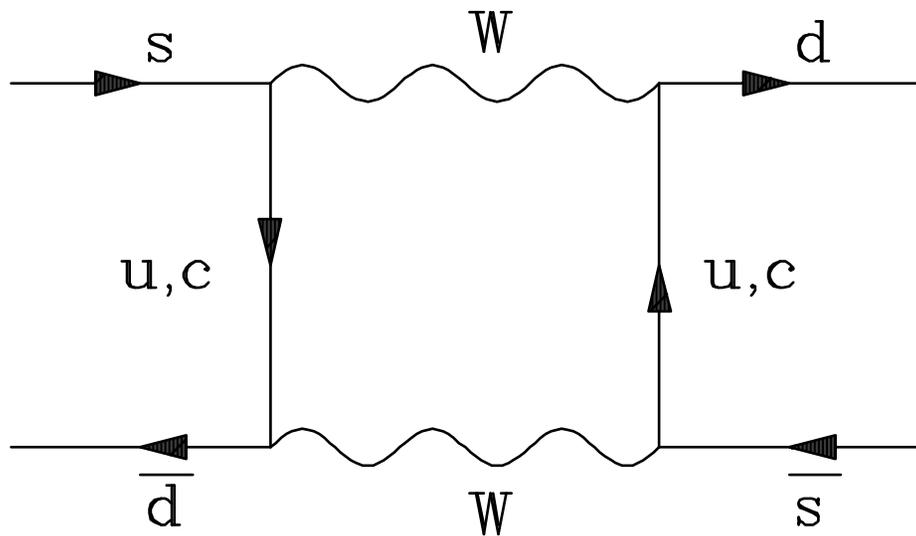
Leptons: $(0)^2 \left(\frac{1}{2}\right) + (-1)^2 \left(-\frac{1}{2}\right) = -\frac{1}{2}$

Quarks: $3 \left[\left(\frac{2}{3}\right)^2 \left(\frac{1}{2}\right) + \left(-\frac{1}{3}\right)^2 \left(-\frac{1}{2}\right) \right] = \frac{1}{2}$

Cancellation within each family requires charmed quark

GAILLARD-LEE CALCULATION ^{6/35}

Example of cancellation: $K^0 - \bar{K}^0$ mixing



plus graph with W^+W^-
in t channel

Contribution $\sim (m_c^2 - m_u^2) / M_W^2$

Either $m_c^2 \sim m_u^2$ (both large)
or $m_c \sim 1.5 - 2 \text{ GeV} \gg m_u$

Prediction $\mathcal{B}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \simeq 10^{-10}$ (now $\times \sim 0.8$)

Current value $(1.7 \pm 1.1) \times 10^{-10}$ (BNL E-949);
CERN NA-62 should collect about 80 such events

Some caution was expressed in relating heavy-quark masses to those of hadrons containing the heavy quark. The link between the two was stronger than anticipated.

WAS CHARM LOST?

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1974: Mary K and Ben invited me to join them; I had been involved with hadron spectroscopy and looking for charm in neutrino interactions: e.g., $\nu_\mu A \rightarrow \mu^- c\bar{s}A$

How could we have missed a whole family of hadrons with masses in the 2 GeV range?

Multiparticle decays of charm were more important than we realized; combinatorial backgrounds needed to be defeated

Charm-anticharm bound state with $J^{PC} = 1^{--}$ should show up in Drell-Yan processes and e^+e^- annihilation; Okubo-Zweig-Iizuka rule suppressing $c\bar{c}$ transitions to light-quark hadrons would make its total width small but also would suppress its production in hadronic reactions

Short-track signatures were crucial in identifying charm

A key signal was missed because of poor mass resolution

GLR CHARM SIGNATURES

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M. K. Gaillard, B. W. Lee, and J. L. Rosner, Rev. Mod. Phys. **47**, 277 (1975) (main text was written before J/ψ)

“Direct” lepton production

Large numbers of strange particles

Narrow peaks in mass spectra of hadrons

Apparent strangeness violations

Short tracks ... lifetimes of order 10^{-13} s

Di-lepton production in neutrino reactions

Narrow peaks in e^+e^- or $\mu^+\mu^-$ mass spectra

Transient threshold ... in deep inelastic leptonproduction

$$\sigma(e^+e^- \rightarrow \text{hadrons})/\sigma(e^+e^- \rightarrow \mu^+\mu^-) \Rightarrow 3\frac{1}{3} \dots$$

GLR CHARM SIGNATURES

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“Direct” lepton production: ✓ but other sources too

Large numbers of strange particles: τ “dilution”

Narrow peaks in mass spectra of hadrons ✓

Apparent strangeness violations: Good particle ID needed

Short tracks ... lifetimes of order 10^{-13} s: Some longer

Di-lepton production in neutrino reactions: ✓

Narrow peaks in e^+e^- or $\mu^+\mu^-$ mass spectra ✓ ($J/\psi, \dots$)!

Transient threshold ... in deep inelastic leptonproduction:

$R = \sigma(e^+e^- \rightarrow \text{hadrons}) / \sigma(e^+e^- \rightarrow \mu^+\mu^-) \Rightarrow 3\frac{1}{3} \dots: \tau$ too

Many of these realized only with measurement progress:
short tracks, particle ID, soft pion in $D^{*+} \rightarrow D^0\pi^+, \dots$

EARLY HINTS OF CHARM

10/30

Christenson +, BNL: 3.5 GeV shoulder in $\mu^+\mu^-$ spectrum;
Lederman at a BNL seminar in July 1971: “Looks like a resonance to me!” (He was talked out of it)

K. Niu and collaborators at Nagoya U.: short tracks in emulsions exposed to cosmic rays. One event consistent with $D^+ \rightarrow \pi^+\pi^0$, mass ~ 2 GeV, lifetime $\sim 10^{-14}$ s.

Dimuon events in neutrino interactions: semileptonic charmed particle decays such as $\nu_\mu d \rightarrow \mu^- c \rightarrow \mu^- s \mu^+ \nu_\mu$ (B. Lee, comment to D. Cundy at London ICHEP, 1974)

Simultaneous announcement of J and ψ , $c\bar{c}$ candidate at 3.1 GeV, and subsequent radial excitation candidate ψ' at 3.7 GeV, greatly narrowed possible charm parameters

Those who were convinced it was charm included Appelquist and Politzer; De Rújula, Georgi, and Glashow (DGG); and Eichten *et al.*; not all (Jan. 6, 1975 PRL)

SORTING OUT CHARM (1975-6)^{11/35}

“Don’t give up the ship. It has just begun to sink.”

$D^0 = c\bar{u} \rightarrow K^-\pi^+$ expected; branching ratio only a few % because of many possible decay modes

Rise in $R = \sigma(e^+e^- \rightarrow \text{hadrons})/\sigma(e^+e^- \rightarrow \mu^+\mu^-)$ was larger than expected: Contribution $\Delta R = 1$ of $\sigma(e^+e^- \rightarrow \tau^+\tau^-)$ needed to be taken into account.

Increase in strange particle production above charm threshold offset by dearth of strange particles in τ decays

Possible charmed baryon production by neutrinos: E. G. Cazzoli +, PRL **34**, 1125 (1975)

Charm discovery at SLAC: $D^0 \rightarrow K^-\pi^+, K^-\pi^+\pi^-\pi^+$ [PRL **37**, 255 (1976)] aided by soft-pion signature in $D^{*+} \rightarrow \pi^+D^0$ [Nussinov, PRL **35**, 1672 (1975)], time-of-flight

SHORT-TRACK DETECTION 12/35

Emulsion: Niu+, PTP **46**, 1644 (1971); FNAL E531, E653

Charm photoproduction: Fermilab E516, E691, E687, E831: short tracks from picosecond lifetimes (1983 –)

Charm hadroproduction a difficult task: J. Ritchie, thesis, University of Rochester, 1983, & references therein

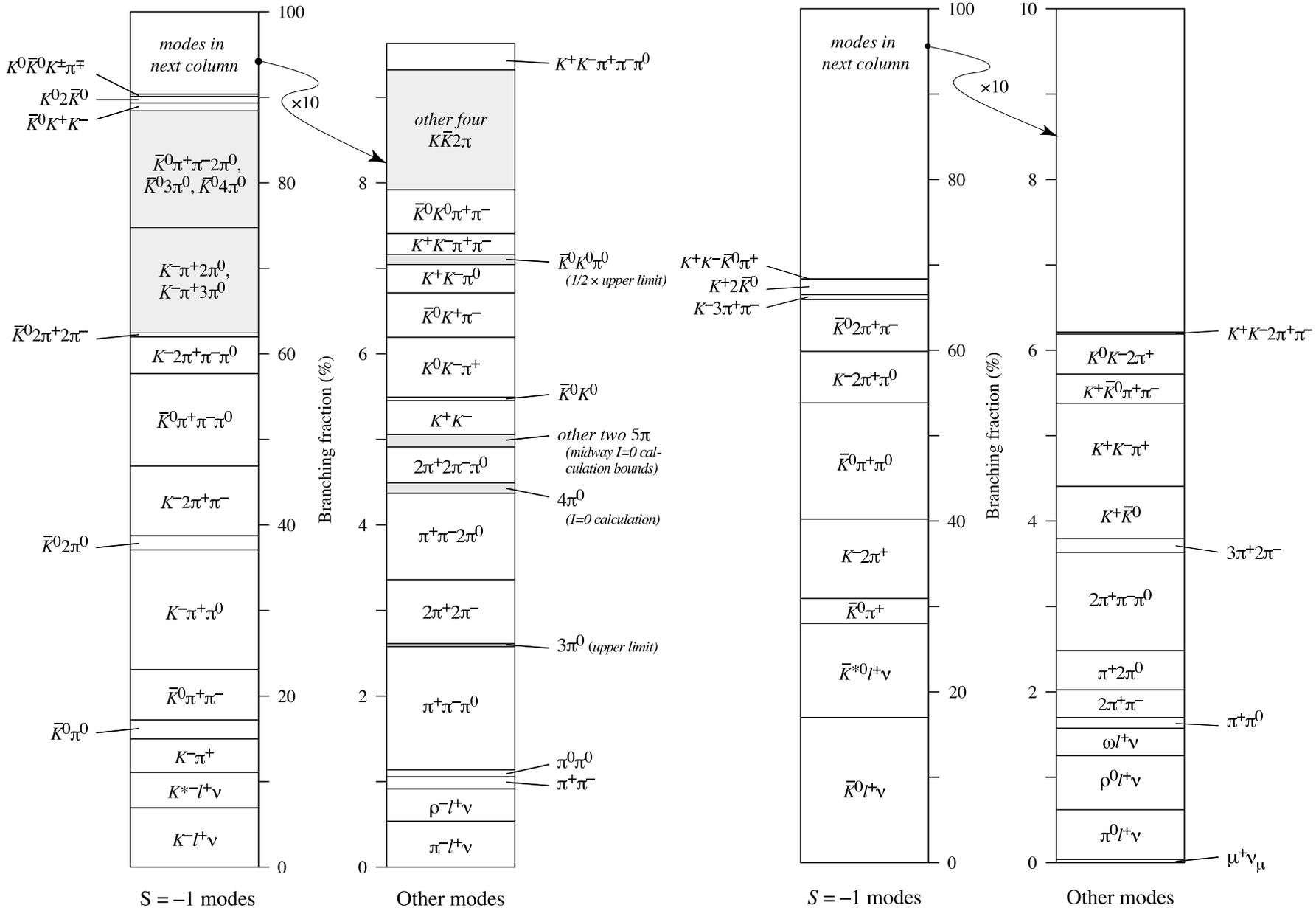
Initially charmed hadrons were not reconstructed but their presence was inferred from direct leptons

Review of early hadroproduction successes: J. Appel, Ann. Rev. Nucl. Part. Sci. **42**, 367 (1992)

Pioneers: e.g., LEBC (rapid-cycling high-resolution bubble chamber); ACCMOR Collab. [R. Butler+, NIM **213**, 201 (1983)], first use of CCDs in a hadron beam

Also, e.g., FNAL E769, E791; CERN WA82, WA89, WA92

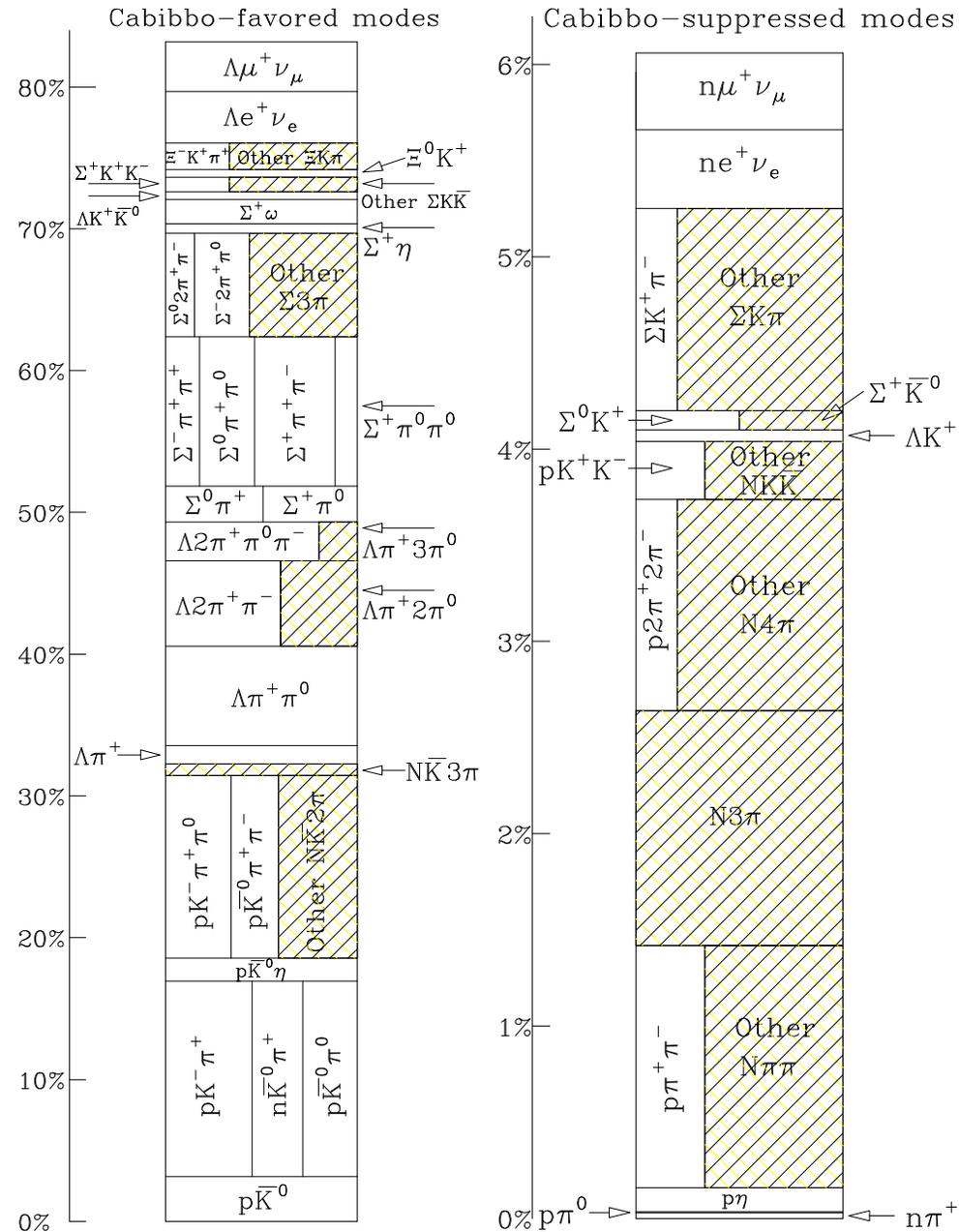
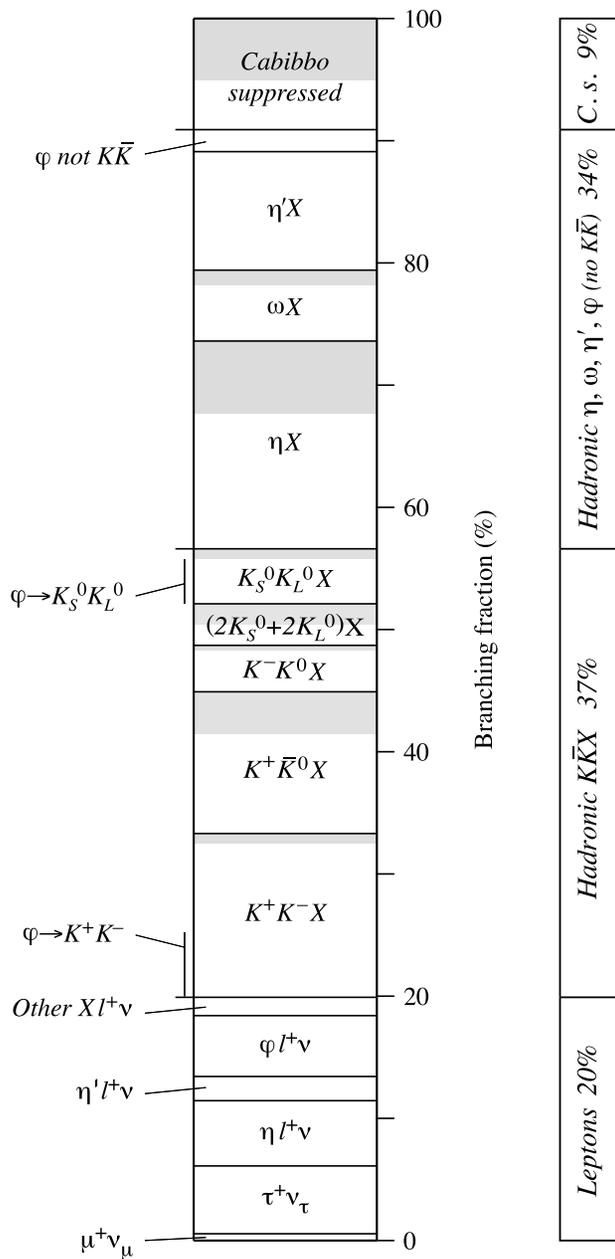
D^0 AND D^+ DECAYS



This and related plots: with M. Gronau, C. Wohl

D_s DECAYS

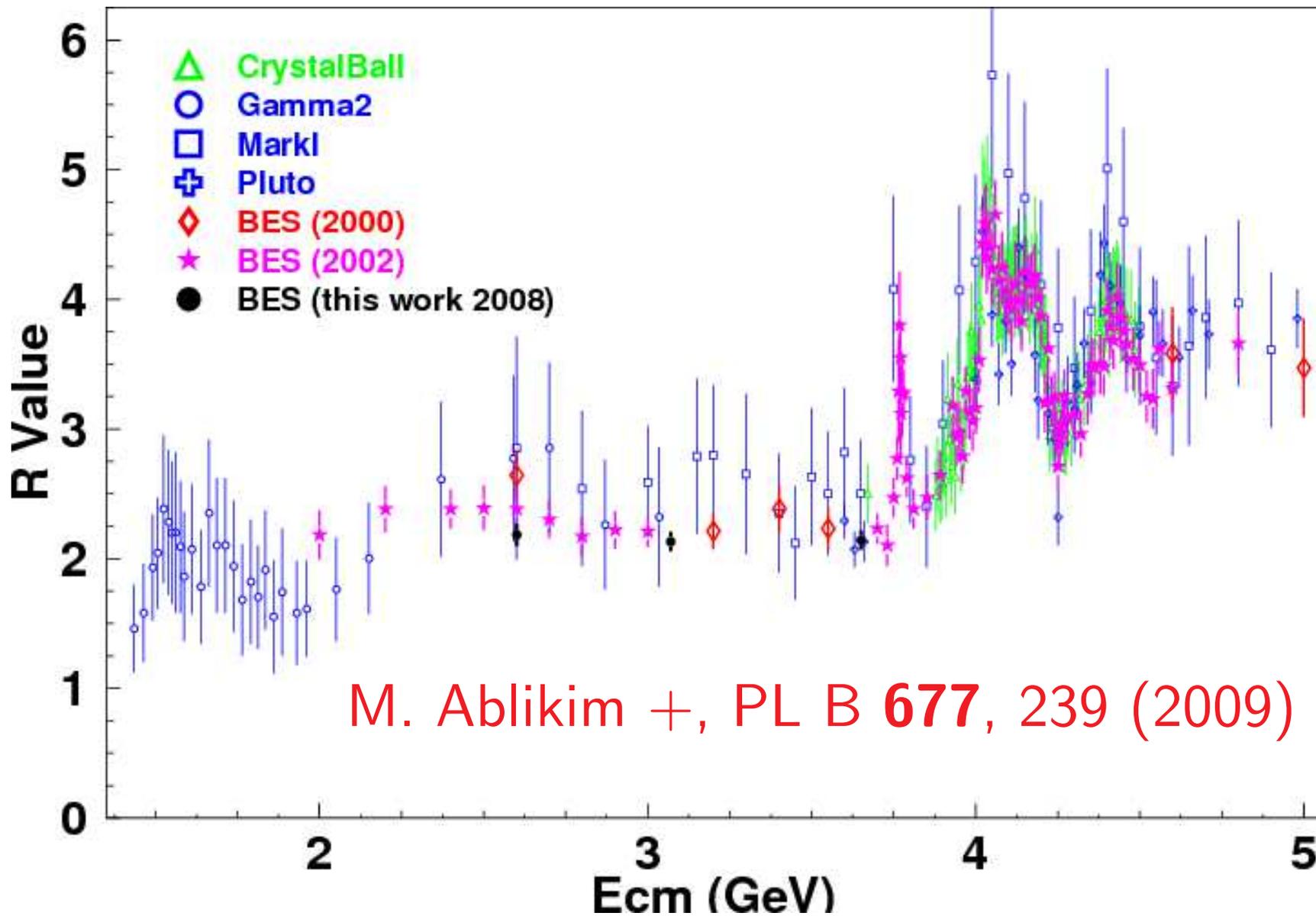
Λ_c DECAYS



Substantial progress in identifying most decay modes

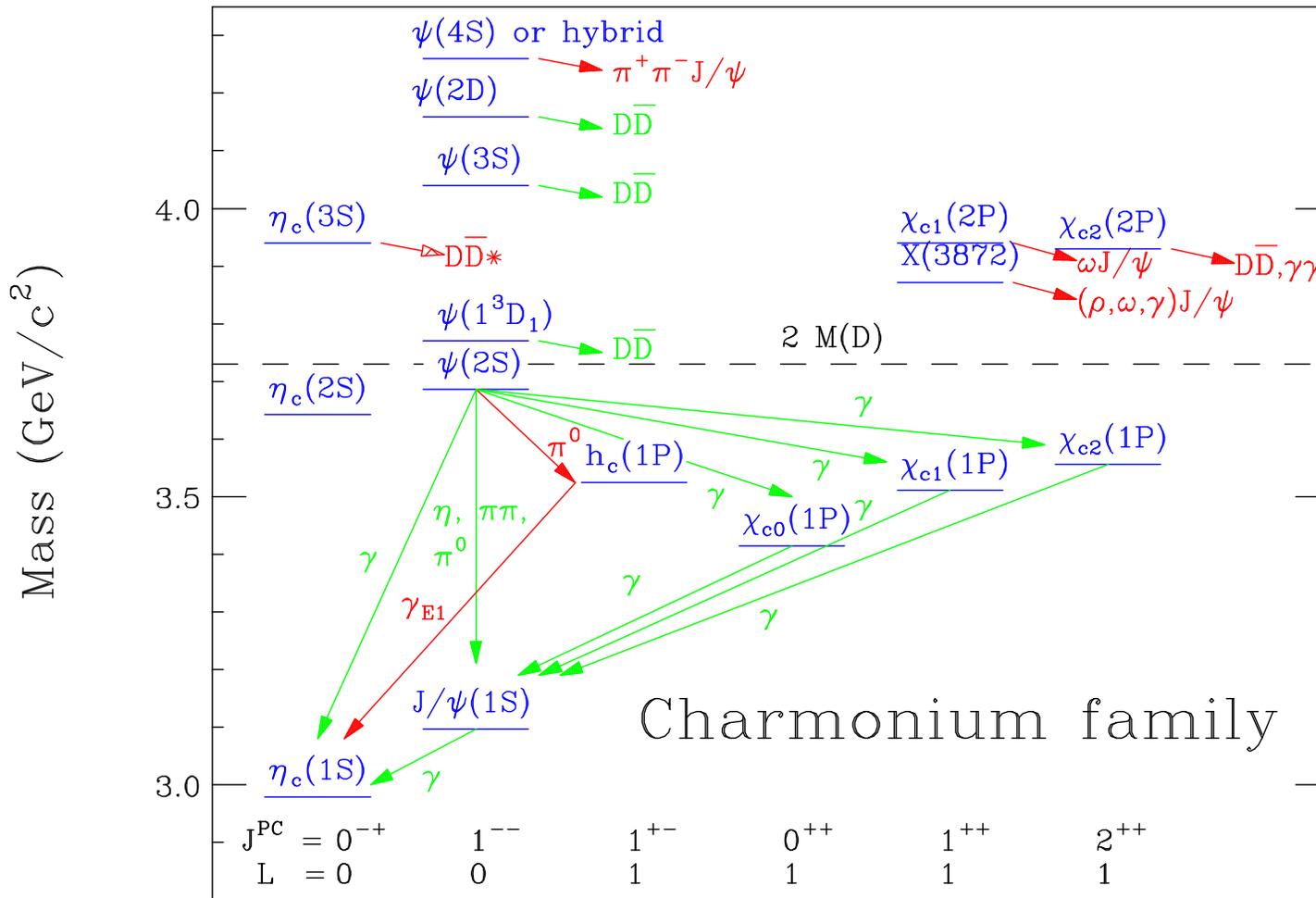
STRUCTURE IN R

15/35



Charm and $\tau^+\tau^-$ thresholds; excited $c\bar{c}$ levels

CHARMONIUM ($c\bar{c}$) SPECTRUM 10/35



I joined CLEO
in 2003 during
sabbatical

Helped find
the $h_c(3525)$

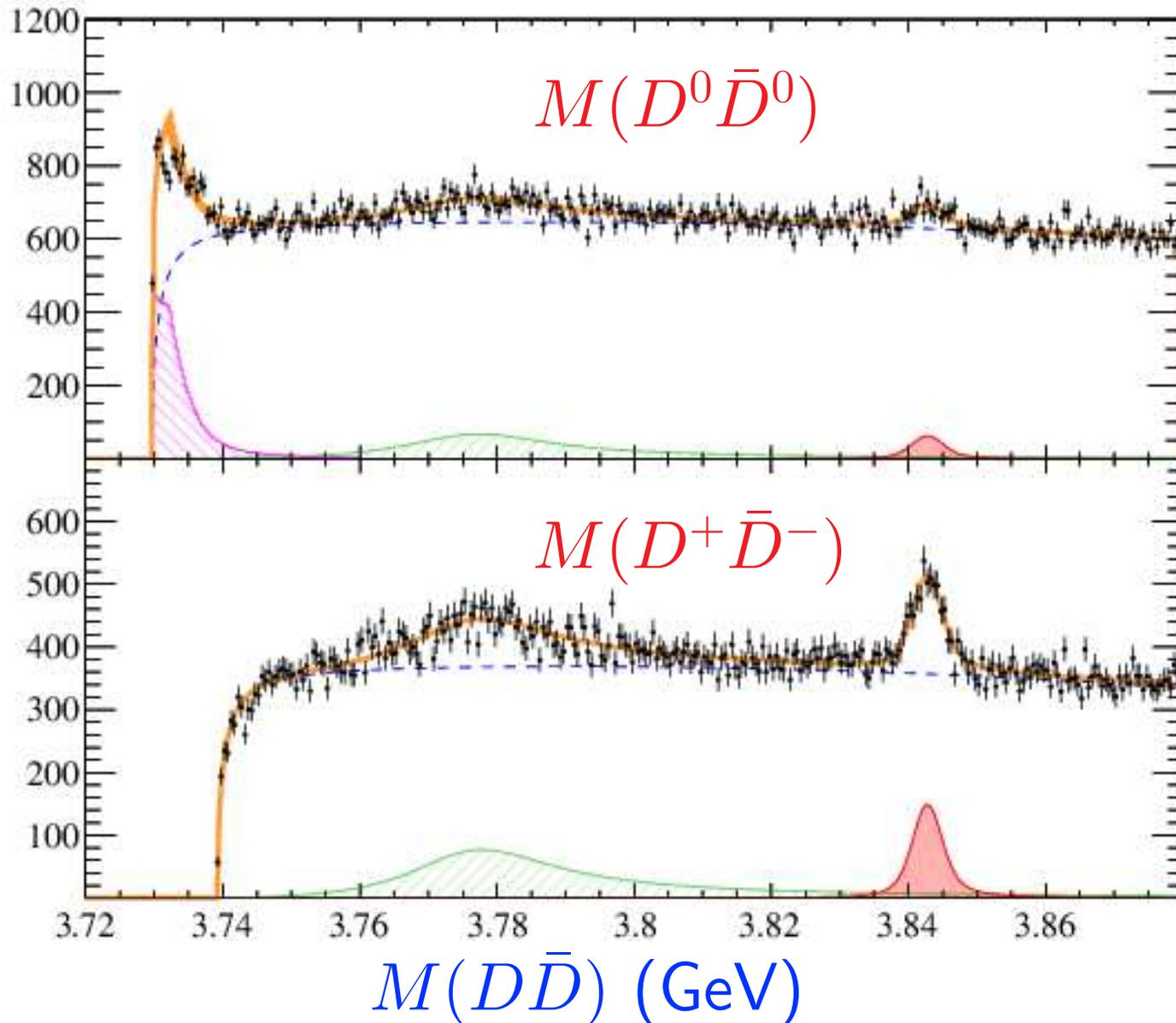
Late addition:
 3D_3 candidate
at 3843 (LHCb)
(1903.12240v2)

$c\bar{c}$ potentials (e.g., Cornell) \Rightarrow cartoon of levels

Linear potential ($\sim r$) at long distance: quark-confining

Gluon exchange (Coulomb-like, $\sim -1/r$) at short distance

3D_3 CANDIDATE AT 3843 MEV ^{17/35}



Hadronic production \Rightarrow wider range of quarkonium states

THE THIRD FAMILY

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Discovery of τ suggested a third neutrino

Quark-lepton symmetry then required a third doublet (t, b)

$$\begin{pmatrix} \nu_e \\ e^- \end{pmatrix} \begin{pmatrix} \nu_\mu \\ \mu^- \end{pmatrix} \begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix} \Leftrightarrow \begin{pmatrix} u \\ d \end{pmatrix} \begin{pmatrix} c \\ s \end{pmatrix} \begin{pmatrix} t \\ b \end{pmatrix}$$

b discovered at Fermilab in 1977: $\Upsilon(1S, 2S, [3S])$ as $b\bar{b}$

Top quark t discovered at Fermilab in 1995: $t \rightarrow Wb$

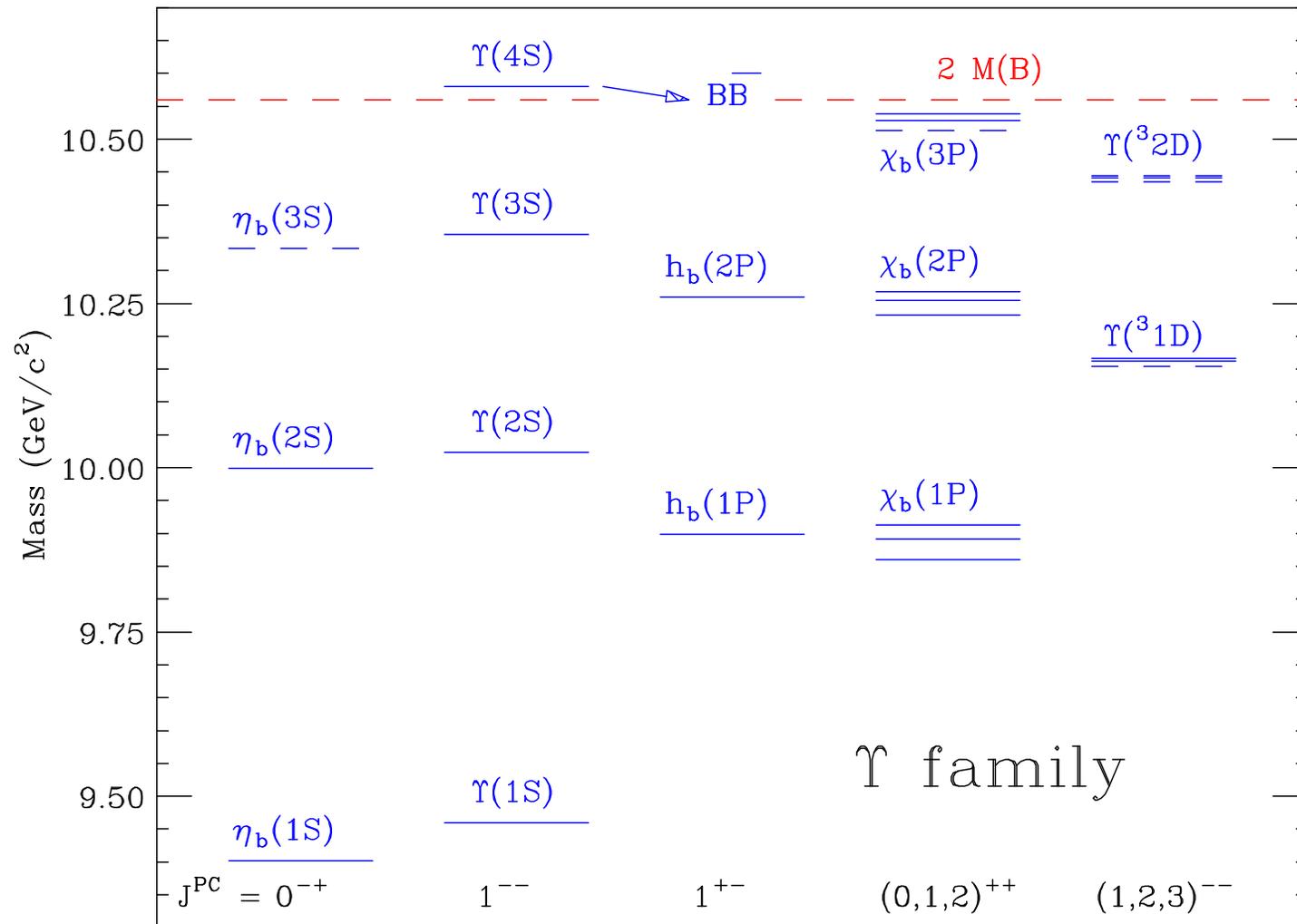
ν_τ discovered at Fermilab in 2000 (emulsion target)

Six quarks anticipated by Kobayashi-Maskawa in 1973; generalized Cabibbo theory to allow CP violation

Weak charge-changing transitions between left-handed fermions now described by $(u, c, t)^T \leftarrow V_{CKM}(d, s, b)^T$

V_{CKM} a 3×3 unitary matrix with three angles and a phase

BOTTOMONIUM ($b\bar{b}$) SPECTRUM 19/25



Level spacings almost independent of reduced mass

C. Quigg + JLR: $\log r$ interpolates between $-1/r$ and r

CARTOON OF HADRON MASSES ^{20/35}

Zeldovich-Sakharov, DGG: Constituent quarks with masses several hundred MeV, hyperfine splittings

M. Karliner + JLR: Additional binding of heavy quarks

Describe ground-state baryons containing u, d, s taking $m_u = m_d \equiv m_q = 363$ MeV, $m_s = 538$ MeV, and hyperfine interaction term $a/(m_q)^2 = 50$ MeV

| State (mass in MeV) | Spin | Expression for mass | Predicted mass (MeV) |
|---------------------|------|---------------------------------------|----------------------|
| $N(939)$ | 1/2 | $3m_q - 3a/(m_q)^2$ | 939 |
| $\Delta(1232)$ | 3/2 | $3m_q + 3a/(m_q)^2$ | 1239 |
| $\Lambda(1116)$ | 1/2 | $2m_q + m_s - 3a/(m_q)^2$ | 1114 |
| $\Sigma(1193)$ | 1/2 | $2m_q + m_s + a/(m_q)^2 - 4a/m_q m_s$ | 1179 |
| $\Sigma(1385)$ | 3/2 | $2m_q + m_s + a/(m_q)^2 + 2a/m_q m_s$ | 1381 |
| $\Xi(1318)$ | 1/2 | $2m_s + m_q + a/(m_s)^2 - 4a/m_q m_s$ | 1327 |
| $\Xi(1530)$ | 3/2 | $2m_s + m_q + a/(m_s)^2 + 2a/m_q m_s$ | 1529 |
| $\Omega(1672)$ | 3/2 | $3m_s + 3a/(m_s)^2$ | 1682 |

CHARMED & BOTTOM BARYONS 21/35

$$M(\Lambda_{c,b}) - M(\Lambda) \Rightarrow m_{c,b}^b = (1710.5, 5043.5) \text{ MeV}$$

Above choices of mass sufficient to describe nonstrange ground state baryons with one c or b quark

When taking account of deeper cs or bs binding in baryons with one or two strange quarks and one charm or bottom fit all baryons with one c or b (with M. Karliner)

| Charmed baryons | | | Bottom baryons | | |
|-----------------------------------|------|------------------------|-----------------------------------|------|------------------------|
| State (M in MeV) | Spin | Predicted M (MeV) | State (M in MeV) | Spin | Predicted M (MeV) |
| $\Lambda_c(2286.5) = cud$ | 1/2 | Input | $\Lambda_b(5619.5) = bud$ | 1/2 | Input |
| $\Sigma_c(2453.4) = cuu, \dots$ | 1/2 | 2444.0 | $\Sigma_b(5814.3) = buu, \dots$ | 1/2 | 5805.1 |
| $\Sigma_c^*(2518.1) = cuu, \dots$ | 3/2 | 2507.7 | $\Sigma_b^*(5833.8) = buu, \dots$ | 3/2 | 5826.7 |
| $\Xi_c(2469.3) = csu, csd$ | 1/2 | 2475.3 | $\Xi_b(5792.7) = bsu, bsd$ | 1/2 | 5801.5 |
| $\Xi_c'(2575.8) = csu, csd$ | 1/2 | 2565.4 | $\Xi_b'(-) = bsu, bsd$ | 1/2 | 5921.3 |
| $\Xi_c^*(2645.9) = csu, csd$ | 3/2 | 2628.6 | $\Xi_b^*(5949.7) = bsu, bsd$ | 3/2 | 5944.1 |
| $\Omega_c(2695.2) = css$ | 1/2 | 2692.1 | $\Omega_b(6046.4) = bss$ | 1/2 | 6042.8 |
| $\Omega_c^*(2765.9) = css$ | 3/2 | 2762.8 | $\Omega_b^*(-) = bss$ | 3/2 | 6066.7 |

Predict masses of baryons with more than one heavy quark?

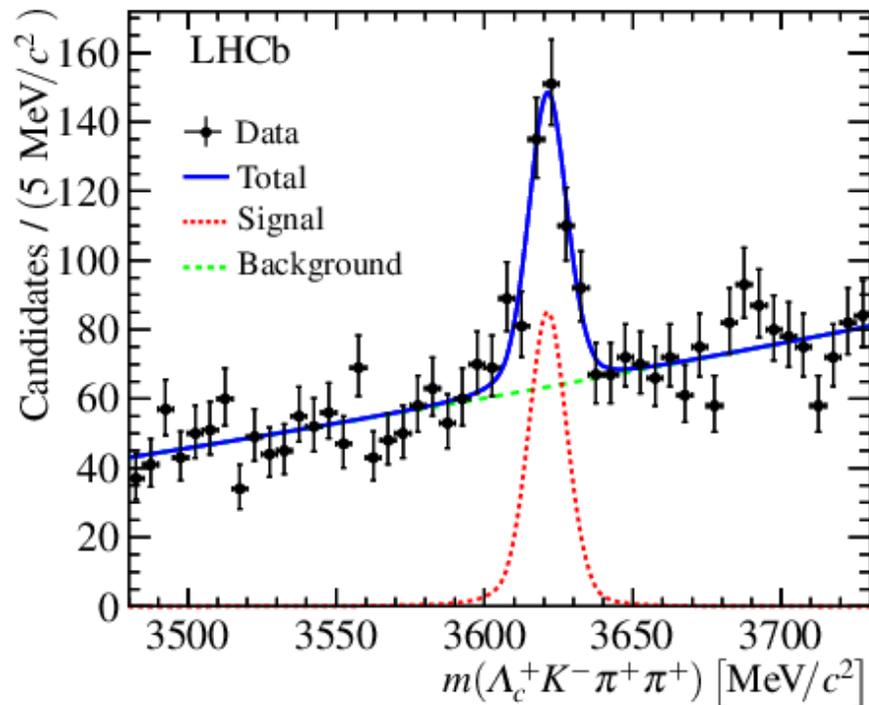
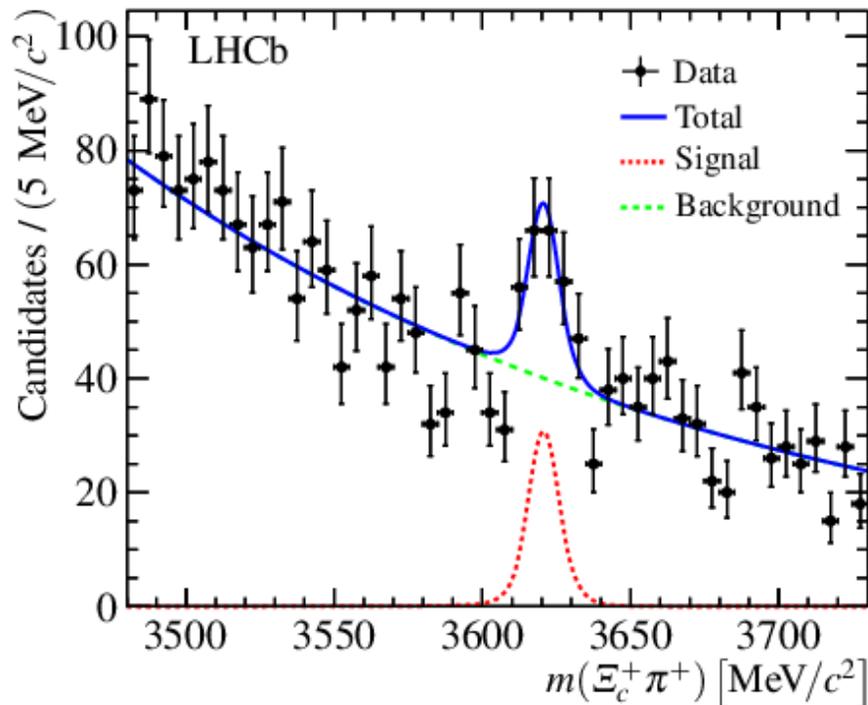
BARYONS WITH > 1 HEAVY QUARK

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SELEX at Fermilab (2002-5) claimed $\Xi_{cc}^{++}(3520) = ccu$ and $\Xi_{cc}^+(3460) = ccd$; not confirmed by others

M. Karliner + JLR (PR D **90**): Constituent-quark masses, hyperfine splittings, estimates of QQ' binding ($q = u, d$):
 $M(\Xi_{cc}, \Xi_{cc}^*) = (3627 \pm 12, 3690 \pm 12)$ MeV

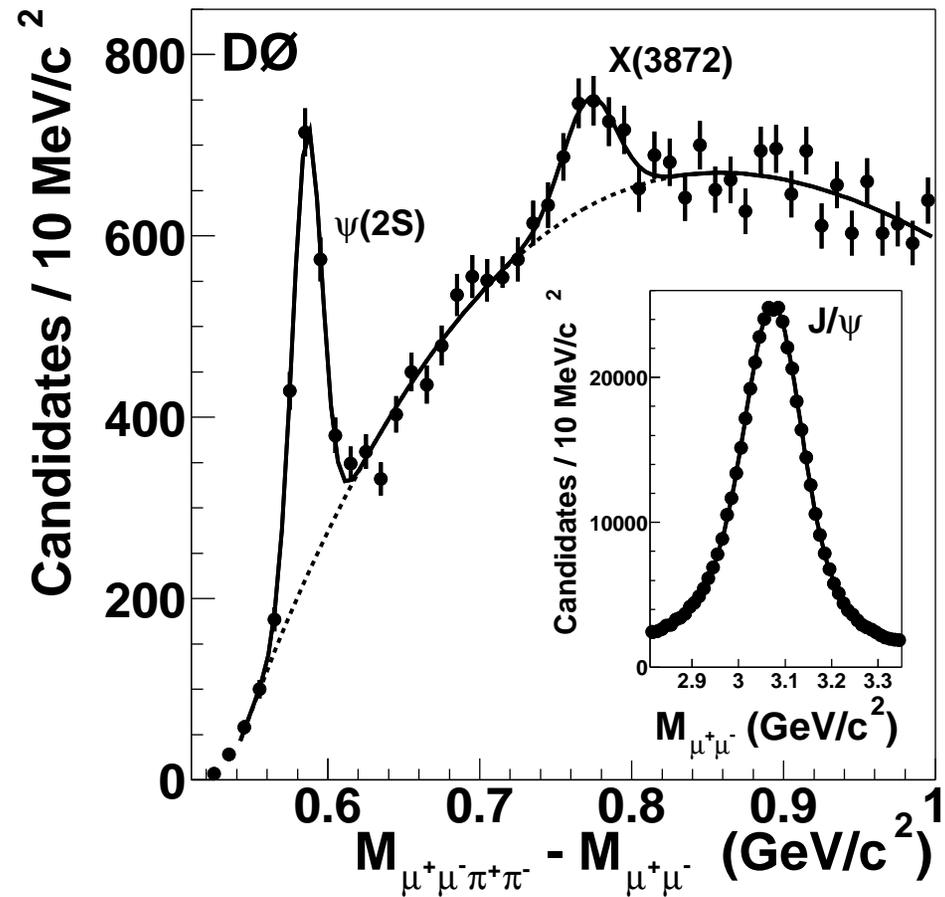
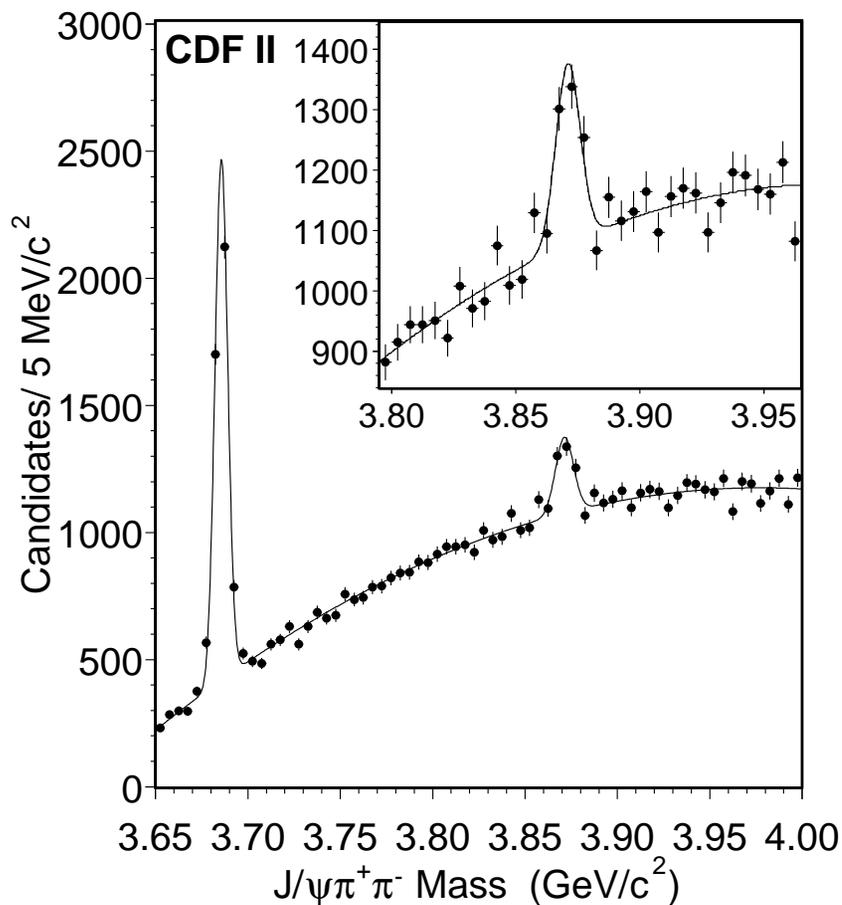
LHCb PRL **119**, 112001: $M(\Xi_{cc}^{++}) = 3621.40 \pm 0.78$ MeV



$X(3872)$: $c\bar{c}u\bar{u}$ STATE

23/35

State decaying to $J/\psi\pi^+\pi^-$ discovered by Belle (2003) at 3872 MeV (shown with $\psi'(3686$ MeV)); also seen by CDF (2004, left), D0 (2004, right), and BaBar (2005)



Within ~ 0.2 MeV of $D^0\bar{D}^{*0}$ threshold (molecule)

$X(3872)$ PROPERTIES

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$M(X) = (3871.69 \pm 0.17) \text{ MeV} \simeq M(D^0) + M(\bar{D}^{*0}) = (3871.68 \pm 0.07) \text{ MeV} \Rightarrow$ key role for that channel

Decay $X \rightarrow \gamma J/\psi$ seen; implies $C(X) = +$ and some admixture of $c\bar{c}$ in its wave function

Angular distribution of decay products implies $J^{PC} = 1^{++}$ as expected for S-wave state of $D^0\bar{D}^{*0} + \text{c.c.}$

C invariance $\Rightarrow C(\pi^+\pi^-) = - \Rightarrow \pi^+\pi^-$ in a ρ meson

Large $M(D^{(*)+} - D^{(*)0}) \Rightarrow$ little $D^{(*)\pm}$ in wave function

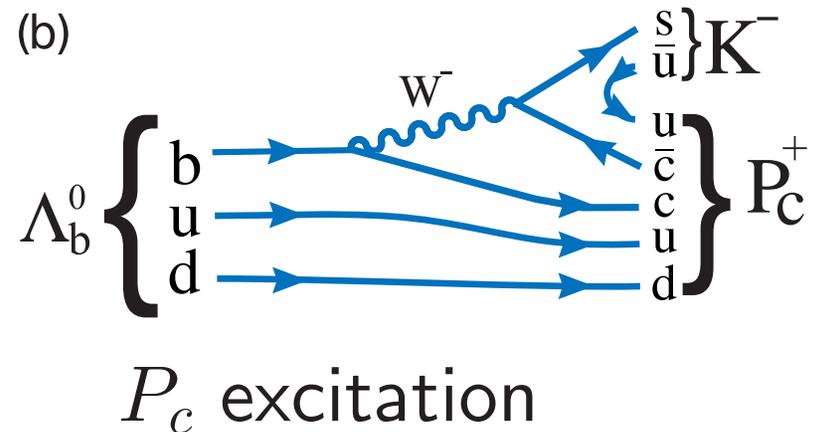
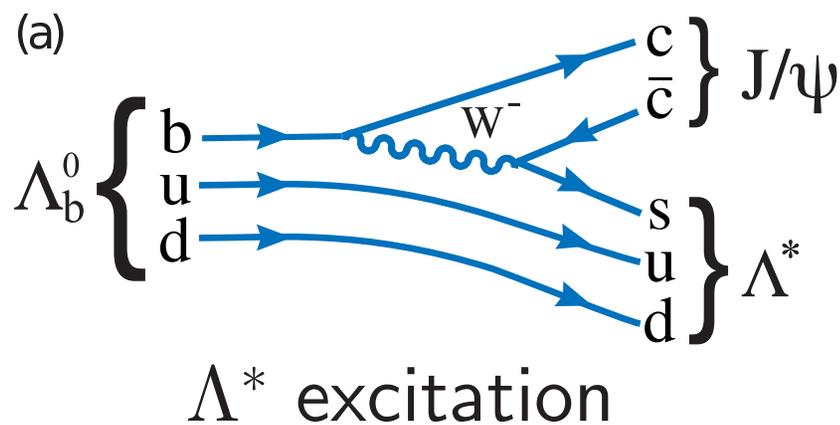
$\Gamma(X \rightarrow \omega J/\psi)$ comparable to $\Gamma(X \rightarrow J/\psi\rho)$, as one would expect for a state with $c\bar{c}u\bar{u}$ admixture

In addition to $X(3872)$ (mixture of 2^3P_1 $c\bar{c}$ state and $J^{PC} = 1^{++}$ $c\bar{c}u\bar{u}$ state) one expects an orthogonal mixture (potential models: probably $> 3900 \text{ MeV}$)

PENTAQUARKS P_c

25/35

LHCb (PRL **115**, 072001) saw bumps in J/ψ p invariant mass in the decay $\Lambda_b \rightarrow K^- J/\psi p$ at 4380 and 4450 MeV



Dalitz plot: many $I = 0$ $K^- p$ states ($\Delta I[b \rightarrow c\bar{c}s] = 0$).

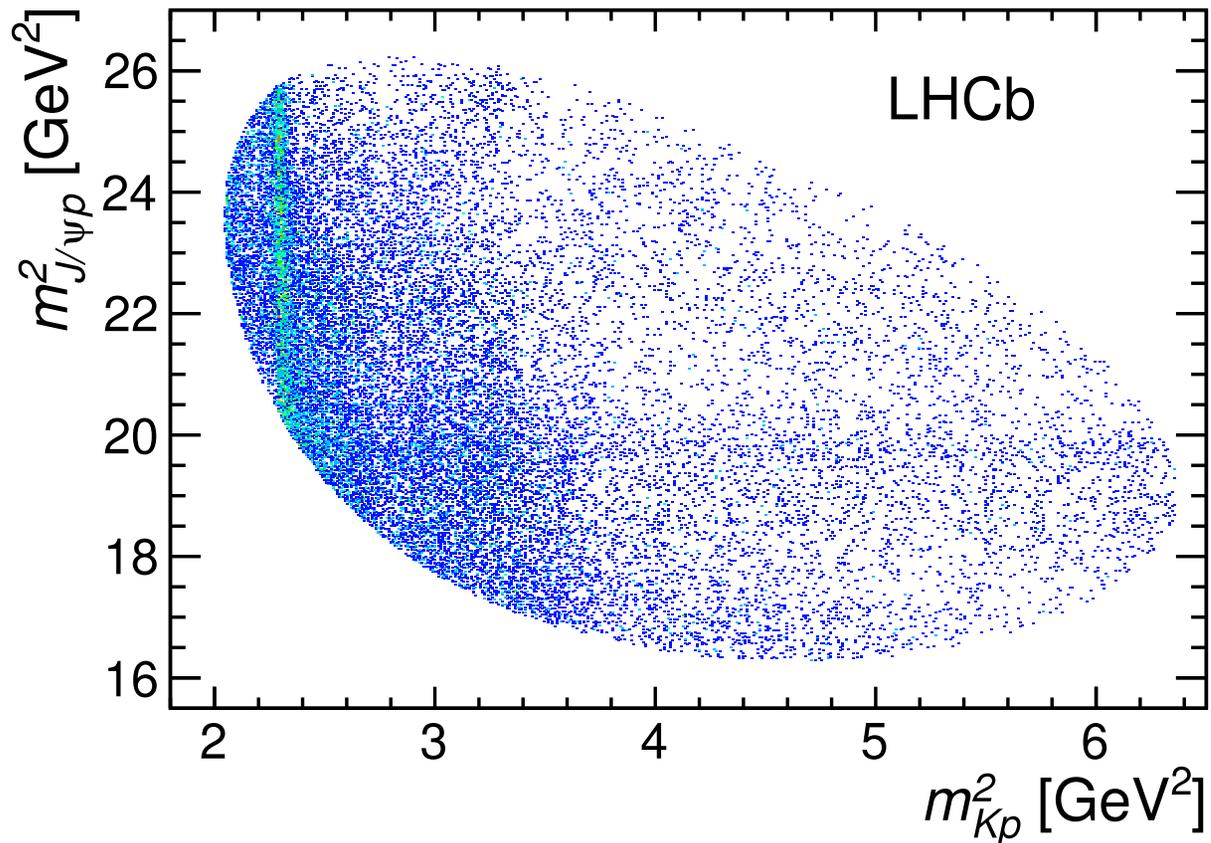
Updated result: *Three* narrow J/ψ resonances: 4311.9, 4440.3, 4457.3 MeV; widths 9.8, 20.6, 6.4 MeV

Near $\Sigma_c \bar{D}$ and $\Sigma_c \bar{D}^*$ thresholds; binding mechanism?

One-pion exchange can't couple to $D\bar{D}$; $\pi^+\pi^-$ may favor $\Sigma_c \bar{D}$ over $D\bar{D}$: lowest intermediate state is $\Lambda_c \bar{D}^*$ vs. $D^* \bar{D}^*$

$K^- p J/\psi$ DALITZ PLOT

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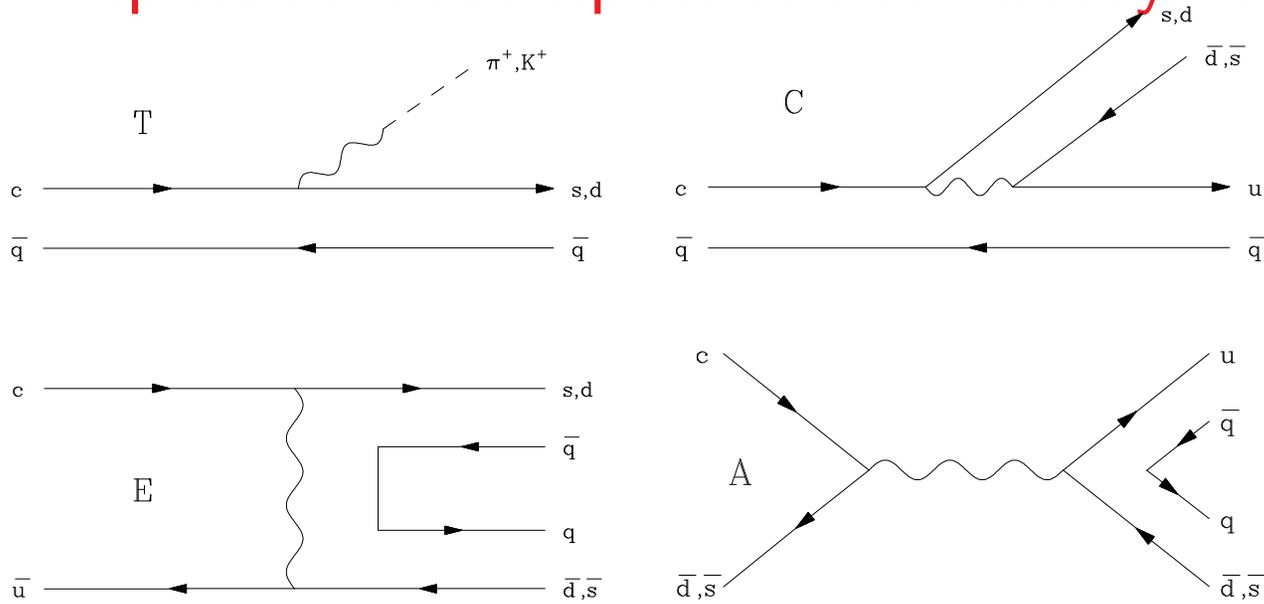
Asymmetric behavior along $M(J/\psi p)$ bands indicates interference with opposite-parity amplitude(s)

arXiv:1903.11503: Possible isospin impurity of $P_c(4457)$

arXiv:1903.11560: Total of 7 $\Sigma_c^{(*)} \bar{D}^{(*)}$ molecules

CHARMED MESON DECAYS 27/35

Amplitude decomposition for decays to two pseudoscalars



With M. Gronau,
B. Bhattacharya

$$T = 2.927,$$

$$C = 2.337 e^{-i 151.66^\circ},$$

$$E = 1.573 e^{i 120.56^\circ},$$

$$A = 0.33 e^{i 70.47^\circ}$$

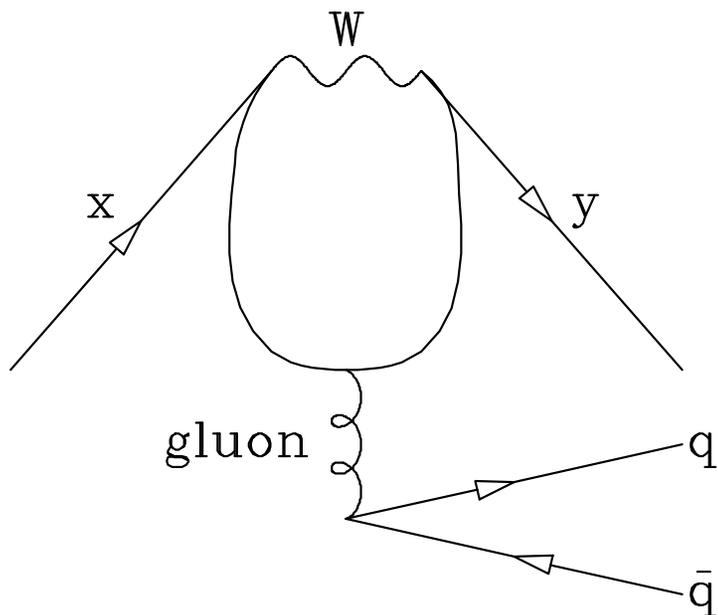
| Meson | Decay mode | \mathcal{B} [PDG] (%) | p^* (MeV) | $ \mathcal{A} $ ($10^{-6} GeV$) | Rep. | Predicted \mathcal{B} (%) |
|---------|-------------------|-------------------------|-------------|-----------------------------------|----------------------|-----------------------------|
| D^0 | $K^- \pi^+$ | 3.891 ± 0.077 | 861.1 | 2.52 ± 0.03 | $T + E$ | 3.905 |
| | $\bar{K}^0 \pi^0$ | 2.380 ± 0.092 | 860.4 | 1.97 ± 0.04 | $(C - E)/\sqrt{2}$ | 2.347 |
| | $\bar{K}^0 \eta$ | 0.962 ± 0.060 | 771.9 | 1.32 ± 0.04 | $C/\sqrt{3}$ | 1.002 |
| | $\bar{K}^0 \eta'$ | 1.900 ± 0.108 | 564.9 | 2.17 ± 0.06 | $-(C + 3E)/\sqrt{6}$ | 1.920 |
| D^+ | $\bar{K}^0 \pi^+$ | 3.074 ± 0.097 | 862.4 | 1.41 ± 0.02 | $C + T$ | 3.090 |
| D_s^+ | $\bar{K}^0 K^+$ | 2.98 ± 0.17 | 850.3 | 2.12 ± 0.06 | $C + A$ | 2.939 |
| | $\pi^+ \eta$ | 1.84 ± 0.15 | 902.3 | 1.62 ± 0.07 | $(T - 2A)/\sqrt{3}$ | 1.810 |
| | $\pi^+ \eta'$ | 3.95 ± 0.34 | 743.2 | 2.61 ± 0.11 | $2(T + A)/\sqrt{6}$ | 3.603 |

CP VIOLATION IN $\Delta S = 0$ CHARM DECAYS ^{28/35}

Leading-order CKM matrix \Rightarrow more CP violation in b decays than in charm decays (phase mainly for 3rd family):

$$V = \begin{bmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{bmatrix} \approx \begin{bmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{bmatrix}$$

LHCb PRL **122**, 211803 (2019)] has seen CP violation in difference between $D^0 \rightarrow K^- K^+$ and $D^0 \rightarrow \pi^- \pi^+$ CP asymmetries: $\Delta A_{CP} = (-15.4 \pm 2.9) \times 10^{-4}$



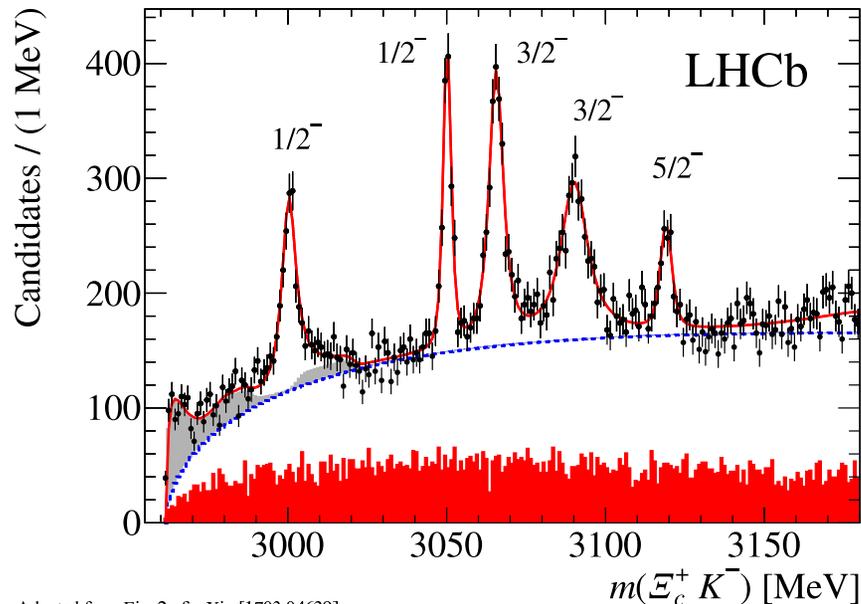
Most likely cause: small $c \rightarrow u$ penguin with b quark in the loop utilizing nonzero phase of V_{ub}
 1903.10490: No; 1903.10(638,952): Yes

If so, expect CP asymmetries in other $\Delta S = 0$ charm decays to two pseudoscalar mesons

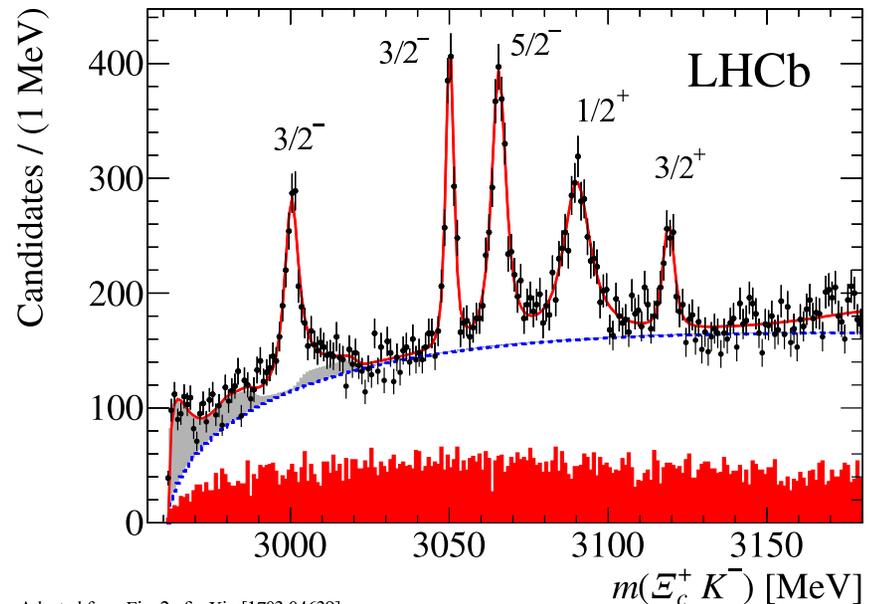
EXCITED $\Omega_c = c s s$ STATES

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LHCb, PRL **118**, 182001 (2017):



Adapted from Fig. 2 of arXiv:[1703.04639]



Adapted from Fig. 2 of arXiv:[1703.04639]

Two possible spin-parity assignments. Narrowness is a consequence of absence of nonstrange (u, d) quarks

If two highest-mass states have positive parity, they are radial excitations of $1/2^+$ and $3/2^+$ ground states, and two more negative-parity states are expected at lower masses [M. Karliner + JLR, PR D **95**, 114012 (2017)]

CHARMED BARYON LIFETIMES 30/35

LHCb: PRL **121**, 092003 (2018); PRD **100**, 032001
Theory: H. Y. Cheng, Front. Phys. **10**, 101406 (2015)

| | PDG 2018 | LHCb | Theory |
|-------------------|-------------------|---------------------------------|--------|
| $\Lambda_c = cud$ | 200 ± 6 | $203.5 \pm 1.0 \pm 1.3 \pm 1.4$ | 264 |
| $\Xi_c^+ = csu$ | 442 ± 26 | $456.8 \pm 3.5 \pm 2.9 \pm 3.1$ | 368 |
| $\Xi_c^0 = csd$ | 112_{-10}^{+13} | $154.5 \pm 1.7 \pm 1.6 \pm 1.0$ | 193 |
| $\Omega_c = css$ | 69 ± 12 | $268 \pm 24 \pm 10 \pm 2$ | 171 |

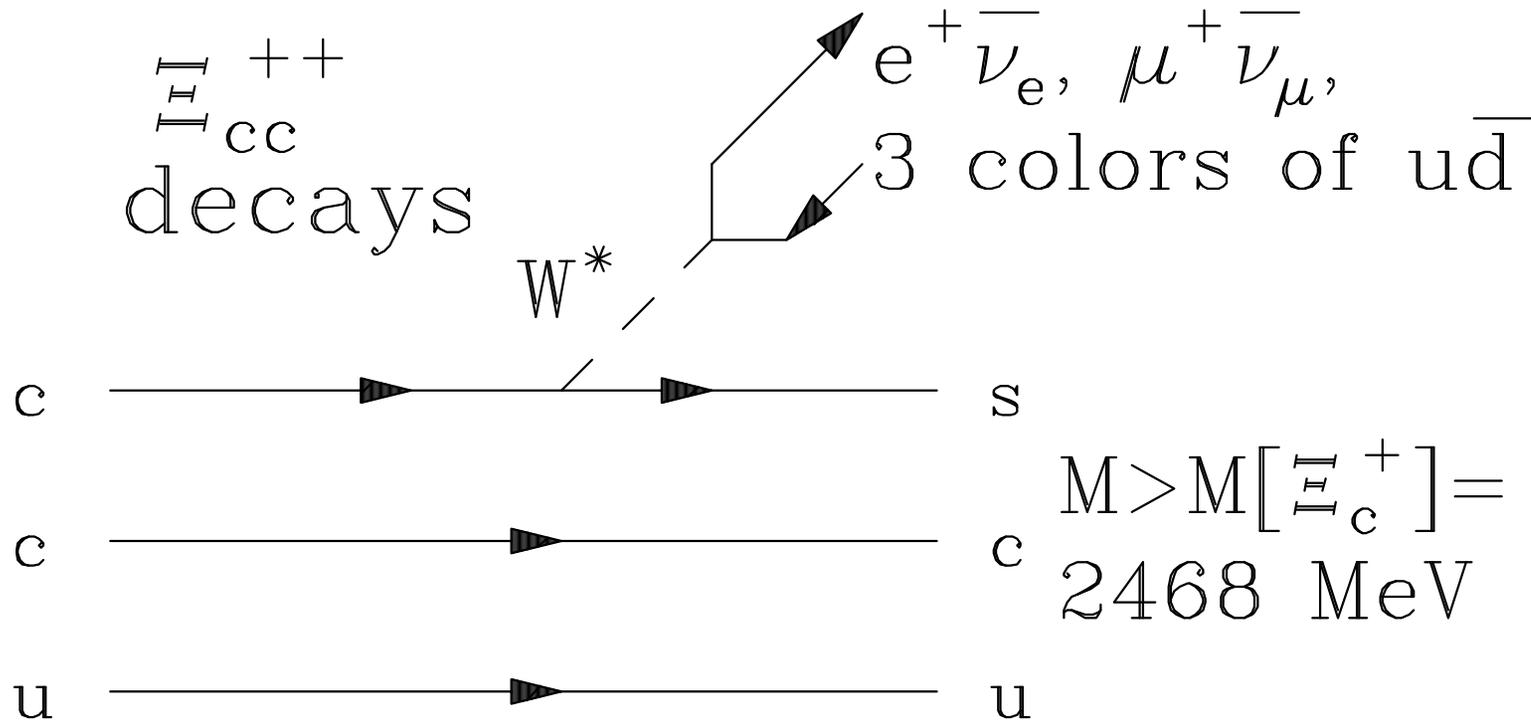
Λ_c and Ξ_c^0 lifetimes short: subprocess $cd \rightarrow su$

Operator product expansion overestimated the enhancement of Ω_c decay rate due to constructive interference of two strange quarks

Still some open questions for theorists!

Ξ_{cc}^{++} LIFETIME

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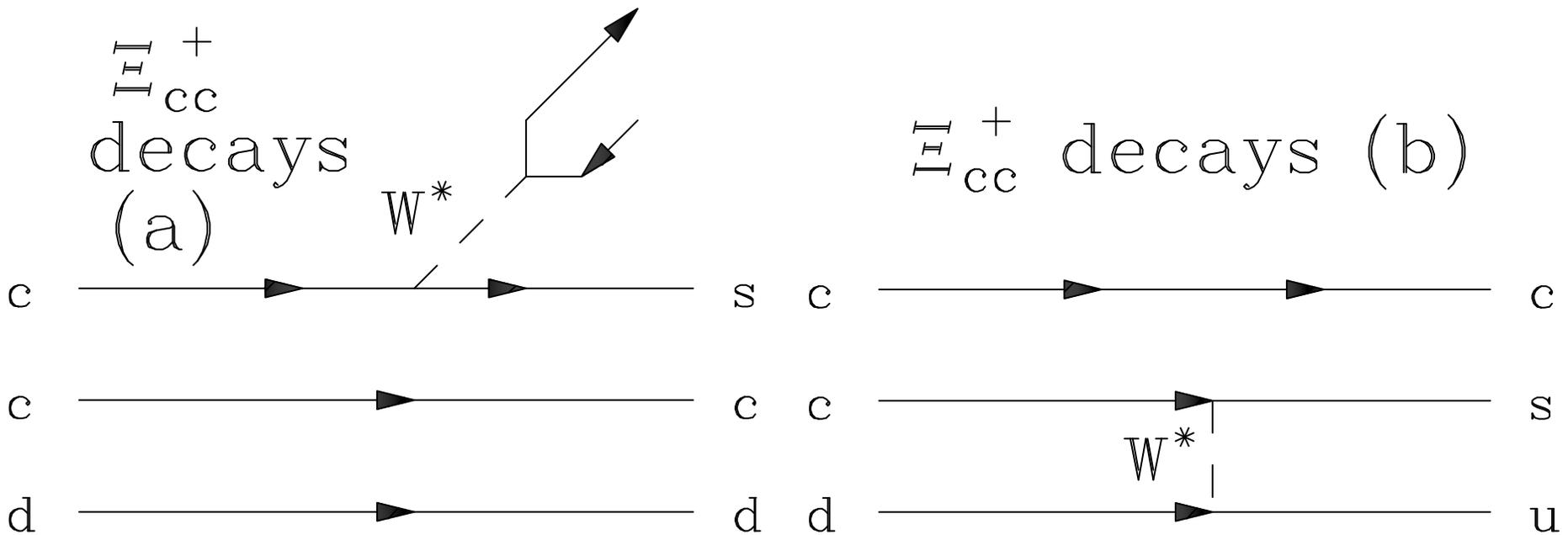
Dominant process: $c \rightarrow s [W^* \rightarrow e^+ \bar{\nu}_e, \mu^+ \bar{\nu}_\mu, 3 \text{ colors of } u\bar{d}]$

Emulate kinematic suppression with $x_{cc} \equiv [M(\Xi_c)/M(\Xi_{cc})]^2$:

$$F(x) = 1 - 8x + 8x^3 - x^4 + 12 \ln(1/x)$$

$$\Gamma(\Xi_{cc}^{++}) = \frac{10G_F^2 M(\Xi_{cc}^{++})^5}{192\pi^3} F(x_{cc}) \Rightarrow \tau(\Xi_{cc}^{++}) = 188 \text{ fs}$$

Ξ_{cc}^+ LIFETIME



(a) "Spectator"

(b) "Exchange"

$$\Gamma_a = \hbar/\tau(\Xi_{cc}^{++}) = \hbar/(256 \text{ fs}) = 2.57 \times 10^{-12} \text{ GeV}$$

$$\Gamma_b = 2[\hbar/\tau(\Xi_c^0) - \hbar/\tau(\Xi_c^+)] = 5.64 \times 10^{-12} \text{ GeV}$$

$$\tau(\Xi_c^0) = 154.5 \pm 1.7 \pm 1.6 \pm 1.0 \text{ fs} ; \quad \text{LHCb arXiv: 1906:08350}$$

$$\tau(\Xi_c^+) = 458.8 \pm 3.6 \pm 2.9 \pm 3.1 \text{ fs}$$

$$\Gamma_a + \Gamma_b = 8.21 \times 10^{-12} \text{ GeV} \Rightarrow \tau(\Xi_{cc}^+) = 80 \text{ fs}$$

STABLE $bb\bar{u}\bar{d}$ TETRAQUARK 33/35

We looked at $QQ'\bar{u}\bar{d}$ systems ($Q = c$ or b) (PRL **119**)

We found $cc\bar{u}\bar{d}$ unbound; it could decay to DD^* or $DD\gamma$

Lowest-lying $bc\bar{u}\bar{d}$ state was near $BD\gamma$ threshold and we could not tell for sure whether it was bound or unbound

Predicted $M(bb\bar{u}\bar{d}) = 10,389 \pm 12$ MeV, 215 MeV below B^-B^{*0} threshold and 170 MeV below $B^-B^0\gamma$ threshold

Regard bb as a color- 3^* diquark (transforming under QCD as an antiquark); fermi statistics require its spin to be 1

Lightest $\bar{q}\bar{q}'$ state ($q, q' = u, d$) is a color-3 $\bar{u}\bar{d}$ state with isospin zero; fermi statistics require its spin to be zero

Mass prediction then relies on accounting for constituent-quark masses, hyperfine interactions, and binding effects

TETRAQUARKS $QQ'\bar{u}\bar{d}$

Contributions (MeV) to mass of lightest tetraquark:

| $cc\bar{u}\bar{d}, J^P = 1^+$ | | $bc\bar{u}\bar{d}, J^P = 0^+$ | | $bb\bar{u}\bar{d}, J^P = 1^+$ | |
|-------------------------------|---------------|-------------------------------|---------------|-------------------------------|----------------|
| Contribution | Value | Contribution | Value | Contribution | Value |
| $2m_c^b$ | 3421.0 | $m_b + m_c$ | 6754.0 | $2m_b^b$ | 10087.0 |
| $2m_q^b$ | 726.0 | $2m_q^b$ | 726.0 | $2m_q^b$ | 726.0 |
| cc hyperfine | 14.2 | bc hyperfine | -25.5 | bb hyperfine | 7.8 |
| $-3a/(m_q^b)^2$ | -150.0 | $-3a/(m_q^b)^2$ | -150.0 | $-3a/(m_q^b)^2$ | -150.0 |
| cc binding | -129.0 | bc binding | -170.8 | bb binding | -281.4 |
| Total | 3882 ± 12 | Total | 7134 ± 13 | Total | 10389 ± 12 |

Spin zero allowed for the $bc\bar{u}\bar{d}$ state, taking advantage of the attractive bc hyperfine interaction

Since $M(cc\bar{u}\bar{d}) > M(D^0) + M(D^+) = 3734$ MeV, it can decay to $D^0 D^+ \gamma$ (decay to $D^0 D^+$ is forbidden)

$M(bc\bar{u}\bar{d}) < M(D^0) + \bar{M}(B^0) = 7144$ MeV?

Estimated lifetime of $bb\bar{u}\bar{d}$ state: 367 fs

THANKS

35/35

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