

Novel approach to measure quark/gluon jets at the LHC

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1. Introduction
2. Theory
3. Results
4. Conclusion

Back-up

- SM searches: often signature for a BSM signals: many quark, backgrounds: QCD gluons
 - 8-jet Gluino event: $pp \rightarrow \tilde{g}\tilde{g}$ and each \tilde{g} decays to 4 quarks:

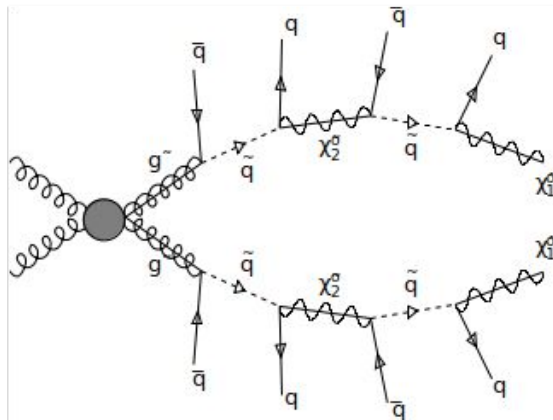
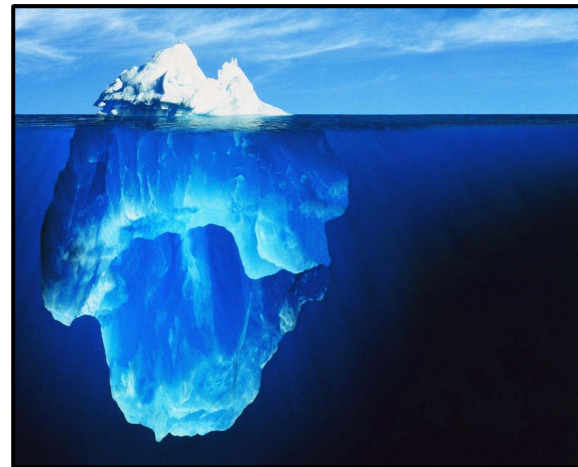
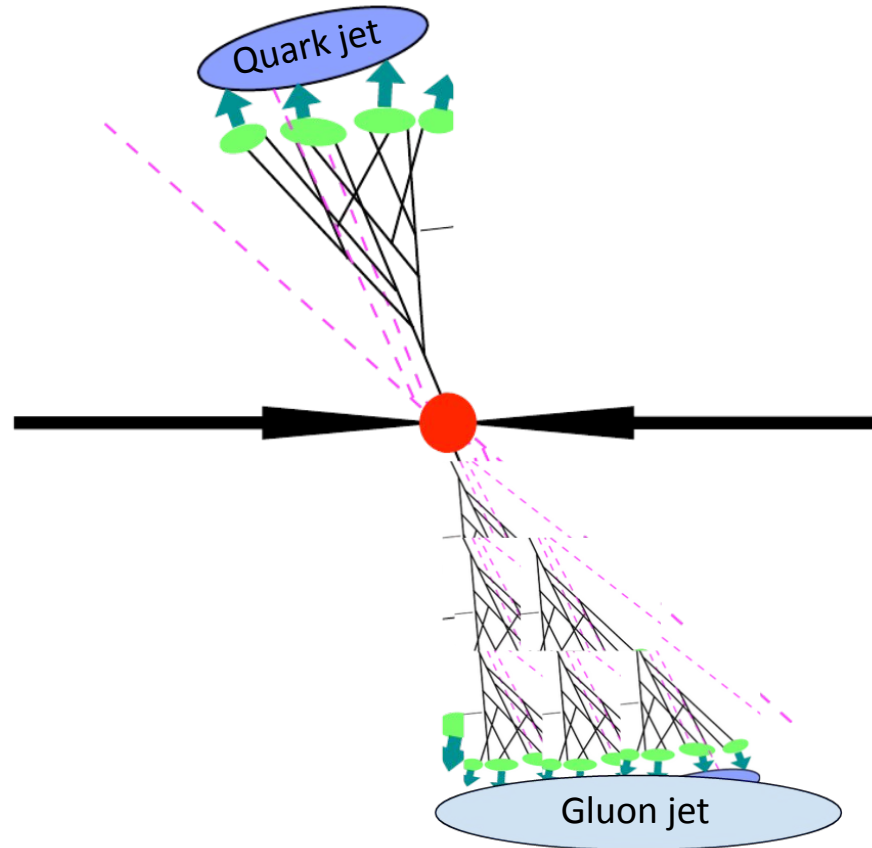


Fig. From J. Gallicchio and M. D. Schwartz, Phys. Rev. Lett.107 (2011)



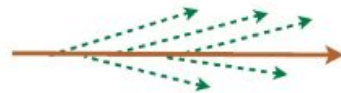
- Higgs $H^+ \rightarrow c\bar{s}$ (for charged Higgs mass between τ and t mass)
- Measure Z' coupling to hadrons (or find a leptophobic Z'/W')



2. Theory

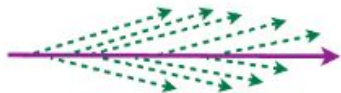
[Gras, Hoeche, Kar, Larkoski, Lönnblad, Plätzer, AS, Skands, Soyez, Thaler, JHEP 1707 (2017) 091]

Cartoon:



Quark: $C_F = 4/3$

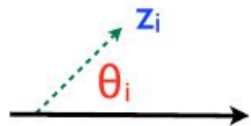
vs.



Gluon: $C_A = 3$

Probe radiation pattern with
e.g. Generalized Angularities

$$\lambda_{\beta}^{\kappa} = \sum_{i \in \text{jet}} z_i^{\kappa} \theta_i^{\beta}$$

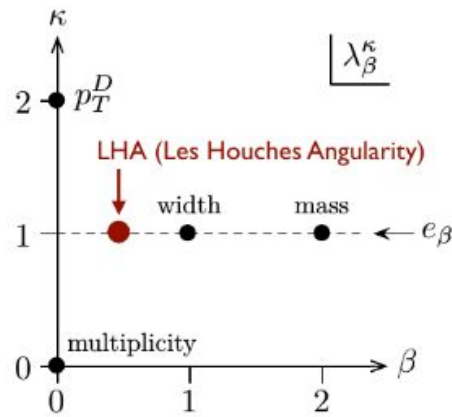


momentum
fraction

angle to
recoil-free axis

$$(\lambda_{\beta}^{\kappa})_{\text{quark}} < (\lambda_{\beta}^{\kappa})_{\text{gluon}}$$

(except p_T^D)

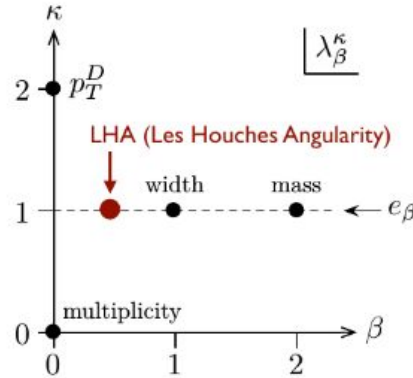


[Larkoski, Salam, Thaler, 13]
[Larkoski, Thaler, Waalewijn, 14]

Probe radiation pattern with
e.g. Generalized Angularities

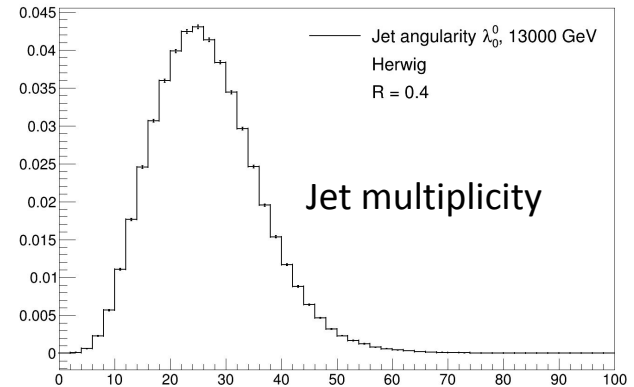
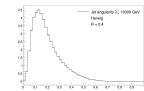
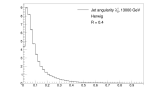
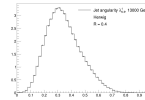
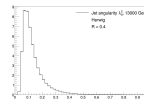
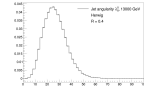
$$\lambda_{\beta}^{\kappa} = \sum_{i \in \text{jet}} z_i^{\kappa} \theta_i^{\beta}$$

z_i : momentum fraction
 θ_i : angle to recoil-free axis
 $(\lambda_{\beta}^{\kappa})_{\text{quark}} < (\lambda_{\beta}^{\kappa})_{\text{gluon}}$



[Larkoski, Salam, Thaler, 13]
[Larkoski, Thaler, Waalewijn, 14]

- λ_0^0 - multiplicity
- λ_0^2 - p_T^D
- $\lambda_{0.5}^1$ - LHA (Les Houches Angularity)
- λ_2^1 - Mass
- λ_1^1 - Width



Each angularity λ is composed of gluon λ_g and quark λ_q angularities

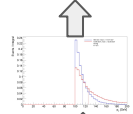
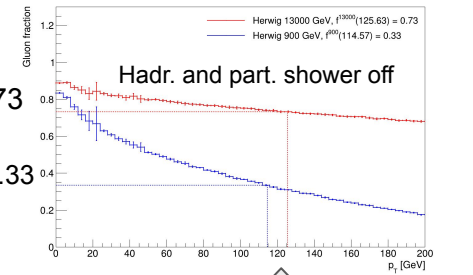
$$\lambda = f \lambda_g + (1 - f) \lambda_q$$

where

f ... gluon fraction

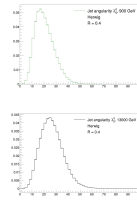
$(1 - f)$... quark fraction

taken from simulation



$p_T^{\text{cut}} = 100 \text{ GeV}$

Let's write equations for measurement at energy 900 GeV and 13 000 GeV (with the same event selection)



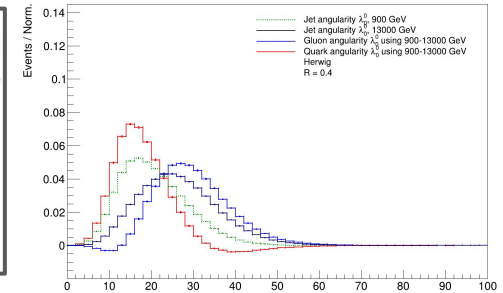
$$\lambda^{900} = f^{900} \lambda_g + (1 - f^{900}) \lambda_q$$

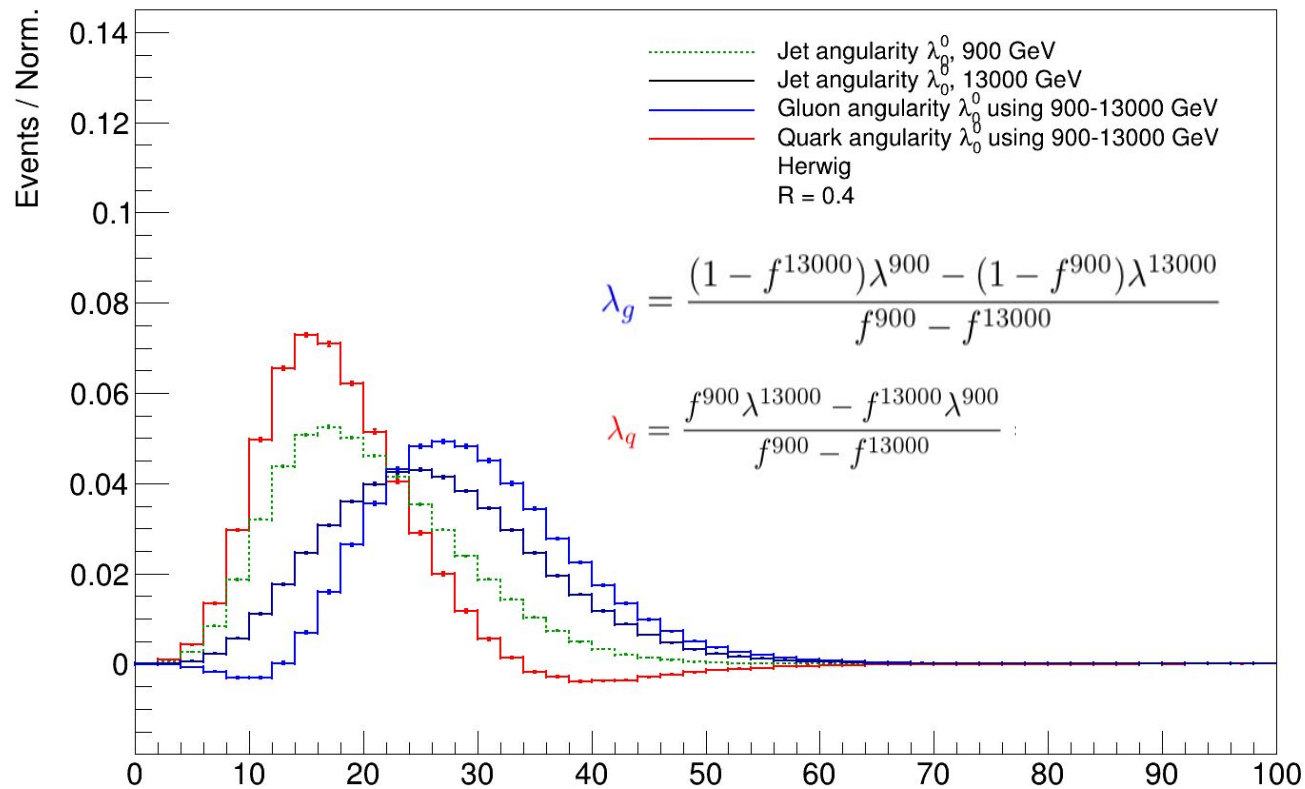
$$\lambda^{13000} = f^{13000} \lambda_g + (1 - f^{13000}) \lambda_q$$

Assuming λ_g and λ_q are energy independent.

$$\lambda_g = \frac{(1 - f^{13000}) \lambda^{900} - (1 - f^{900}) \lambda^{13000}}{f^{900} - f^{13000}}$$

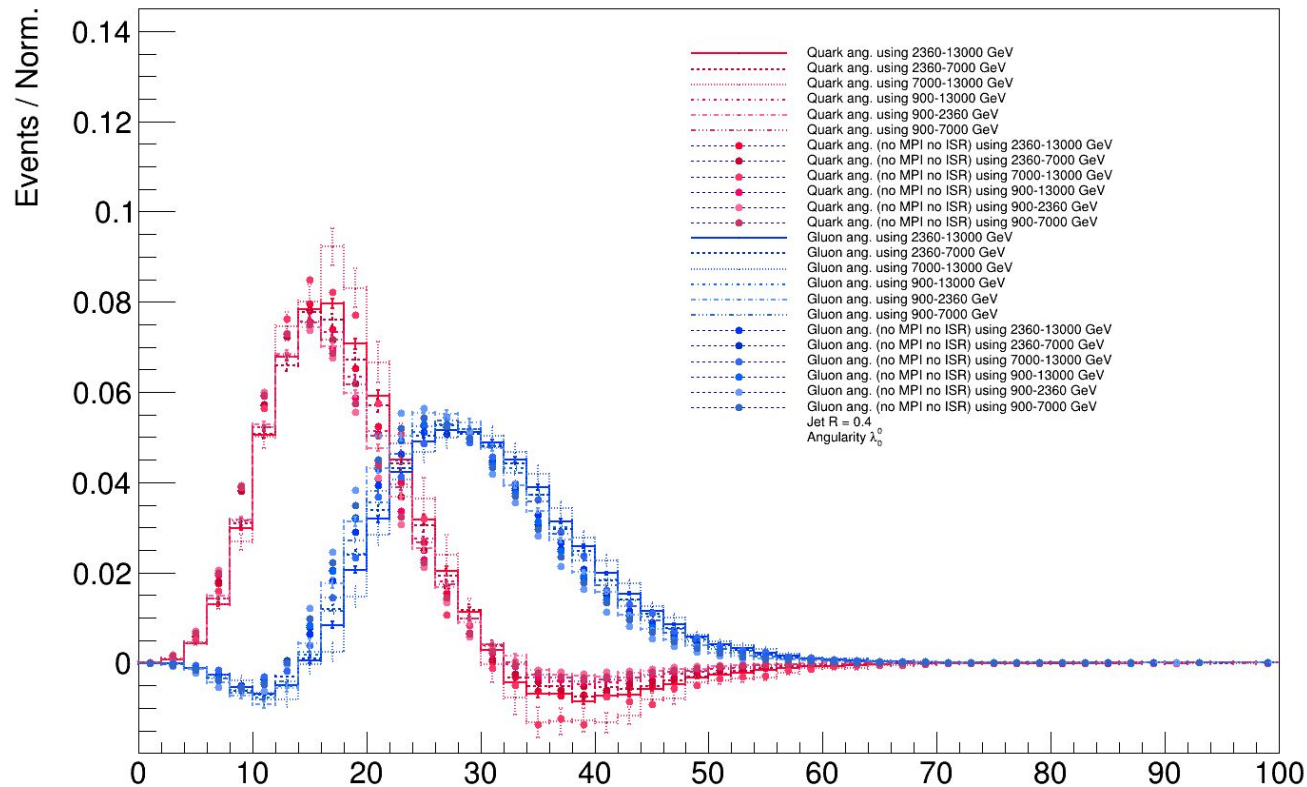
$$\lambda_q = \frac{f^{900} \lambda^{13000} - f^{13000} \lambda^{900}}{f^{900} - f^{13000}}$$





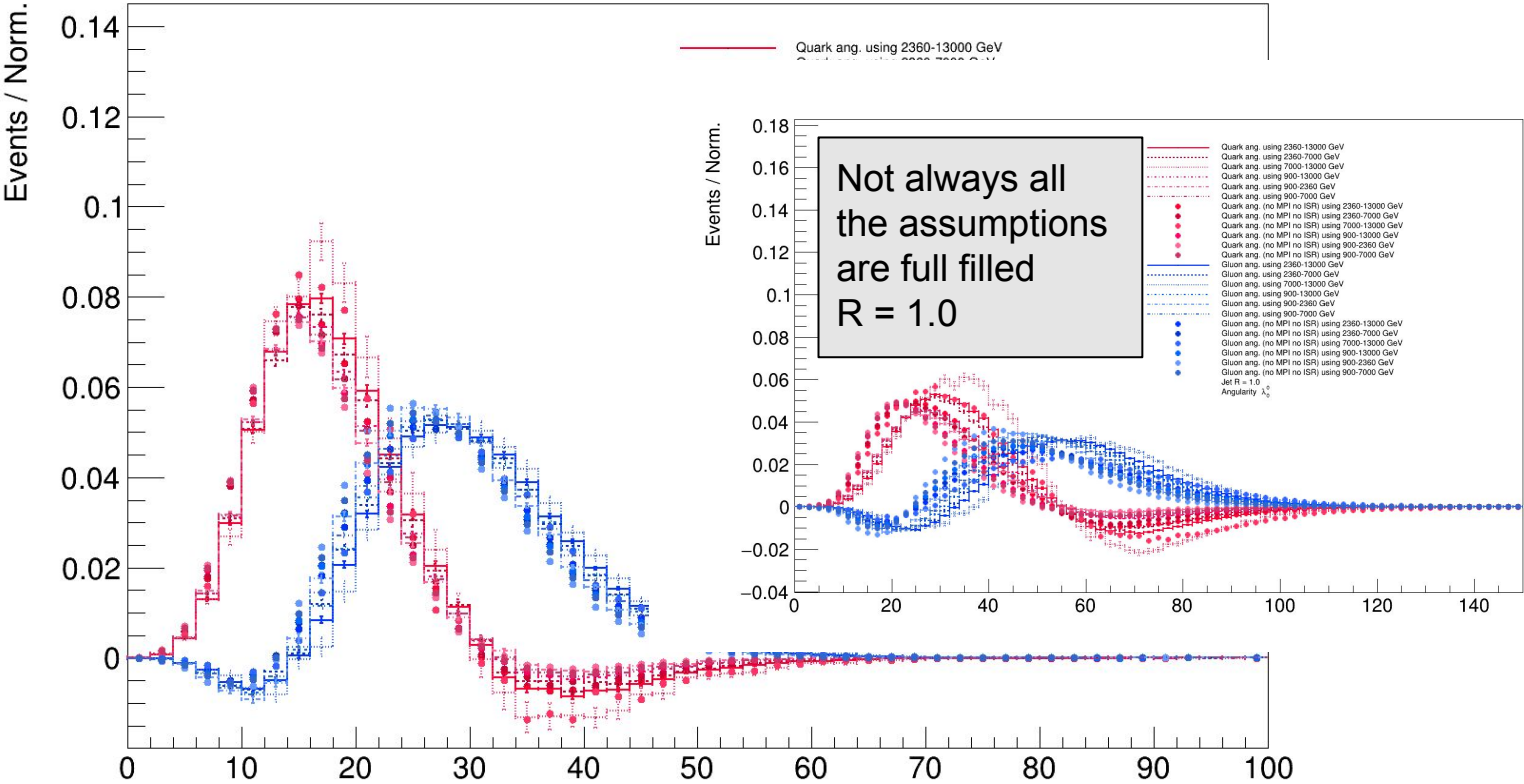
Let's use more 6 energy combinations:

900-2360, 900-7000, 900-13000, 2360-7000, 2360-13000, 7000-13000 GeV



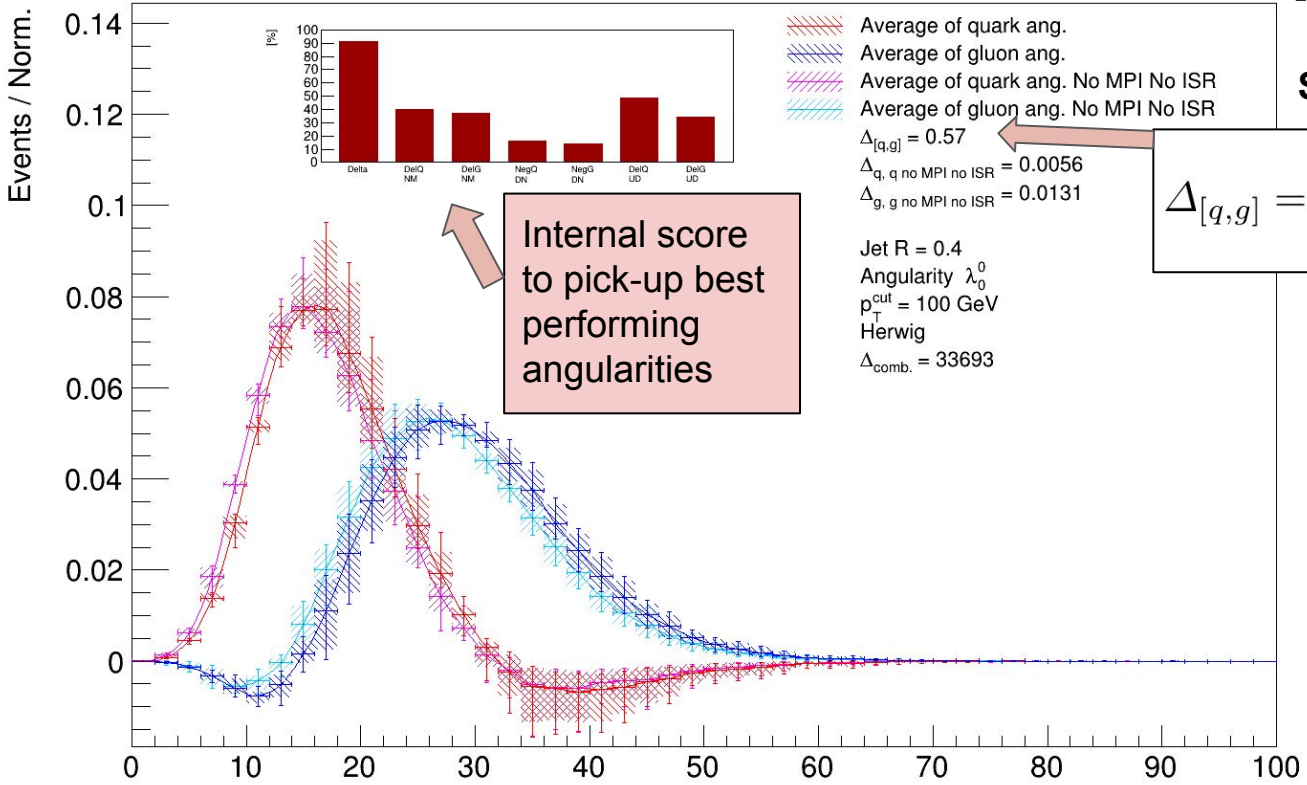
Dotted lines test the robustness to Multi Parton Interactions MPI and Initial State Radiation ISR

Let's use more 6 energy combinations:
900-2360, 900-7000, 900-13000, 2360-7000, 2360-13000, 7000-13000 GeV



Dotted lines test the robustness to Multi Parton Interactions MPI and Initial State Radiation ISR

Simplified averaged plot over 6 energy combinations: - filled area (energy comb. variation),
- ticks - stat. Unc.



Separation power:

$$\Delta_{[q,g]} = \frac{1}{2} \sum_{i=1}^N \frac{(\lambda_{q_i} - \lambda_{g_i})^2}{\lambda_{q_i} + \lambda_{g_i}}$$

$\Delta_{[q,g]} = 0.57$
 $\Delta_{q, q \text{ no MPI no ISR}} = 0.0056$
 $\Delta_{g, g \text{ no MPI no ISR}} = 0.0131$

Jet R = 0.4
Angularity λ_0^0
 $p_T^{\text{cut}} = 100 \text{ GeV}$
Herwig
 $\Delta_{\text{comb.}} = 33693$

3. Results

Selection of dijet events:

$N_{\text{jets}} = 2$
veto neutrinos

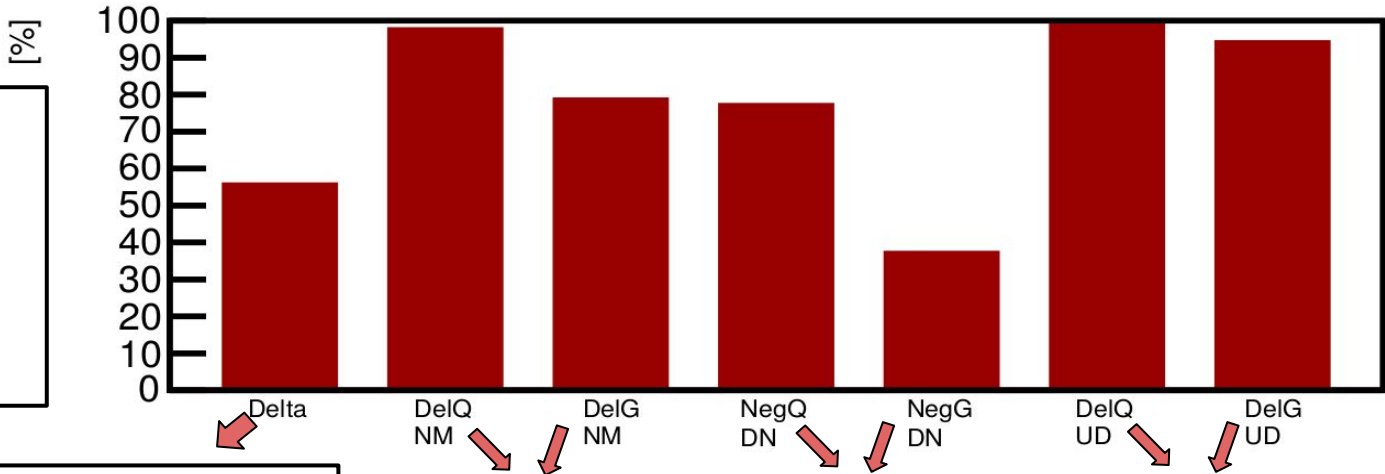
$$p_{T \text{ sublead}}/p_{T \text{ lead}} > 0.8$$

We considered all combinations of:

- 5 – angularities $\lambda_0^0, \lambda_{0.5}^1, \lambda_1^1, \lambda_0^2, \lambda_2^1$
- 2 – using groomed (MMDT) / not groomed jets
- 5 – jet radii $R = 0.2, 0.4, 0.6, 0.8, 1.0$
- 4 – regions - dijet average $p_T^{\text{cut}} = 50 \text{ GeV}, 100, 200,$
and 400 GeV $(p_{T \text{ lead}} + p_{T \text{ sublead}})/2 > p_T^{\text{cut}}$
- 2 – quark/gluon
- 2 – MPI and ISR switched on/off
- 6 – energy combinations: 900–2360, 900–7000, 900–
13000, 2360–7000, 2360–13000, 7000–13000 GeV
- 2 – event generators HERWIG and PYTHIA

Results:

Each column represents percentile of given feature of all our studied plots.



In **separation between Q/G jets** about 55 % of variations have lower performance

Separation power between noMPI variations - **robustness to noMPI no ISF**

Negativity (percentage of negative area to whole area) **negative bins**

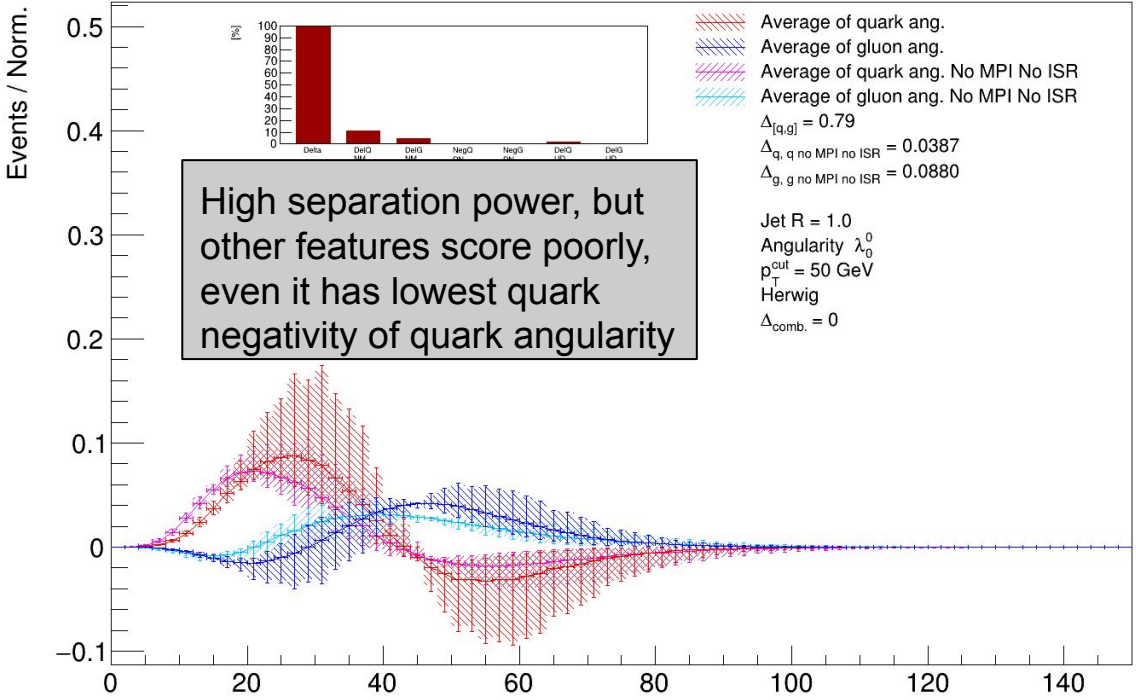
Separation power to UP and DOWN energy combination variations **robustness to different energy combination used**

Combining all columns gives us our internal score :

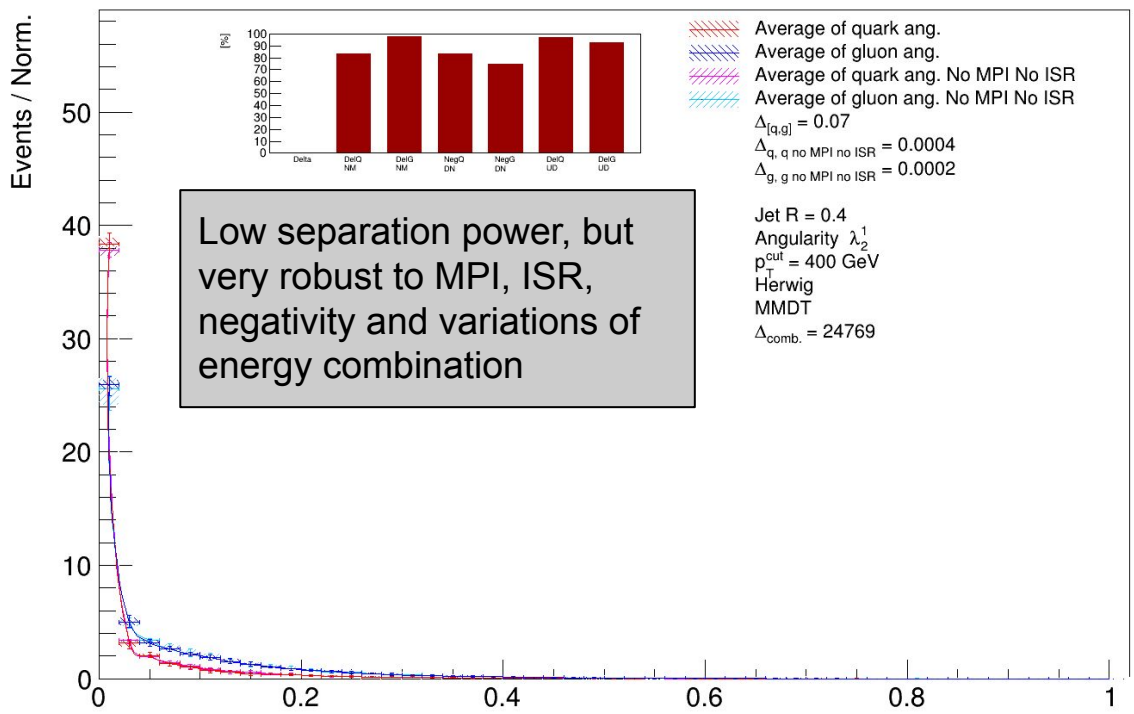
$$\Delta_{\text{comb}} = 1000 \cdot \ln \left[1 + (\text{Delta})^3 \cdot (\text{DelQ NM}) \cdot (\text{DelG NM}) \cdot (\text{NegQ DN}) \cdot (\text{NegG DN}) \cdot (\text{DelQ UD}) \cdot (\text{DelG UD}) \right]$$

min-max
0-41447

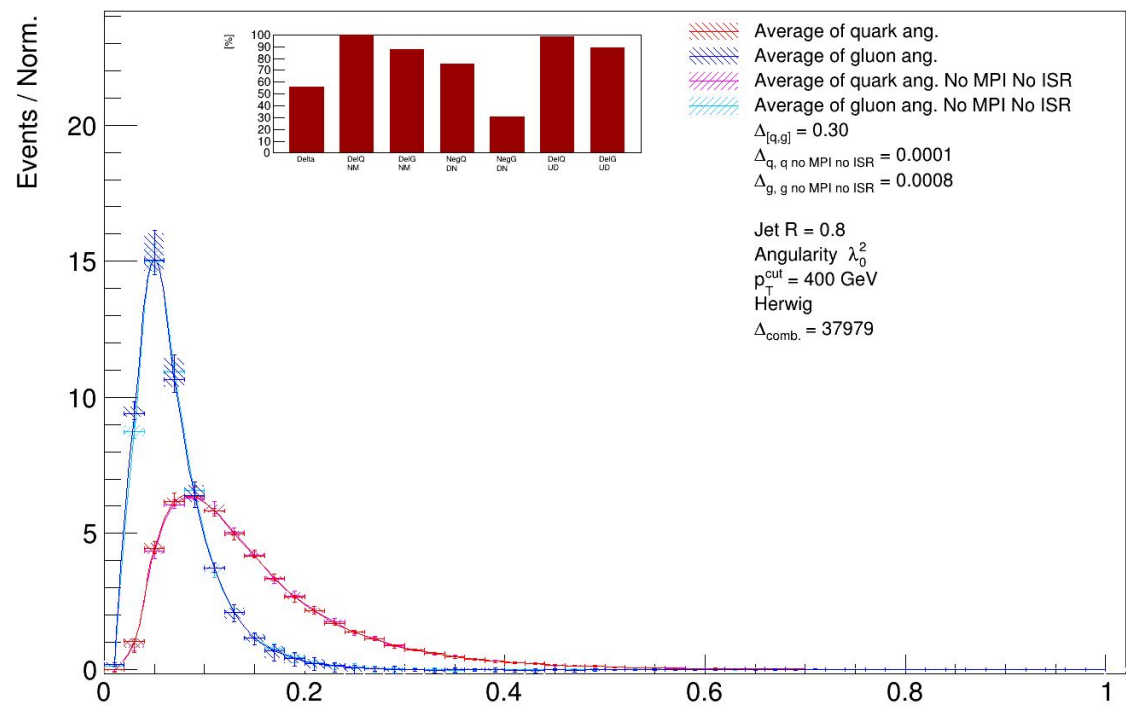
Why we looking into other features then separation power (bad examples):



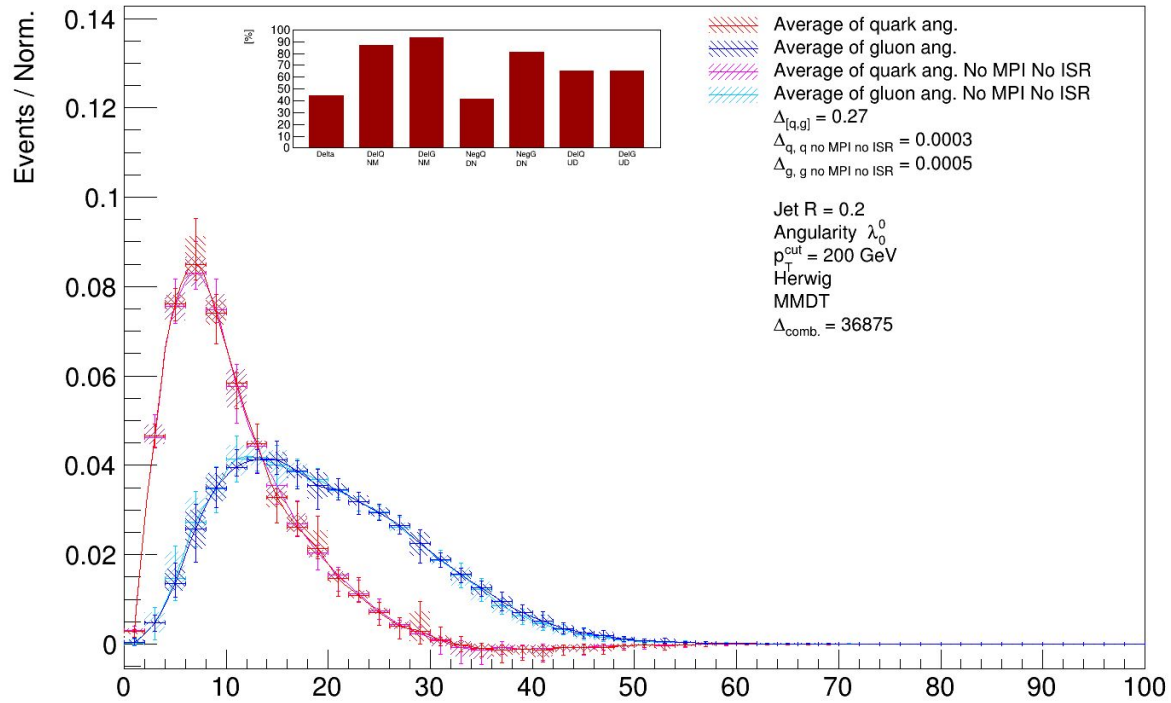
Why we looking into other features then separation power (bad examples):



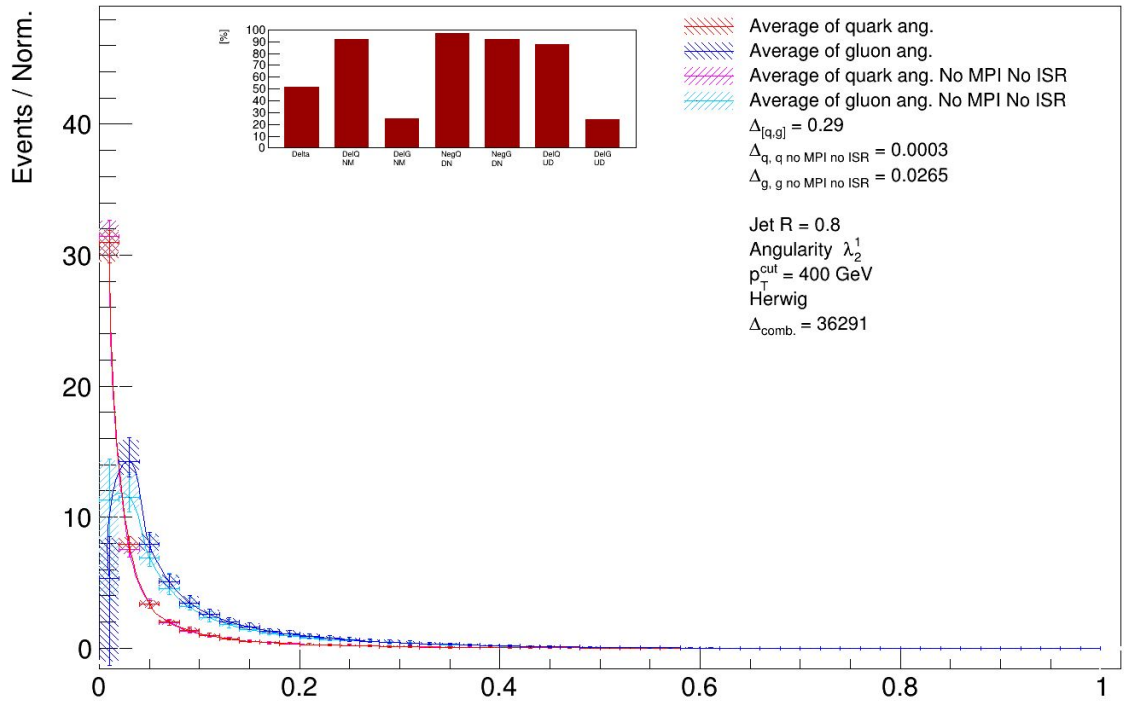
Best performing angularities: p_T^D



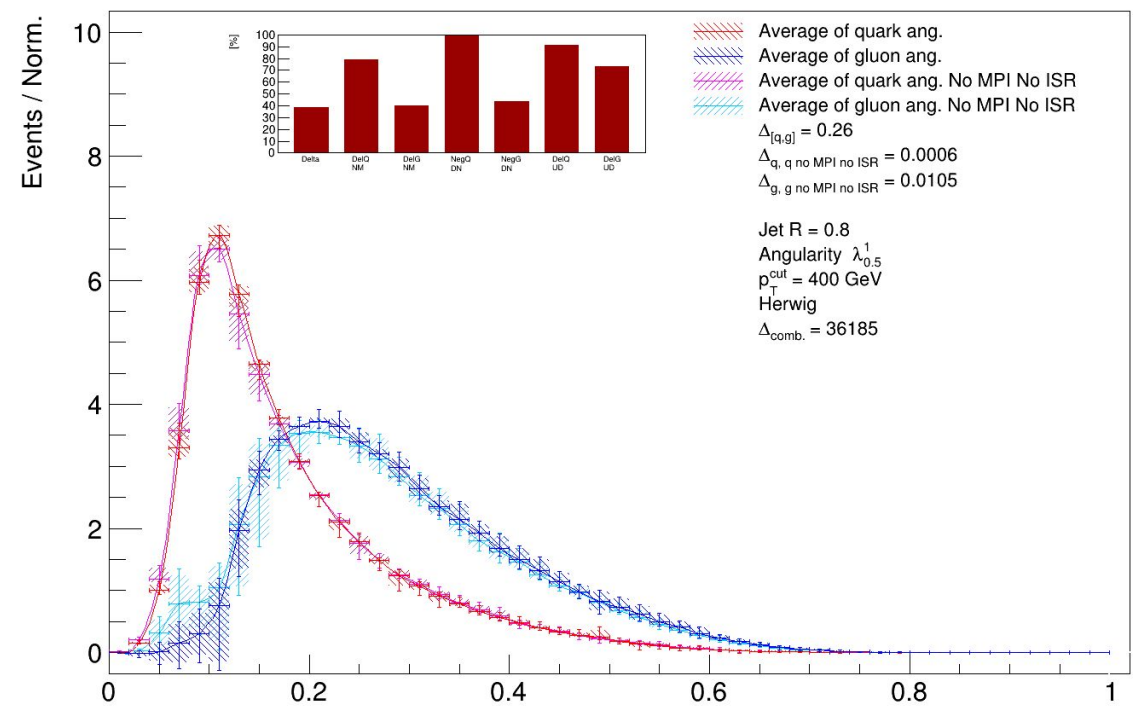
Best performing angularities: Multiplicity



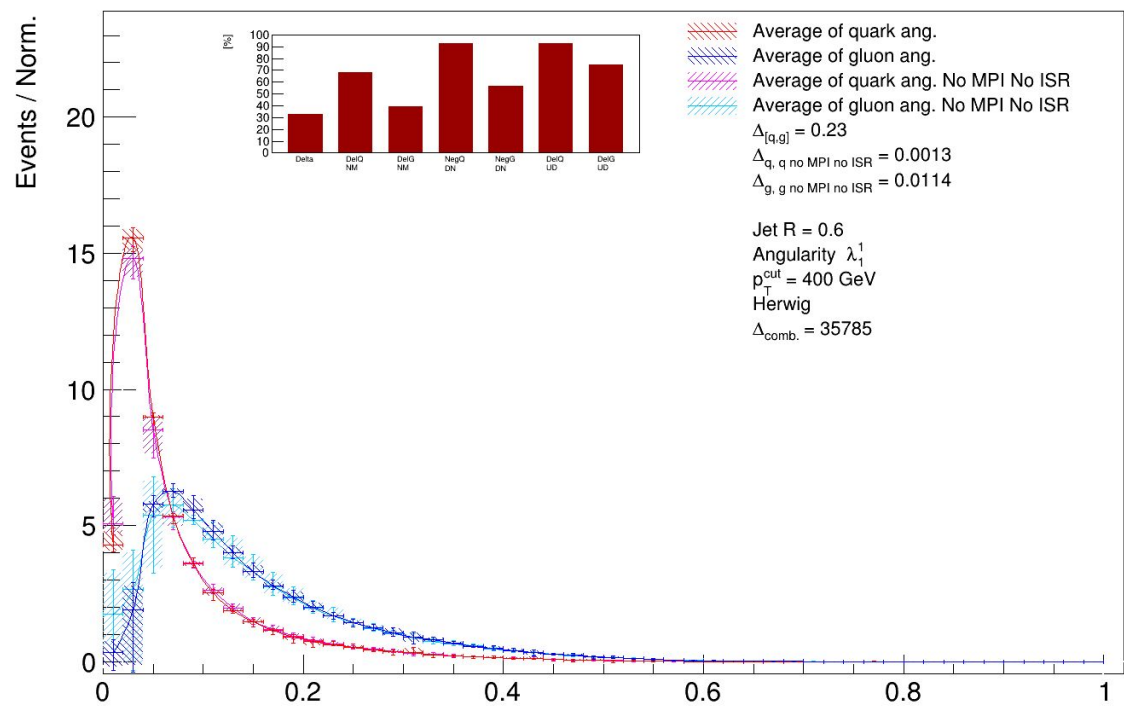
Best performing angularities: Mass



Best performing angularities: LHA



Best performing angularities: Width



4. Conclusion

Conslution:

- Main idea it that properties of jets of a given flavour and transverse momentum, are almost entirely independent of the jet's production mechanism.

Thus, the energy-dependence can be used to extract the flavour-dependent properties on a statistical basis.

- We proposed selection of best performing angularities
- More details: <https://arxiv.org/abs/2307.15378>
- Plans:
 - Multidim angularities (2D, 3D) with machine learning approach to enhance separation power
 - Derive jet topics: <https://arxiv.org/abs/1802.00008> since all ingredients should be in place
 - Perform measurement - we will be happy to contribute:-)

Novel approach to measure quark/gluon jets at the LHC

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²Lancaster-Manchester-Sheffield Consortium for Fundamental Physics, Department of Physics and Astronomy, University of Manchester, M13 9PL, U.K.
³Jagiellonian University, ul. prof. Stanisława Łojasiewicza 11, 30-348 Kraków, Poland

15378v1 [hep-ph] 28 Jul 2023

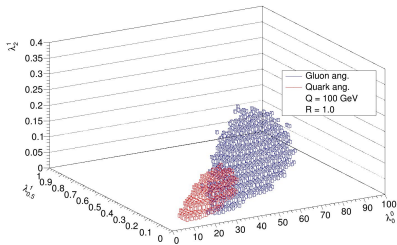
Abstract In this paper, we present a new proposal on how to measure quark/gluon jet properties at the LHC. The measurement strategy takes advantage of the fact that the LHC has collected data at different energies. Measurements at two or more energies can be combined to yield distributions of any jet property separated into quark and gluon jet samples on a statistical basis, without the need for an independent event-by-event tag. We illustrate our method with a variety of different angularity observables, and discuss how to narrow down the search for the most useful observables.

1 Introduction

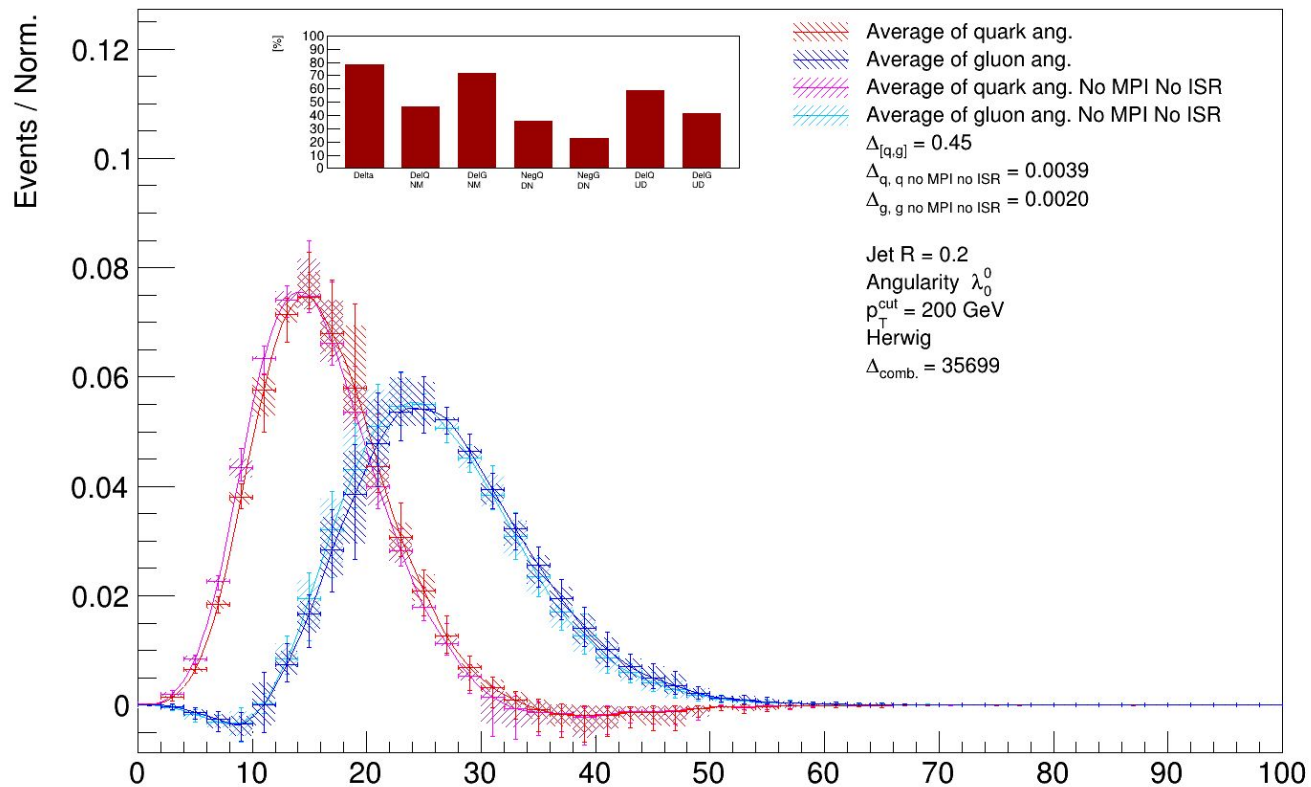
Experimentally, we can study partons (quarks and gluons) by analyzing jets (narrow, energetic sprays of particles) whose kinematic characteristics mirror those of

As well as proposing an observable that can distinguish quark jets and gluon jets [7–16], any quantitative analysis must also propose how to calibrate that observable by independently tagging quark and gluon jet samples. In some studies, this has been done by calibrating against Monte Carlo samples in which the “truth” flavour of the jet is known. However, one might worry about whether event generators make sufficiently reliable predictions of these flavour-dependent properties [17–19] and, indeed, this is something one would like to test against the data. In other studies, another method is used to tag the jet flavour, for example the hard process dependence [20,21], and used to calibrate the measurement of the proposed observable. Here, one would worry that the two tagging methods are correlated, yielding a biased measurement of the jet property.

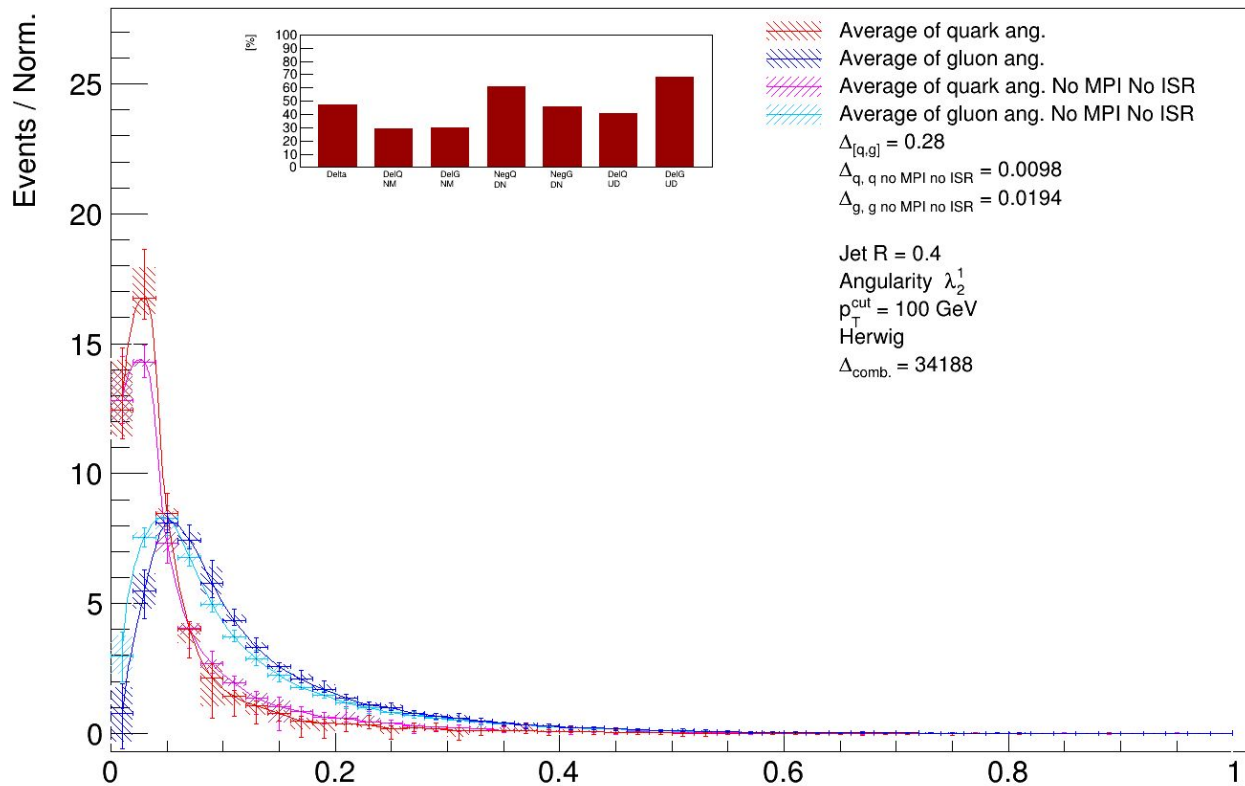
In this paper, we study a variety of angularity ob-



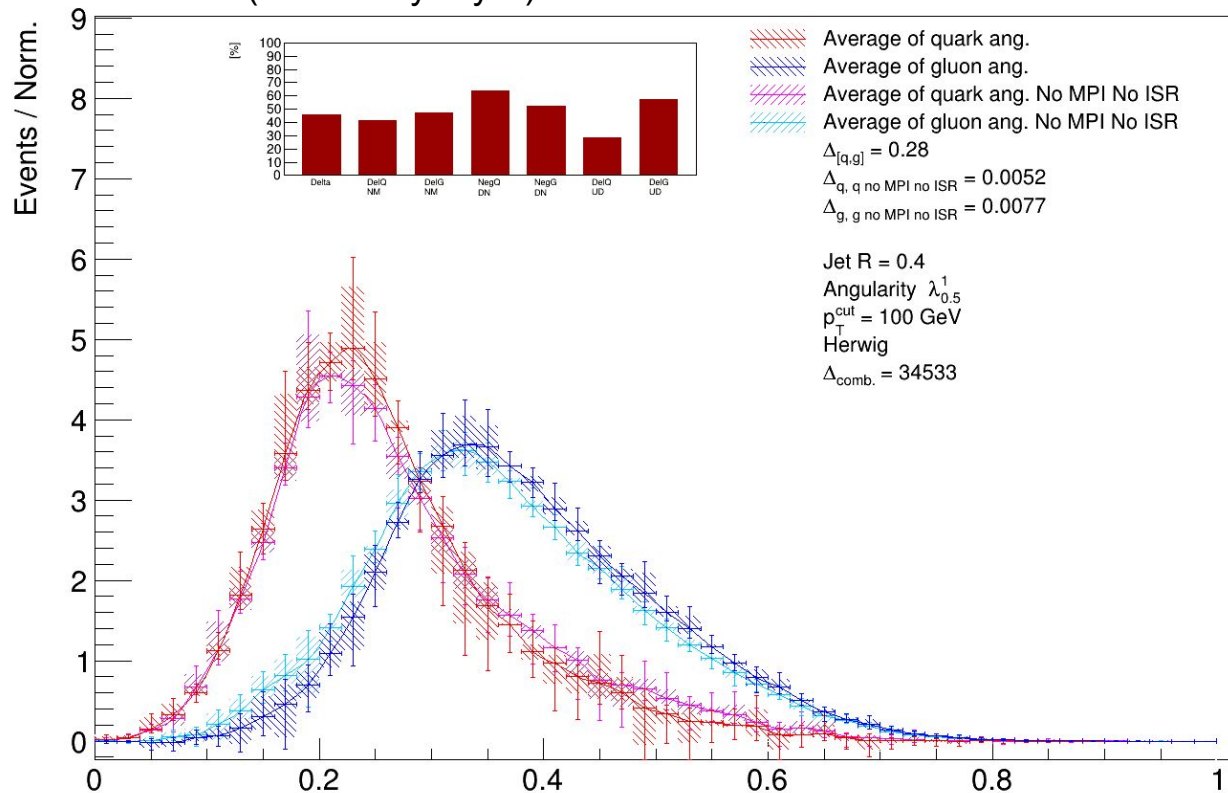
Wild cards (chosen by “eye”): multiplicity



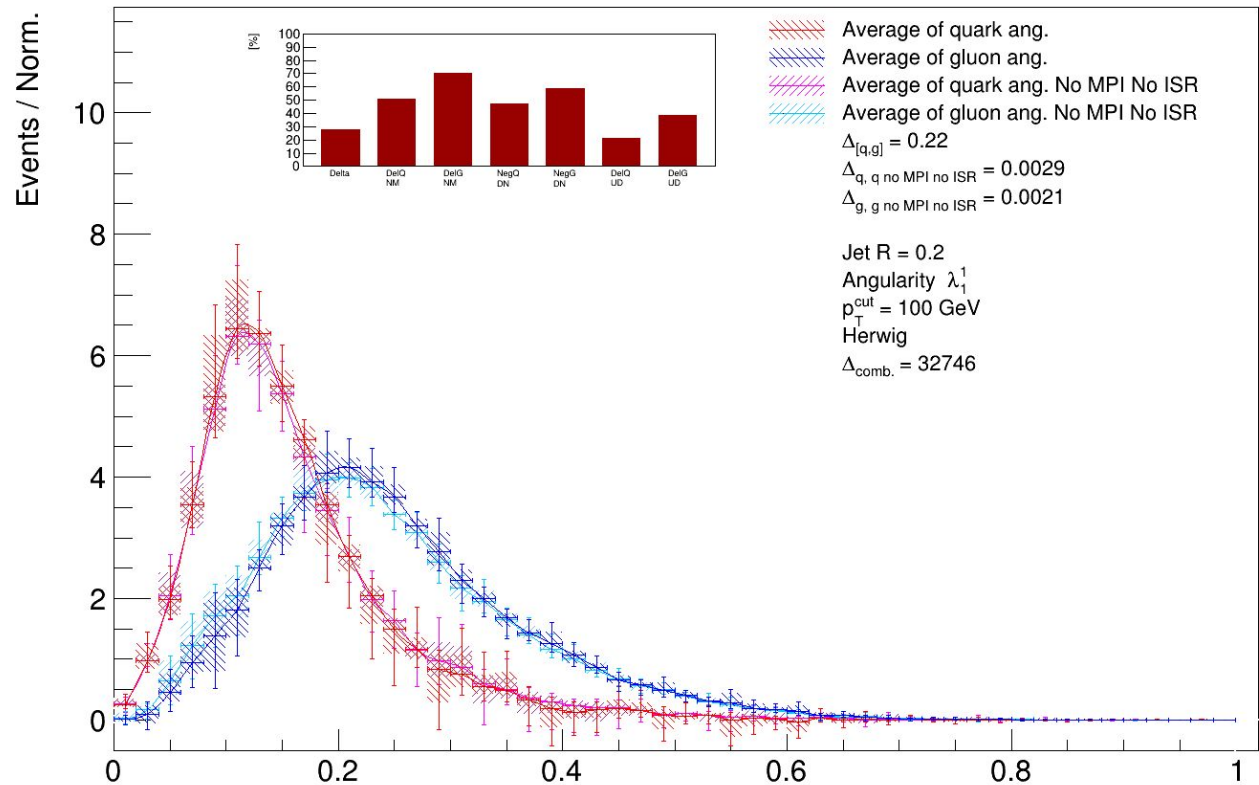
Wild cards (chosen by “eye”): mass



Wild cards (chosen by “eye”): LHA



Wild cards (chosen by “eye”): width



Herwig vs Pythia

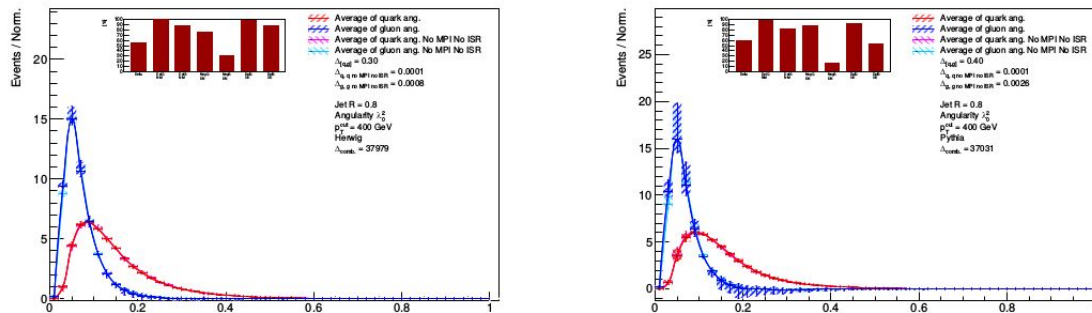


Fig. 17 Quark and gluon averaged angularities λ_0^2 , $R = 0.8$ with highest score $\Delta_{\text{comb}} = 37979$ using HERWIG event generator (left) and $\Delta_{\text{comb}} = 37031$ using PYTHIA event generator (right), with $p_T^{\text{cut}} = 400$ GeV, using the average of 6 energy combinations 900–2360, 900–7000, 900–13000, 2360–7000, 2360–13000, 7000–13000 GeV.

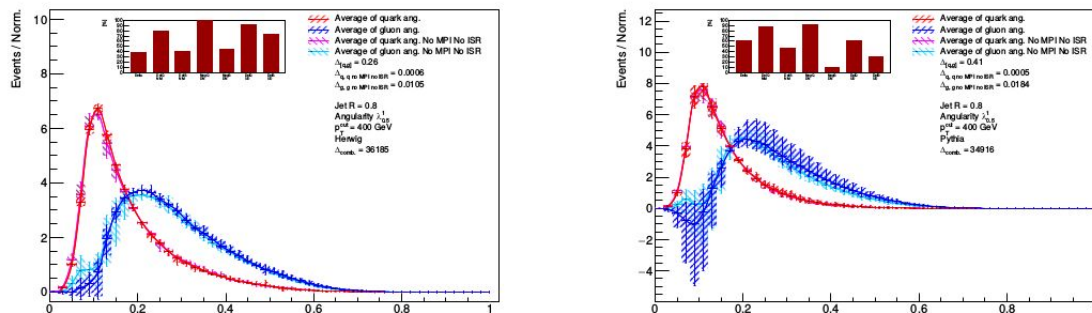
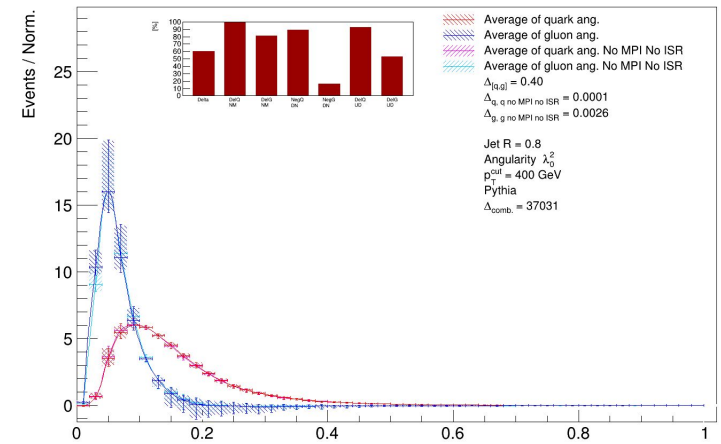
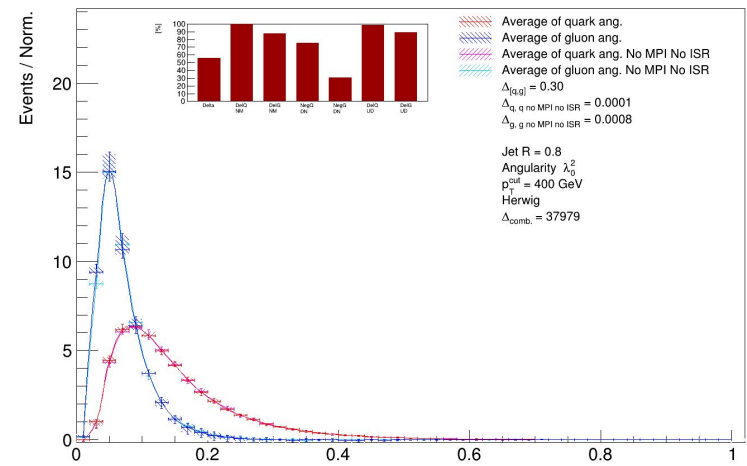
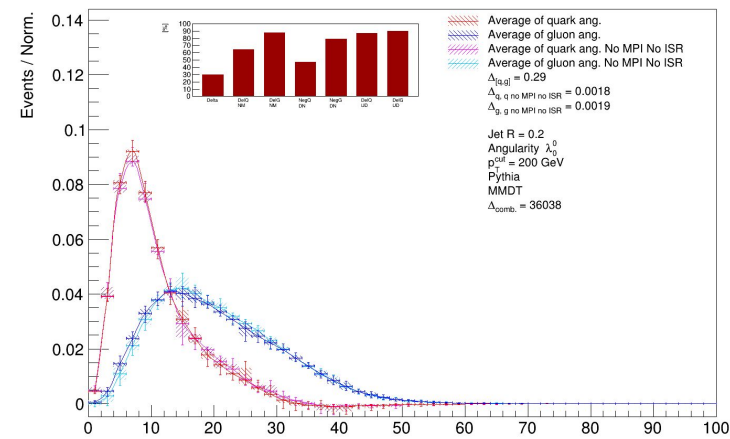
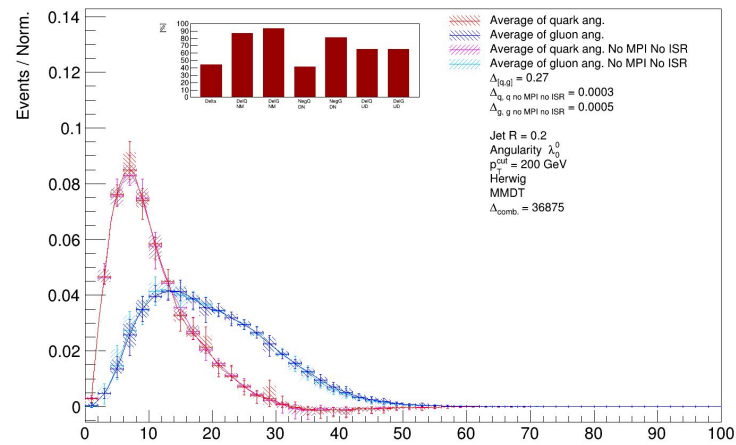


Fig. 18 Quark and gluon averaged angularities MMDT $\lambda_{0.5}^1$, $R = 0.8$ with score $\Delta_{\text{comb}} = 36185$ using HERWIG event generator (left) and $\Delta_{\text{comb}} = 34916$ using PYTHIA event generator (right), with $p_T^{\text{cut}} = 400$ GeV, using the average of 6 energy combinations 900–2360, 900–7000, 900–13000, 2360–7000, 2360–13000, 7000–13000 GeV.

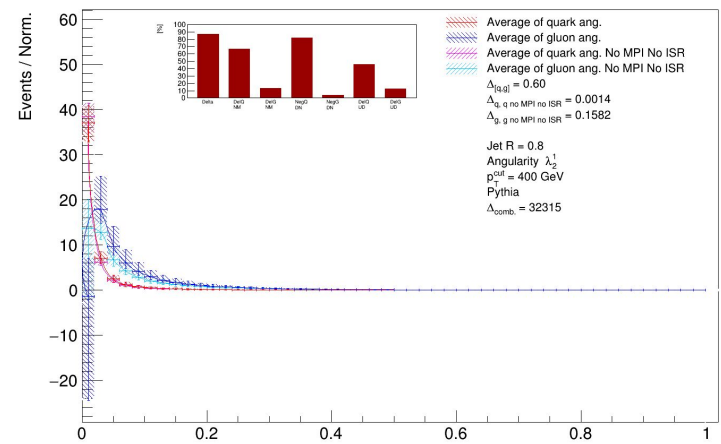
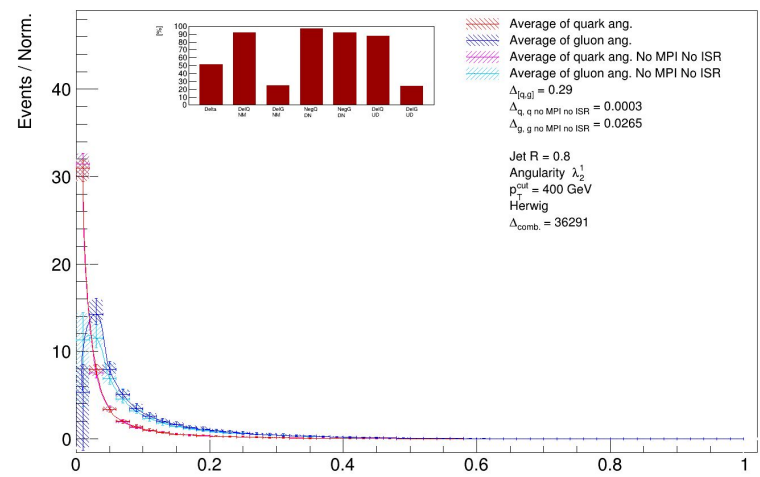
Best performing angularities: p_T^D



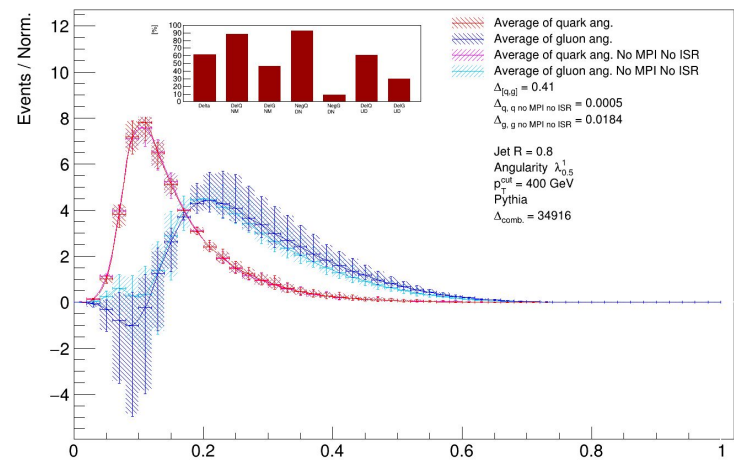
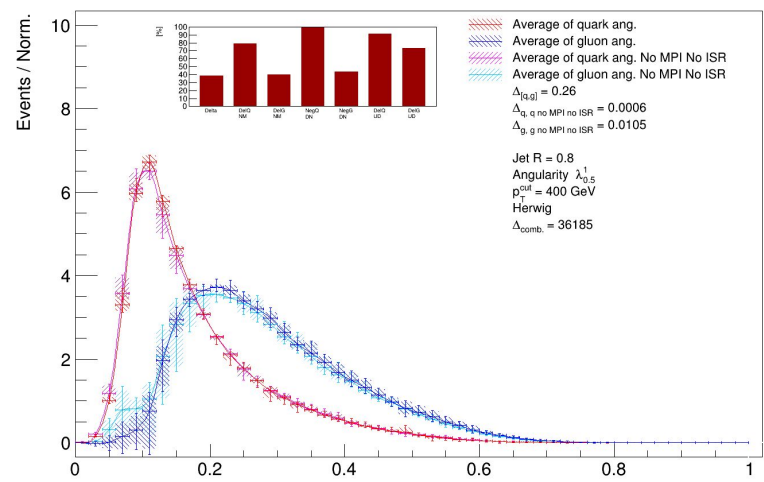
Best performing angularities: Multiplicity



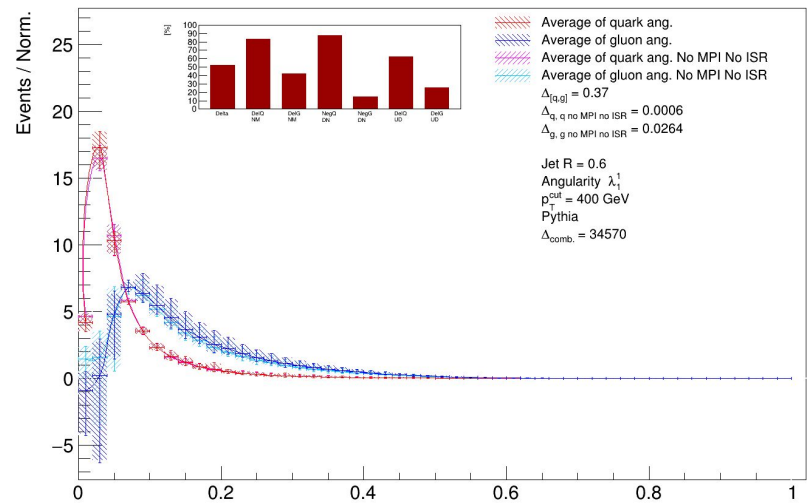
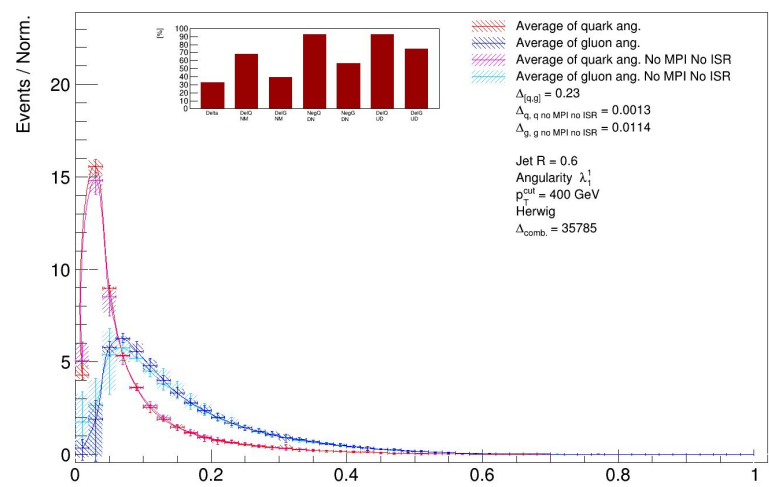
Best performing angularities: Mass



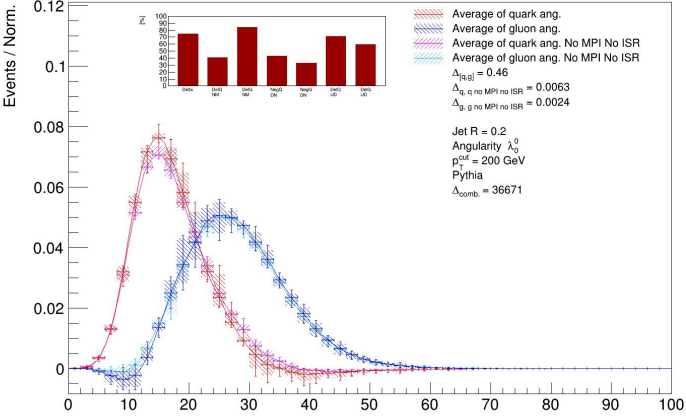
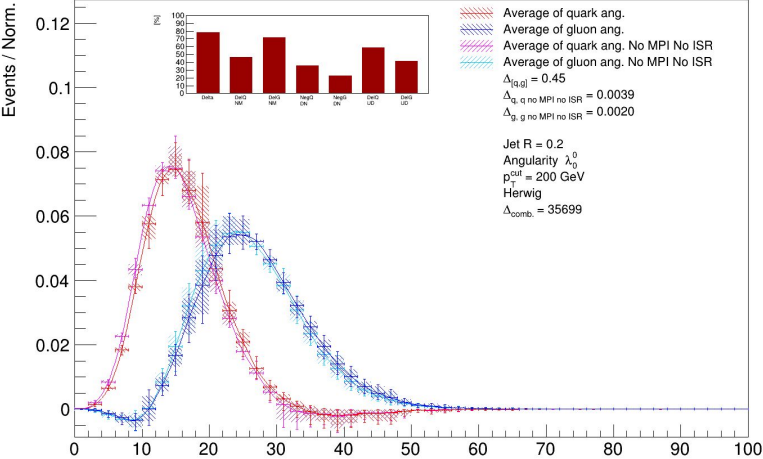
Best performing angularities: LHA



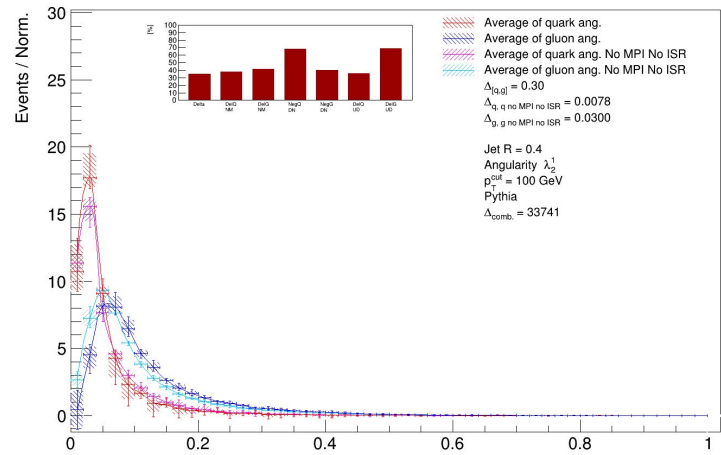
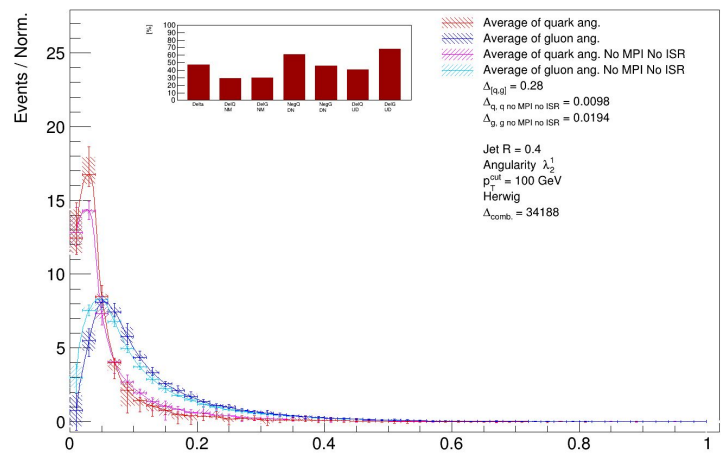
Best performing angularities: Width



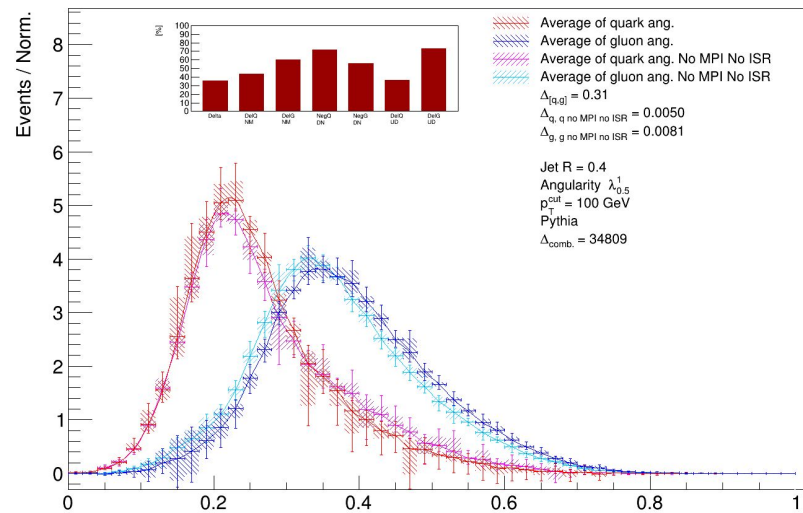
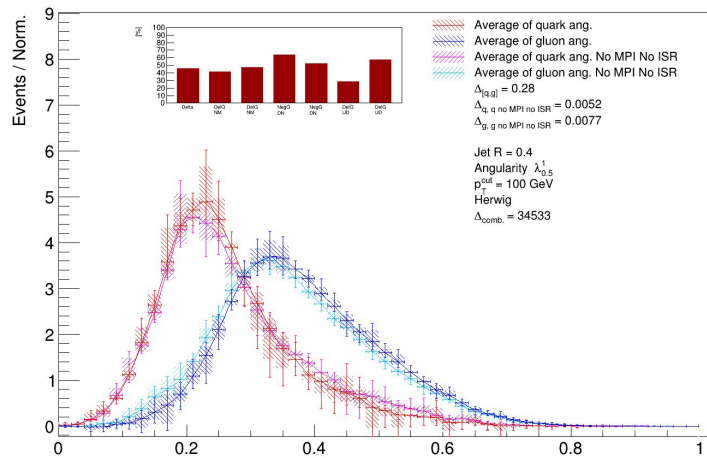
Wild cards (chosen by “eye”): multiplicity



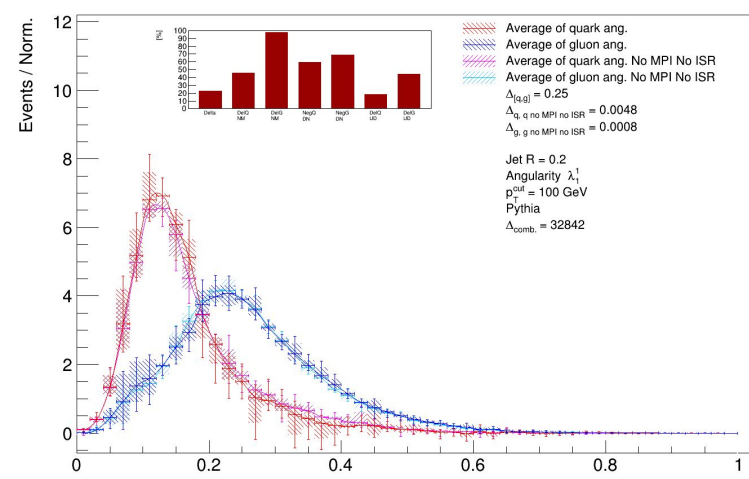
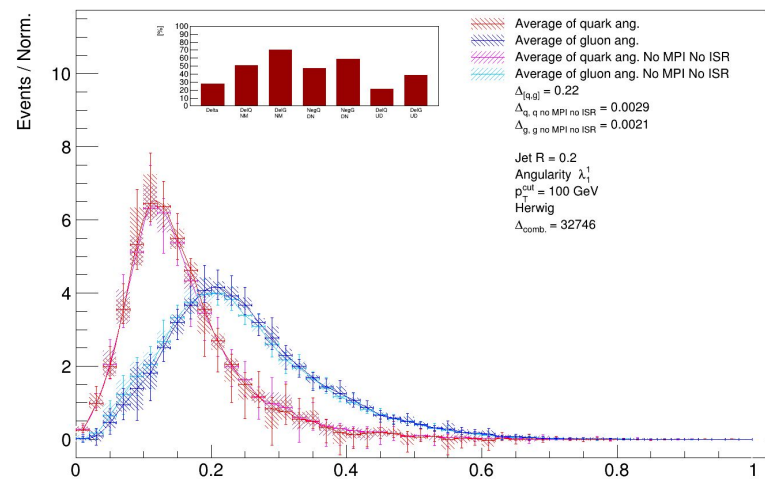
Wild cards (chosen by “eye”): mass

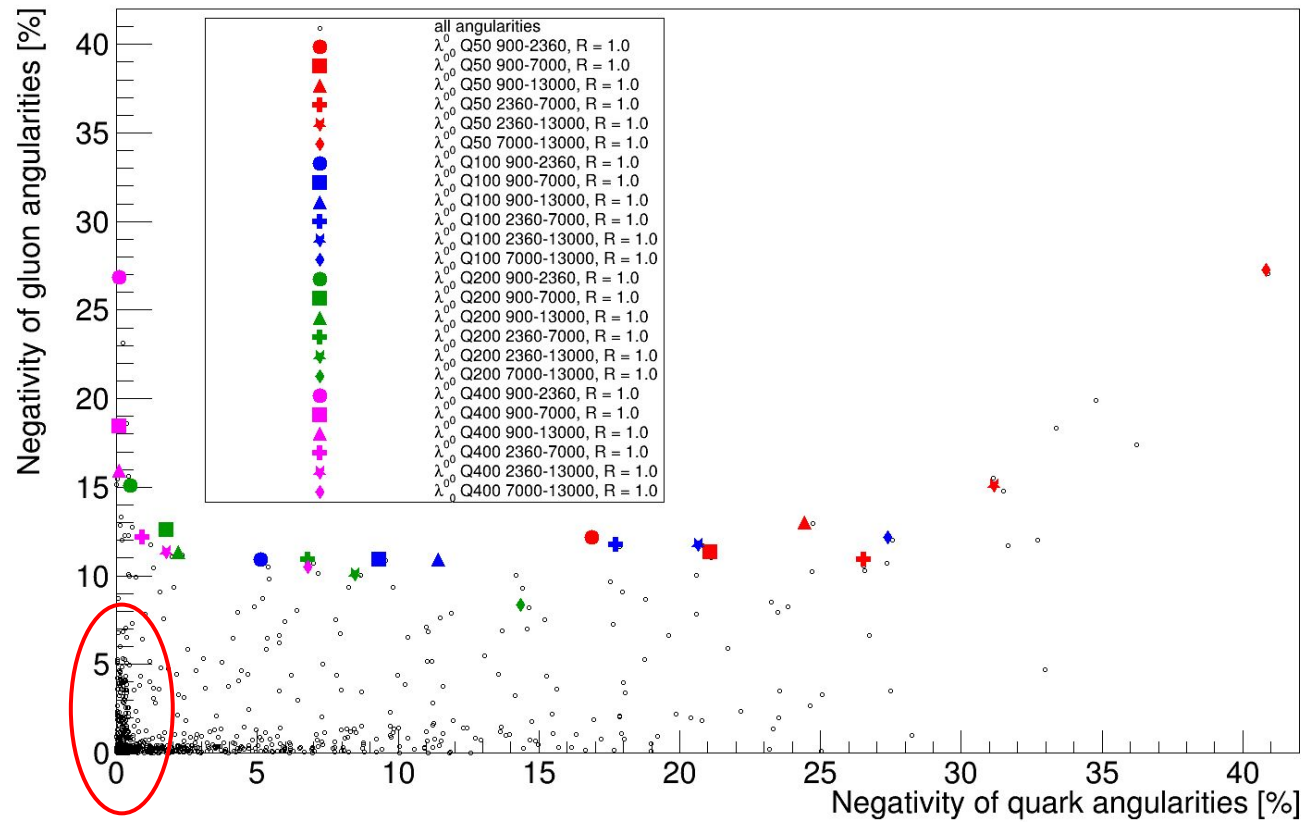


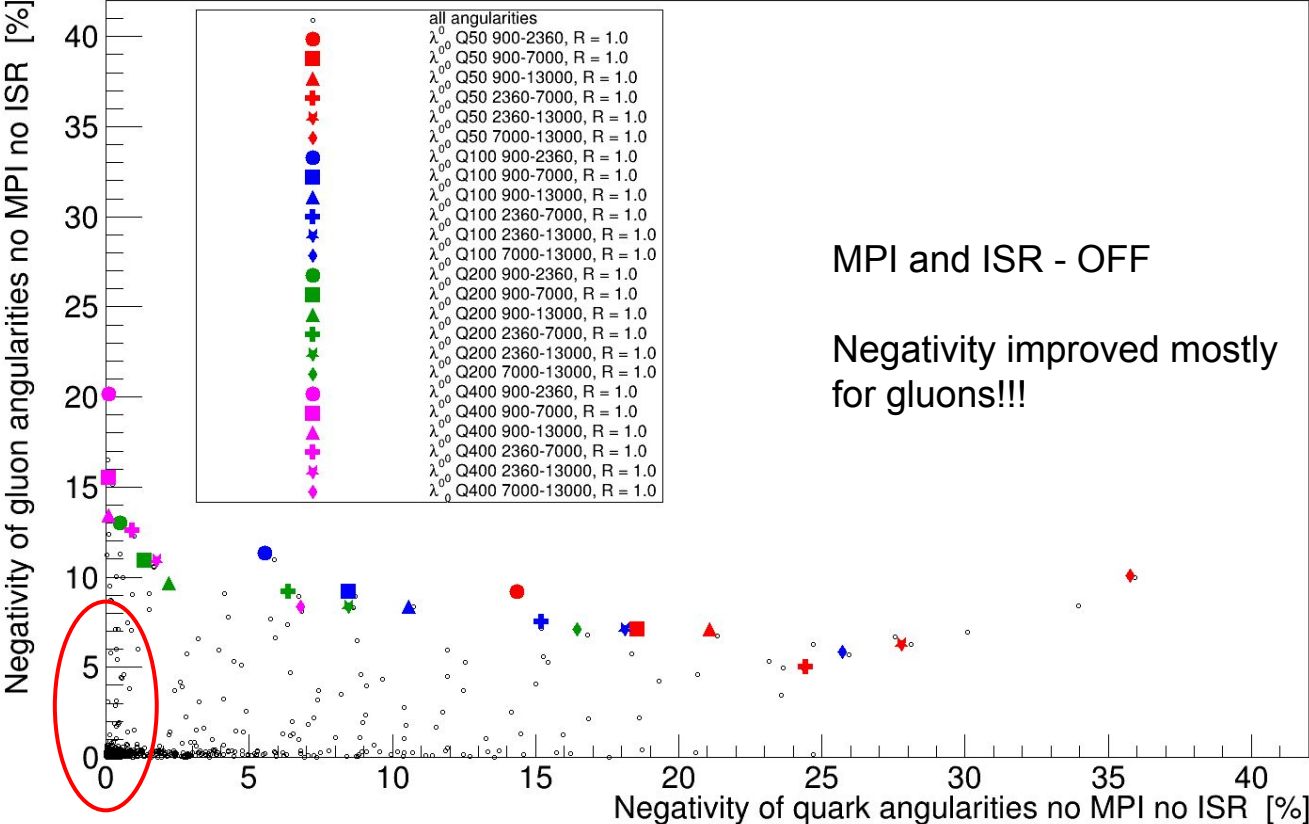
Wild cards (chosen by “eye”): LHA

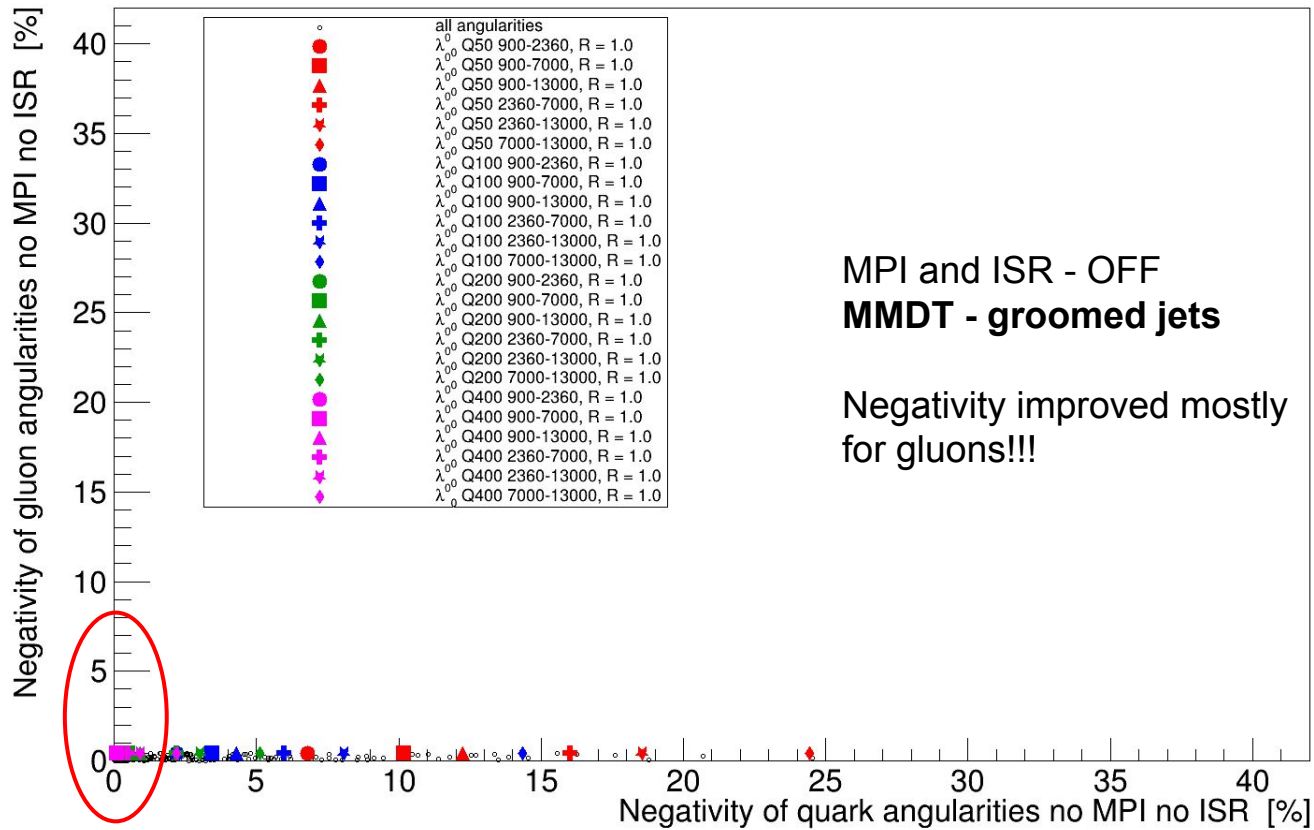


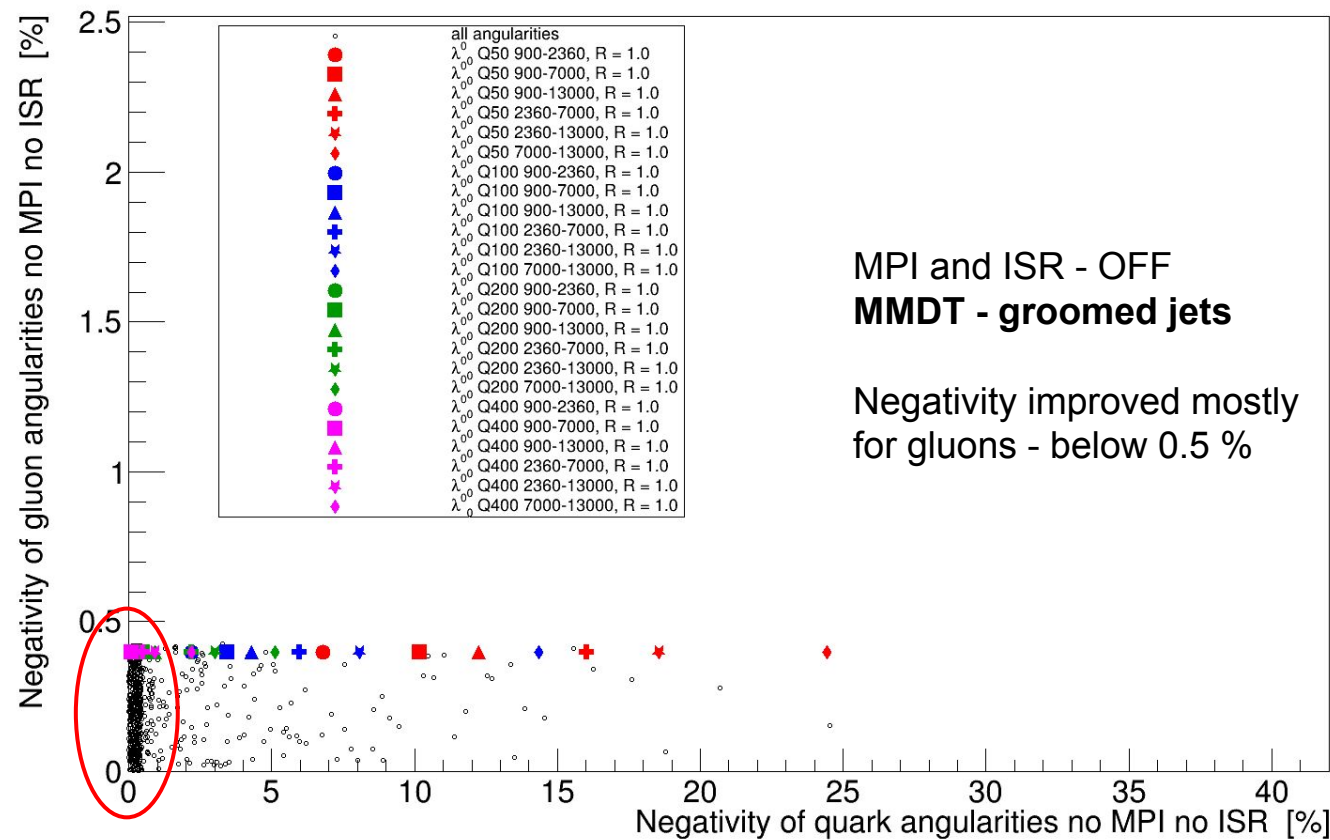
Wild cards (chosen by “eye”): width











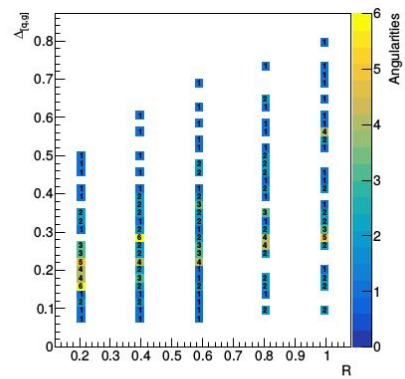


Fig. 19 First column scatter plot of $\Delta_{[q,g]}$ as a function of jet radius.

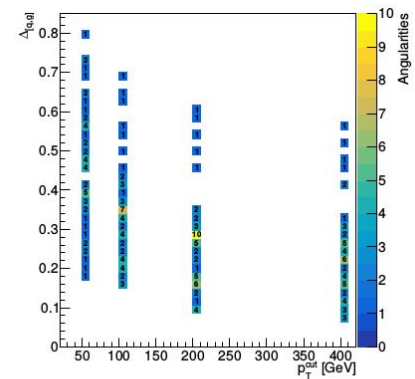


Fig. 21 First column scatter plot of $\Delta_{[q,g]}$ as a function of p_T^{cut} .

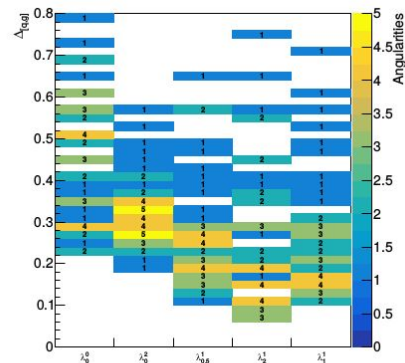


Fig. 20 First column scatter plot of $\Delta_{[q,g]}$ as a function of jet angularity.

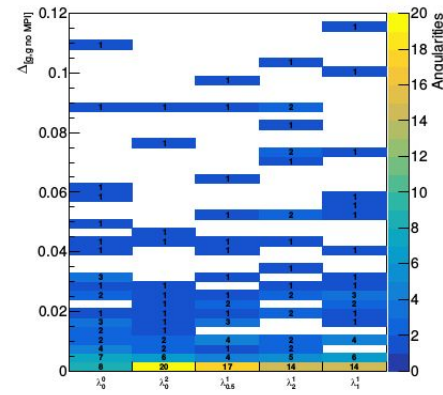
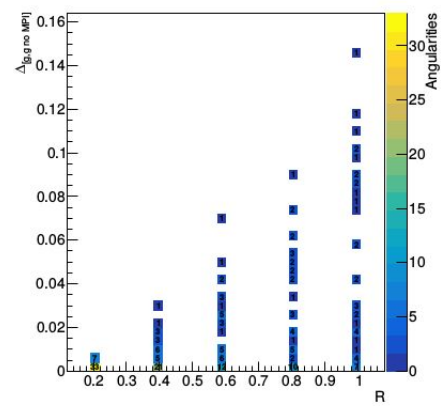
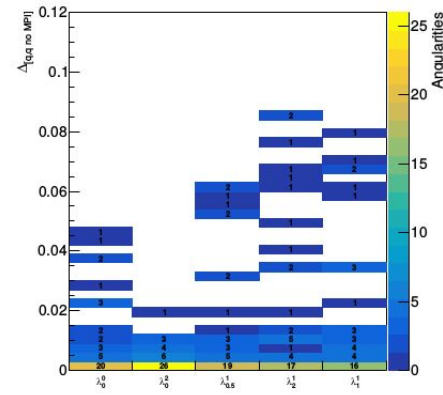
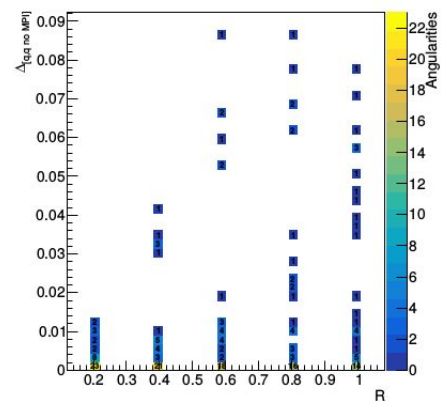


Fig. 22 Second/third column quark $\Delta_{[q,q] \text{ noMPI}}$ (top) and gluon $\Delta_{[g,g] \text{ noMPI}}$ (bottom) as a function of jet radius.

Fig. 23 Second/third column quark $\Delta_{[q,q] \text{ noMPI}}$ (top) and gluon $\Delta_{[g,g] \text{ noMPI}}$ (bottom) as a function of angularities.

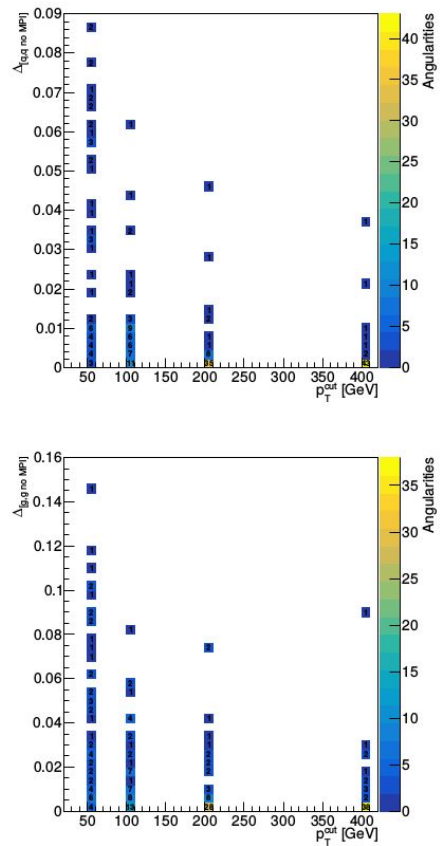


Fig. 24 Second/third column quark $\Delta_{[q,q] \text{ no MPI}}$ (top) and gluon $\Delta_{[g,g] \text{ no MPI}}$ (bottom) as a function of p_T^{jet} .

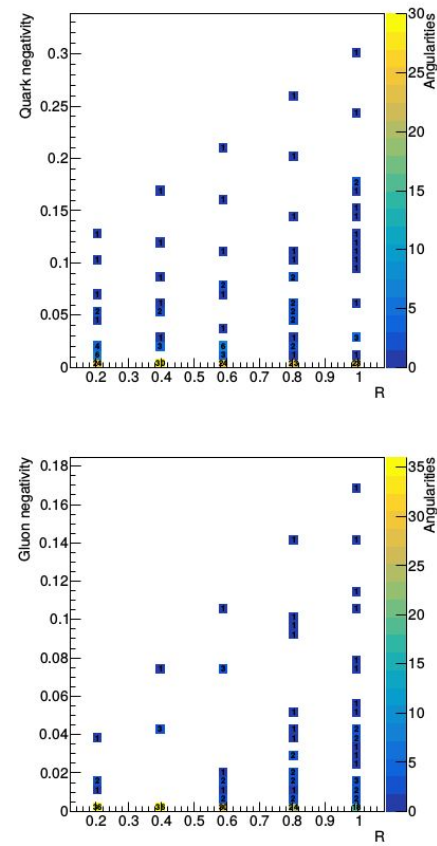


Fig. 25 Fourth/fifth column quark (top) and gluon negativity (bottom) as a function of jet radius.

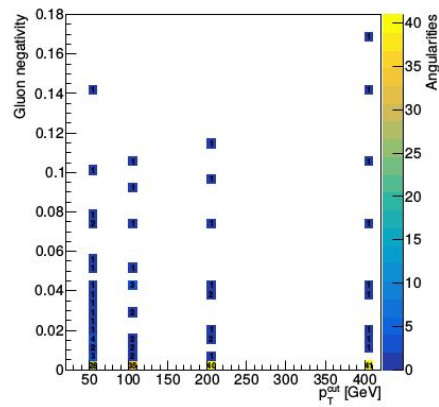
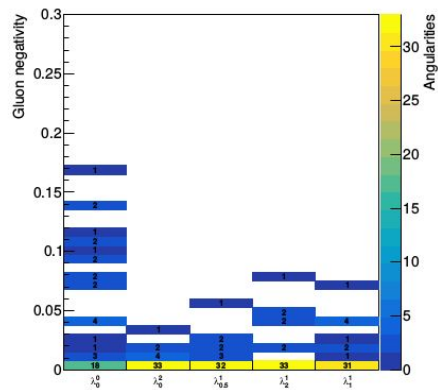
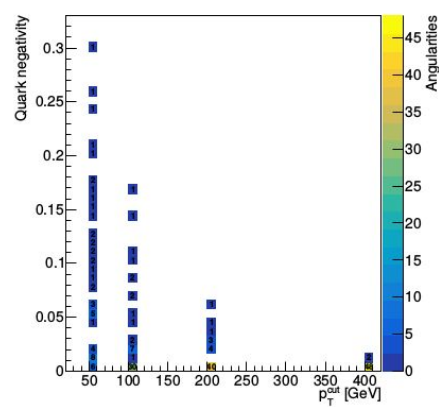
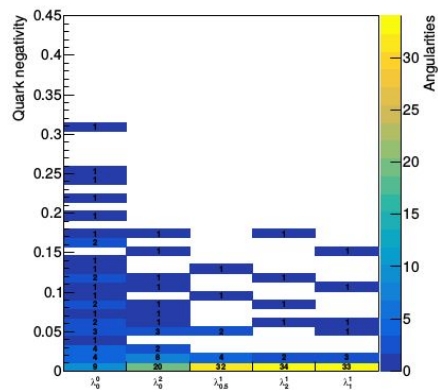


Fig. 26 Fourth/fifth column quark (top) and gluon negativity (bottom) as a function of angularities.

Fig. 27 Fourth/fifth column quark (top) and gluon negativity (bottom) as a function of p_T^{cut} .

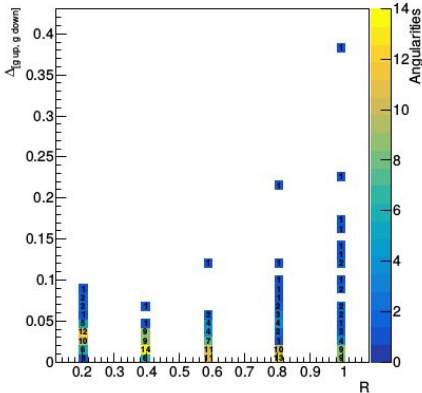
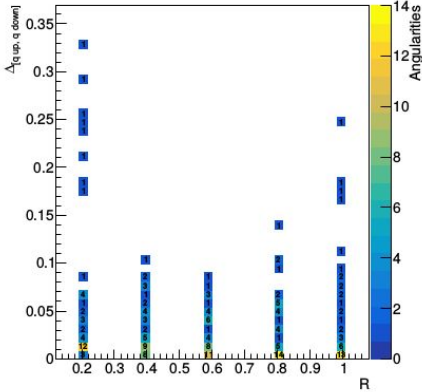


Fig. 28 Sixth/Seventh column quark $\Delta_{[q \text{ down}, q \text{ up}]}$ (top) and gluon $\Delta_{[g \text{ down}, g \text{ up}]}$ (bottom) as a function of jet radius.

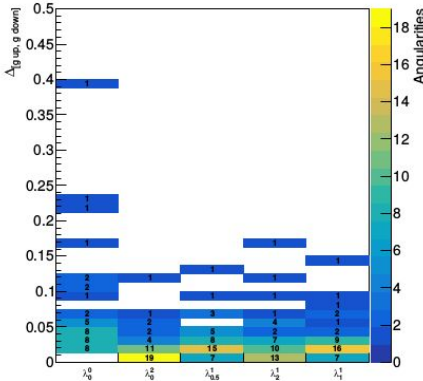
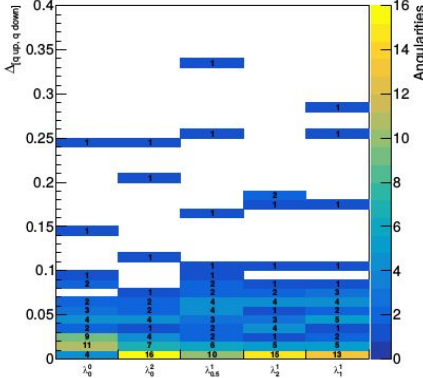


Fig. 29 Sixth/Seventh column quark $\Delta_{[q \text{ down}, q \text{ up}]}$ (top) and gluon $\Delta_{[g \text{ down}, g \text{ up}]}$ (bottom) as a function of angularities.

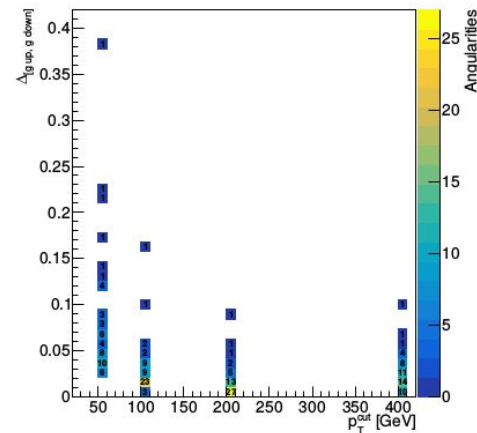
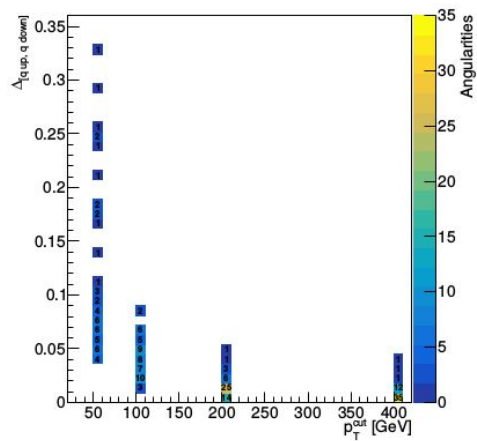


Fig. 30 Sixth/Seventh column quark $\Delta_{[q \text{ down}, q \text{ up}]}$ (top) and gluon $\Delta_{[g \text{ down}, g \text{ up}]}$ (bottom) as a function of p_T^{cut} .