Resonances in the Dijet Mass

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Overview

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- Example Models
 - \circ Excited Quarks
 - $\circ \quad \ \ {\rm Quantum \ Black \ Holes}$
 - $\circ \quad \ \ Z' \ DM \ \, Mediator$
- Modeling the Background
- Selection Criteria
- Results

Motivation

- New particles that directly interact with partons of the proton can product partons when they decay
- Common feature in many BSM models accessible at the LHC
- A new resonant state decaying to two jets may introduce an excess in m_{jj} distribution, near the mass of the resonance
- The production rates for BSM dijet final states can be large, allowing searches at high masses
- Excited quarks, quantum black holes, and W', W*, and Z' bosons would produce peaks in the invariant mass distribution
- Contact interactions would introduce smooth changes in the high-mass tail of the m_{jj} distribution that could be detected in the analysis of the χ distributions (Recall Alex's Talk)

Bump Hunting 101



Excited Quarks (q*)

- Composite models predict them
- Excited state should not be much lighter than the substructure scale → predicted to be pretty heavy
- Predicted to decay predominantly into ordinary quarks and gluons
- They assume spin-1/2 excited quarks with same coupling constants as SM quarks
- Only decay of excited quark \rightarrow up or down quark simulated (BR 85%)

Quantum Black Holes

- Not accessible at LHC based on limits on the Planck scale, smallest black hole mass $\sim 2 {\rm x} 10^{-8} ~{\rm kg}$
- If extra spatial dimensions exist \rightarrow Planck scale could be smaller \rightarrow the fundamental scale of gravity could be lowered
- Most optimistic scenarios of extra dimensions \rightarrow lowest Planck scale \rightarrow LHC is a black hole factory
- In this model the branching ratio to dijets is 96%
- ADD scenario and dimensions n=6
- QBH would appear as in excess in $\rm m_{jj}$ distribution localized near the threshold mass for QBH production $\rm M_{th}$

Z' DM Mediator

- Z' dark matter mediator model
- Z' can in principle have large couplings to both DM and SM fermions
- Any mediator produced from SM particles in the initial state can also decay back into SM states
- Searches for a DM mediator complementary to usual missing ET searches. Resonance searches complementary to direct detection constraints as well.
- The model assumes axial-vector couplings to all SM quarks and to a Dirac fermion dark matter candidate. Leptophobic.
- In this scenario the Z' BR to DM is negligible, so that production rate and resonance width depends only on g_q and the mass of the resonance

Modeling the QCD Background

- Background dominated by $2 \rightarrow 2$ scattering of partons predicted by perturbative QCD
- No searches use MC simulation to model the QCD background
- Theoretical uncertainties in QCD (PDFs, renormalization and factorization scales, strong coupling constant, non-perturbative effects, ...)





Data-Driven BG Estimate: Sliding-Window Fit

- Limitations at the high mass tail
- Before: $f(z) = p_1(1-z)^{p_2} z^{p_3} z^{p_4 \log z}$ where $z = m_{jj} / s^{1/2}$
- Increases in luminosity \rightarrow cannot rely in $f(z) \rightarrow$ opportunity to develop new techniques...
- Traditional global fit with p_4 is still viable for this analysis \rightarrow opportunity to use new method
- Fit data in each window \rightarrow fit value at center of window is BG description
- Window size chosen such that 3-parameter f(z) describes data well in window
- Nominal window size $\sim \frac{1}{2}$ of total bins
- Nominal window size wide enough for all considered models to fit within one window.

Selection Criteria

	$p_{\mathrm{T}}^{\mathrm{leading}}$	$p_{\mathrm{T}}^{\mathrm{subleading}}$	$ y^* $	$ y_{ m B} $	m_{jj}
Resonance	> 0.44 TeV	> 0.06 TeV	< 0.6	_	> 1.1 TeV
W^*	> 0.44 TeV	$> 0.06 { m TeV}$	< 1.2	2	> 1.7 TeV
Angular	> 0.44 TeV	$> 0.06 { m TeV}$	< 1.7	< 1.1	> 2.5 TeV

 p_{T} requirement ensures a trigger efficiency of at least 99.5% for collisions that enter into the analysis

Rapidity of an outgoing parton: $y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right)$ Rapidity difference between the two jets: $y * = \frac{y_1 - y_2}{2}$

Wider $|y^*| < 1.2$ region optimized for signals produced at forward angles (W* peaks at $|y^*| > 1.0$)

All other benchmark models peak at $y^*=0$

The selection is only fully efficient for $m_{ii} > 1.1 \text{ TeV} (1.7 \text{ TeV})$

Reconstructed Dijet Mass Distribution



 $\begin{array}{l} \text{BumpHunter algorithm} \\ \text{quantifies statistical} \\ \text{significance of localized} \\ \text{excess in } m_{jj} \end{array}$

Model Specific Limits







At high coupling $g_q > 0.6$ the width of the Z' mass increases to 15% and beyond, loss of sensitivity. Search limited to $g_q < 0.5$.

Model Independent Limits

Resonance Shapes:

- Narrow width: width dominated by experimental resolution
- Difference in widths dominated by difference in # of decay channels
- Limits worsen for wider signals



95% CL Lower Limits on Masses

Model	95% CL exclusion limit			
	Observed	Expected		
Quantum black hole	$8.9~{ m TeV}$	8.9 TeV <		
W'	$3.6 { m ~TeV}$	$3.7~{\rm TeV}$		
W^*	3.4 TeV 3.77 TeV – 3.85 TeV	3.6 TeV		
Excited quark	$6.0 { m TeV}$	$5.8 { m TeV}$		
$Z' \ (g_q = 0.1)$	$2.1 { m ~TeV}$	$2.1 { m ~TeV}$		
$Z' \ (g_q = 0.2)$	$2.9~{ m TeV}$	$3.3~{\rm TeV}$		
Contact interaction $(\eta_{\rm LL} = -1)$	$21.8~{\rm TeV}$	$28.3 { m TeV}$		
Contact interaction $(\eta_{LL} = +1)$	$\begin{array}{c} 13.1 {\rm TeV} \\ 17.4 {\rm TeV} - 29.5 {\rm TeV} \end{array}$	$15.0 { m TeV}$		

 $\begin{array}{c} \text{Largest } \text{m}_{\text{jj}} \text{ detected } 8.12 \\ \text{TeV} \end{array}$

Larger than largest mass detected!

Extrapolating into region with no data.

Conclusion

- Limits set on the masses of excited quarks, qbh, W' and Z' bosons, excited chiral W* bosons, on contact interaction scales, and generic Gaussian-shaped signal production.
- Improvements on previous limits from 7% for QBH masses to 40% for W' masses
- Any new model predicting excesses in the dijet mass has generic limits set on that excess as a function of width



Table 3. Searches for dijet resonances at hadron colliders. For each search we list the experiment, year of data publication, center-of-mass energy, integrated luminosity, techniques used to define the resonance shape and the background, dijet mass range of the data, the cut applied on the center of mass scattering angle, and the primary reference for the search.

Expt.	Yr.	\sqrt{s} (TeV)	$\int L dt \\ (pb^{-1})$	Resonance Shape	Background Shape	m_{JJ} (TeV)	$\cos \theta^*$ Cut	Ref. #
UA1	86	0.63	0.26	$BW \oplus Gaussian$	LO QCD	.07 - 0.3	-	12
UA1	88	0.63	0.49	$BW \oplus Gaussian$	LO QCD	.11 - 0.3	bins	25
CDF	90	1.8	0.026	$BW \oplus Gaussian$	LO QCD	.06 - 0.5	-	26
UA2	90	0.63	4.7	Gaussian	Fit Func.	.05 - 0.3	-	27
CDF	93	1.8	4.2	$\mathrm{BW} \oplus \mathrm{Resolution}$	LO QCD	.14 - 1.0	-	28
UA2	93	0.63	11	Gaussian	Fit Func.	.05 - 0.3	.60	29
CDF	95	1.8	19	Pythia \oplus Sim.	Fit Func.	.15 - 0.9	.67	30
CDF	97	1.8	106	Pythia \oplus Sim.	Fit Func.	.18 - 1.0	.67	31
D0	04	1.8	109	Pythia \oplus Sim.	NLO QCD	.18 - 1.2	.67	32
CDF	09	1.96	1130	Pythia \oplus Sim.	Fit Func.	.18 - 1.3	-	33
ATLAS	10	7	0.32	Pythia \oplus Sim.	Fit Func.	.20 - 1.7	.57	34
CMS	10	7	2.9	Pythia \oplus Sim.	Fit Func.	.22 - 2.1	.57	35
ATLAS	11w	7	36	PYTHIA	Fit Func.	.50 - 2.8	.57	36
CMS	11	7	1000	Pythia \oplus Sim.	Fit Func.	.84 - 3.7	.57	37
ATLAS	11s	7	1000	Pythia \oplus Sim.	Fit Func.	.72 - 4.1	.54	38