

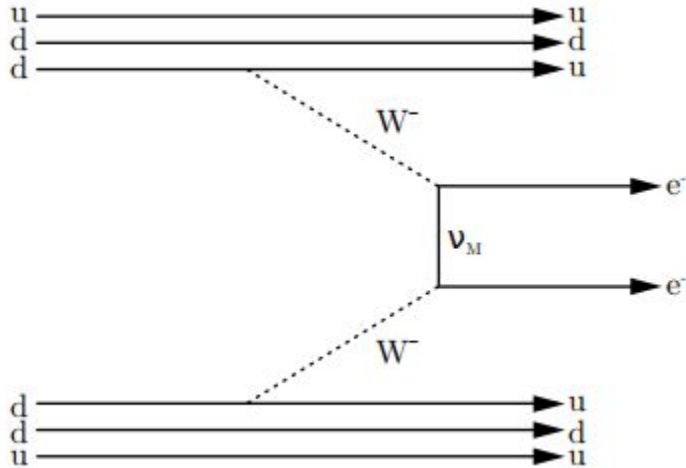
Neutrinoless Double Beta Decay

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Overview

- Theoretical background
- Experimental design
- Selection of results

Why Should I Care?



- $0\nu\beta\beta$: nuclear double beta decay where no neutrinos are in the final state (left)
- $(A, Z) \rightarrow (A, Z + 2) + 2e^-$
- SM conserves total lepton number L
- $0\nu\beta\beta$ would be direct evidence of BSM processes
- Premised on neutrinos being Majorana particles

Majorana what nows?



Ettore Majorana

- All SM particles are (so far) all Dirac
- Majorana particles are massive particles which are their own antiparticles
- Neutrinos are only SM candidates for Majorana particles (massive and neutral)
- Appear in Lagrangians as terms like

$$\mathcal{L}_{\text{mass}} = \frac{1}{2} \sum_{\ell, \ell' = e, \mu, \tau} \nu_{\ell}^t C^{-1} M_{\ell\ell'} \nu_{\ell'} + h.c.$$

- Most interesting man alive?

Right Handed Neutrinos and the Seesaw Mechanism

How can I shove these cool new mass terms into the standard model?

Kill two birds with one stone!

$$\mathcal{L}_{\text{mass}} = -\bar{\nu}_{Ri} m_{\ell i}^{\text{Dirac}} \nu_{Ll} + \frac{1}{2} \bar{\nu}_{Ri} M_i C \bar{\nu}_{Ri}^t + h.c.$$

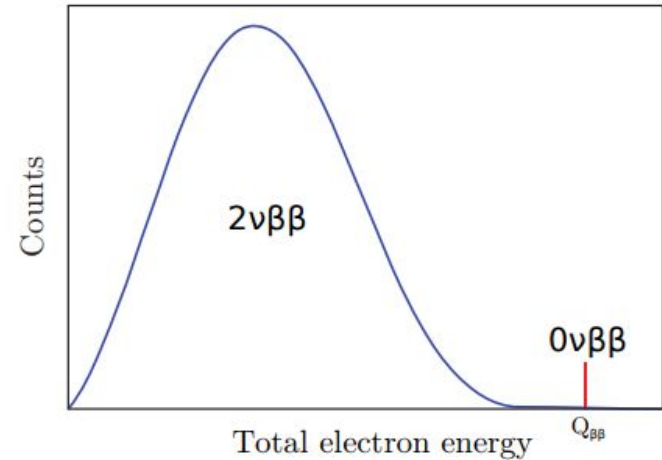
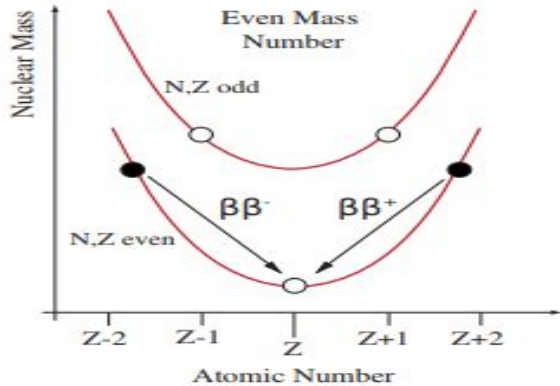
Heavy neutrinos have Majorana mass and at low energies, light neutrinos end up with effective Majorana mass inversely proportional to heavy neutrino mass

$$A = \begin{pmatrix} 0 & M \\ M & B \end{pmatrix}, \quad \lambda_{\pm} = \frac{B \pm \sqrt{B^2 + 4M^2}}{2}, \quad \lambda_{-} \approx -\frac{M^2}{B}.$$

Why haven't we already done this?

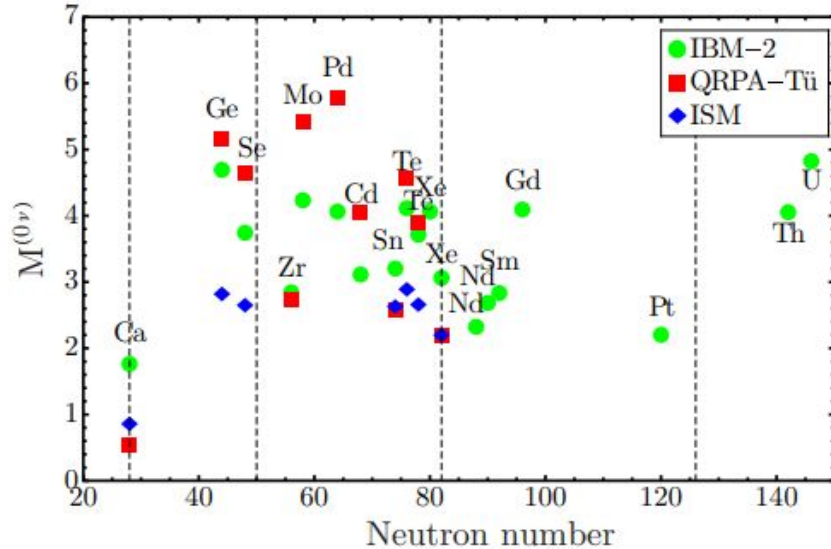
Now we want to observe $0\nu\beta\beta$, what's stopping us?

1. Nuclear matrix element calculations
2. Tiny rate
3. Energy resolution



$$\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 \left(\frac{\langle m_{\beta\beta} \rangle}{m_e} \right)^2$$

NME: Why You Shouldn't Try to Do Nuclear Physics

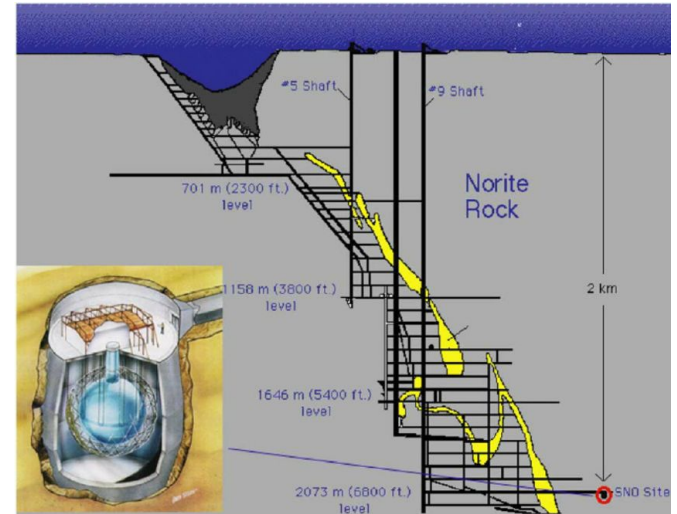
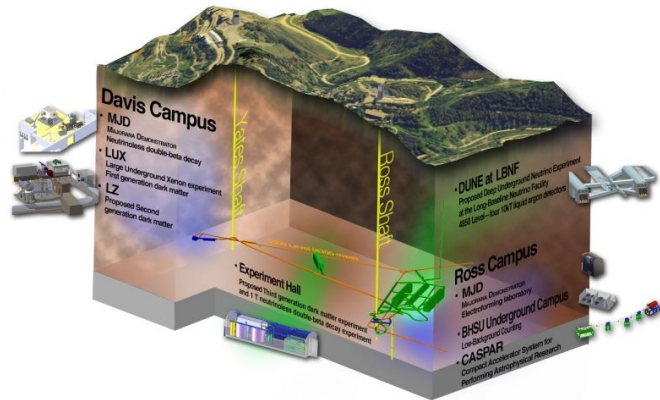


- Experiments are large piles of material in which this decay is possible.
- Simply measuring the energy of the decays
- Nuclear Matrix Element Calculations are very difficult
- Note that other processes give an idea of the error in $0\nu\beta\beta$, its higher than it seems

Tiny Rate: Let's Check Out This Abandoned Mine Scooby

NME is a theoretical issue, so why can't we measure it and figure out the rest later? $2\nu\beta\beta$ and consequently $0\nu\beta\beta$ have very low probabilities of occurring, leading to very low rates.

One crucial requirement: low background rate



Energy Resolution: Squint, Alot

There are two broad classes of experiment

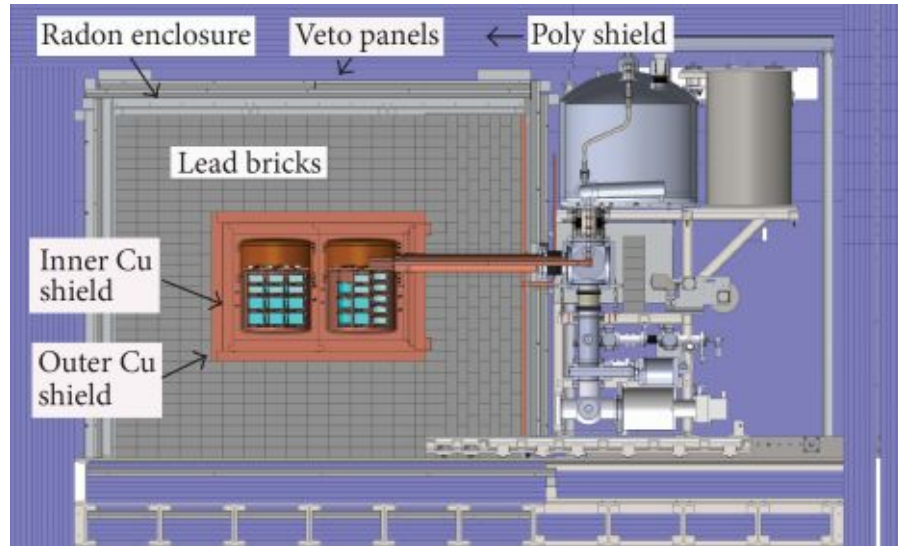
Detector + Source

- Fully reconstructible event topology enabling background reduction (good)
- Low energy resolution and intrinsic limits on the size of the source material (bad)

Detector = Source

- Good energy resolution (good)
- Scalable (good)
- Not generally able to reconstruct the event topology (bad)
- Limited in choices of material (bad)

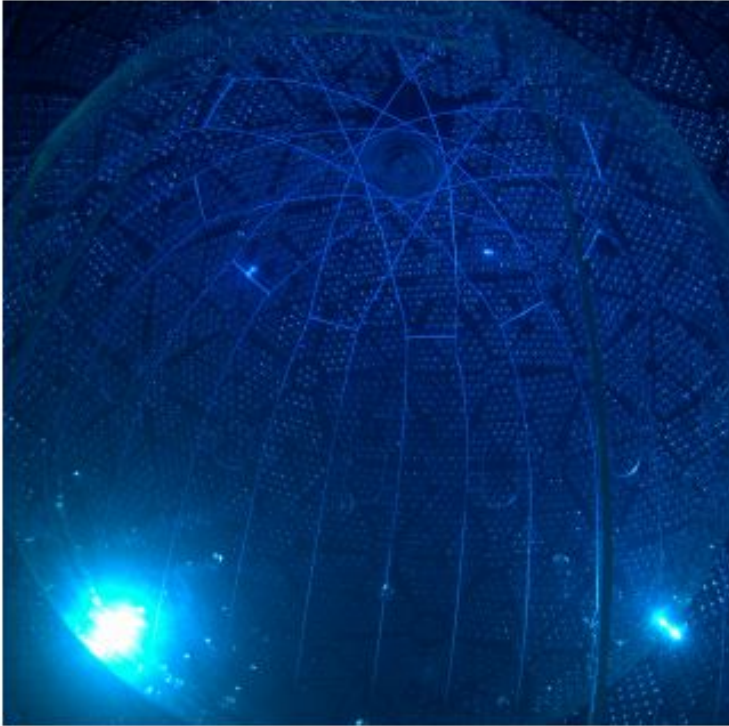
Experiment Example: Majorana Demonstrator



- Ge 76 (enriched and natural)
- Calorimetric approach (source coincides with detector)
- Competitive predictive limits for Ge with full lifetime
- Ge Diode detection principle

arXiv:1711.11145

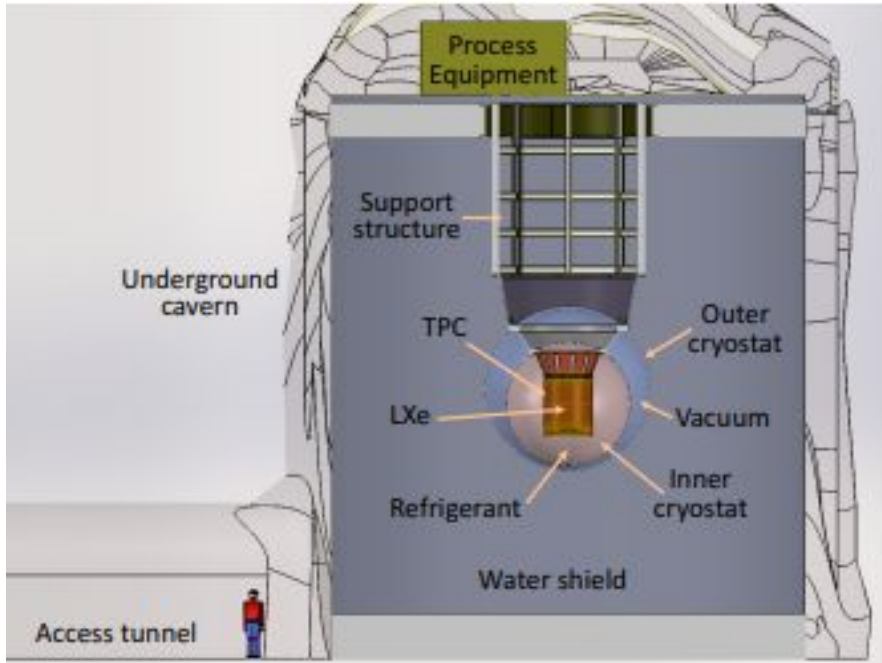
Experiment Example: SNO+



- Te based experiment
- The isotope is mixed with water (liquid scintillator)
- Read out by pmts
- Water veto on the outside

arXiv:1711.11094

Experiment Example: (n)EXO



- Liquid Xenon based detector
- Uses both scintillation light and a TPC to detect events
- Full event topology is reconstructed with good energy resolution
- Able to tag background processes for veto to further suppress backgrounds
- Predicted best limits

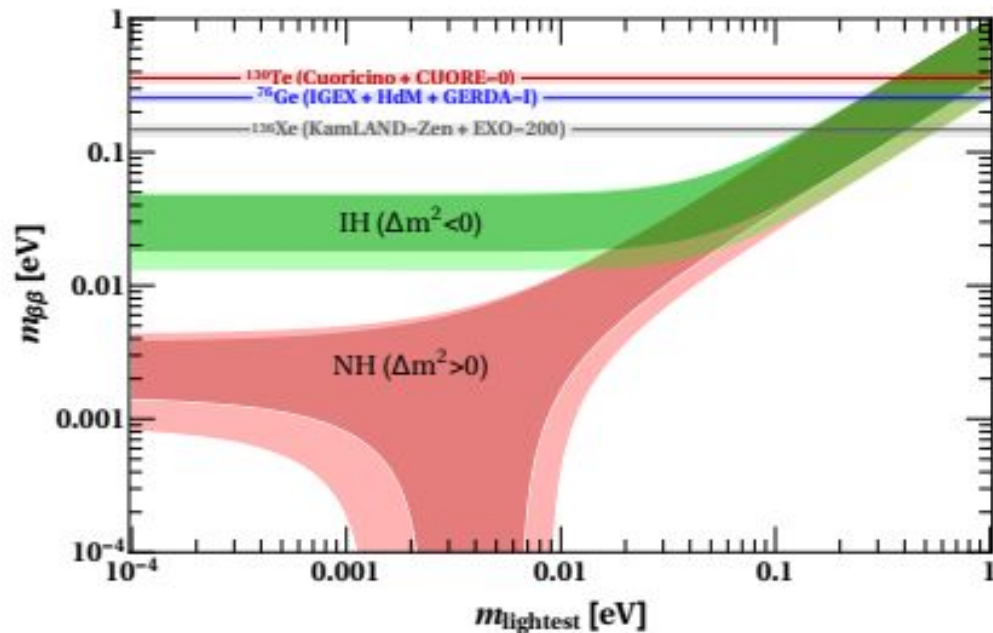
arXiv:1710.05075

Table of Results

Experiment	Isotope	Technique	Total mass [kg]	Exposure [kg yr]	FWHM @ $Q_{\beta\beta}$ [keV]	Background [counts/keV/kg/yr]	$S^{0\nu}$ (90% C.L.) [10^{25} yr]
<i>Past</i>							
Cuoricino, [179]	^{130}Te	bolometers	40.7 (TeO_2)	19.75	5.8 ± 2.1	0.153 ± 0.006	0.24
CUORE-0, [180]	^{130}Te	bolometers	39 (TeO_2)	9.8	5.1 ± 0.3	0.058 ± 0.006	0.29
Heidelberg-Moscow, [181]	^{76}Ge	Ge diodes	11 ($^{\text{enr}}\text{Ge}$)	35.5	4.23 ± 0.14	0.06 ± 0.01	1.9
IGEX, [182, 183]	^{76}Ge	Ge diodes	8.1 ($^{\text{enr}}\text{Ge}$)	8.9	~ 4	$\lesssim 0.06$	1.57
GERDA-I, [167, 184]	^{76}Ge	Ge diodes	17.7 ($^{\text{enr}}\text{Ge}$)	21.64	3.2 ± 0.2	~ 0.01	2.1
NEMO-3, [185]	^{100}Mo	tracker + calorimeter	6.9 (^{100}Mo)	34.7	350	0.013	0.11
<i>Present</i>							
EXO-200, [186]	^{136}Xe	LXe TPC	175 ($^{\text{enr}}\text{Xe}$)	100	89 ± 3	$(1.7 \pm 0.2) \cdot 10^{-3}$	1.1
KamLAND-Zen, [187, 188]	^{136}Xe	loaded liquid scintillator	348 ($^{\text{enr}}\text{Xe}$)	89.5	244 ± 11	~ 0.01	1.9
<i>Future</i>							
CUORE, [189]	^{130}Te	bolometers	741 (TeO_2)	1030	5	0.01	9.5
GERDA-II, [174]	^{76}Ge	Ge diodes	37.8 ($^{\text{enr}}\text{Ge}$)	100	3	0.001	15
LUCIFER, [190]	^{82}Se	bolometers	17 (Zn^{82}Se)	18	10	0.001	1.8
MAJORANA D., [191]	^{76}Ge	Ge diodes	44.8 ($^{\text{enr}/\text{nat}}\text{Ge}$)	100 ^a	4	0.003	12
NEXT, [192, 193]	^{136}Xe	Xe TPC	100 ($^{\text{enr}}\text{Xe}$)	300	12.3 – 17.2	$5 \cdot 10^{-4}$	5
AMoRE, [194]	^{100}Mo	bolometers	200 ($\text{Ca}^{\text{enr}}\text{MoO}_4$)	295	9	$1 \cdot 10^{-4}$	5
nEXO, [195]	^{136}Xe	LXe TPC	4780 ($^{\text{enr}}\text{Xe}$)	12150 ^b	58	$1.7 \cdot 10^{-5}$ ^b	66
PandaX-III, [196]	^{136}Xe	Xe TPC	1000 ($^{\text{enr}}\text{Xe}$)	3000 ^c	12 – 76	0.001	11 ^c
SNO+, [197]	^{130}Te	loaded liquid scintillator	2340 ($^{\text{nat}}\text{Te}$)	3980	270	$2 \cdot 10^{-4}$	9
SuperNEMO, [198, 199]	^{82}Se	tracker + calorimeter	100 (^{82}Se)	500	120	0.01	10

Current Limits

- Common to consider intersection of oscillations and $0\nu\beta\beta$
- Different $0\nu\beta\beta$ limits lead to different implications for mass splitting scenarios
- Cosmological constraints and different models, also play a role in the full physics of this search



Questions

Reference

1. S. Dell'Oro, S. Marcocci, M. Viel, F. Vissani. Neutrinoless double beta decay: 2015 review. *Advances in High Energy Physics*, Volume 2016 (2016).