### Double differential factorization in $\tau_0$ and $\tau_1$

### Laís Schunk Work in progress with Goutam Das and Frank Tackmann

DESY

SCET 25 March 2019



Laís Schunk Work in progress with Goutam I

Factorization in  $au_0$  and  $au_1$ 

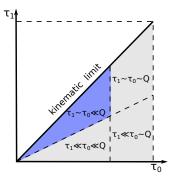
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## Introduction

- We want the double differential factorization for  $\tau_0$  and  $\tau_1 \rightarrow$  measures 2 consecutive resolved emissions
- Some kinematic regions are already known
  - $\tau_1 \ll \tau_0 \ll Q$  : known with SCET+ approach  $\rightarrow$  strongly ordered emissions Pietrulewicz, Tackmann , Waalewijn 2016
  - $au_1 \ll au_0 \sim Q$  : only  $au_1$  resummation

 $\rightarrow$  one hard emission

•  $\tau_1 \sim \tau_0 \sim Q$  : purely fixed-order  $\rightarrow$  two hard emissions

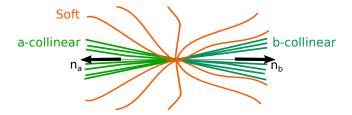


•  $\tau_1 \sim \tau_0 \ll Q$  is the missing piece in the  $\tau_0 \times \tau_1$  space  $\rightarrow$  region with two unordered (but resolved) emissions

- Prototype for factorization beyond strongly ordered regime
  - $\rightarrow$  using N-jettiness as a "simple" resolution variable
  - $\rightarrow$  treatment of unordered emissions is currently unknown
  - $\rightarrow$  relevant to increase accuracy of Parton Showers
- Immediate application: Geneva framework at NNLL' + NNLO  $\rightarrow$  Needs full 0/1-jet and 1/2-jet separation
- Jet substructure : understanding cases with multiple observables with no hierarchy between them  $\rightarrow$  Example : C- and D-parameters in the region  $D \sim C^2 \ll 1$ Larkoski, Procita 2019 – See Larkoski's talk later today  $\rightarrow$ Finite N-subjettiness ratios  $\tau_{21} = \tau_2/\tau_1$  and  $\tau_{32} = \tau_3/\tau_2$ Napoletano, Soyez 2018

## Notation and Setup

• Because  $au_0 \ll Q 
ightarrow$  power counting is defined by  $au_0$  only



• Definition of observables here

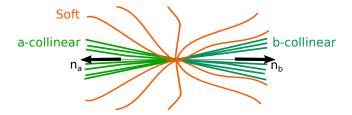
$$\tau_0 = \frac{1}{Q} \sum_{i \in \text{event}} \min\left(n_a \cdot p_i, n_b \cdot p_i\right), \quad \tau_1 = \frac{1}{Q} \sum_{i \in \text{event}} \min\left(n_1 \cdot p_i, n_a \cdot p_i, n_b \cdot p_i\right)$$

• Decompose N-jet as

$$\tau_{N} = \tau_{s,N} + \tau_{a,N} + \tau_{b,N}$$

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$$\tau_{N} = \tau_{s,N} + \tau_{a,N} + \tau_{b,N}$$

#### How to write $\tau_1$ with these modes?

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Factorization in  $au_0$  and  $au_1$ 

## "Inspiration" : Fully Recursive N-jettiness

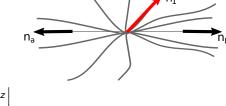
- Fully recursive  $\tau_0^{FR}$  is given by
  - Step 1 Compute  $\tau_1$ , with the resp. axis  $n_1$
  - Step 2 Emissions divided in subgroups – depends on which axis the have "chosen"

Step 3 Add contribution

$$\Delta^{\mathsf{FR}} = \left|\sum \vec{p}\right| - \left|\sum p^{z}\right|$$

 $\bullet~$  Corollary: we can always write  $\tau_0^{FR}=\tau_1+\Delta^{FR}$ 

Definition of  $\tau_0$  used for *pp* collision in GENEVA (Alioli and al. 2015)



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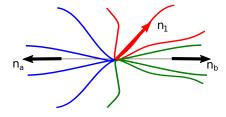
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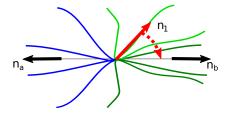
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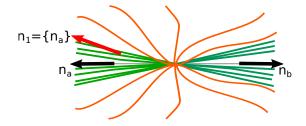
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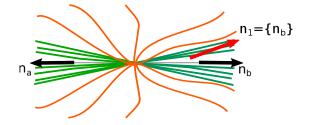
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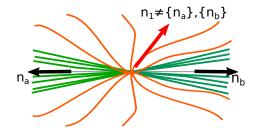
- We select the axis  $n_1$  that minimizes the 1-jettiness
- ullet Possibilities for the axis choice are  $\rightarrow$  in the a-collinear sector



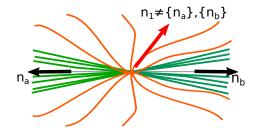
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- We select the axis  $n_1$  that minimizes the 1-jettiness
- Possibilities for the axis choice are ightarrow none of the above



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• From the power-counting point of view

$$\{n_a\} \equiv \left(1, \frac{\vec{p_a}}{|\vec{p_a}|}\right), \qquad p_a \sim (p_a^+, p_a^-, p_{a\perp}) \sim Q^2(\lambda^2, 1, \lambda)$$

• Decompose axis minimization as

$$\tau_{1} = \min\left[\min_{n_{1}=\{n_{a}\}}\left(\sum_{i\in event}\min\left(n_{a}\cdot p_{i}, n_{b}\cdot p_{i}, n_{1}\cdot p_{i}\right)\right),$$
$$\min_{n_{1}=\{n_{b}\}}\left(\sum_{i\in event}\min\left(n_{a}\cdot p_{i}, n_{b}\cdot p_{i}, n_{1}\cdot p_{i}\right)\right),$$
$$\min_{n_{1}\neq\{n_{a}\},\{n_{b}\}}\left(\sum_{i\in event}\min\left(n_{a}\cdot p_{i}, n_{b}\cdot p_{i}, n_{1}\cdot p_{i}\right)\right)\right]$$

 $\bullet$  Divide and conquer  $\rightarrow$  prove factorization for each piece

• Decompose axis minimization as

$$\begin{aligned} \tau_1 &= \min\left[\min_{n_1 \in \{n_a\}} \left(\sum_{i \in event} \min\left(n_a \cdot p_i, n_b \cdot p_i, n_1 \cdot p_i\right)\right), \\ &\min_{n_1 \in \{n_b\}} \left(\sum_{i \in event} \min\left(n_a \cdot p_i, n_b \cdot p_i, n_1 \cdot p_i\right)\right), \\ &\min_{n_1 \neq \{n_a\}, \{n_b\}} \left(\sum_{i \in event} \min\left(n_a \cdot p_i, n_b \cdot p_i, n_1 \cdot p_i\right)\right)\right] \end{aligned}$$

 $\bullet$  Divide and conquer  $\rightarrow$  prove factorization for each piece

Start with  $n_1 = \{n_a\}$  case

- **→ →** →

## Factorization for $n_1 = \{n_a\}$

• Reminder : we can write  $\tau_N = \tau_{s,N} + \tau_{a,N} + \tau_{b,N}$ 

For a-collinear piece

$$\tau_{1,a} = \min\left(n_a \cdot b_a, n_b \cdot b_a, n_1 \cdot b_a\right) = b_a^+ - \delta \tau_a$$

Por b-collinear piece

$$\tau_{1,b} = \min\left(n_a \cdot b_b, n_b \cdot b_b, n_1 \cdot b_b\right) = b_b^+$$

So For soft piece:  

$$\{n_a\} = n_a + \lambda^2 \vec{n} \rightarrow n_1 \cdot k_s = n_a \cdot k_s (1 + \mathcal{O}(\lambda^2))$$

$$\tau_{1,s} = \min(n_a \cdot k_s, n_b \cdot k_s, n_1 \cdot k_s) = k_s$$

If  $n_1 \in \{n_a\}$ , then  $\tau_1 = \tau_0 - \delta \tau_a$ Only emissions in the a-collinear sector matter

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Factorization in  $au_0$  and  $au_1$ 

# Factorization for $n_1 \neq \{n_a\}, \{n_b\}$

Suppose n<sub>1</sub> ≠ {n<sub>a</sub>}, {n<sub>b</sub>}
Collinear pieces → same reasoning as before

$$\tau_{1,a} = \boldsymbol{b}_a^+, \qquad \tau_{1,b} = \boldsymbol{b}_b^+$$

Soft emissions

$$\tau_{1,s} = \min\left(n_a \cdot k_s, n_b \cdot k_s, n_1 \cdot k_s\right) = k_s - \delta \tau_s$$

# Factorization for $n_1 \neq \{n_a\}, \{n_b\}$

Suppose n<sub>1</sub> ≠ {n<sub>a</sub>}, {n<sub>b</sub>}
Collinear pieces → same reasoning as before

$$\tau_{1,a} = b_a^+, \qquad \tau_{1,b} = b_b^+$$

Soft emissions

$$\tau_{1,s} = \min\left(n_a \cdot k_s, n_b \cdot k_s, n_1 \cdot k_s\right) = k_s - \delta \tau_s$$

$$\begin{array}{ll} \text{If } n_1 = \{n_a\}, & \text{then } \tau_1 = \tau_0 - \delta \tau_a \\ \text{If } n_1 = \{n_b\}, & \text{then } \tau_1 = \tau_0 - \delta \tau_b \\ \text{If } n_1 \neq \{n_a\}, \{n_b\}, & \text{then } \tau_1 = \tau_0 - \delta \tau_s \end{array}$$

Axis in a sector implies  $\delta \tau_i$  in the same sector

$$\tau_{1} = \min \left[ b_{b}^{+} + k_{s} + \min_{n_{1} = \{n_{a}\}} \left( b_{a}^{+} - \delta \tau_{a} \right), \\ b_{a}^{+} + k_{s} + \min_{n_{1} = \{n_{b}\}} \left( b_{b}^{+} - \delta \tau_{b} \right), \\ b_{a}^{+} + b_{b}^{+} + \min_{n_{1} \neq \{n_{a}\}, \{n_{b}\}} \left( k_{s} - \delta \tau_{s} \right) \right]$$

$$\tau_{1} = \min \left[ \tau_{0} + \min_{n_{1} = \{n_{a}\}} \left( -\delta \tau_{a} \right), \\ \tau_{0} + \min_{n_{1} = \{n_{b}\}} \left( -\delta \tau_{b} \right), \\ \tau_{0} + \min_{n_{1} \neq \{n_{a}\}, \{n_{b}\}} \left( -\delta \tau_{s} \right) \right]$$

$$\tau_{1} = \tau_{0} - \max\left[\max_{n_{1} = \{n_{a}\}} (\delta\tau_{a}), \\ \max_{n_{1} = \{n_{b}\}} (\delta\tau_{b}), \\ \max_{n_{1} \neq \{n_{a}\}, \{n_{b}\}} (\delta\tau_{s})\right]$$

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$$\tau_1 = \tau_0 - \max\left[\Delta_a, \Delta_b, \Delta_s\right]$$

Complete separation of contributions from different modes

$$\tau_1 = \tau_0 - \max\left[\Delta_a, \Delta_b, \Delta_s\right]$$

#### Complete separation of contributions from different modes

- The rest of factorization is similar to pure  $\tau_0$  case
- Need new Beam and Soft functions that depend on  $\Delta_i$  measurements

$$\frac{d^2\sigma}{d\tau_0 d\tau_1} = \sigma_0 \int db_a^+ db_b^+ d\Delta_a d\Delta_b d\Delta_s B_i^{(a)}(b_a^+, \Delta_a, x_a) B_j^{(b)}(b_b^+, \Delta_b, x_b) \\ \times \frac{S(\tau_0 - b_a^+ - b_b^+, \Delta_s)}{\delta(\tau_1 - \tau_0 + \max(\Delta_i))} \delta(\tau_1 - \tau_0 + \max(\Delta_i))$$

### Features

$$\frac{d^2\sigma}{d\tau_0 d\tau_1} = \sigma_0 \int db_a^+ db_b^+ d\Delta_a d\Delta_b d\Delta_s B_i^{(a)}(b_a^+, \Delta_a, x_a) B_j^{(b)}(b_b^+, \Delta_b, x_b) \\ \times \frac{S(\tau_0 - b_a^+ - b_b^+, \Delta_s)}{\delta(\tau_1 - \tau_0 + \max(\Delta_i))} \delta(\tau_1 - \tau_0 + \max(\Delta_i))$$

• We must recover  $\tau_0$  SCET factorization

$$\int d\tau_1 \frac{d^2\sigma}{d\tau_0 d\tau_1} = \frac{d\sigma}{d\tau_0}$$

- Integration over  $\delta( au_1 \dots) 
  ightarrow$  integration over  $\Delta_i$  measurements
- Implies condition on Beam and Soft functions

$$\int d\Delta_a B_i^{(a)}(b_a^+, \Delta_a, x_a) = B_i^{(a)}(b_a^+, x_a)$$
$$\int d\Delta_s S(k_s, \Delta_s) = S(k_s)$$

• At NLO  $\mathcal{O}(\alpha_s)$ , we know that there is no 1-jettiness contribution

$$\frac{d^2\sigma^{NLO}}{d\tau_0 d\tau_1} = \frac{d\sigma}{d\tau_0}\delta(\tau_1)$$

• The 1-loop Beam and Soft functions are given by

$$B_i^{1-\text{loop}}(b_a^+, x_a, \Delta_a) = B_i^{1-\text{loop}}(b_a^+, x_a)\delta(\Delta_a - b_a)$$
  
$$S^{1-\text{loop}}(k_s, \Delta_s) = S^{1-\text{loop}}(k_s)\delta(\Delta_s - k_s)$$

• Auxiliary measurement  $\delta(\tau_1)$  does not change structure of RGEs  $\rightarrow$  resummation up to NLL' is unchanged

### Features

- First unkown order is at  $\mathcal{O}(\alpha_s^2) \to$  needs 2-loop functions Work in progress
- Cross-terms are known

$$\begin{aligned} \frac{d\sigma}{d\tau_{0}d\tau_{1}} &= \sigma_{0}\int db_{a}^{+} db_{b}^{+} dk_{s} \ B_{i}^{1-loop}(b_{a}^{+}, x_{a})B_{i}^{tree}(b_{b}^{+}, x_{b}) \\ &\times \ S^{1-loop}(k_{s})\delta(\tau_{0} - b_{a}^{+} - b_{b}^{+} - k_{s})\delta(\tau_{1} - \min(b_{a}^{+}, k_{s})) \\ &= \sigma_{0}\frac{f_{j}(x_{b})}{x_{b}E_{CM}^{2}}\Theta\left(\frac{\tau_{0}}{2} - \tau_{1}\right) \\ &\times \ \left[B_{i}^{1-loop}(\tau_{1}, x_{a})S^{1-loop}(\tau_{0} - \tau_{1}) + B_{i}^{1-loop}(\tau_{0} - \tau_{1}, x_{a})S^{1-loop}(\tau_{1})\right] \end{aligned}$$

• We recover the 
$$\tau_1 < \tau_0/2$$
 limit  
 $\rightarrow$  intrinsic property of the  $\tau_1 \sim \tau_0$  regime

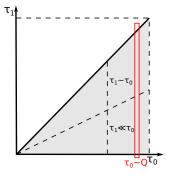
## Matching : reminder for large- $\tau_0$

 $au_1 \ll au_0 \sim Q \rightarrow au_1 \sim au_0 \sim Q$ Match  $au_1$  factorization with fixed-order calculation

• For a given  $\tau_0$ :

$$d\sigma^{match} = d\sigma^{(1)} + \left( d\sigma^{FO} - \left[ d\sigma^{(1)} \right]_{FO} \right)$$

• With  $[d\sigma^{(1)}]_{FO}$  such that •  $[d\sigma^{(1)}]_{FO} \rightarrow d\sigma^{(1)}$  when  $\tau_1 \sim \tau_0$ •  $[d\sigma^{(1)}]_{FO} \rightarrow d\sigma^{FO}$  when  $\tau_1 \ll \tau_0$ • Set scales to turn off resummation



 $\mu_{(1)} \rightarrow \mu_{FO}$ 

# Matching with SCET+

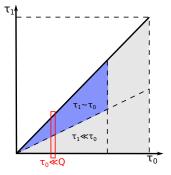
#### $au_1 \ll au_0 \ll Q \to au_1 \sim au_0 \ll Q$ Match our factorization with existing SCET+ factorization

• For a given  $\tau_0$ :

1

$$d\sigma^{match} = d\sigma^+ + \left( d\sigma^{(0,1)} - \left[ d\sigma^+ 
ight]_{"FO"} 
ight)$$

- "FO" has to include resummation in  $\tau_{\rm 0}$
- With  $[d\sigma^+]_{"FO"}$  such that
  - $[d\sigma^+]_{"FO"} \rightarrow d\sigma^+$  when  $\tau_1 \sim \tau_0$
  - $[d\sigma^+]_{"FO"}^{"} \rightarrow d\sigma^{(0,1)}$  when  $\tau_1 \ll \tau_0$
- Set scales to "turn off"  $\log(\tau_1/\tau_0)$  resummation  $\mu_+ \to \mu_{(0,1)}$

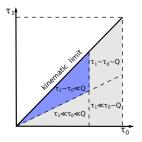


### Work in progress

## Conclusion

#### Recap

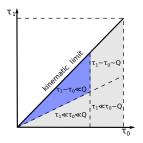
- Presented the factorization of the double differential  $\frac{d\sigma}{d\tau_0 d\tau_1}$  in the region  $\tau_0 \sim \tau_1$
- Completes the description of the two resolved emissions phase space
   → including unordered emissions



## Conclusion

#### Recap

- Presented the factorization of the double differential  $\frac{d\sigma}{d\tau_0 d\tau_1}$  in the region  $\tau_0 \sim \tau_1$
- Completes the description of the two resolved emissions phase space
   → including unordered emissions



#### Future

- Fully understand details of matching to SCET+
- Calculate soft  $S(k_s, \Delta_s)$  and beam  $B(x_i, b_i, \Delta_i)$  functions at  $\mathcal{O}(\alpha_s^2)$