

# **Beam-beam effects in plasma acceleration-based linear collider**

**Carl B. Schroeder**

Lawrence Berkeley National Laboratory

**ICFA Mini-Workshop on Beam-Beam Effects  
in Circular Colliders**

LBL, Berkeley, CA

6 Feb 2018

Work supported by the U.S. DOE, Office of Science, Office of High Energy Physics,  
under Contract No. DE-AC02-05CH11231

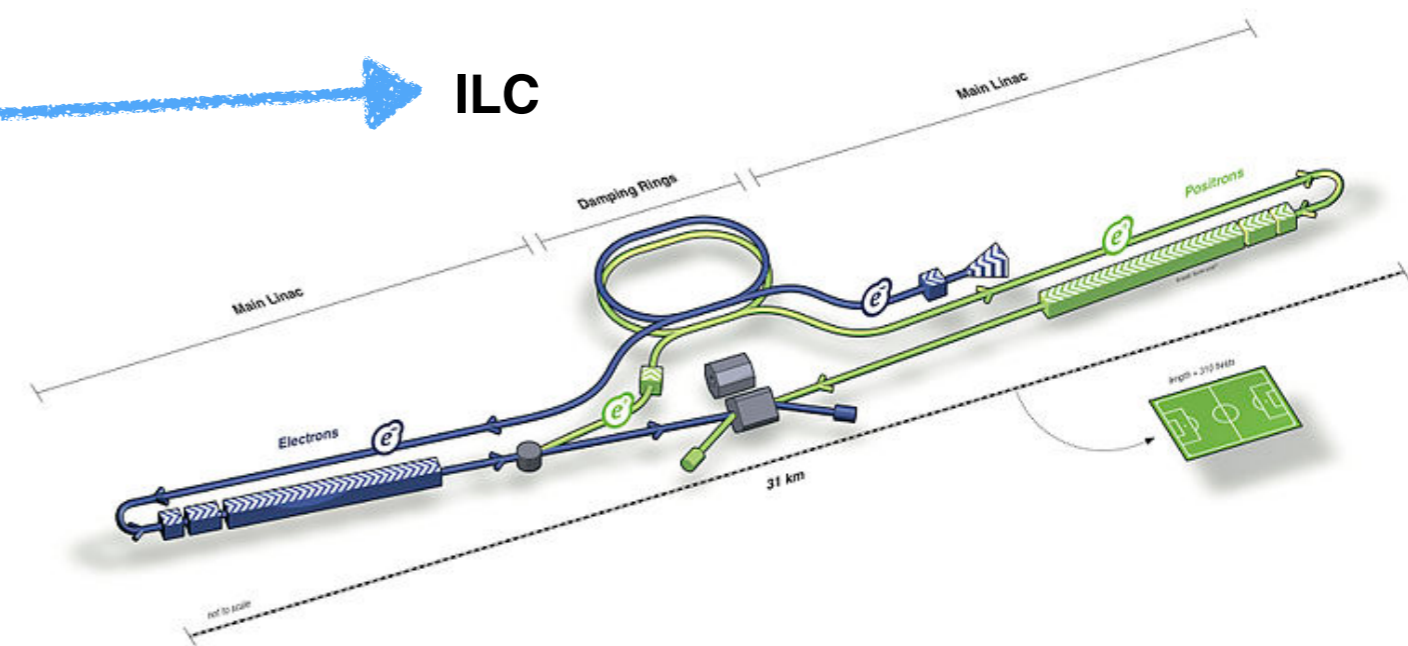
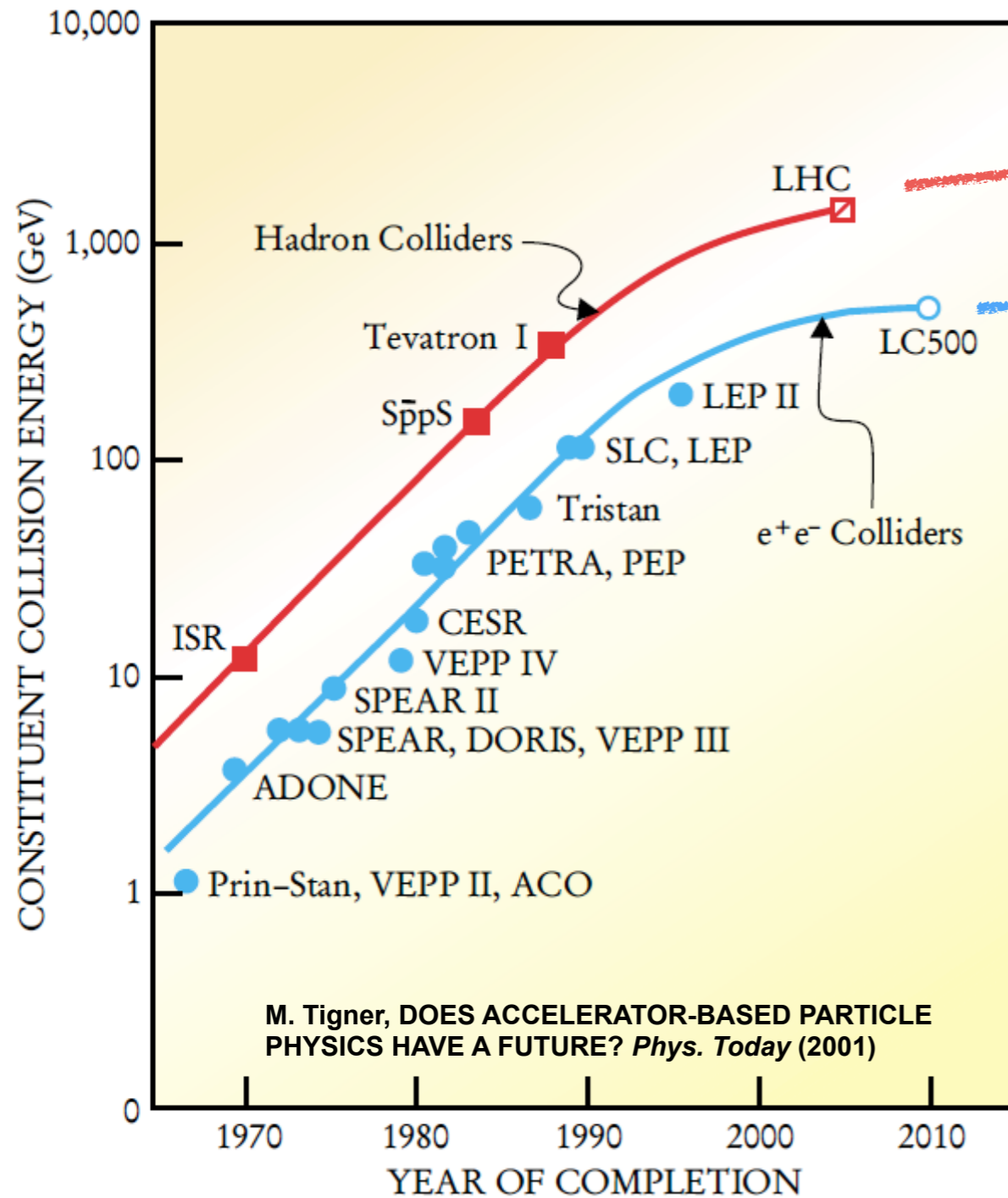
# Outline

- ▶ **Laser-driven plasma-based accelerators (LPAs):** ultra-high gradient (tens of GV/m) provides compact source of energetic, fs, kA electron beams
  - GeV electron beams in cm-scale plasmas using LPAs
- ▶ **LPA-based linear collider design considerations**
  - operational plasma density
  - beam distribution
  - LPA staging toward high energy
- ▶ **Beam-beam effects in LPA-based linear colliders**
  - Beam-beam interaction effects determine the operational plasma density (determine required laser technology)



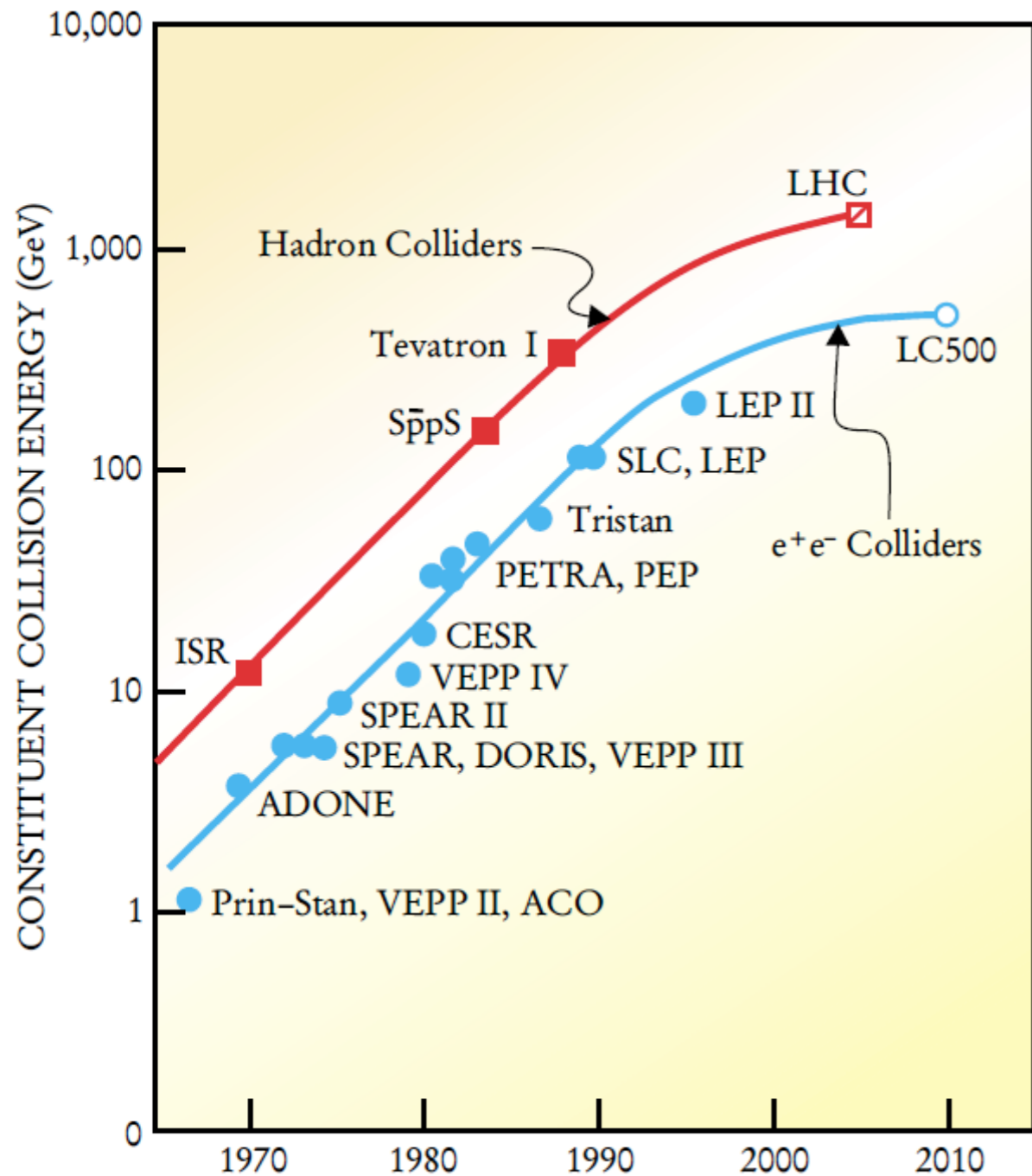
# High-energy physics application of LPAs

- “Livingston Plot” Saturation of accelerator tech.
  - Practical limit reached for conventional accelerator technology (RF metallic structures)
  - Gradient limited by material breakdown
    - e.g., X-band demonstration  $\sim 100$  MV/m



- Largest cost driver is acceleration
  - $\sim 50$  MV/m implies  $\sim 20$  km/TeV
  - Facility costs scale roughly with facility size (and power consumption)
  - $>50\%$  cost in main linacs (e.g., ILC)

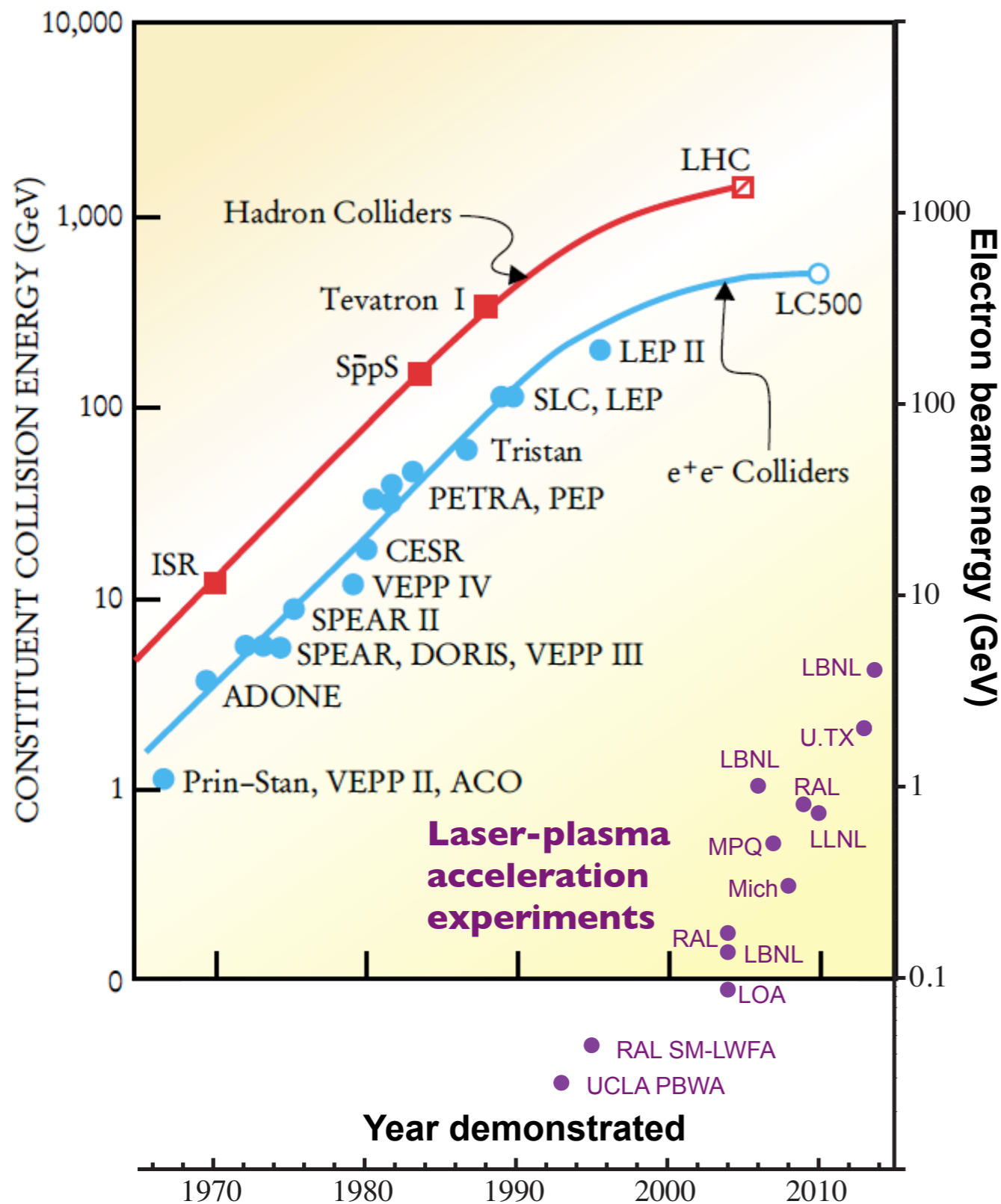
# LPA application: Lepton Collider



- Any future linear TeV ( $> \text{TeV}$ ) collider is a massive (ultra-massive) project
  - require  $>$ order of magnitude increase in acceleration gradient
- Ultra-high gradient requires structures to sustain high fields:
  - Dielectric structures:  $\sim 1 \text{ GV/m}$
  - Plasmas:  $\sim 10 \text{ GV/m}$
- High gradients require high peak power:
  - Beam driven
  - Laser driven
- Significant progress worldwide in LPAs in the last 20+ years
- Critical developments:
  - Better understanding of LPA physics
  - Development of laser technology (CPA) for high peak power delivery



# LPA application: Lepton Collider



- Any future linear TeV (>TeV) collider is a massive (ultra-massive) project
  - require >order of magnitude increase in acceleration gradient
- Ultra-high gradient requires structures to sustain high fields:
  - Dielectric structures: ~1 GV/m
  - Plasmas: ~10 GV/m
- High gradients require high peak power:
  - Beam driven
  - Laser driven
- Significant progress worldwide in LPAs in the last 20+ years
- Critical developments:
  - Better understanding of LPA physics
  - Development of laser technology (CPA) for high peak power delivery

# Laser-plasma accelerators: compact sources (>10 GV/m) of fs e-beams

*Esarey, Schroeder, Leemans, Rev. Mod. Phys. (2009)*

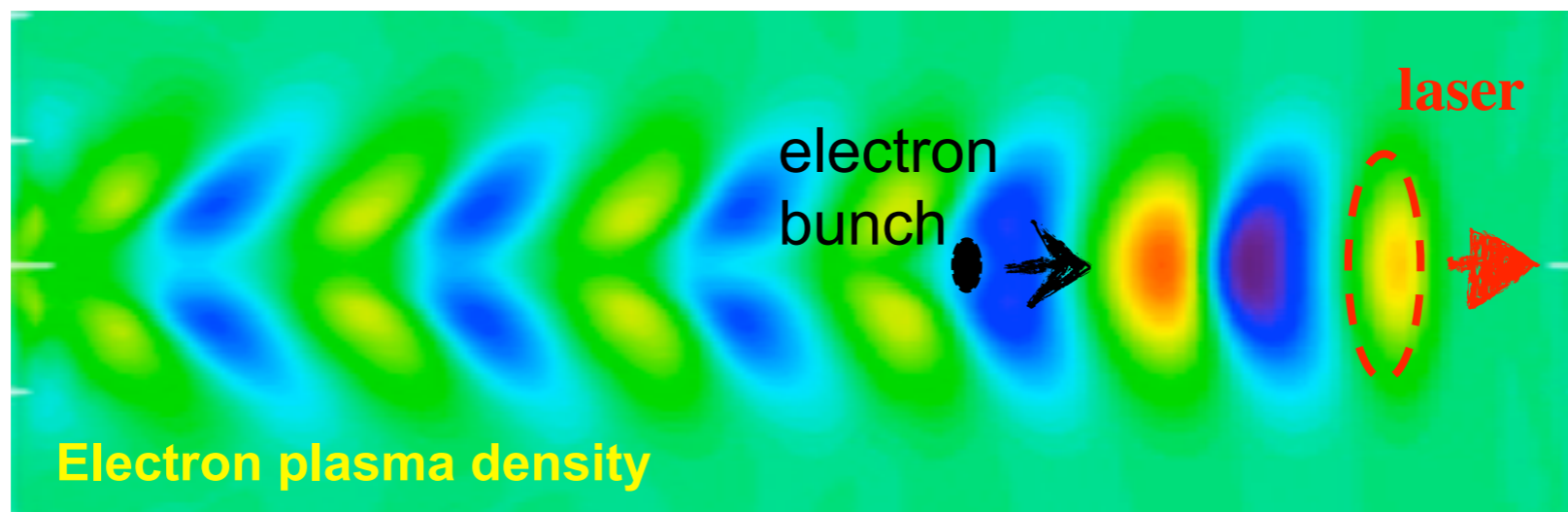
► **Accelerating field:** characteristic size of plasma wave field, driven by the ponderomotive force of a resonant laser pulse with relativistic intensity ( $>10^{18}$  W/cm<sup>2</sup>) propagating in underdense plasma:

$$E \sim \left( \frac{m c \omega_p}{e} \right) \approx (96 \text{ V/m}) \sqrt{n_0 [\text{cm}^{-3}]}$$

- E.g., for  $\sim 10^{18}$  cm<sup>-3</sup>, gradient  $\sim 100$  GV/m

► **Bunch duration:**

- Accelerating bucket  $\sim$  plasma wavelength



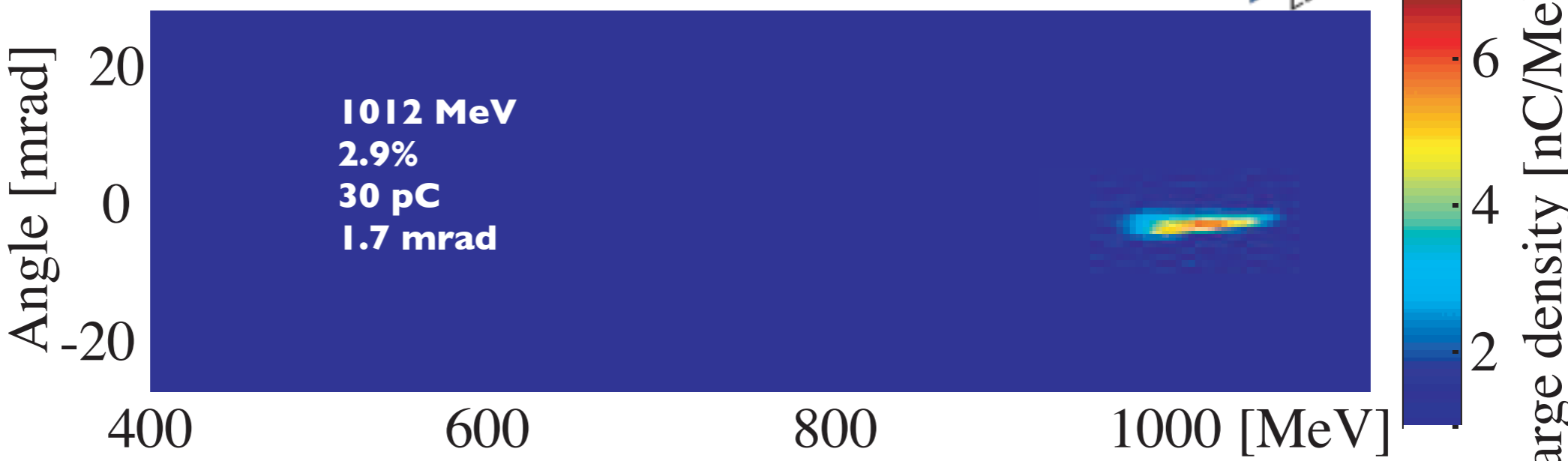
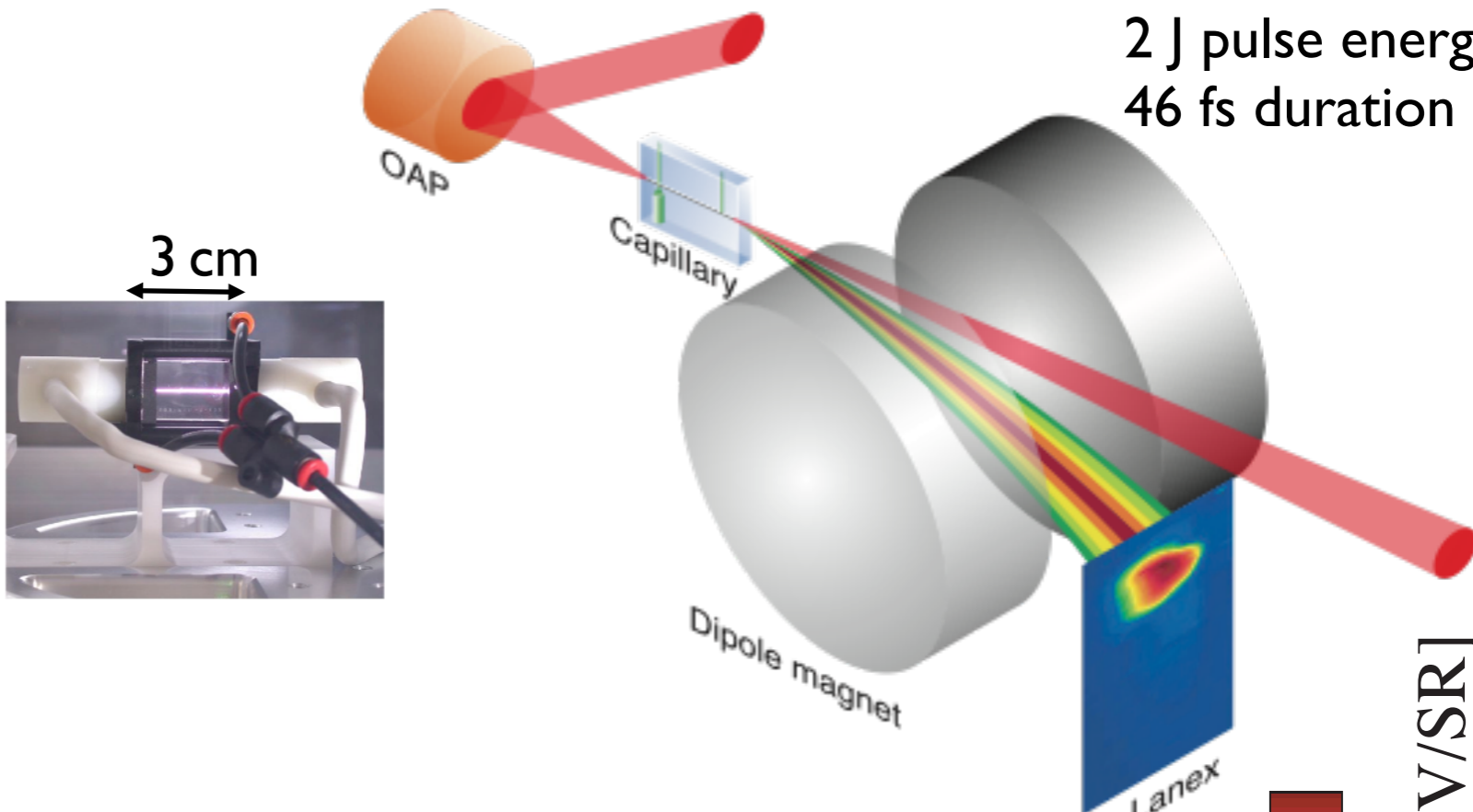
- Beam charge (set by beam loading, plasma density):  $\sim 1$ - $100$  pC
  - Beam duration (set by trapping physics and density):  $\sim 1$ - $10$  fs
- $\sim 1$ - $10$  kA peak current



# Laser-plasma accelerator (LPA) experimental demonstration of GeV electron beam at LBNL

Ti:Sapphire laser:  
0.8 micron wavelength  
2 J pulse energy  
46 fs duration

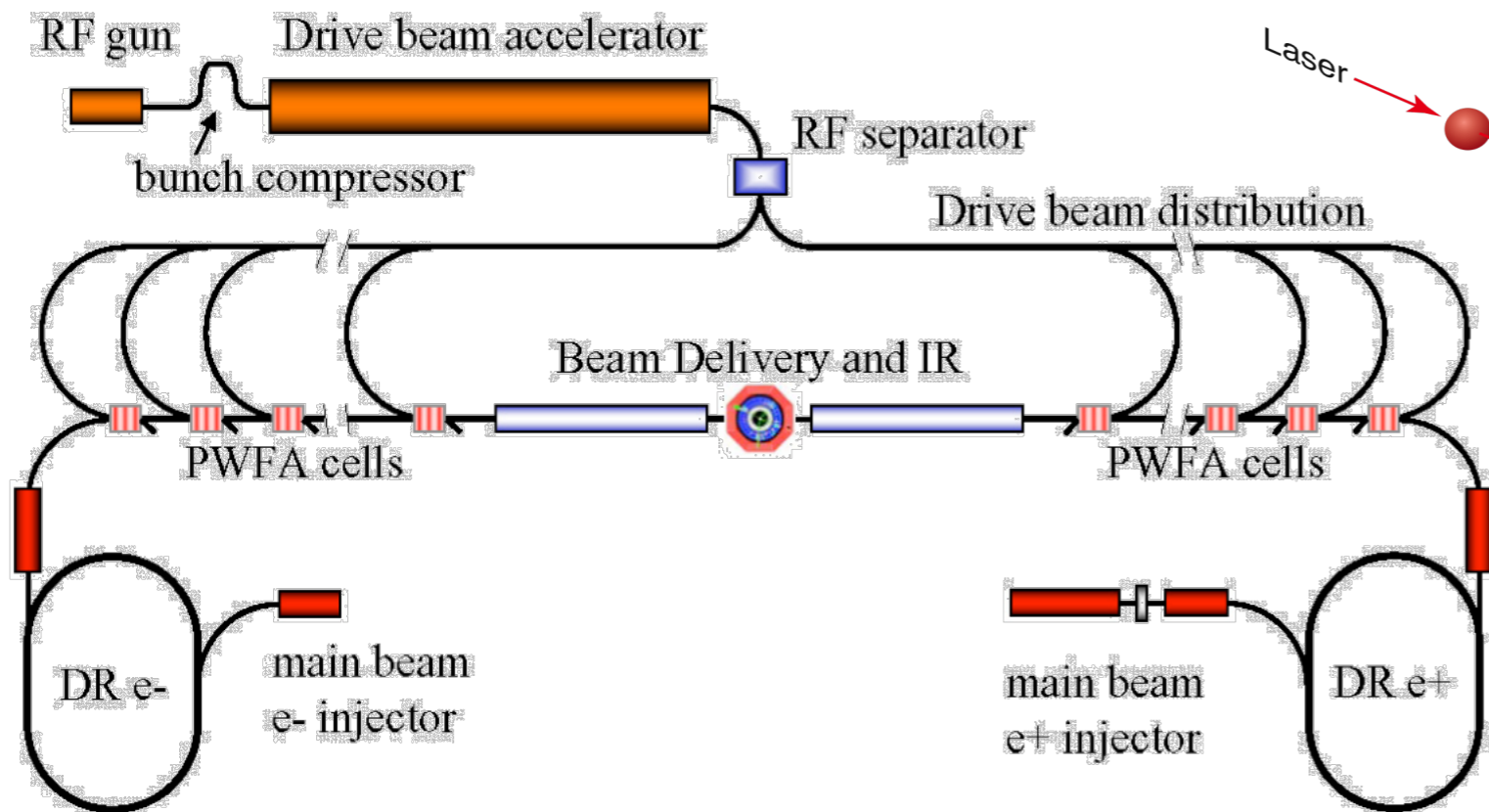
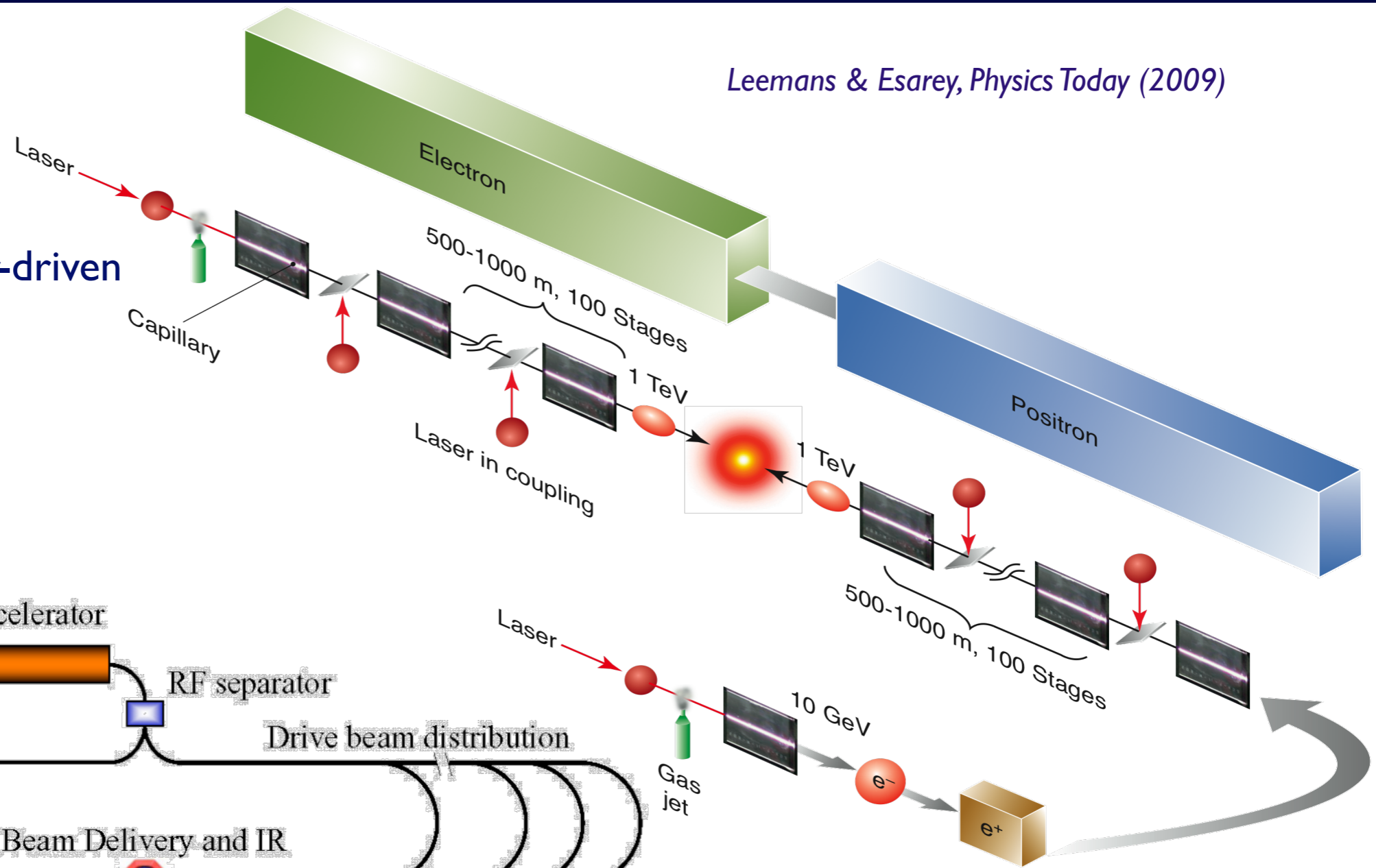
Plasma source:  
H-discharge capillary  
number density =  $3 \times 10^{18} \text{ cm}^{-3}$



Leemans et al., Nature Phys. (2006); Nakamura et al., Phys. Plasmas (2007)

# Plasma-based linear collider concepts

► Collider based on laser-driven plasma acceleration:



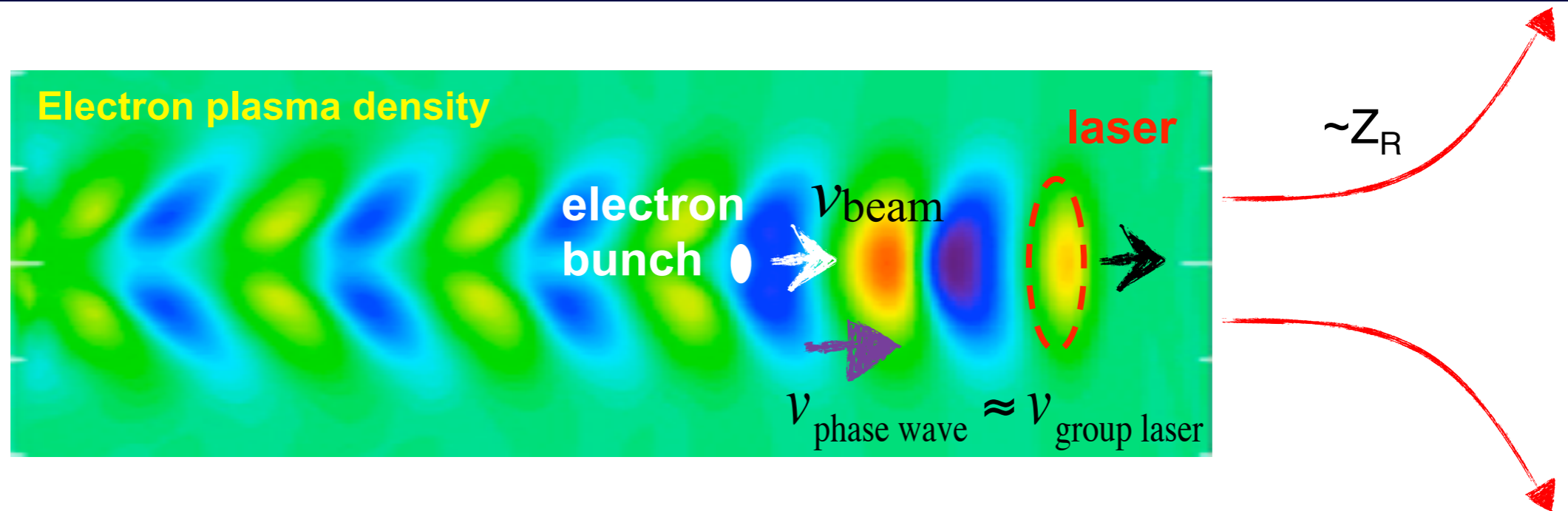
► Collider based on beam-driven plasma acceleration (PWFA):

Figure 1: Concept for a multi-stage PWFA-based Linear Collider.

*Seryi et al., PAC Proc (2009)*



# Limits to energy gain in laser-plasma accelerator (LPA): diffraction, dephasing, depletion



**Limits to single stage energy gain:**  $mc^2 \Delta\gamma \sim q(mc\omega_p/e)L_{\text{int}}$

- **Laser Diffraction:**  $\sim$  Rayleigh range (typically most severe)
  - Controlled by transverse plasma density tailoring (plasma channel) and/or relativistic self-guiding and ponderomotive self-channeling
- **Beam-Wave Dephasing:** Slippage between e-beam and plasma wave
  - Controlled by longitudinal plasma density tailoring (plasma tapering)
- **Laser Energy Depletion:** Rate of laser energy deposition of into plasma wave

$$L_{\text{deplete}} \propto n^{-3/2} \lambda^{-2}$$

# Plasma density scalings

- Laser-plasma interaction (depletion) length:

$$L_{\text{acc}} \sim \lambda_p^3 / \lambda_L^2 \propto n^{-3/2}$$

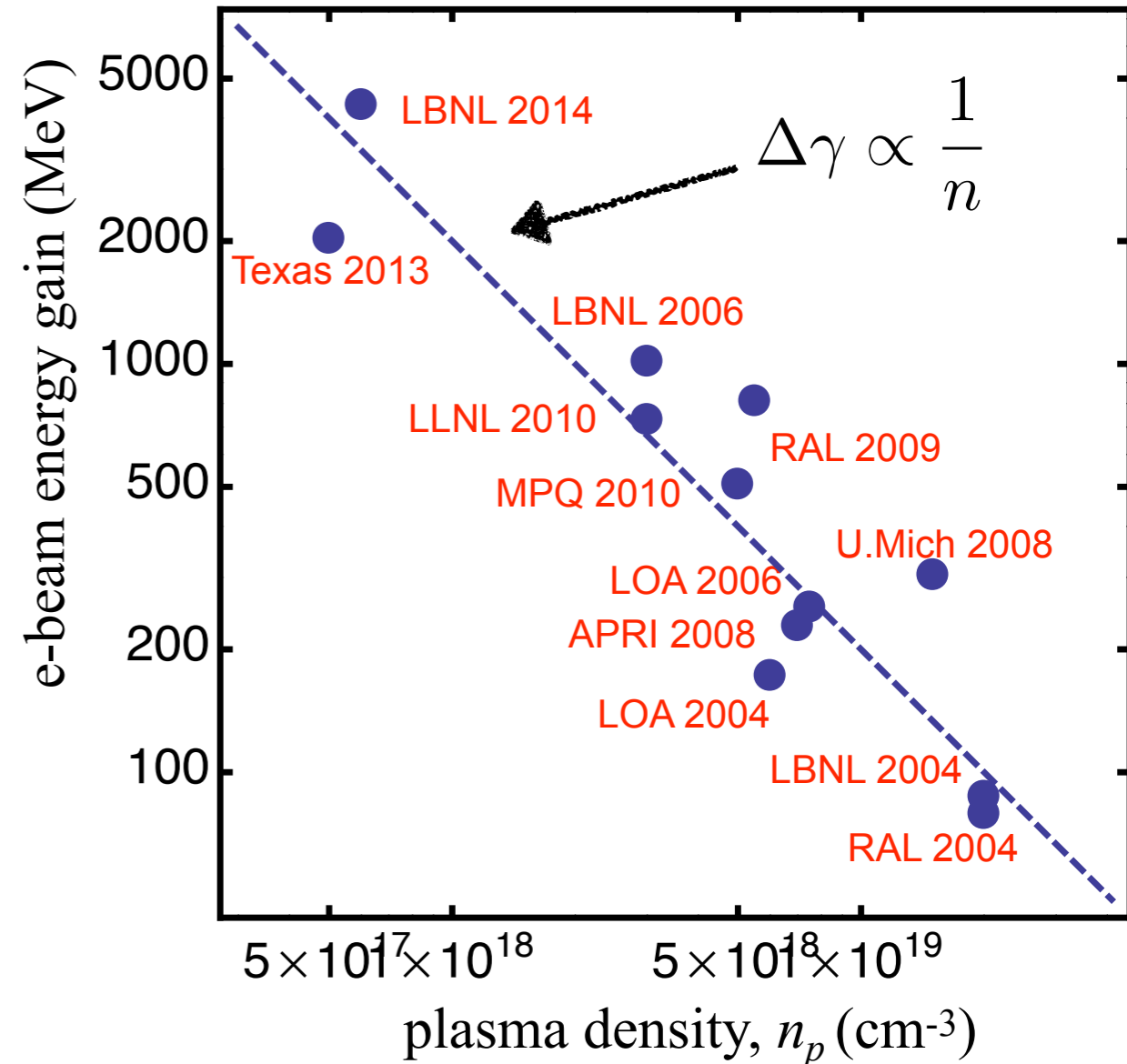
- Accelerating gradient:

$$E \sim E_0 = (m_e c \omega_p / e) \propto \sqrt{n}$$

- Energy gain:

$$W \sim (m c \omega_p / e) L_{\text{acc}} \propto 1/n$$

For high-energy applications, laser depletion (and reasonable gradient) necessitates staging laser-plasma accelerators





# Plasma density scalings

- Laser-plasma interaction (depletion) length:

$$L_{\text{acc}} \sim \lambda_p^3 / \lambda_L^2 \propto n^{-3/2}$$

- Accelerating gradient:

$$E \sim E_0 = (m_e c \omega_p / e) \propto \sqrt{n}$$

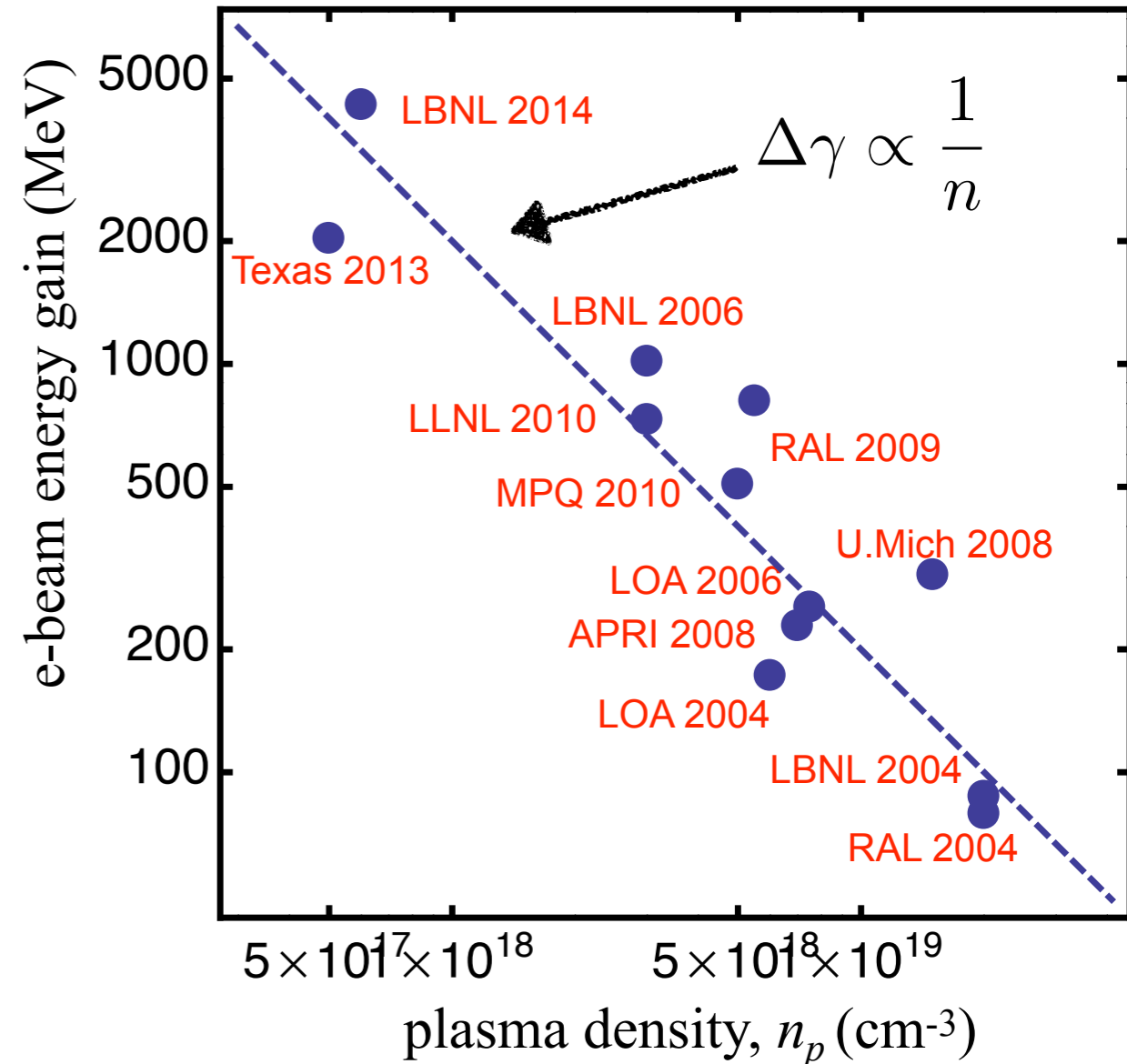
- Energy gain:

$$W \sim (m c \omega_p / e) L_{\text{acc}} \propto 1/n$$

For high-energy applications, laser depletion (and reasonable gradient) necessitates staging laser-plasma accelerators

## Laser requirements:

- Laser energy:  $U_{\text{laser}} \propto \lambda_p^3 \propto n^{-3/2}$
- Laser duration:  $\tau_{\text{laser}} \propto \lambda_p \propto n^{-1/2}$
- Laser peak power:  $P_{\text{laser}} \propto n^{-1}$



# Plasma density scalings

- Laser-plasma interaction (depletion) length:

$$L_{\text{acc}} \sim \lambda_p^3 / \lambda_L^2 \propto n^{-3/2}$$

- Accelerating gradient:

$$E \sim E_0 = (m_e c \omega_p / e) \propto \sqrt{n}$$

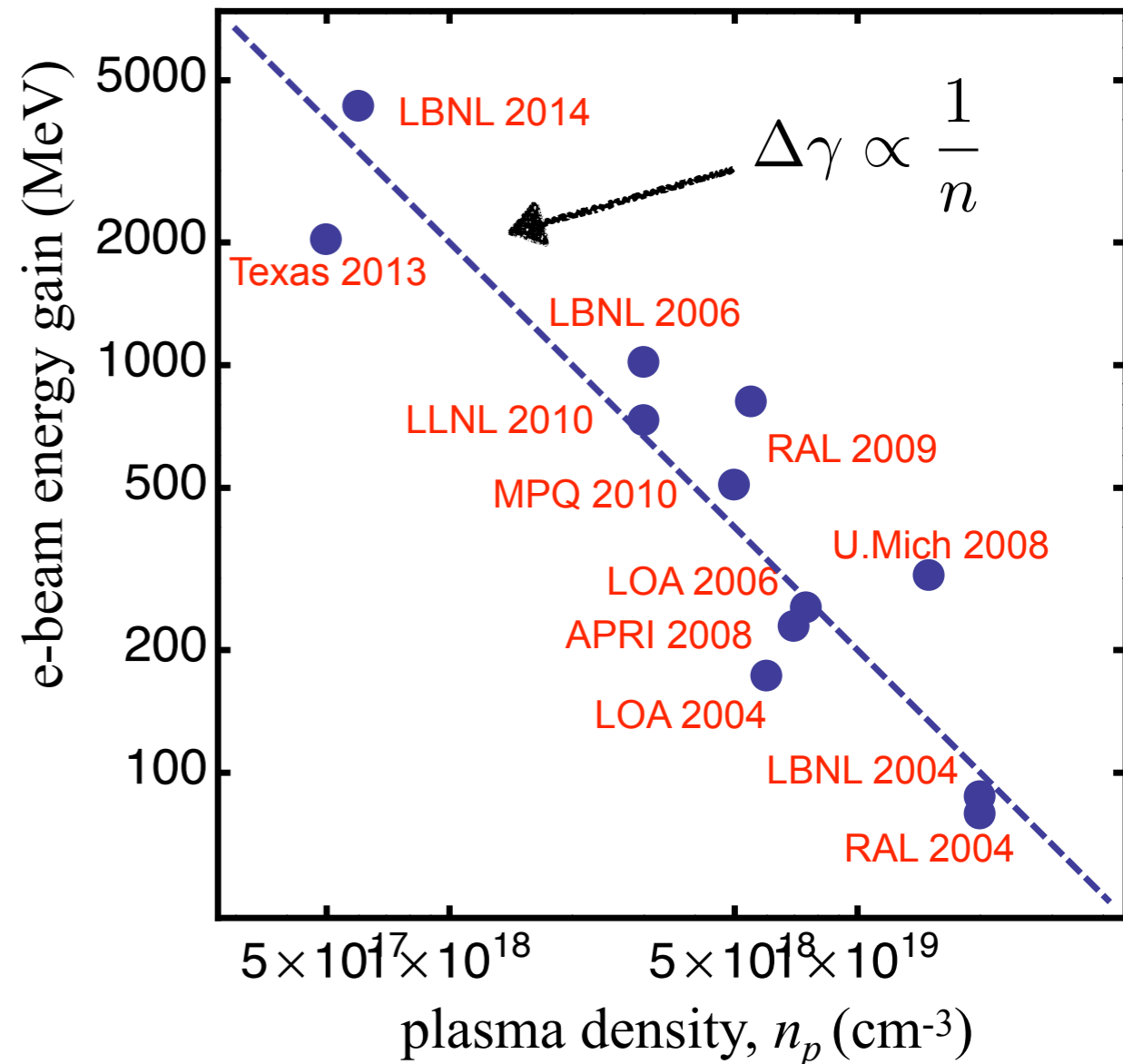
- Energy gain:

$$W \sim (m c \omega_p / e) L_{\text{acc}} \propto 1/n$$

For high-energy applications, laser depletion (and reasonable gradient) necessitates staging laser-plasma accelerators

## Laser requirements:

- Laser energy:  $U_{\text{laser}} \propto \lambda_p^3 \propto n^{-3/2}$
- Laser duration:  $\tau_{\text{laser}} \propto \lambda_p \propto n^{-1/2}$
- Laser peak power:  $P_{\text{laser}} \propto n^{-1}$



**W ~ 1 GeV**

**n ~ 3x10<sup>18</sup> cm<sup>-3</sup>**

**L<sub>acc</sub> ~ 1 cm**

**U<sub>laser</sub> ~ 2 J**

**P<sub>laser</sub> ~ 100 TW**

**W ~ 4 GeV**

**n ~ 7x10<sup>17</sup> cm<sup>-3</sup>**

**L<sub>acc</sub> ~ 8 cm**

**U<sub>laser</sub> ~ 16 J**

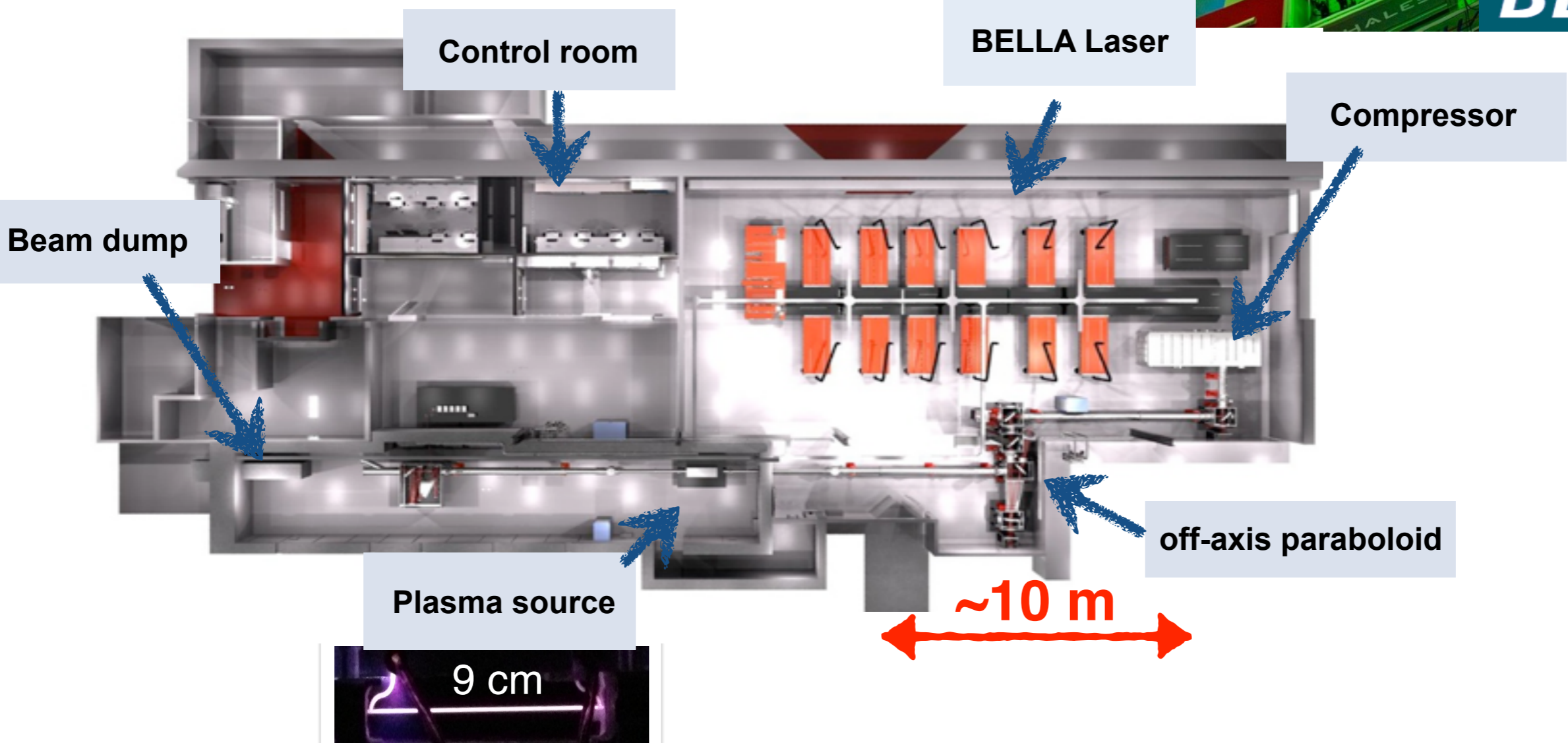
**P<sub>laser</sub> ~ 0.4 PW**

# BELLA Laser Facility

Leemans et al., PAC Proc. (2013)

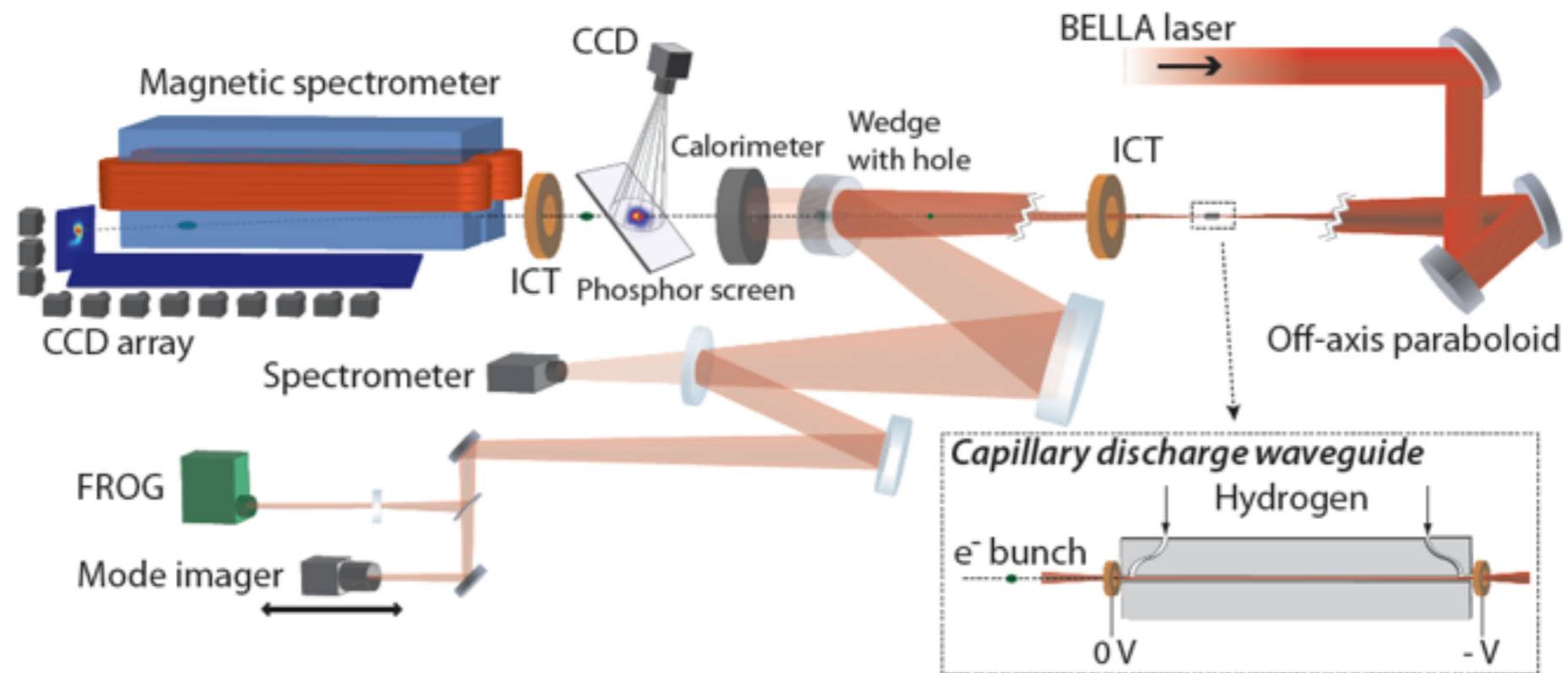
## BELLA (BERkeley Lab Laser Accelerator) Facility:

- PW-laser for laser-accelerator science
- **>42 J in <40 fs (> IPW) at 1 Hz laser** and supporting infrastructure at LBNL
- commissioned in 2013

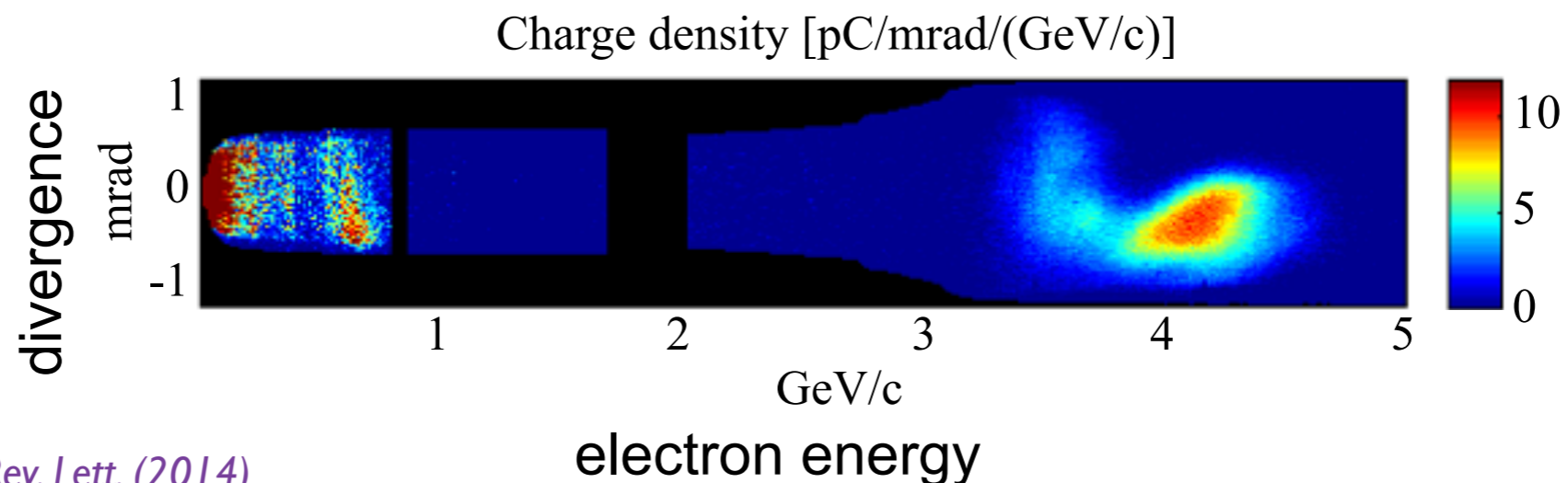




# Multi-GeV electron acceleration using BELLA

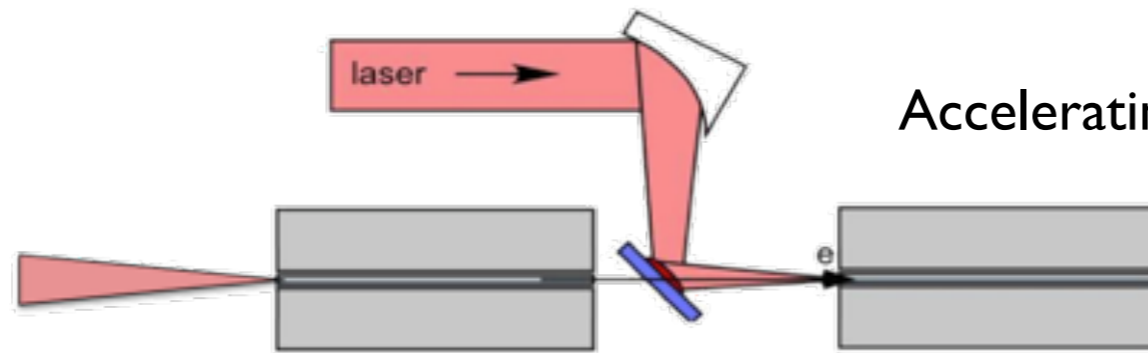


- ▶ plasma: 9 cm H-discharge capillary, on-axis density  $7 \times 10^{17}/\text{cc}$
- ▶ laser: 16 J, 40 fs
- ▶ 4.25 GeV e-beam, 6% energy spread, 0.3 mrad divergence, 10 pC



# Staged LPAs: average gradient determined by driver in-coupling distance

Schroeder et al., PR STAB (2010)



Accelerating gradient:  $E \sim E_0 = (m_e c \omega_p / e) \propto \sqrt{n}$

number of stages:  $N_{\text{stage}} \propto n \lambda^2$

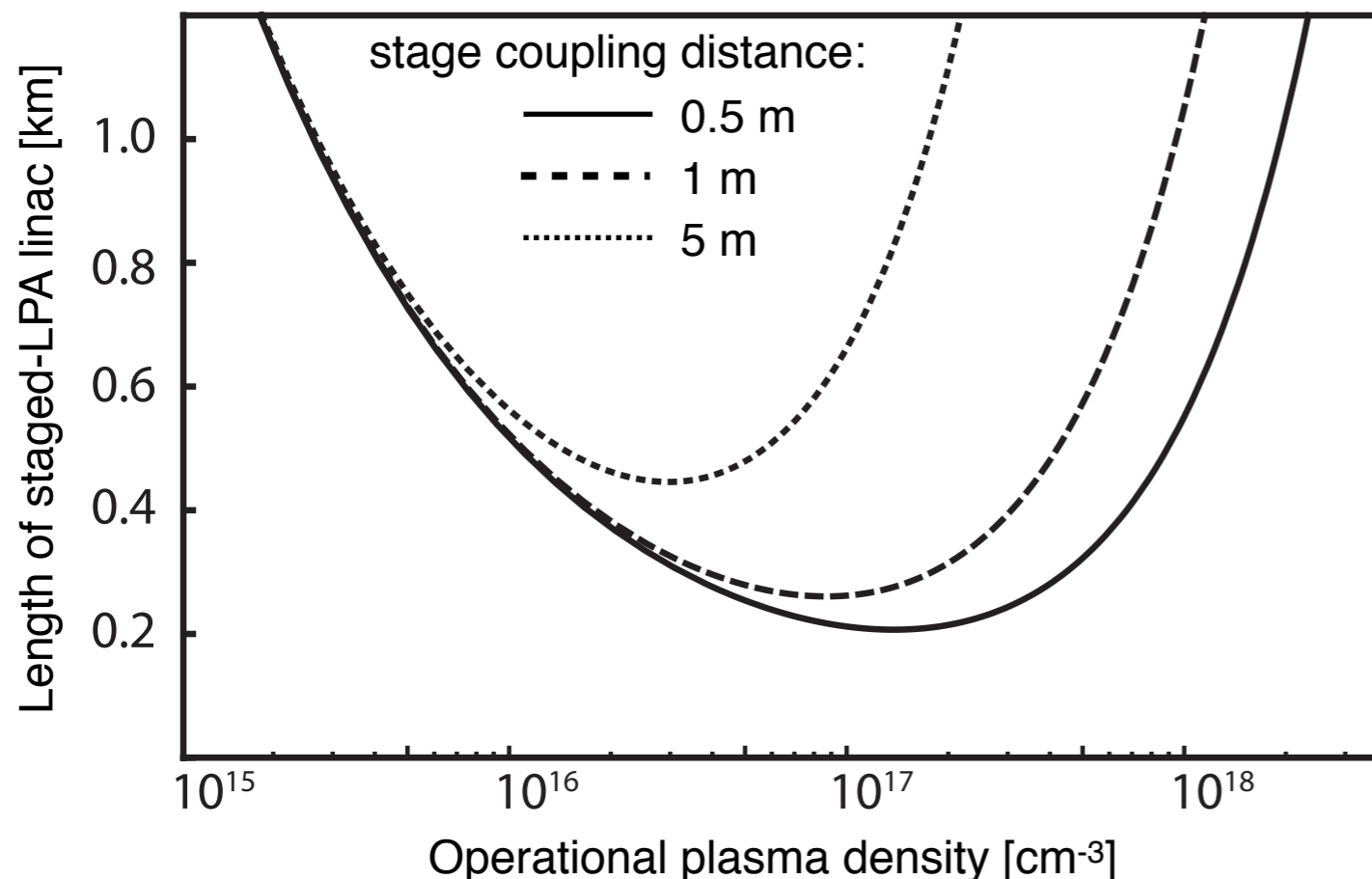
Compact laser in-coupling distance enables high **average** accelerating gradient:

- **Conventional laser optics:** requires many Rayleigh ranges to reduce fluency on optic (avoid damage)

$$L_{\text{coupling}} \propto n^{-5/4} \lambda^{-1}$$

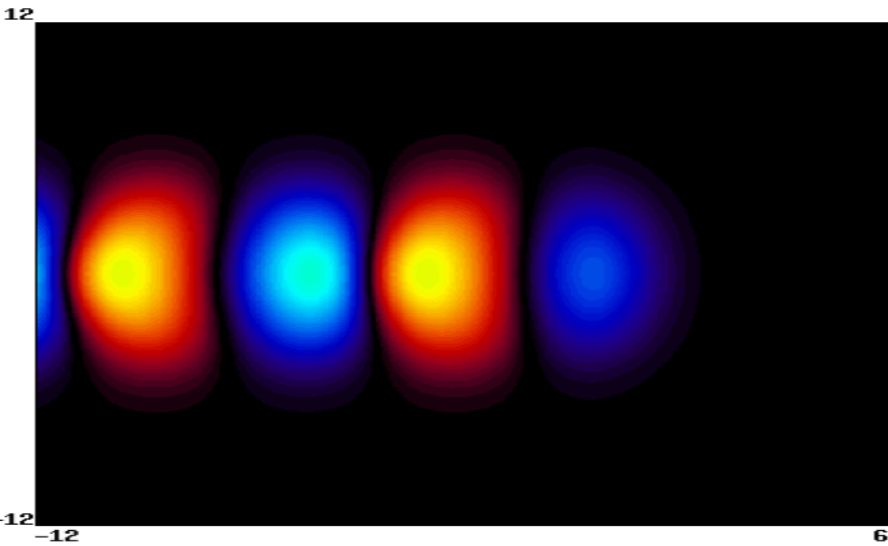
- **Plasma mirror:** relies on critical density plasma production (high laser intensity): laser coupling < 1 m

Length of 1 TeV staged-LPA linac

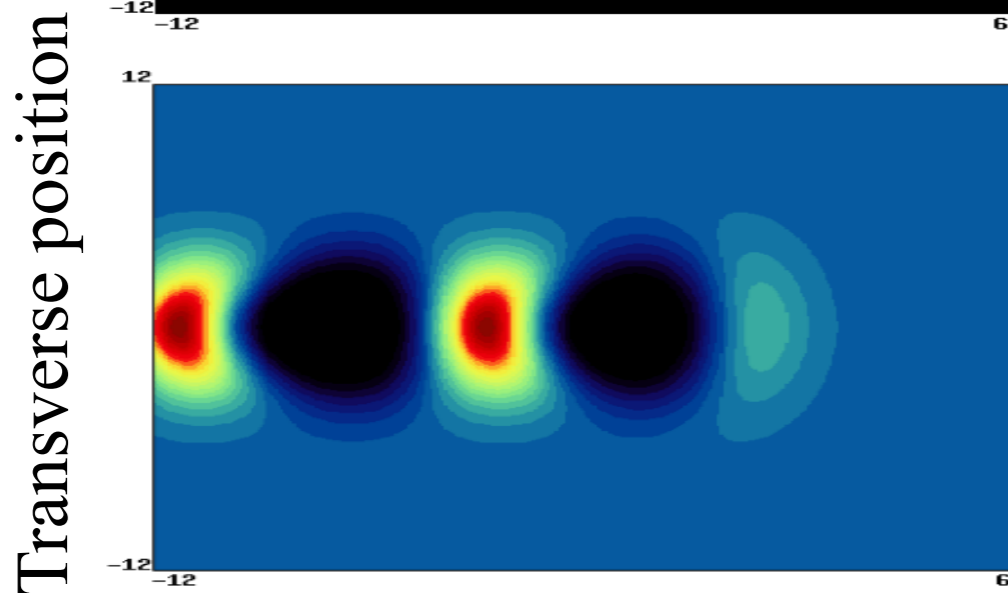


# Quasi-linear regime: positron focusing & independent control of acceleration and focusing

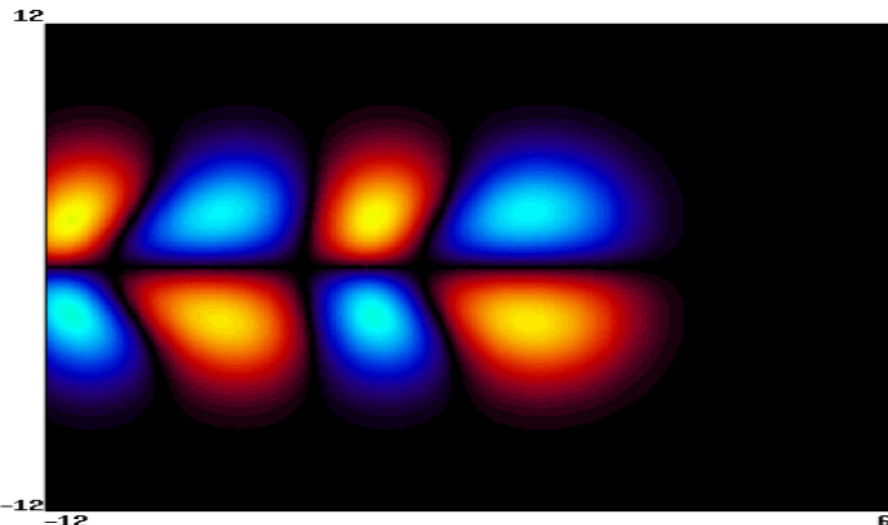
Accelerating field



Plasma density



Focusing field



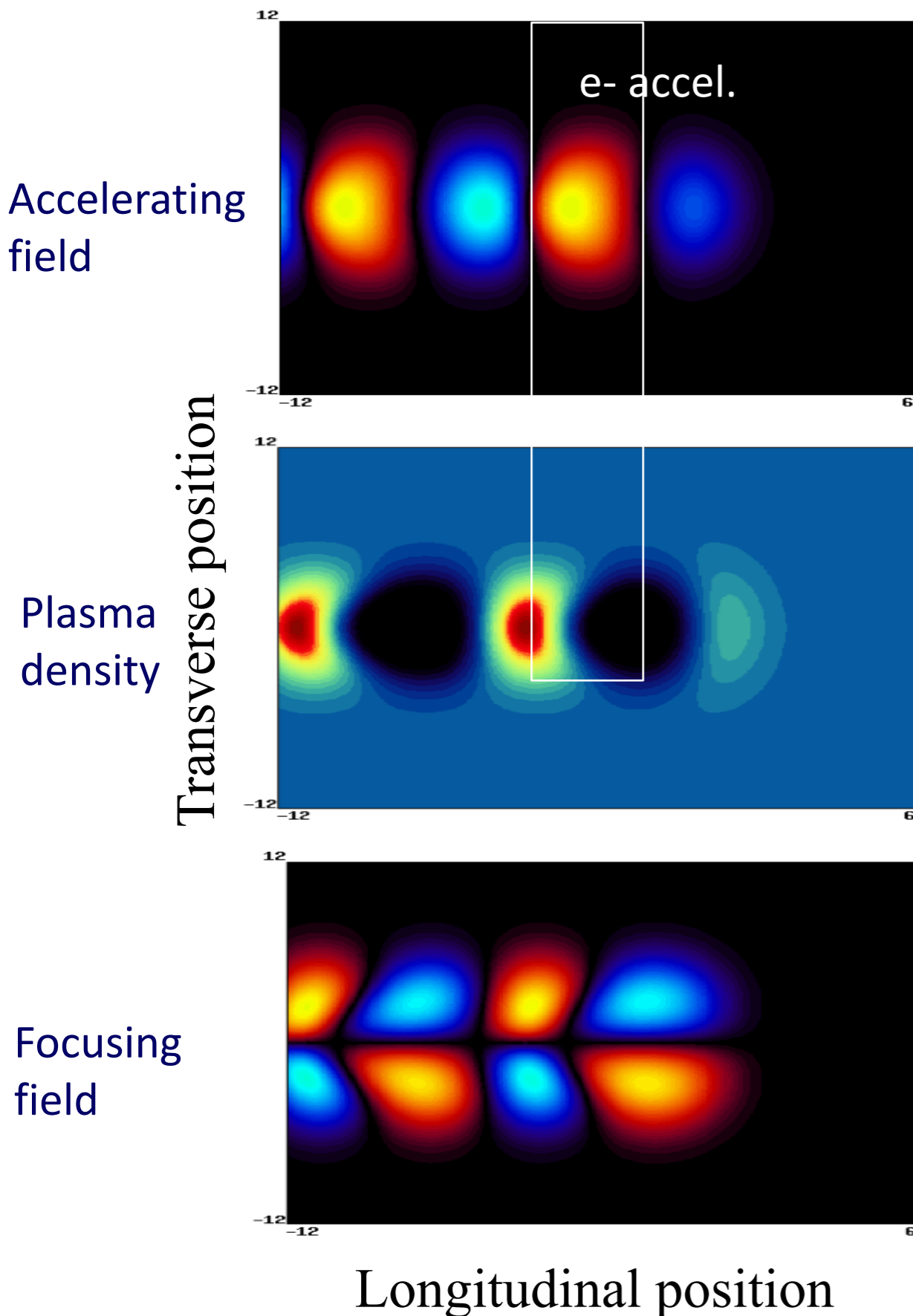
Longitudinal position

► Operate in “quasi-linear” regime:

- Quiver momentum weakly-relativistic  $a \sim 1$   
(Intensity  $\sim 10^{18}$  W/cm<sup>2</sup>)
- Region of acceleration/focusing for both electrons and positrons
- Stable laser propagation in plasma channel
- Independent control of accelerating and focusing forces



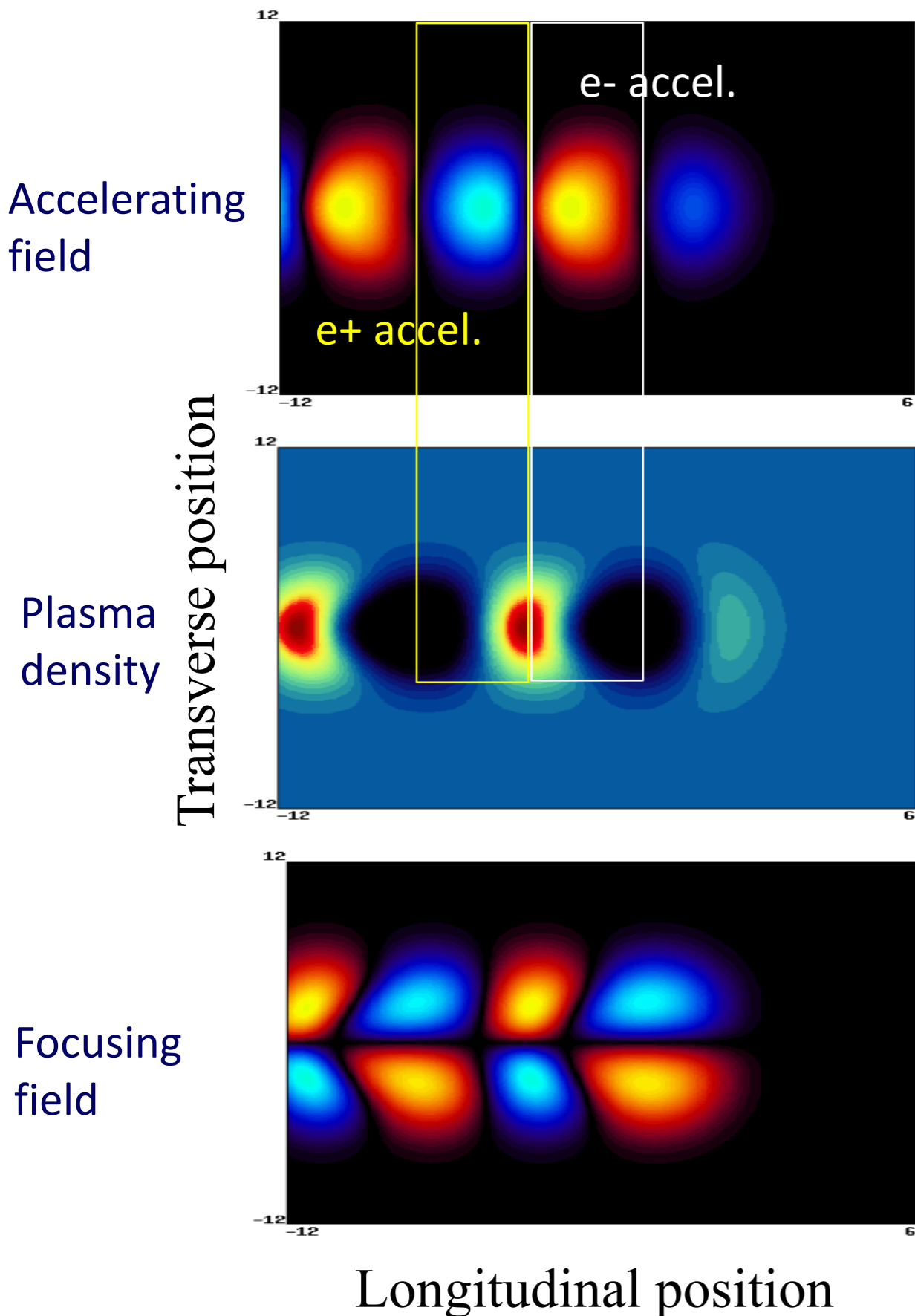
# Quasi-linear regime: positron focusing & independent control of acceleration and focusing



## ► Operate in “quasi-linear” regime:

- Quiver momentum weakly-relativistic  $a \sim 1$   
(Intensity  $\sim 10^{18}$  W/cm<sup>2</sup>)
- Region of acceleration/focusing for both electrons and positrons
- Stable laser propagation in plasma channel
- Independent control of accelerating and focusing forces

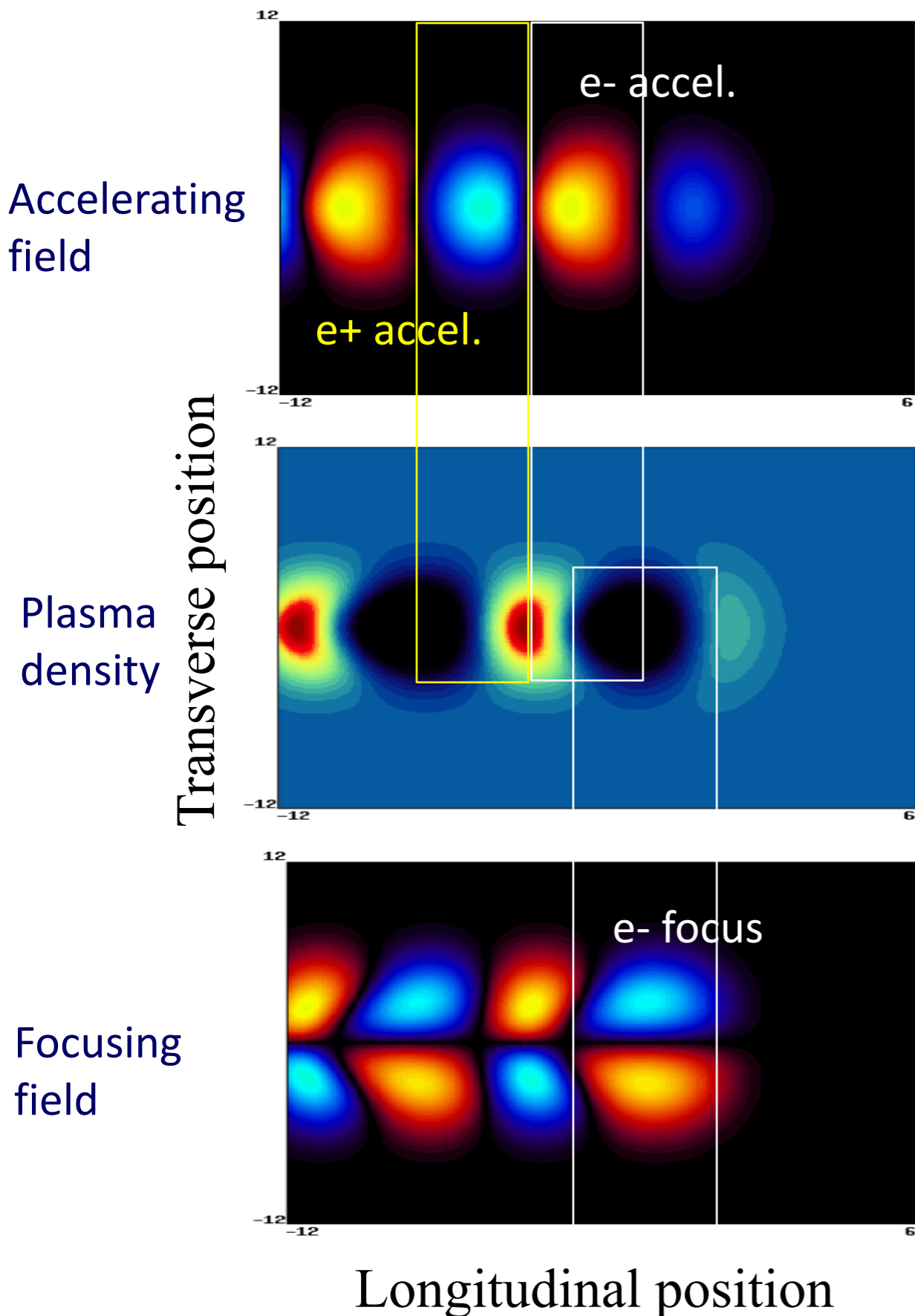
# Quasi-linear regime: positron focusing & independent control of acceleration and focusing



## ▶ Operate in “quasi-linear” regime:

- Quiver momentum weakly-relativistic  $a \sim 1$   
(Intensity  $\sim 10^{18}$  W/cm<sup>2</sup>)
- Region of acceleration/focusing for both electrons and positrons
- Stable laser propagation in plasma channel
- Independent control of accelerating and focusing forces

# Quasi-linear regime: positron focusing & independent control of acceleration and focusing

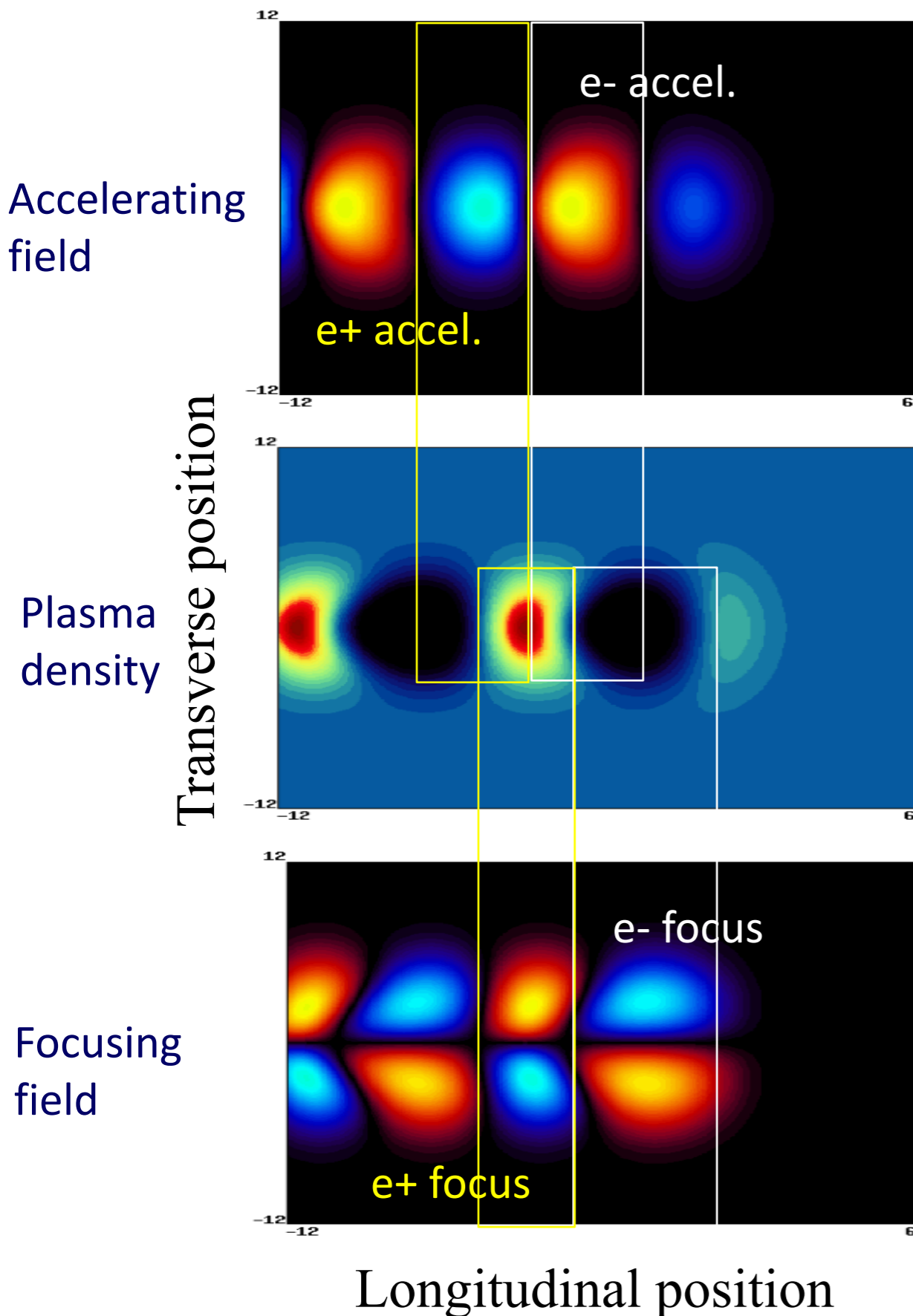


## ► Operate in “quasi-linear” regime:

- Quiver momentum weakly-relativistic  $a \sim 1$   
(Intensity  $\sim 10^{18}$  W/cm<sup>2</sup>)
- Region of acceleration/focusing for both electrons and positrons
- Stable laser propagation in plasma channel
- Independent control of accelerating and focusing forces



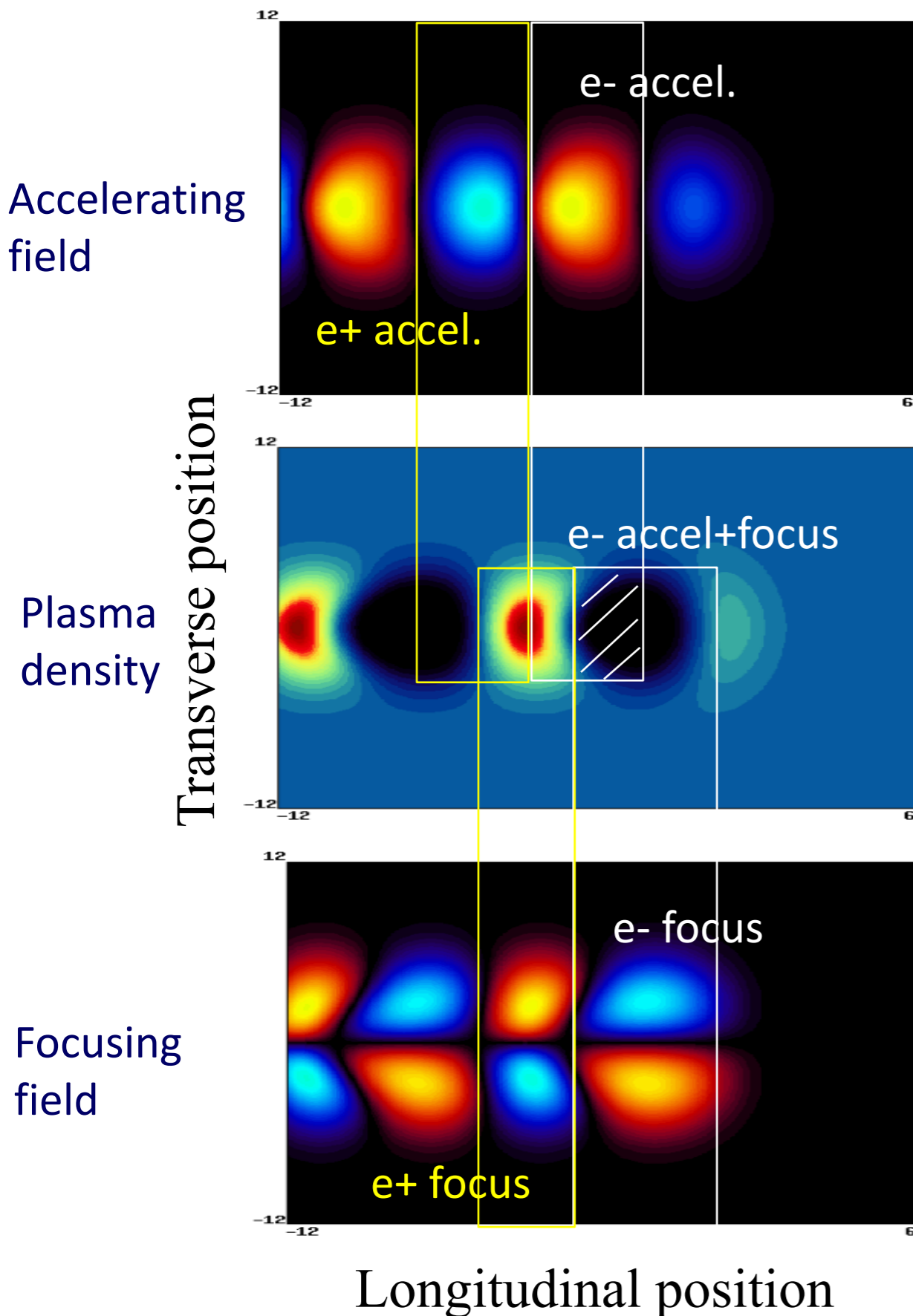
# Quasi-linear regime: positron focusing & independent control of acceleration and focusing



## ▶ Operate in “quasi-linear” regime:

- Quiver momentum weakly-relativistic  $a \sim 1$   
(Intensity  $\sim 10^{18}$  W/cm<sup>2</sup>)
- Region of acceleration/focusing for both electrons and positrons
- Stable laser propagation in plasma channel
- Independent control of accelerating and focusing forces

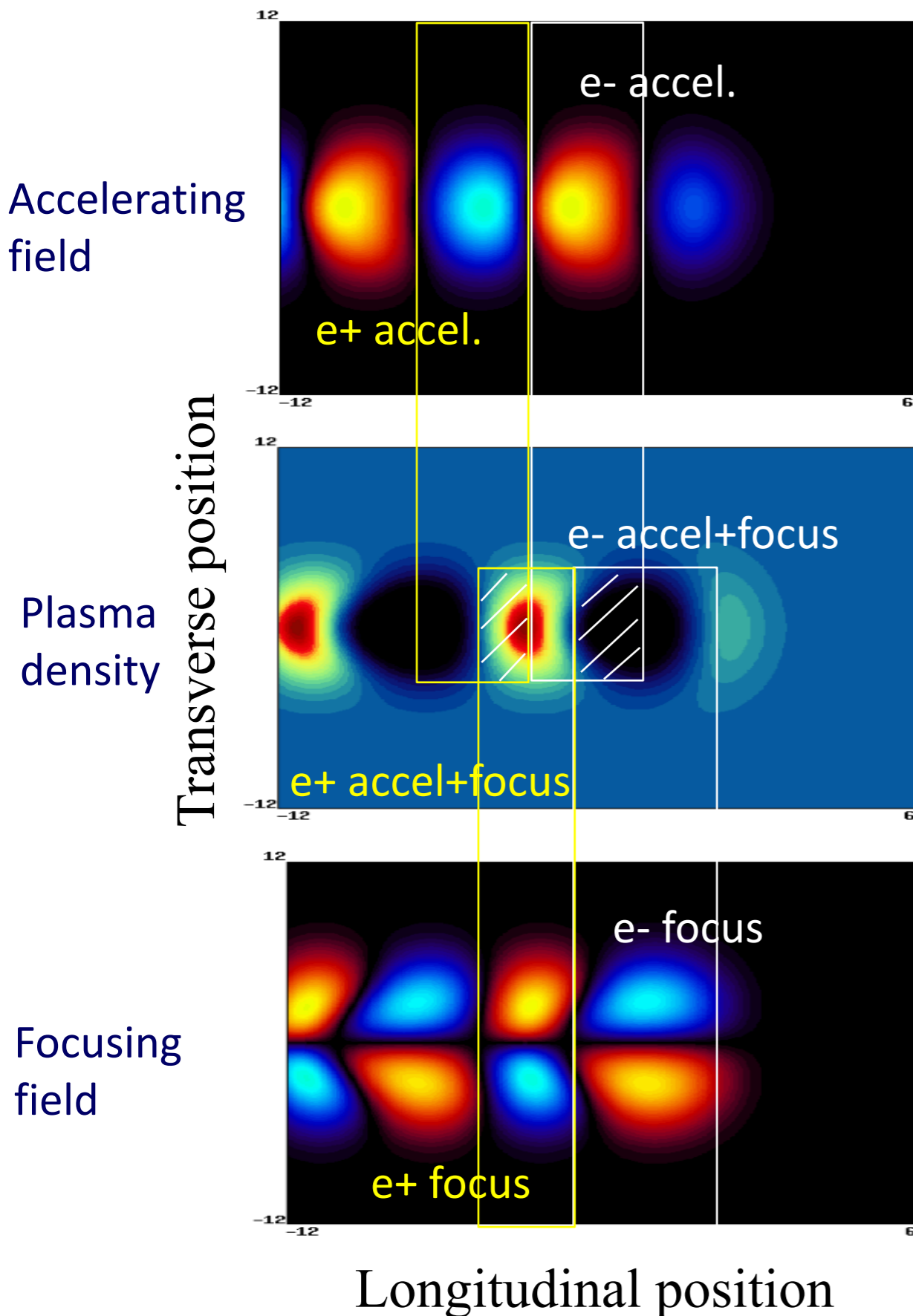
# Quasi-linear regime: positron focusing & independent control of acceleration and focusing



## ► Operate in “quasi-linear” regime:

- Quiver momentum weakly-relativistic  $a \sim 1$   
(Intensity  $\sim 10^{18}$  W/cm<sup>2</sup>)
- Region of acceleration/focusing for both electrons and positrons
- Stable laser propagation in plasma channel
- Independent control of accelerating and focusing forces

# Quasi-linear regime: positron focusing & independent control of acceleration and focusing

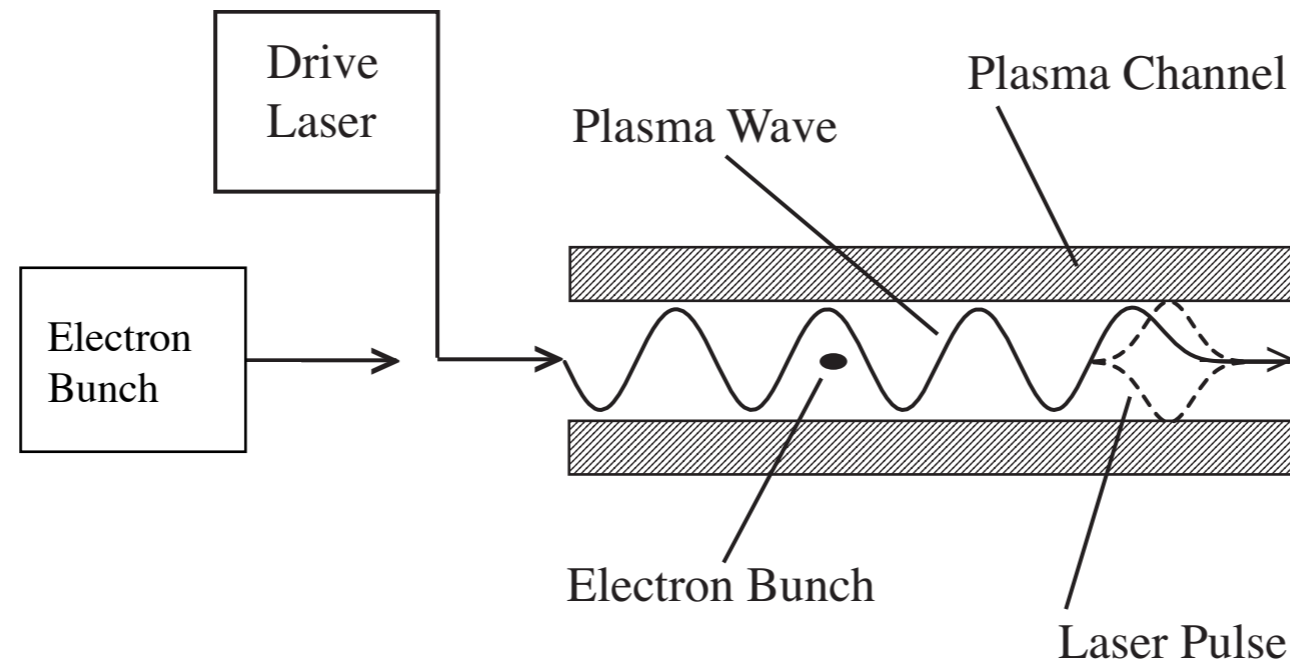


## ▶ Operate in “quasi-linear” regime:

- Quiver momentum weakly-relativistic  $a \sim 1$   
(Intensity  $\sim 10^{18}$  W/cm<sup>2</sup>)
- Region of acceleration/focusing for both electrons and positrons
- Stable laser propagation in plasma channel
- Independent control of accelerating and focusing forces



# hollow plasma channels: Excellent wakefield properties for ultra-low emittance preservation



► Excellent wakefield properties in plasma channel and *independent* control over accelerating and focusing forces:

- Accelerating wakefield uniform in radial position
- Focusing wakefield uniform in longitudinal position (no head-to-tail variation in focusing force) and linear in radial position
- Ion motion negligible
- Mitigates Coulomb scattering

# Beam loading in near-hollow plasma channels: use shaped bunches to eliminate energy spread

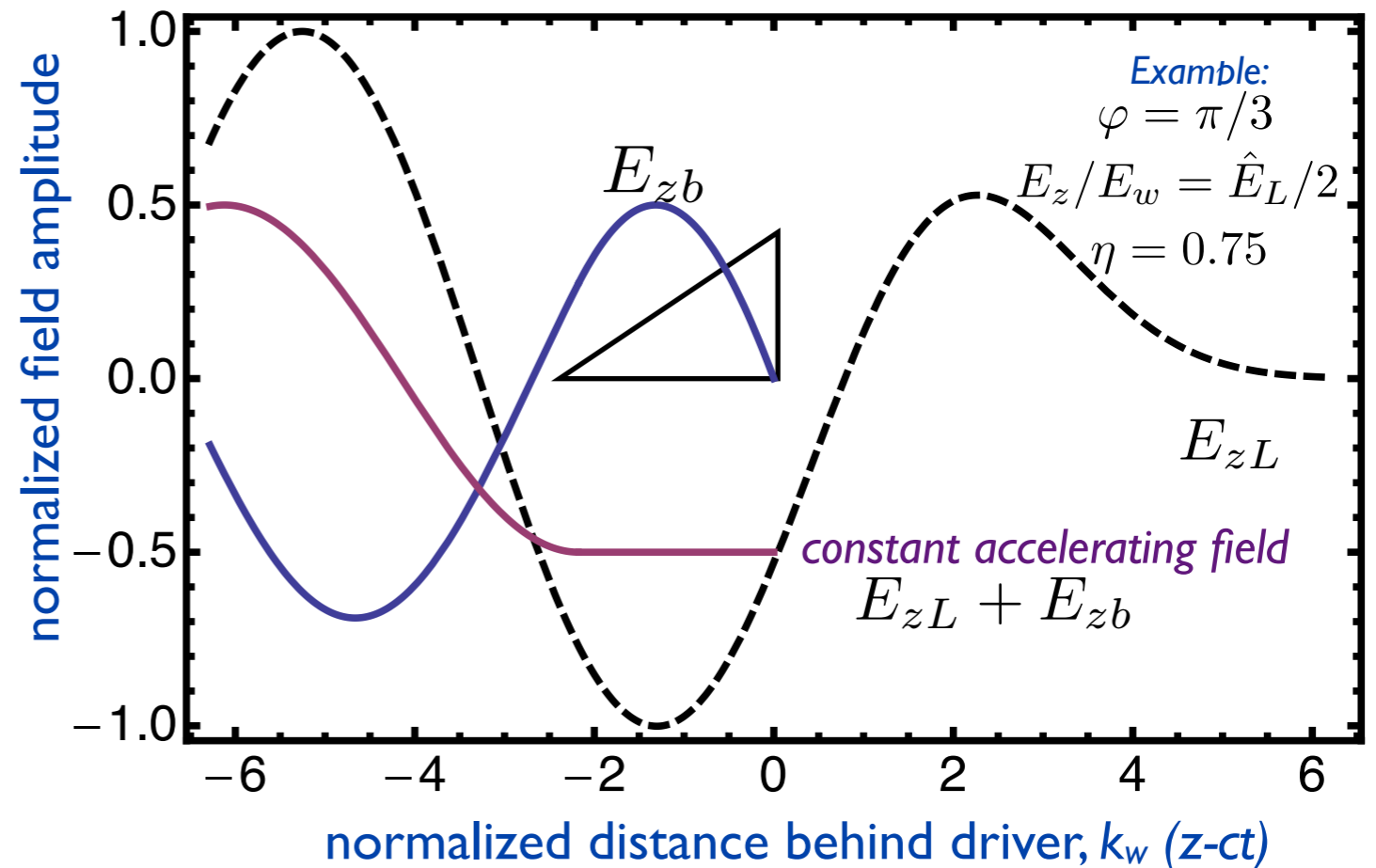
- ▶ Energy spread minimized using shaped beams

# Beam loading in near-hollow plasma channels: use shaped bunches to eliminate energy spread

- ▶ Energy spread minimized using shaped beams

Ramped triangular current distribution :

$$I = (1 + \zeta/L_b)I_b$$



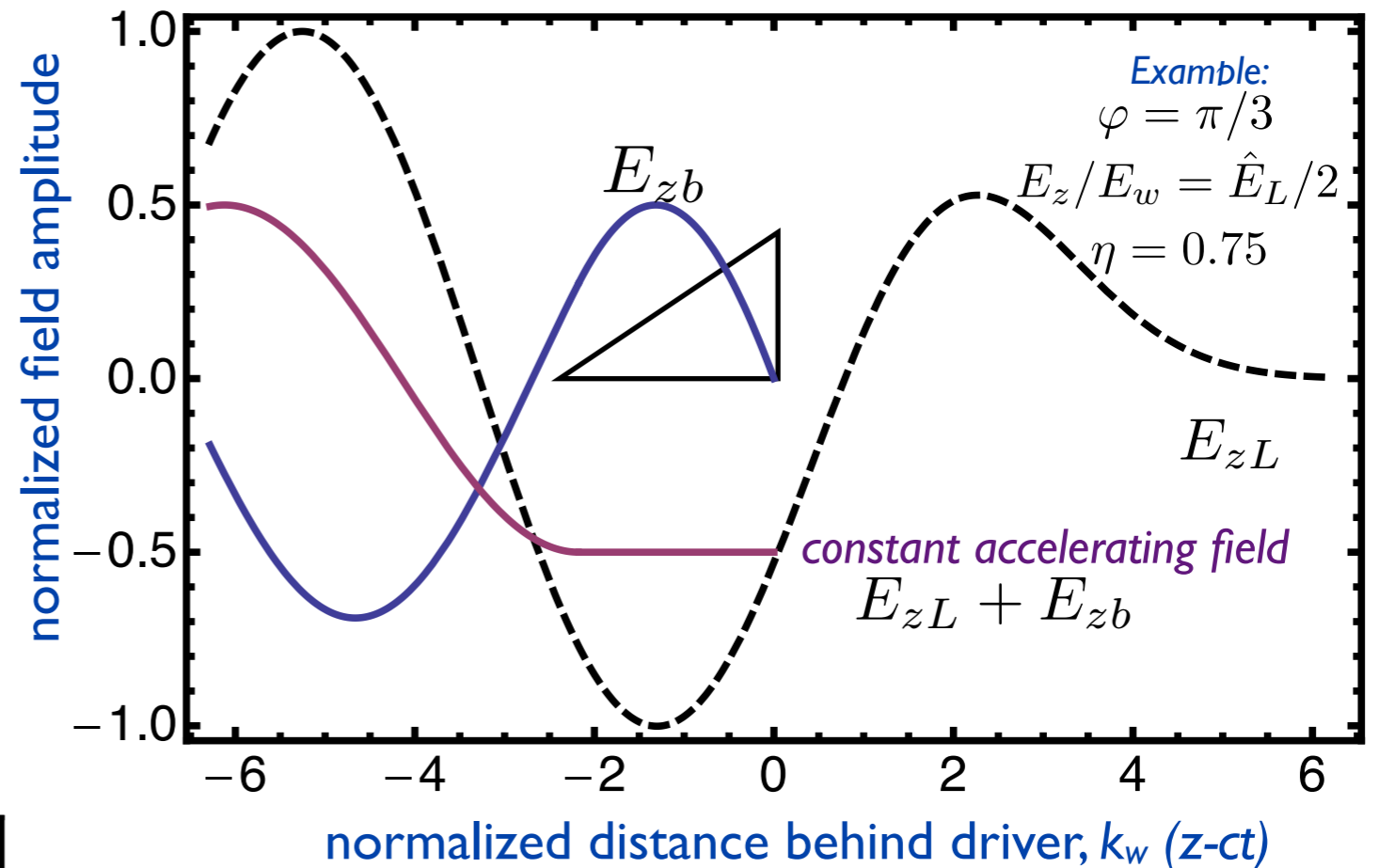
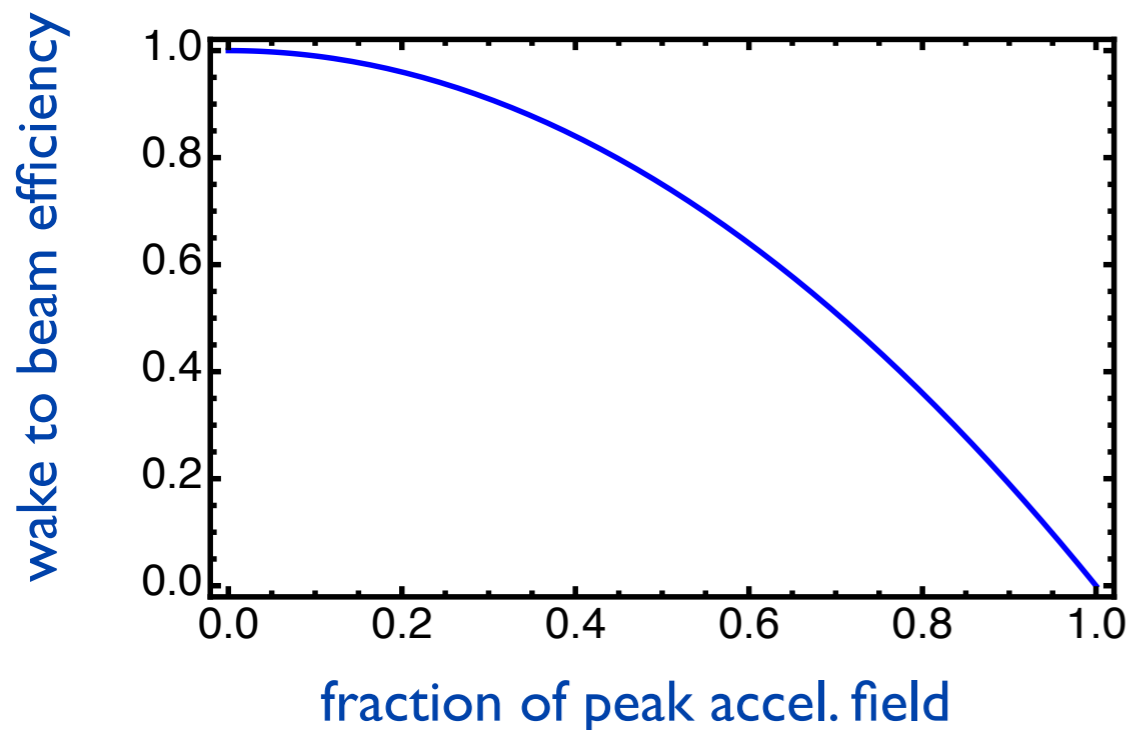
# Beam loading in near-hollow plasma channels: use shaped bunches to eliminate energy spread

- ▶ Energy spread minimized using shaped beams

Ramped triangular current distribution :

$$I = (1 + \zeta/L_b)I_b$$

- ▶ Trade-off between gradient and efficiency (for no induced energy spread)





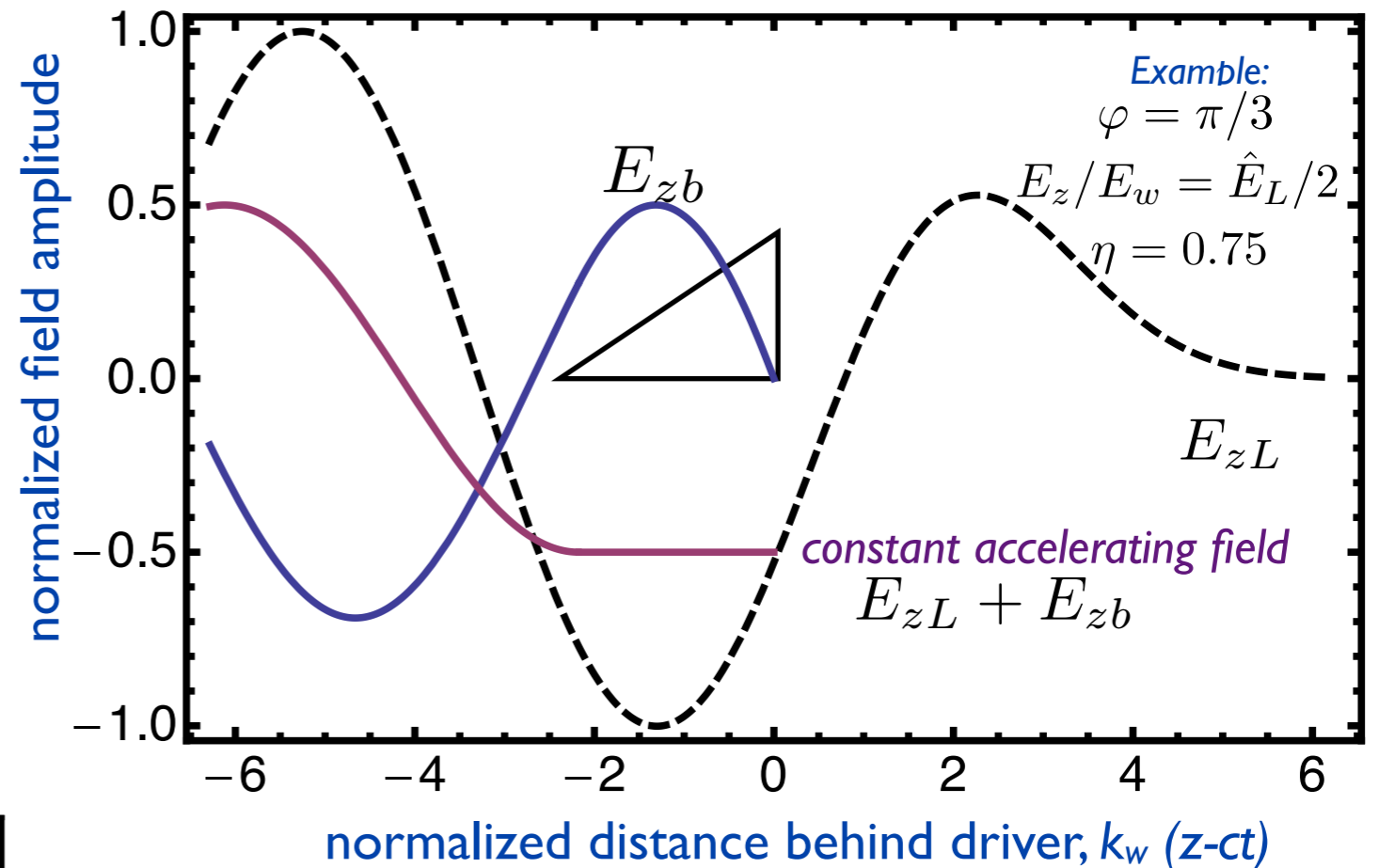
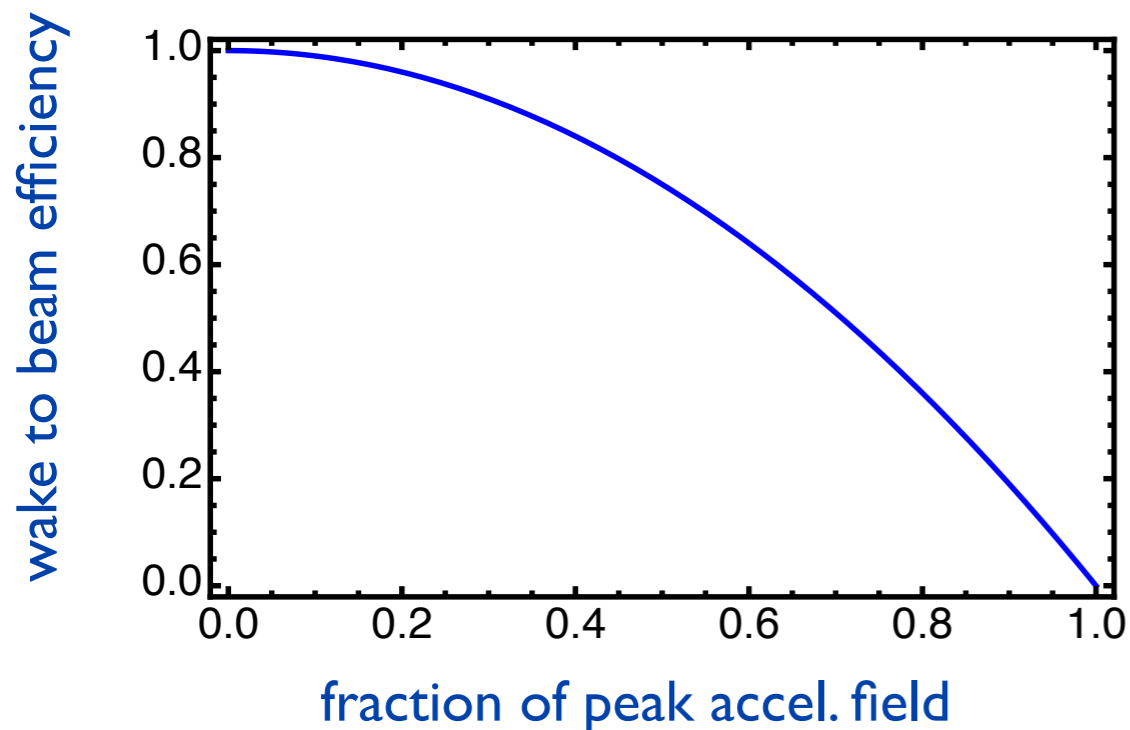
# Beam loading in near-hollow plasma channels: use shaped bunches to eliminate energy spread

- ▶ Energy spread minimized using shaped beams

Ramped triangular current distribution :

$$I = (1 + \zeta/L_b)I_b$$

- ▶ Trade-off between gradient and efficiency (for no induced energy spread)

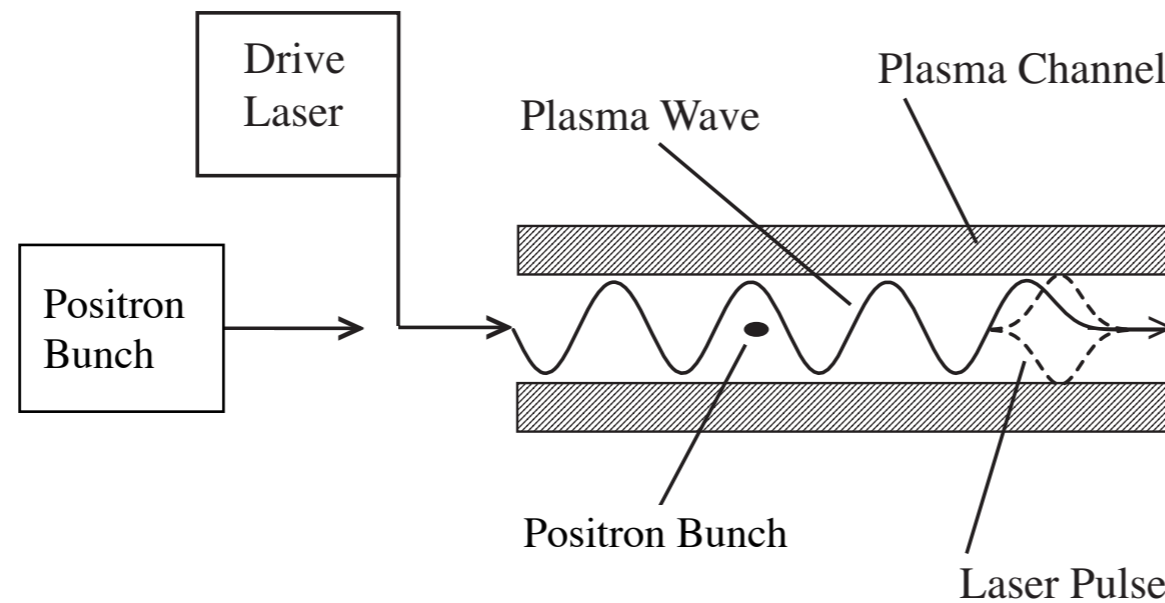


beam charge:  $N_b \propto n_w^{-1/2}$

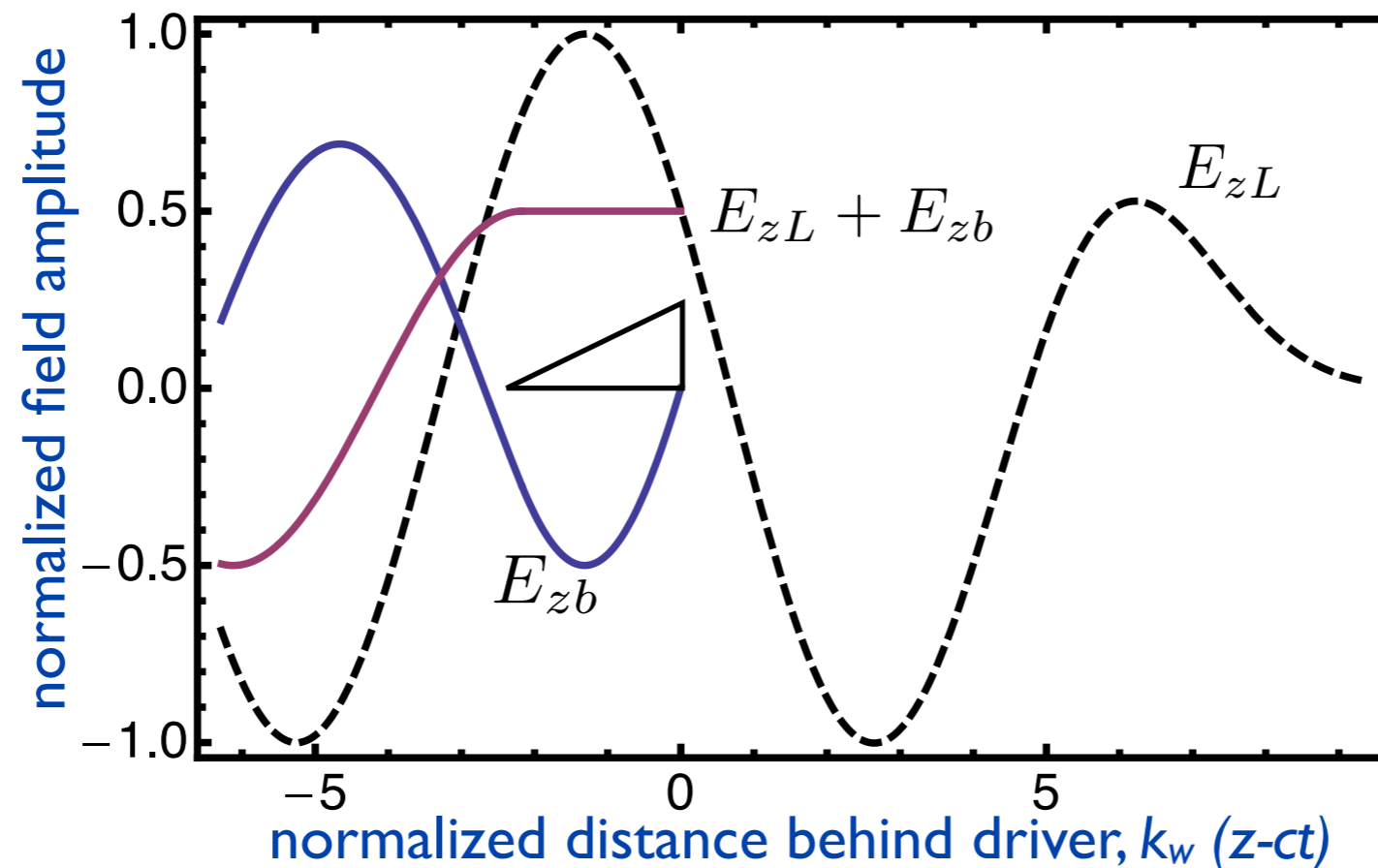
lower plasma density,  
higher beam charge

beam length:  $L_b \sim k_p^{-1} \propto n_w^{-1/2}$

# Positron beams accelerated in linear regime in hollow plasma channel



- In quasi-linear regime, acceleration of positrons is symmetric to electrons



# Basic scalings indicate disruption small for plasma accelerated (ultra-short) beams

- ▶ Disruption parameter:

*P. Chen and K. Yokoya, Phys. Rev. D 38, 987 (1988)*

$$D_x = \frac{2r_e N \sigma_z}{\gamma \sigma_x^* (\sigma_x^* + \sigma_y^*)}$$

- ▶ Plasma-based accelerators intrinsically produce short (< plasma skin depth) beams
- ▶ Luminosity enhancement effects from beam disruption are typically weak  $D \ll 1$  for LPA-based colliders owing to the ultrashort LPA bunch length ( $\ll$  beam-beam lens)

# Beamstrahlung minimization: basic scalings

▶ Beamstrahlung parameter:  $\Upsilon = \gamma \langle E + B \rangle / E_c \approx \frac{5r_e^2 \gamma}{6\alpha(\sigma_x^* + \sigma_y^*)} \frac{N}{\sigma_z}$

▶ average number of beamstrahlung photons/lepton:

*P. Chen and K. Yokoya*  $n_\gamma \approx 2.54 \left( \alpha^2 / r_e \gamma \right) \sigma_z \Upsilon (1 + \Upsilon^{2/3})^{-1/2}$



# Beamstrahlung minimization: basic scalings

▶ Beamstrahlung parameter:  $\Upsilon = \gamma \langle E + B \rangle / E_c \approx \frac{5r_e^2 \gamma}{6\alpha(\sigma_x^* + \sigma_y^*)} \frac{N}{\sigma_z}$

▶ average number of beamstrahlung photons/lepton:

*P. Chen and K. Yokoya*  $n_\gamma \approx 2.54 \left( \alpha^2 / r_e \gamma \right) \sigma_z \Upsilon (1 + \Upsilon^{2/3})^{-1/2}$

▶ TeV-scale colliders will operate in the high-beamstrahlung regime:  $\Upsilon \gg 1$

▶ Beamstrahlung suppressed by using short beams (in limit  $\Upsilon \gg 1$ ):

$$n_\gamma \approx 2.54 \left( \alpha^2 / r_e \gamma \right) \Upsilon^{2/3} \sigma_z$$

$$\delta_\gamma \approx 0.722 \left( \alpha^2 / r_e \gamma \right) \Upsilon^{2/3} \sigma_z$$

# Beamstrahlung minimization: basic scalings

▶ Beamstrahlung parameter:  $\Upsilon = \gamma \langle E + B \rangle / E_c \approx \frac{5r_e^2 \gamma}{6\alpha(\sigma_x^* + \sigma_y^*)} \frac{N}{\sigma_z}$

▶ average number of beamstrahlung photons/lepton:

*P. Chen and K. Yokoya*  $n_\gamma \approx 2.54 (\alpha^2 / r_e \gamma) \sigma_z \Upsilon (1 + \Upsilon^{2/3})^{-1/2}$

▶ TeV-scale colliders will operate in the high-beamstrahlung regime:  $\Upsilon \gg 1$

▶ Beamstrahlung suppressed by using short beams (in limit  $\Upsilon \gg 1$ ):

$$n_\gamma \approx 2.54 (\alpha^2 / r_e \gamma) \Upsilon^{2/3} \sigma_z$$

$$\delta_\gamma \approx 0.722 (\alpha^2 / r_e \gamma) \Upsilon^{2/3} \sigma_z$$

▶ Short beams save power ( $\Upsilon \gg 1$ ):

$$\frac{\mathcal{L}}{U_{\text{cm}}^2} \propto \frac{n_\gamma^{3/2} \eta_{wb} P_{\text{wall}}}{\sigma^* \gamma^{5/2} \sigma_z^{1/2}}$$

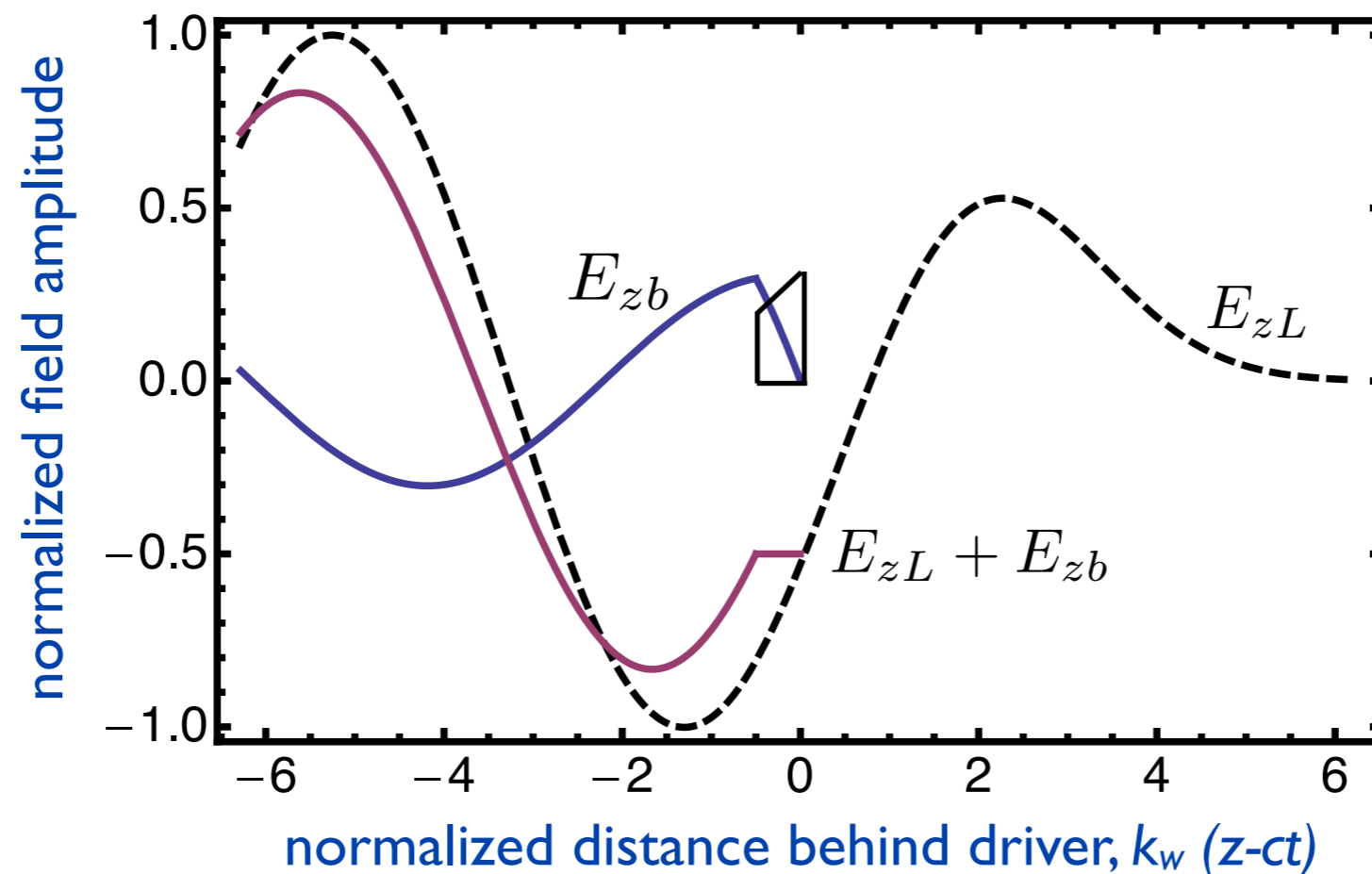
▶ Plasma-based accelerators produce short ( $<$  plasma skin depth) beams:  $\sigma_z < k_p^{-1}$

# Ultra-short beams can be accelerated without energy spread growth

- ▶ Ultra-short beams are desirable, e.g., beamstrahlung suppression in colliders.

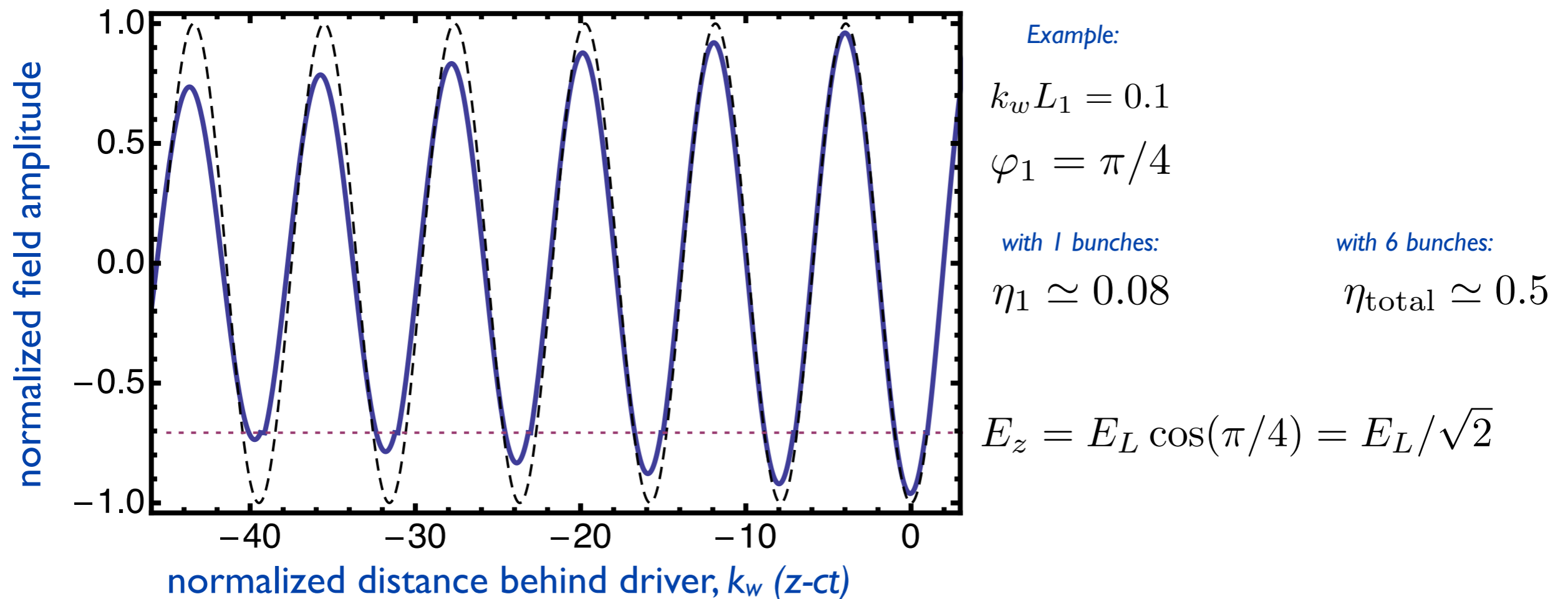
beamstrahlung photons/electron:  $n_\gamma \propto \mathcal{L}^{1/3} \sigma_z^{1/3}$

- ▶ Trapezoidal current distribution:



- ▶ Improved efficiency achieved using bunch trains

# Bunch trains allows ultra-short bunches with high efficiency, gradient, and no energy spread



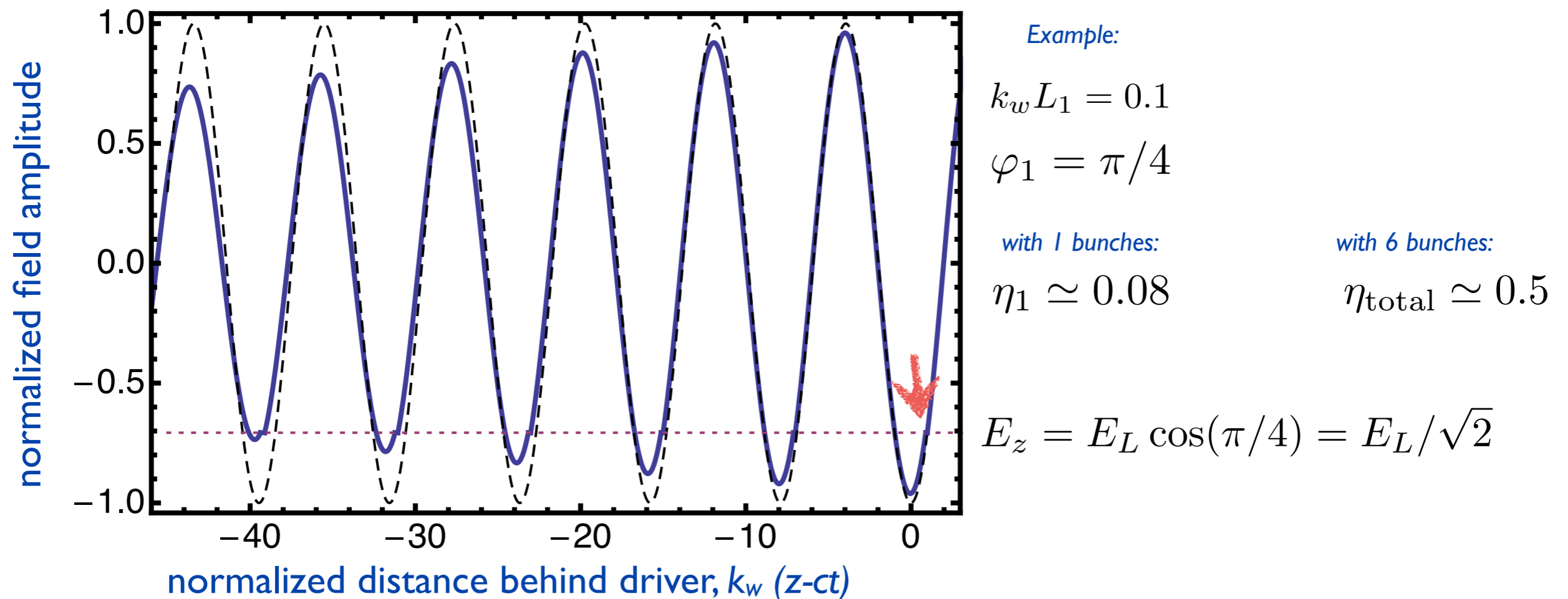
## ► Improved efficiency achieved using bunch trains:

- Each bunch has same charge
- Experiences same (constant across bunch) accelerating gradient
- Efficiency increased by number of bunches in train

► Trade-off between gradient and efficiency, with no energy spread growth using ultra-short beams



# Bunch trains allows ultra-short bunches with high efficiency, gradient, and no energy spread

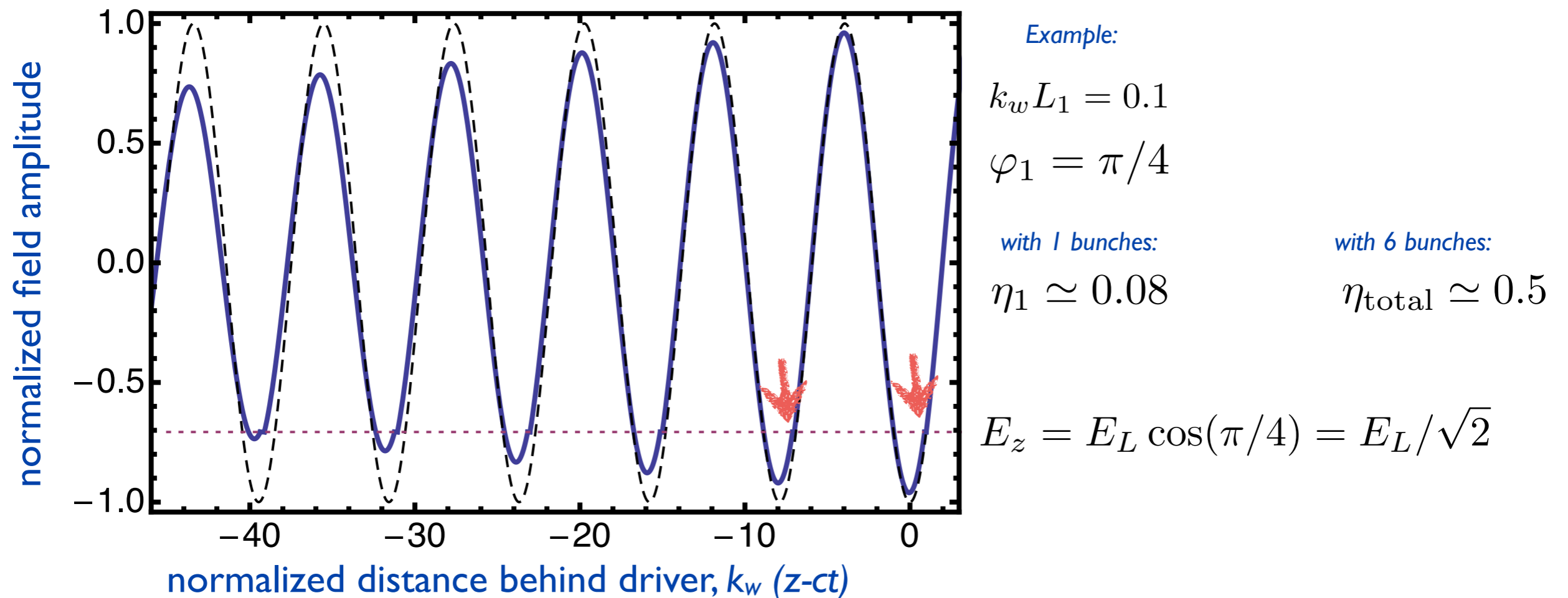


## ► Improved efficiency achieved using bunch trains:

- Each bunch has same charge
- Experiences same (constant across bunch) accelerating gradient
- Efficiency increased by number of bunches in train

► Trade-off between gradient and efficiency, with no energy spread growth using ultra-short beams

# Bunch trains allows ultra-short bunches with high efficiency, gradient, and no energy spread

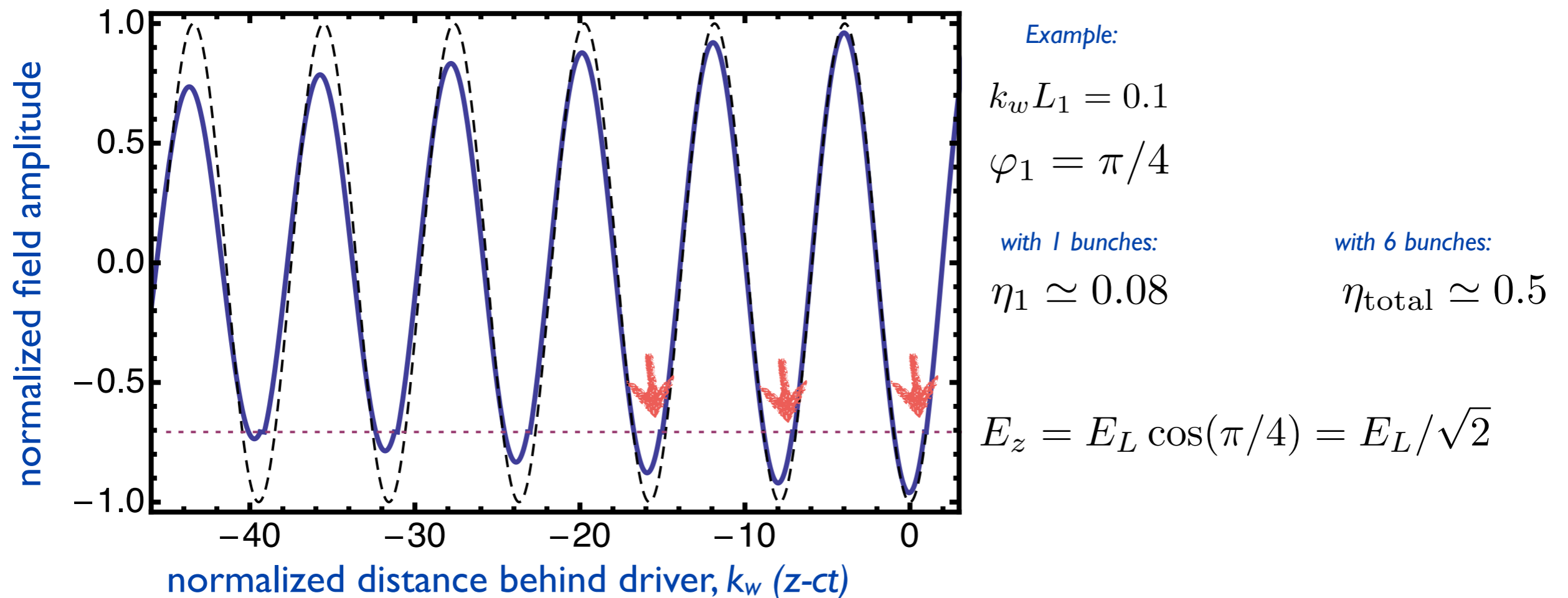


## ► Improved efficiency achieved using bunch trains:

- Each bunch has same charge
- Experiences same (constant across bunch) accelerating gradient
- Efficiency increased by number of bunches in train

► Trade-off between gradient and efficiency, with no energy spread growth using ultra-short beams

# Bunch trains allows ultra-short bunches with high efficiency, gradient, and no energy spread

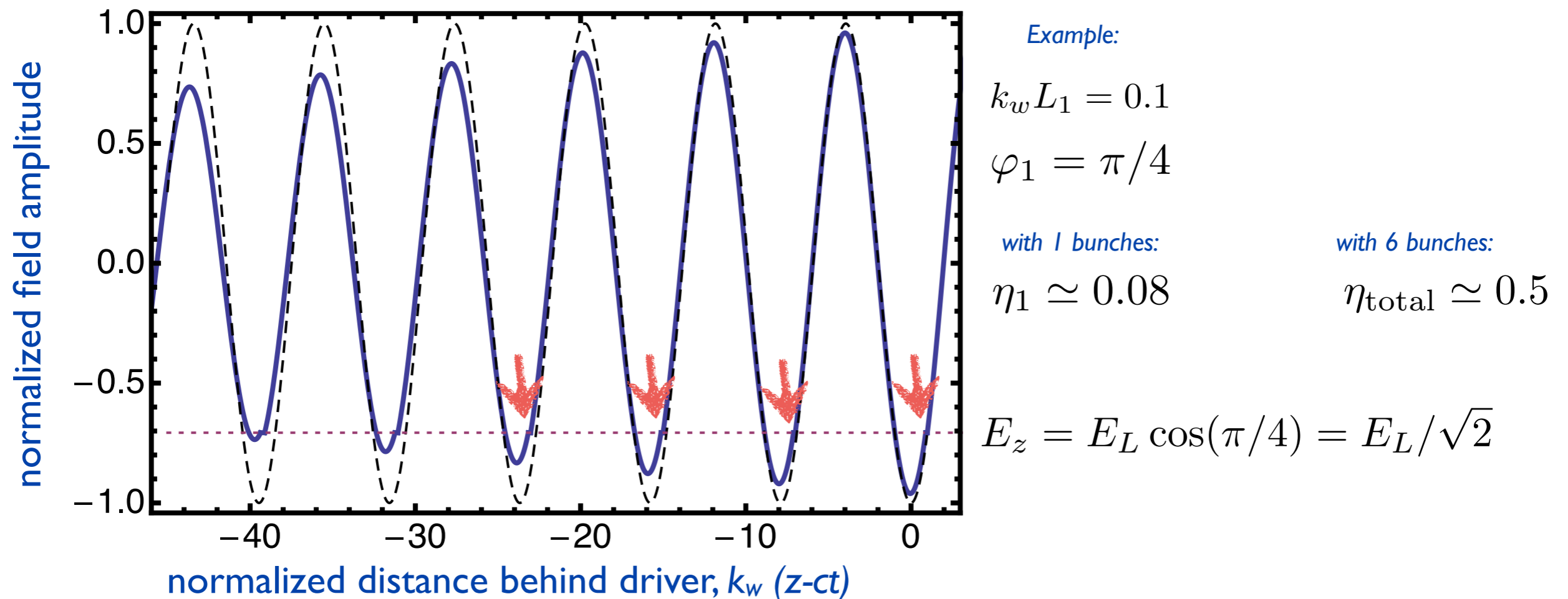


## ► Improved efficiency achieved using bunch trains:

- Each bunch has same charge
- Experiences same (constant across bunch) accelerating gradient
- Efficiency increased by number of bunches in train

► Trade-off between gradient and efficiency, with no energy spread growth using ultra-short beams

# Bunch trains allows ultra-short bunches with high efficiency, gradient, and no energy spread



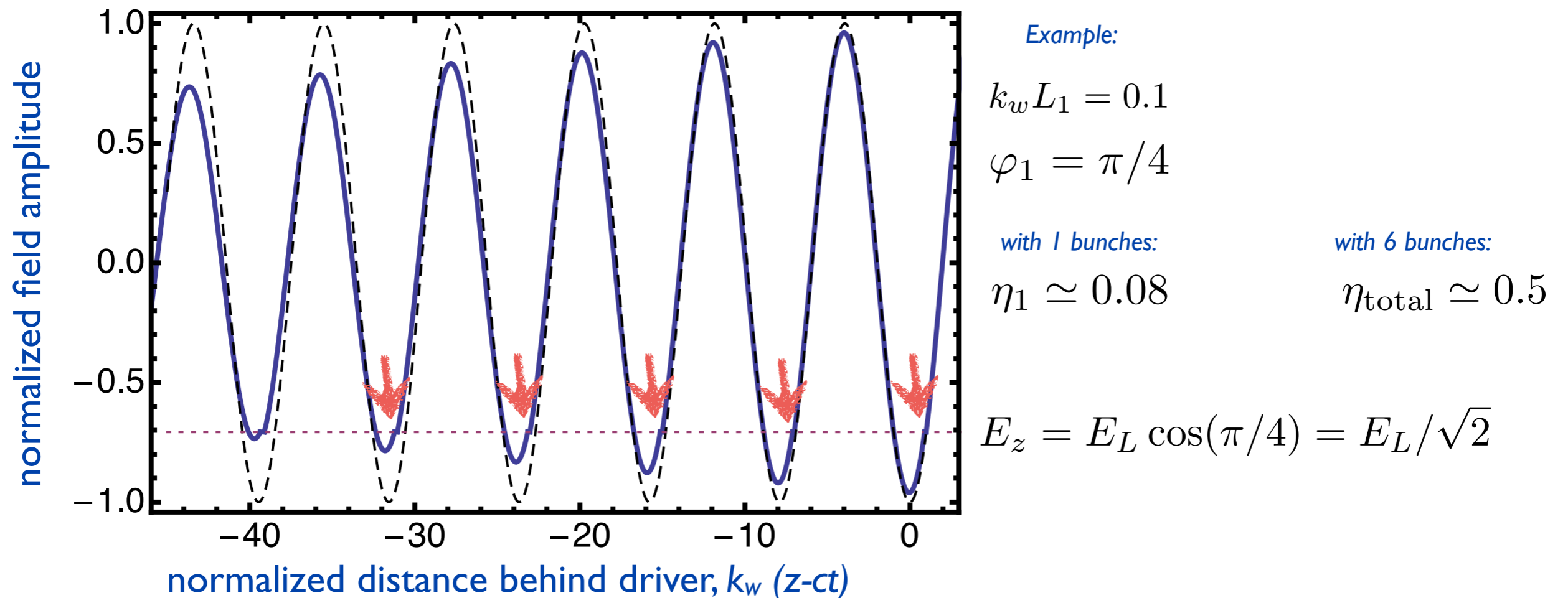
## ► Improved efficiency achieved using bunch trains:

- Each bunch has same charge
- Experiences same (constant across bunch) accelerating gradient
- Efficiency increased by number of bunches in train

► Trade-off between gradient and efficiency, with no energy spread growth using ultra-short beams



# Bunch trains allows ultra-short bunches with high efficiency, gradient, and no energy spread

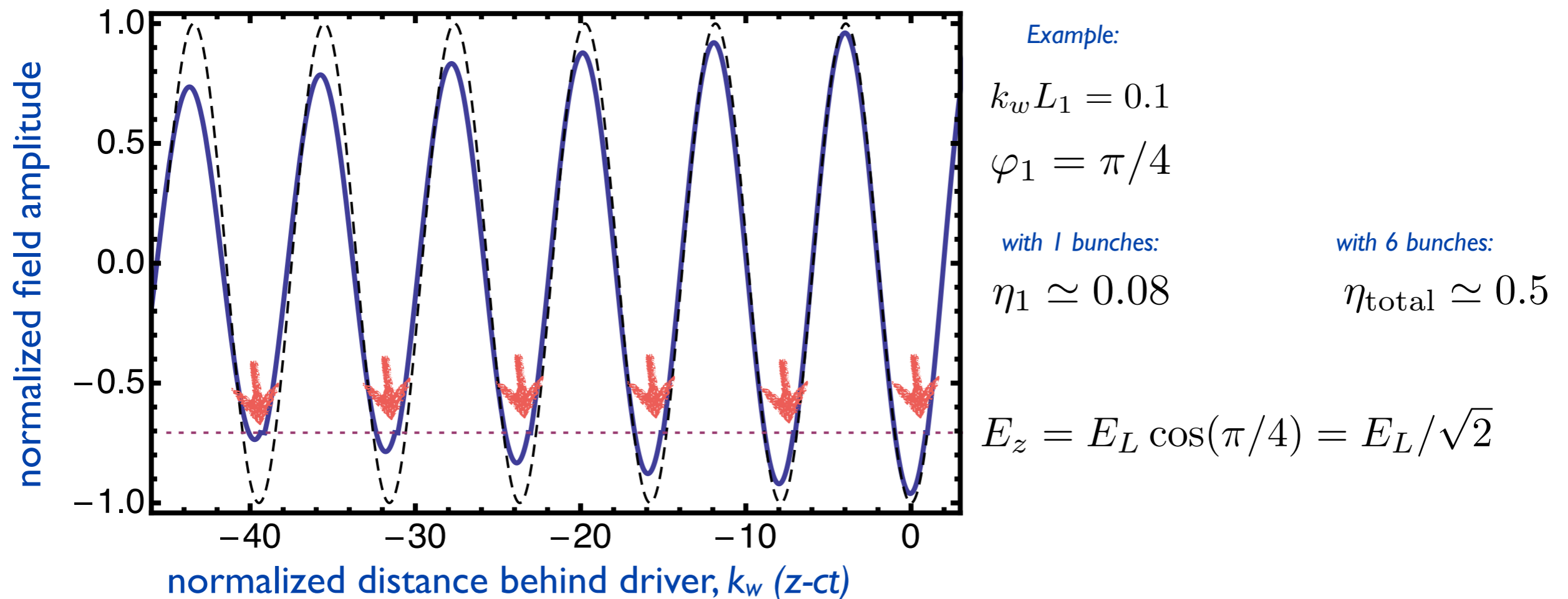


## ► Improved efficiency achieved using bunch trains:

- Each bunch has same charge
- Experiences same (constant across bunch) accelerating gradient
- Efficiency increased by number of bunches in train

► Trade-off between gradient and efficiency, with no energy spread growth using ultra-short beams

# Bunch trains allows ultra-short bunches with high efficiency, gradient, and no energy spread

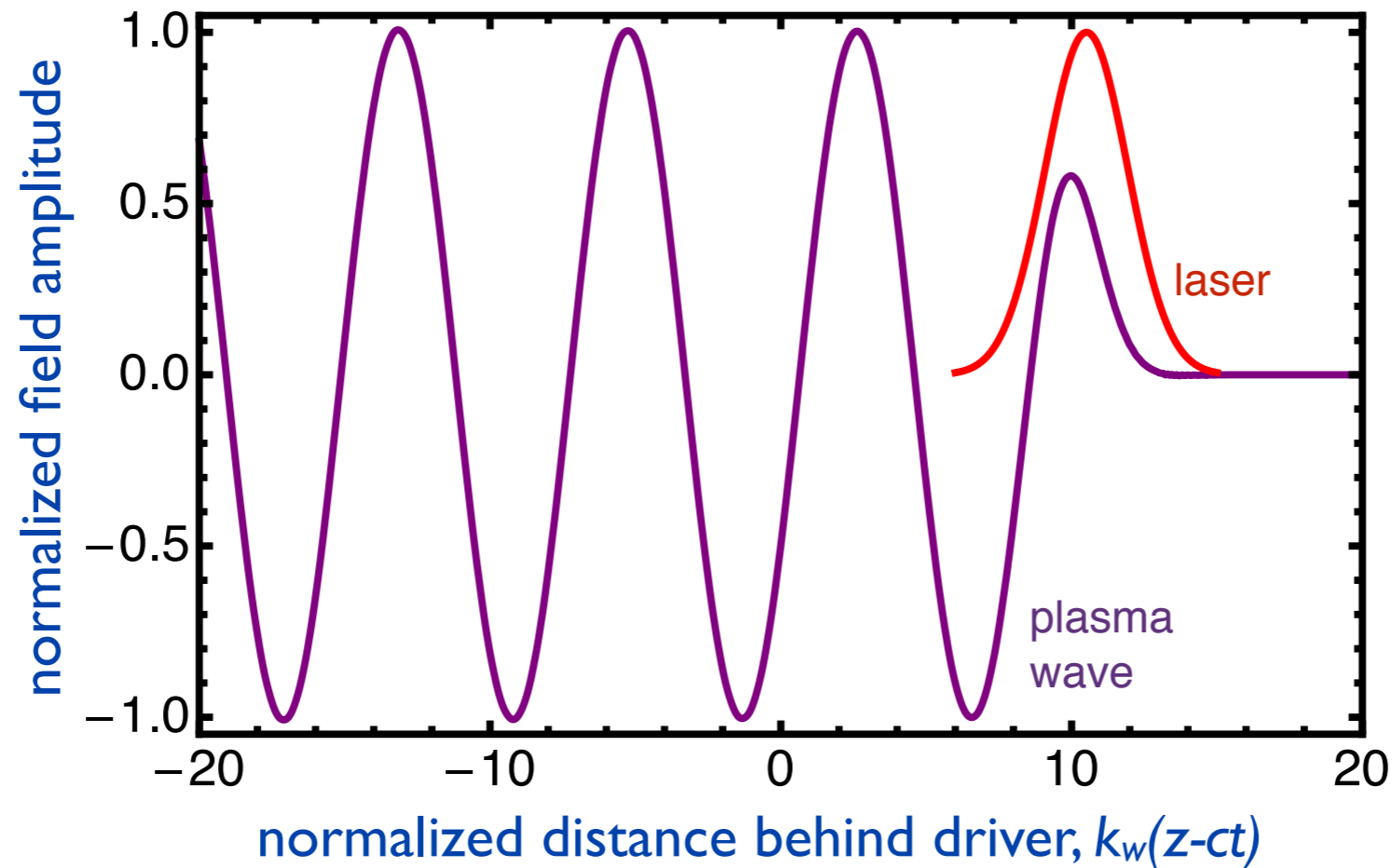


## ► Improved efficiency achieved using bunch trains:

- Each bunch has same charge
- Experiences same (constant across bunch) accelerating gradient
- Efficiency increased by number of bunches in train

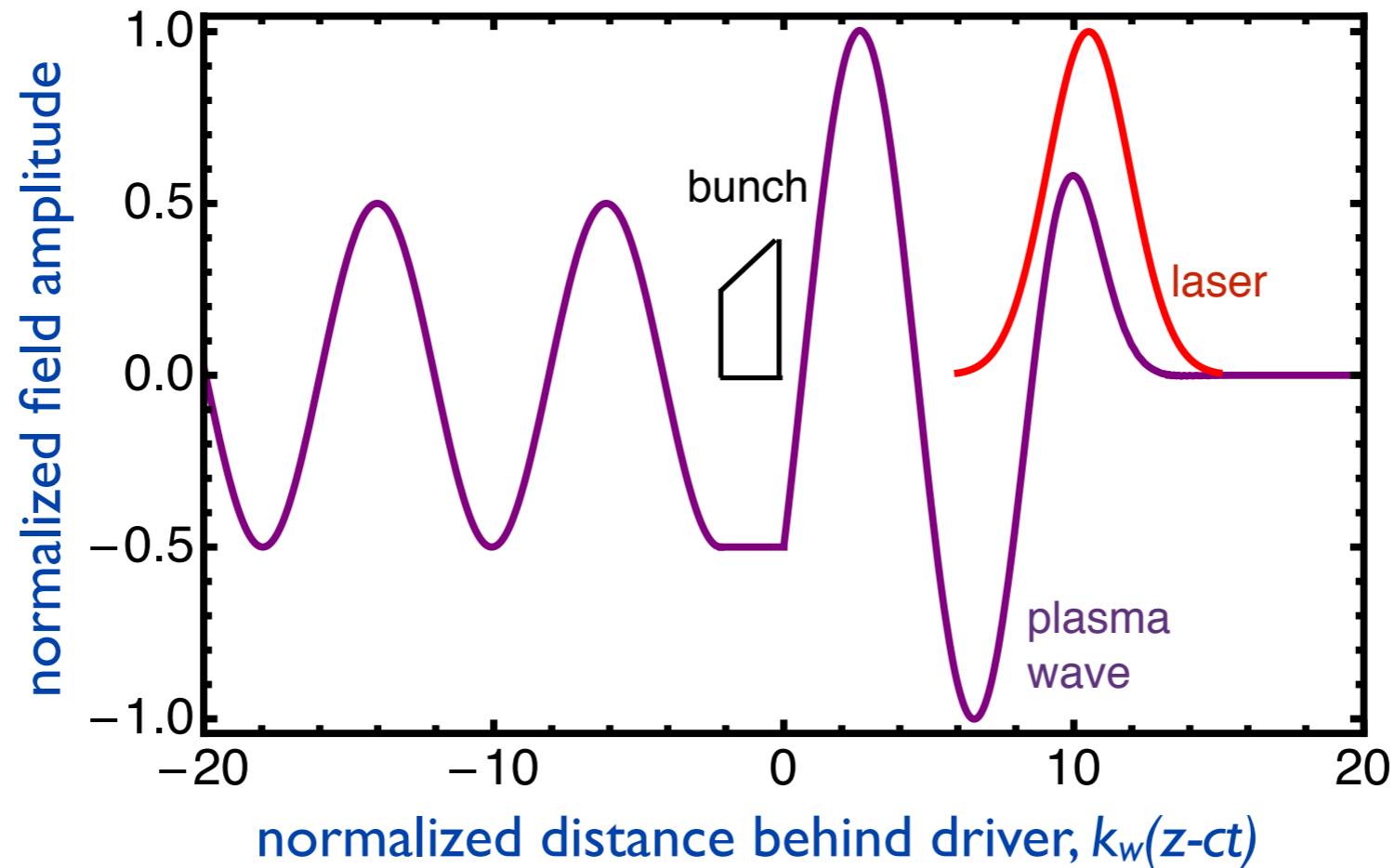
► Trade-off between gradient and efficiency, with no energy spread growth using ultra-short beams

# Improved efficiency using additional laser pulses to absorb remaining plasma wave energy



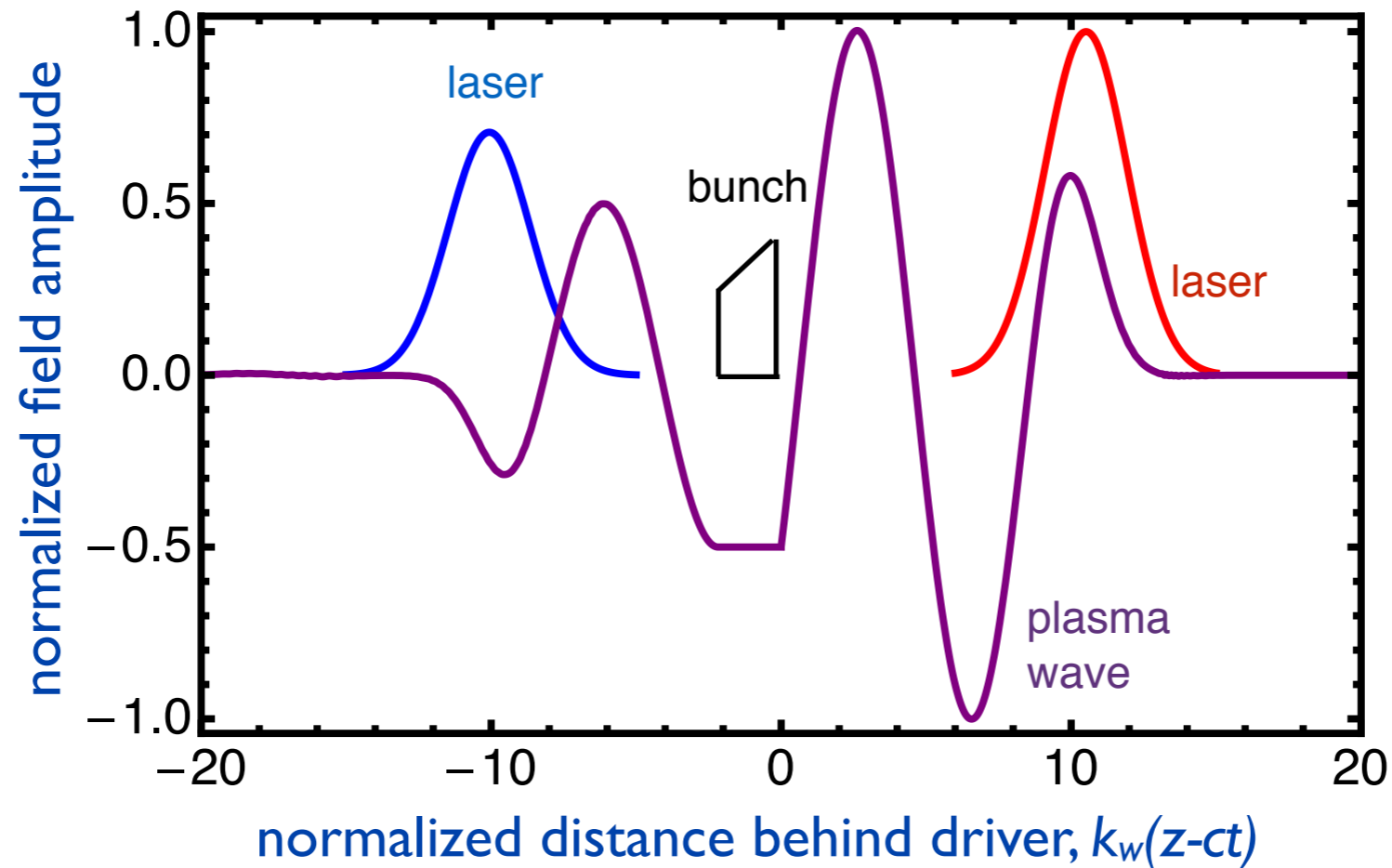
# Improved efficiency using additional laser pulses to absorb remaining plasma wave energy

$\eta(\text{plasma to beam}) = 0.75$   
gradient = 0.5(peak field)



# Improved efficiency using additional laser pulses to absorb remaining plasma wave energy

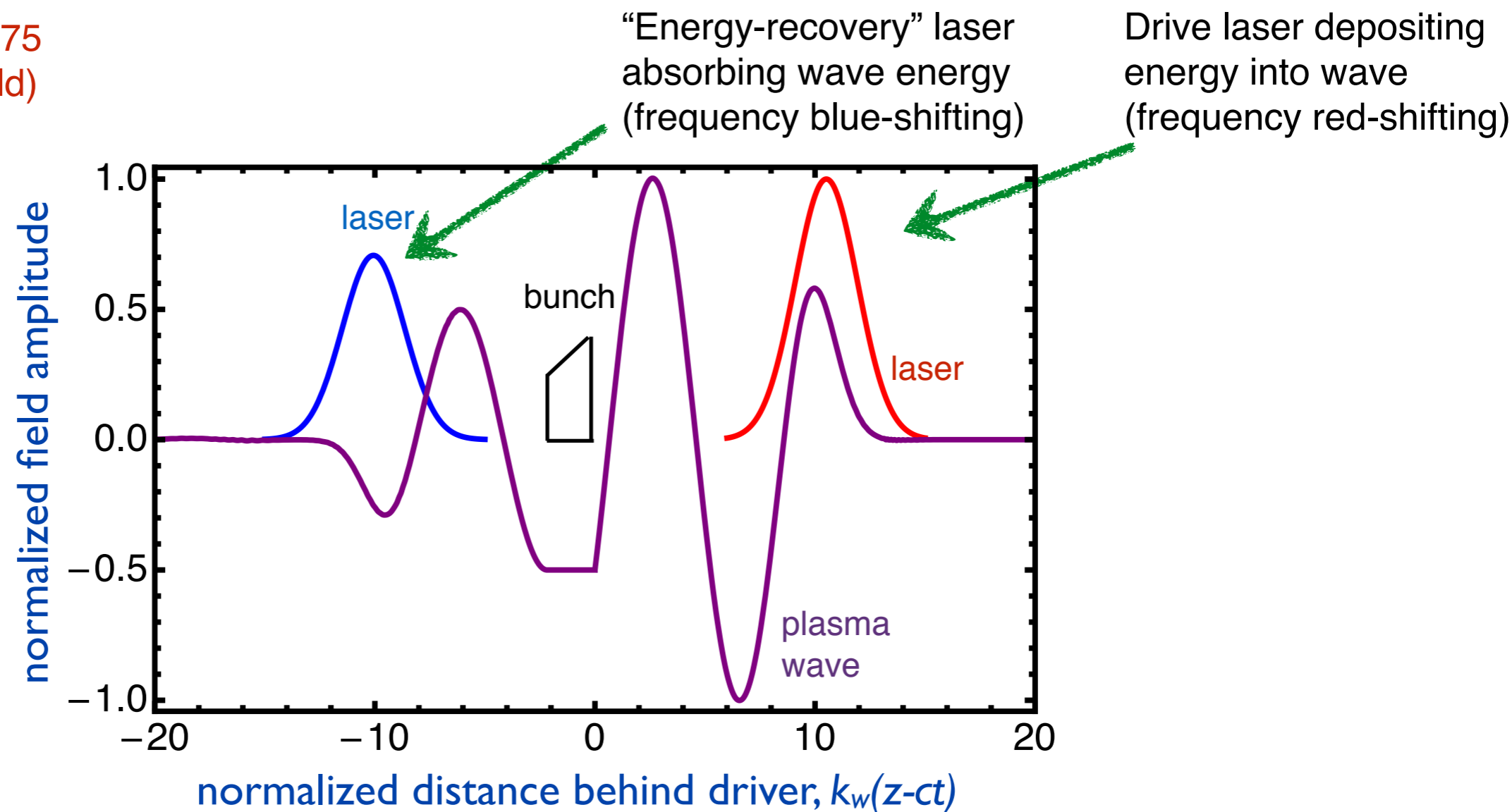
$\eta(\text{plasma to beam}) = 0.75$   
gradient = 0.5(peak field)





# Improved efficiency using additional laser pulses to absorb remaining plasma wave energy

$\eta(\text{plasma to beam}) = 0.75$   
gradient = 0.5(peak field)



- ▶ Additional laser pulse allows for no energy to remain in coherent plasma oscillations after energy transferred to particle beam

# Average power reduced at lower plasma density (Beamstrahlung limits charge/bunch)

Charge/bunch:

$$N \propto \frac{U_L}{\Delta\gamma} \propto n^{-1/2}$$



Laser rep rate (for fixed luminosity goal):

$$f \propto n$$



Wall-plug power:

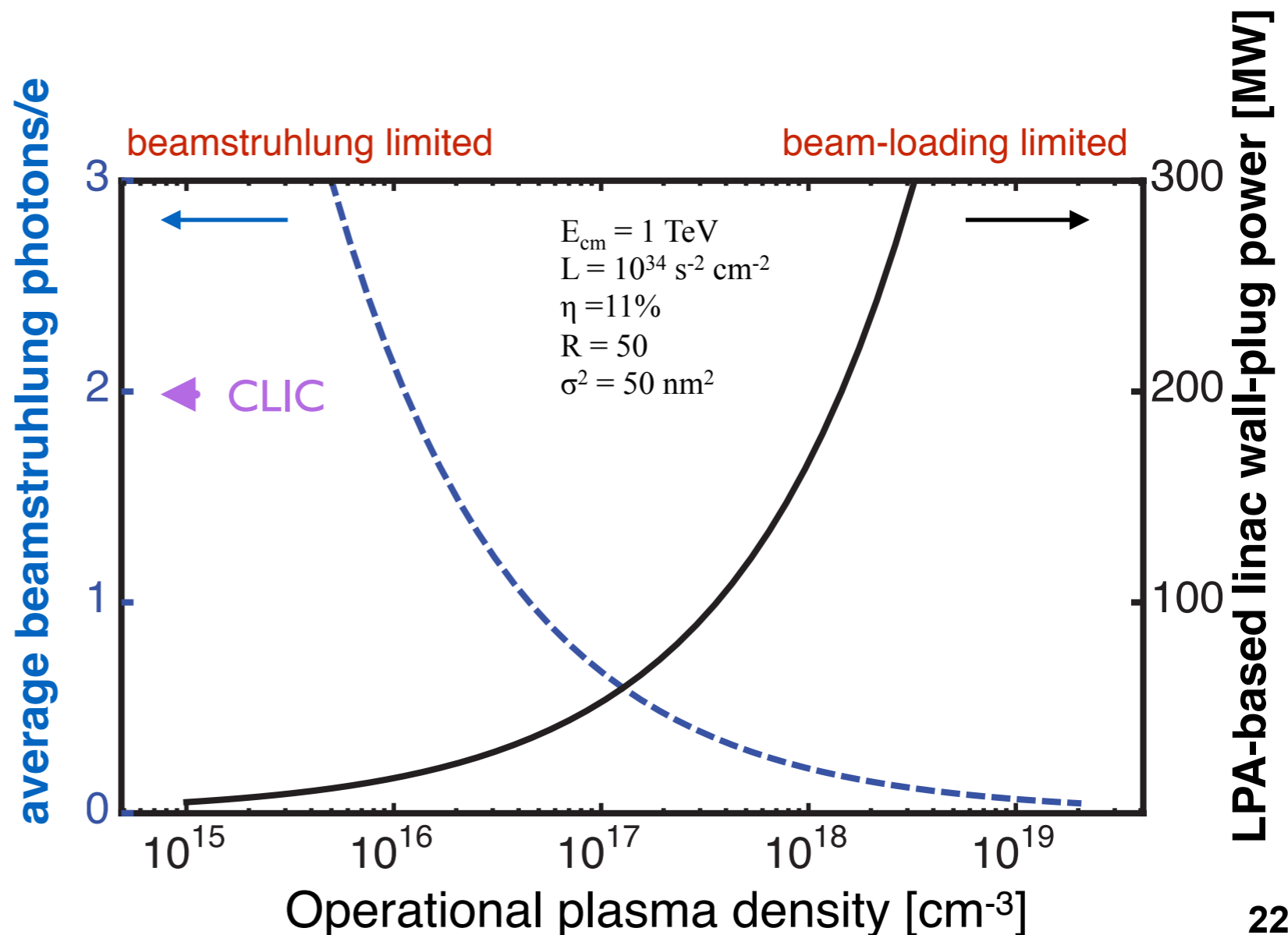
$$P \propto \sqrt{n}$$

*Schroeder et al., PR ST-AB (2010)*

Average beamstrahlung photons/e at IP:

*Schroeder et al., PR ST-AB (2012)*

$$n_\gamma \propto N^{2/3} \sigma_z^{1/3} \propto n^{-1/2}$$



# 3 TeV laser-plasma collider example

LPA stage laser-plasma parameters.

Plasma density (wall), $n_0$ [ $\text{cm}^{-3}$ ]	$10^{17}$
Plasma wavelength, $\lambda_p$ [mm]	0.1
Channel radius, $r_c$ [ $\mu\text{m}$ ]	22
Laser wavelength, $\lambda$ [ $\mu\text{m}$ ]	1
Normalized laser strength, $a_0$	1.2
Peak laser power, $P_L$ [TW]	50
Laser pulse duration (FWHM), $\tau_L$ [fs]	130
Laser energy, $U_L$ [J]	6.5
Normalized accelerating field, $E_L/E_0$	0.2
Peak accelerating field, $E_L$ [GV/m]	6
Laser depletion length, $L_{pd}$ [m]	8.7
Plasma channel length, $L_c$ [m]	1.7
Laser depletion, $\eta_{pd}$ [%]	20

operational plasma density

Shaped electron/positron beam parameters.

Bunch phase (relative to peak field), $\varphi$	$\pi/3$
Loaded gradient, $E_z$ [GV/m]	3
Beam beam current, $I$ [kA]	3
Charge/bunch, $eN_b = Q$ [nC]	0.19
Length (triangular shape), $L_b$ [ $\mu\text{m}$ ]	36

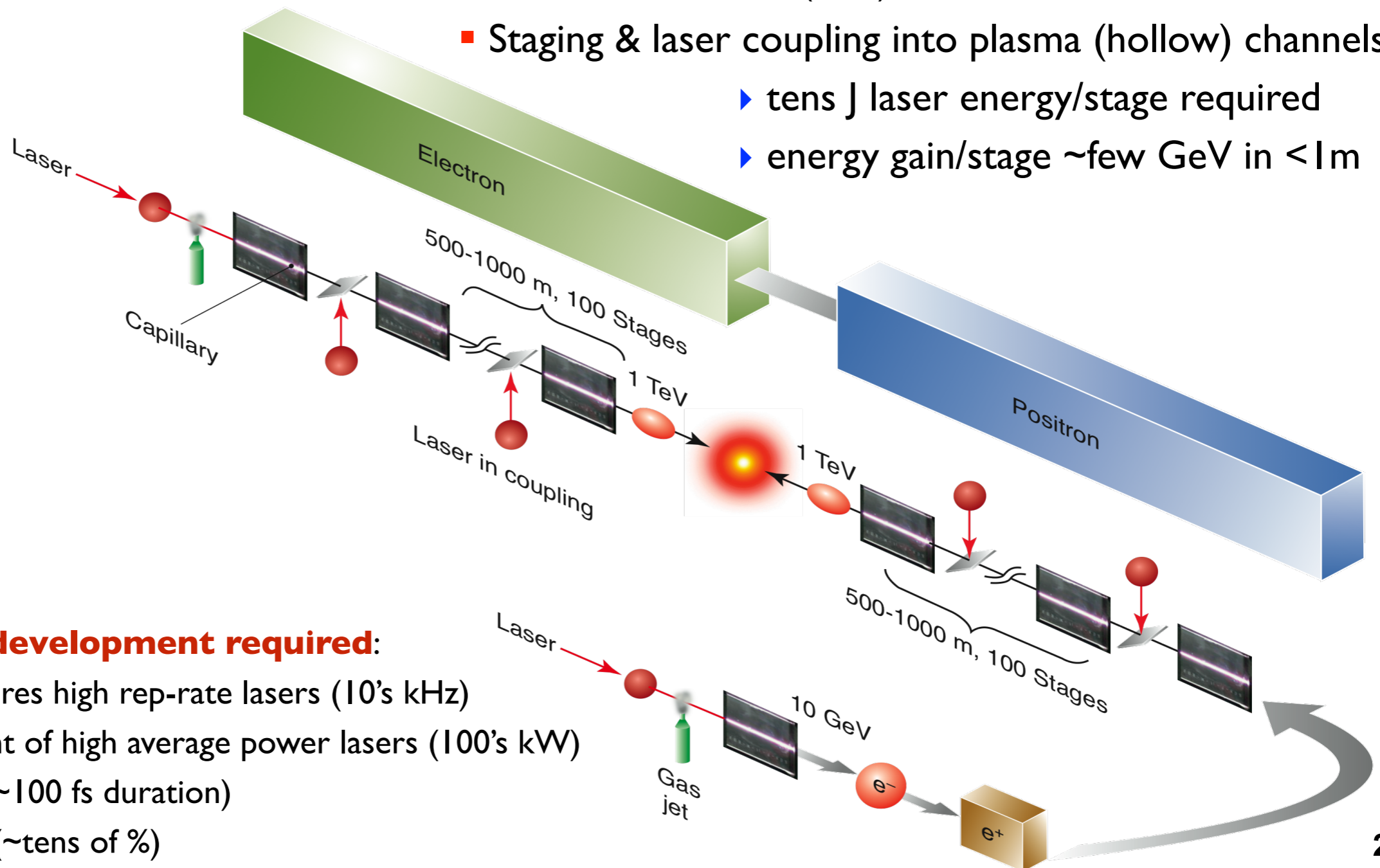
beam energy	1.5 TeV
energy gain/stage	5 GeV
RMS $\sigma_z$ [micron]	8.5
$\sigma_x$ [nm]/ $\sigma_y$ [nm]	18 / 0.5
Luminosity [ $\text{s}^{-1}\text{cm}^{-2}$ ]	1E+35
laser rep rate [kHz]	84
Beamstrahlung parameter	16
Beamstrahlung photons/e	0.81
Beamstrahlung energy spread	0.20
Disruption parameter	0.046
Power (LPA linacs)	407 MW

# Laser-plasma accelerator-based collider concept

Schroeder et al., PR STAB (2010)

Leemans & Esarey, Physics Today (2009)

- Plasma density scalings [minimize **construction** (max. average gradient) and **operational** (min. wall power) costs] indicates:  $n \sim 10^{17} \text{ cm}^{-3}$
- Quasi-linear wake ( $a \sim 1$ ):  $e^-$  and  $e^+$
- Staging & laser coupling into plasma (hollow) channels:
  - ▶ tens J laser energy/stage required
  - ▶ energy gain/stage  $\sim$  few GeV in  $< 1 \text{ m}$



## Laser technology development required:

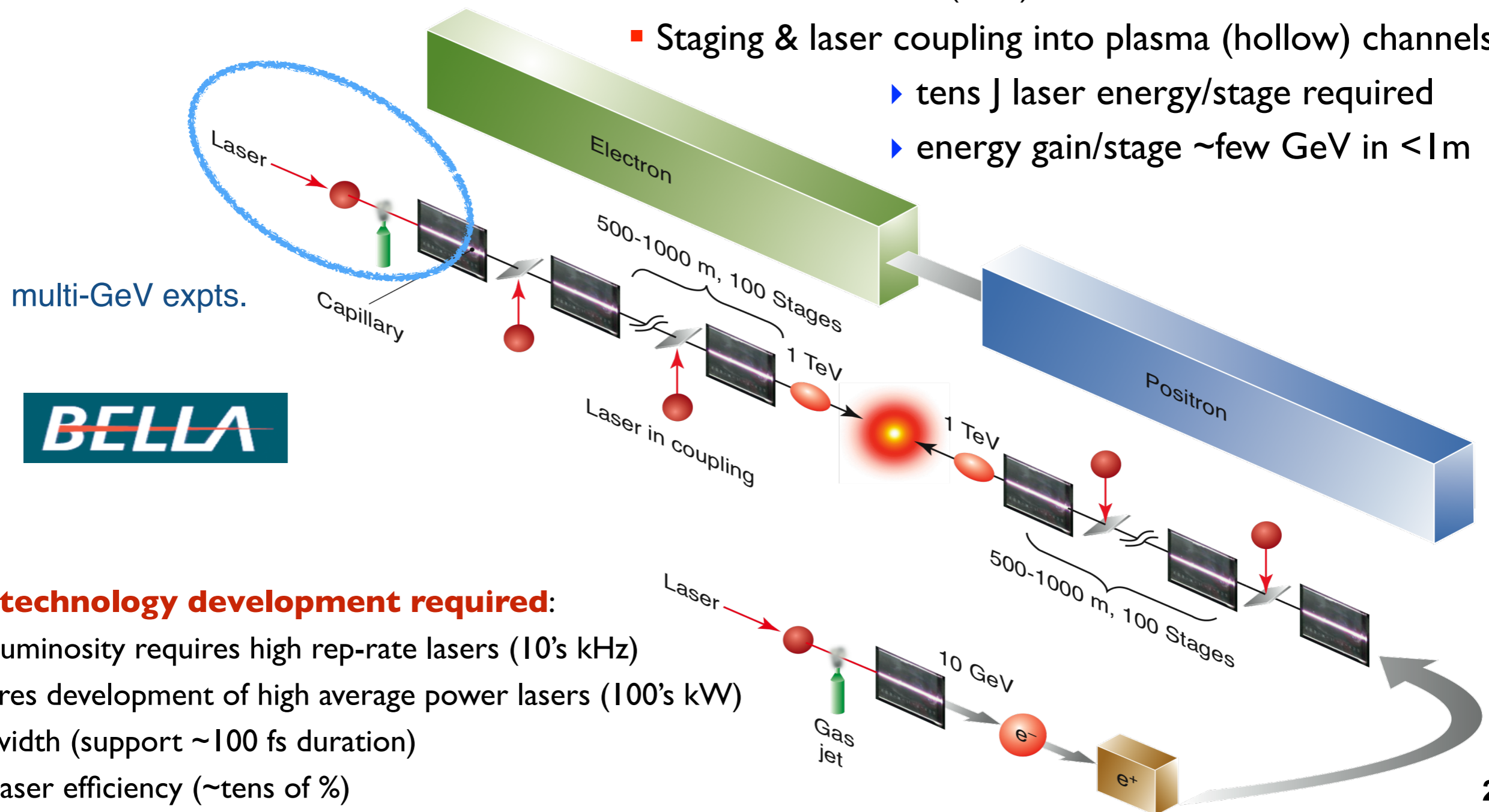
- High luminosity requires high rep-rate lasers (10's kHz)
- Requires development of high average power lasers (100's kW)
- Bandwidth (support  $\sim 100 \text{ fs}$  duration)
- High laser efficiency ( $\sim$  tens of %)

# Laser-plasma accelerator-based collider concept

Schroeder et al., PR STAB (2010)

Leemans & Esarey, Physics Today (2009)

- Plasma density scalings [minimize **construction** (max. average gradient) and **operational** (min. wall power) costs] indicates:  $n \sim 10^{17} \text{ cm}^{-3}$
- Quasi-linear wake ( $a \sim 1$ ):  $e^-$  and  $e^+$
- Staging & laser coupling into plasma (hollow) channels:
  - ▶ tens J laser energy/stage required
  - ▶ energy gain/stage  $\sim$  few GeV in  $< 1 \text{ m}$



## Laser technology development required:

- High luminosity requires high rep-rate lasers (10's kHz)
- Requires development of high average power lasers (100's kW)
- Bandwidth (support  $\sim 100 \text{ fs}$  duration)
- High laser efficiency ( $\sim$  tens of %)

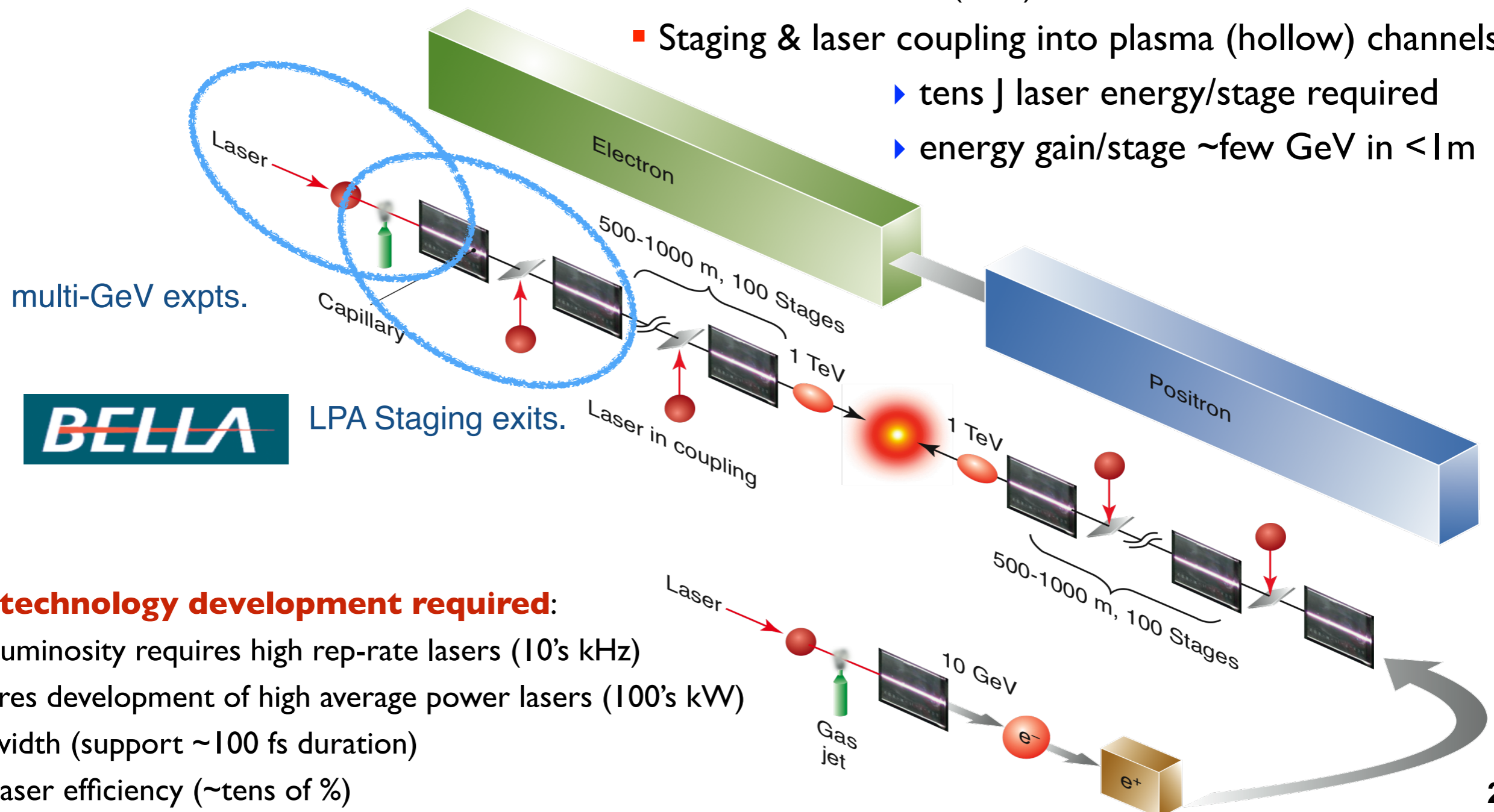


# Laser-plasma accelerator-based collider concept

Schroeder et al., PR STAB (2010)

Leemans & Esarey, Physics Today (2009)

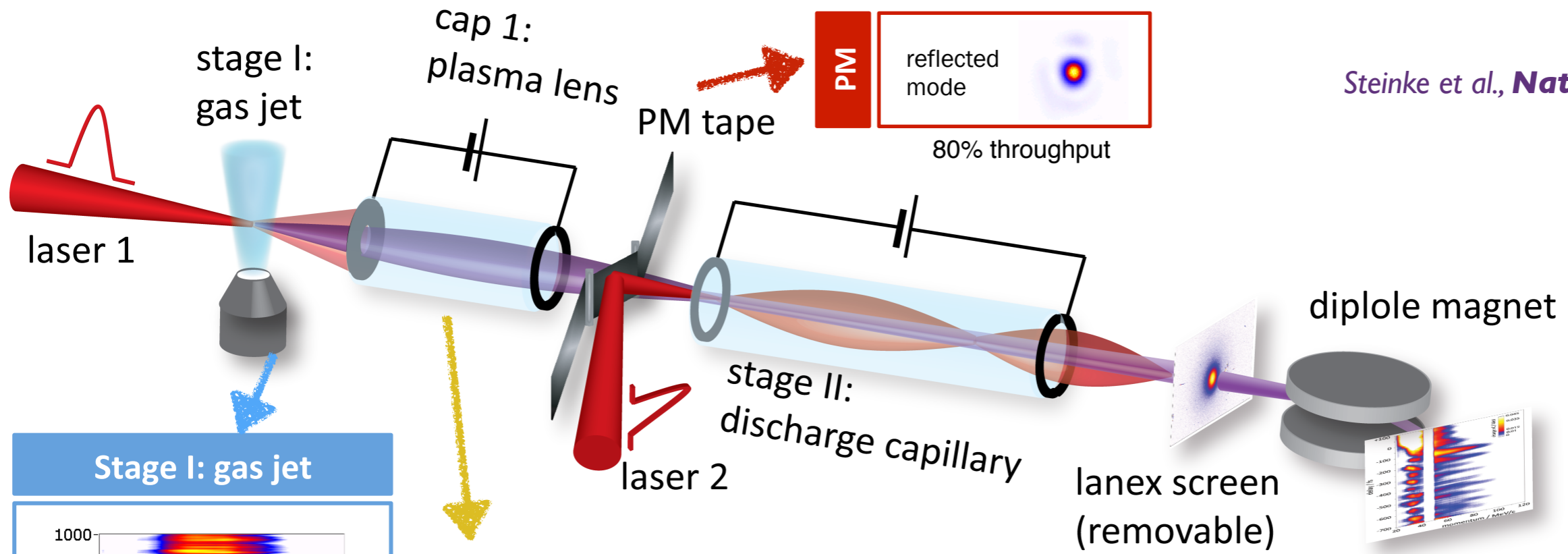
- Plasma density scalings [minimize **construction** (max. average gradient) and **operational** (min. wall power) costs] indicates:  $n \sim 10^{17} \text{ cm}^{-3}$
- Quasi-linear wake ( $a \sim 1$ ):  $e^-$  and  $e^+$
- Staging & laser coupling into plasma (hollow) channels:
  - ▶ tens J laser energy/stage required
  - ▶ energy gain/stage  $\sim$  few GeV in  $< 1 \text{ m}$



## Laser technology development required:

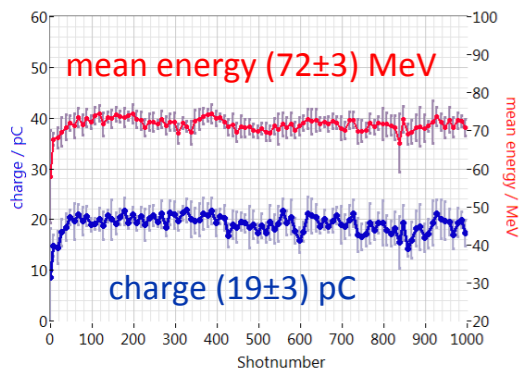
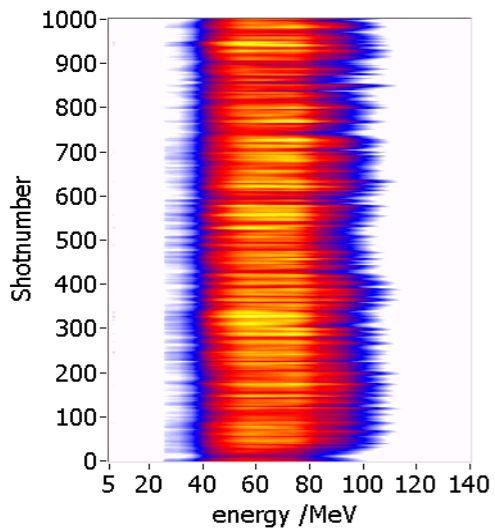
- High luminosity requires high rep-rate lasers (10's kHz)
- Requires development of high average power lasers (100's kW)
- Bandwidth (support  $\sim 100 \text{ fs}$  duration)
- High laser efficiency ( $\sim$  tens of %)

# Demonstration of acceleration in second independently-powered laser-plasma accelerator at LBNL

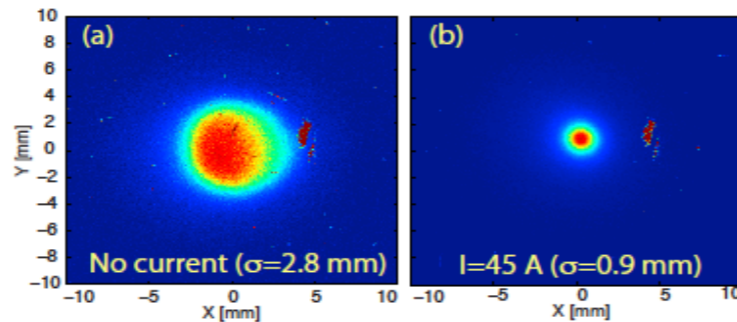
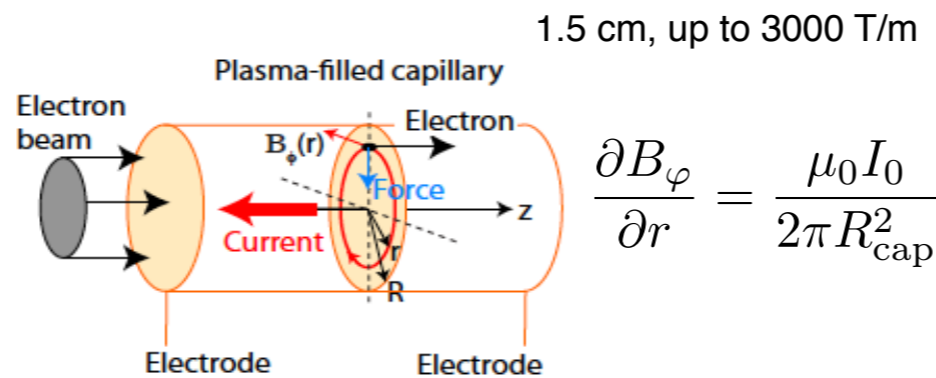


Steinke et al., *Nature* (2016)

## Stage I: gas jet

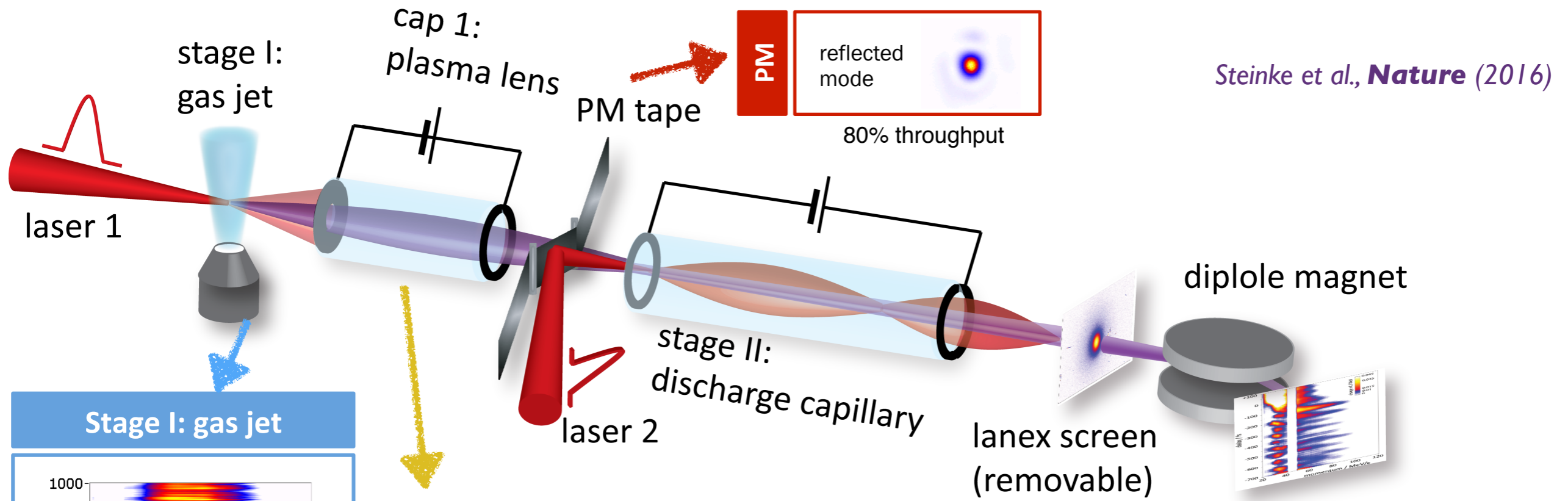


## Cap 1: active plasma lens

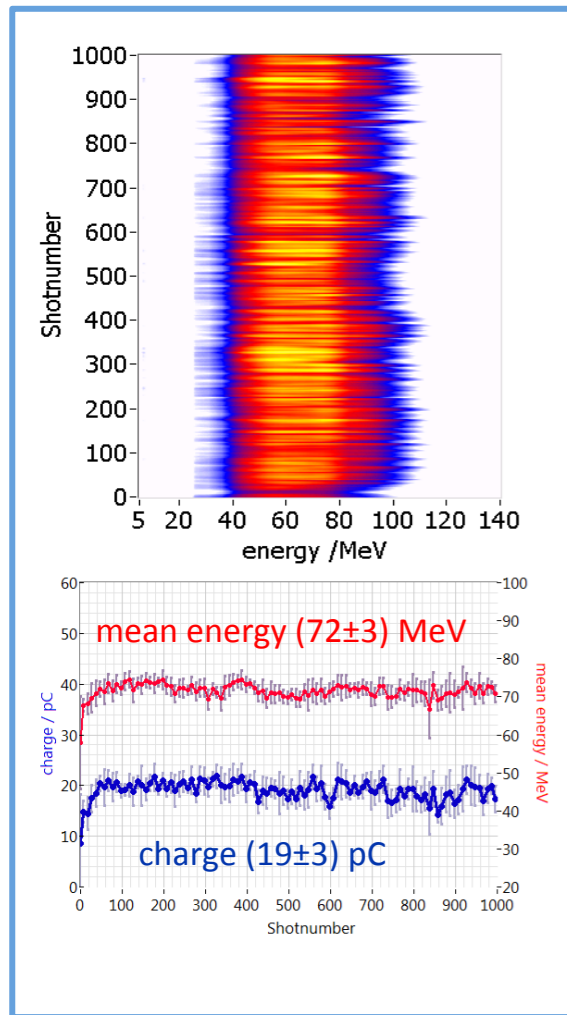


J. van Tilborg et al., *PRL* (2015)

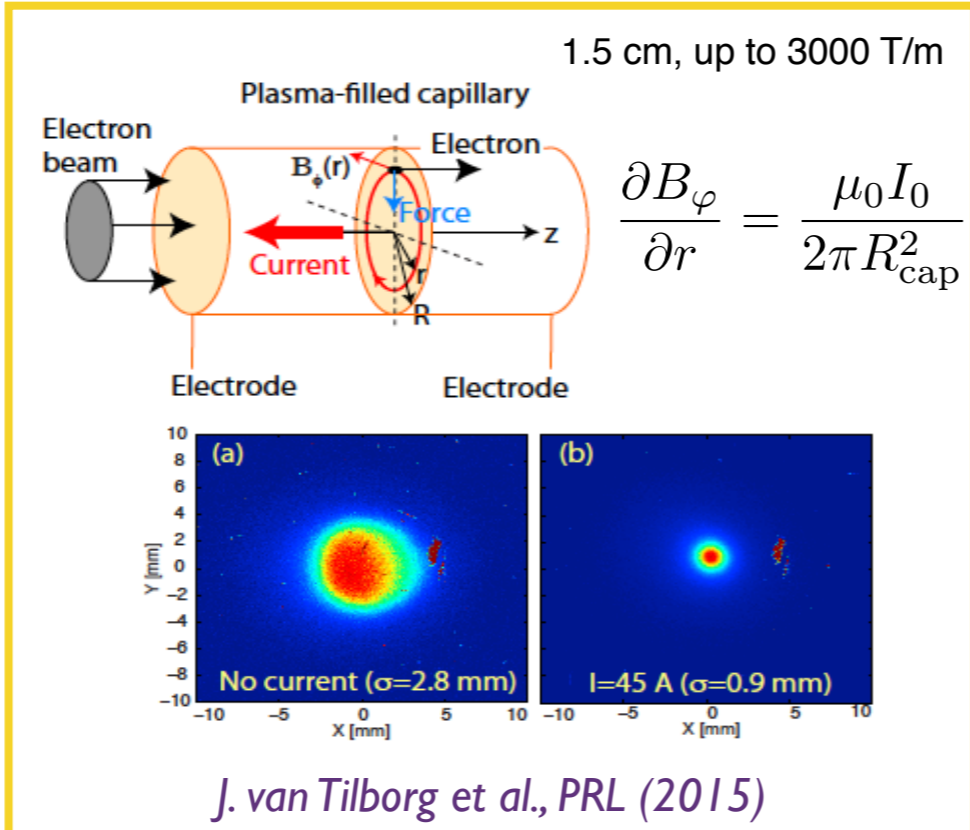
# Demonstration of acceleration in second independently-powered laser-plasma accelerator at LBNL



Stage I: gas jet

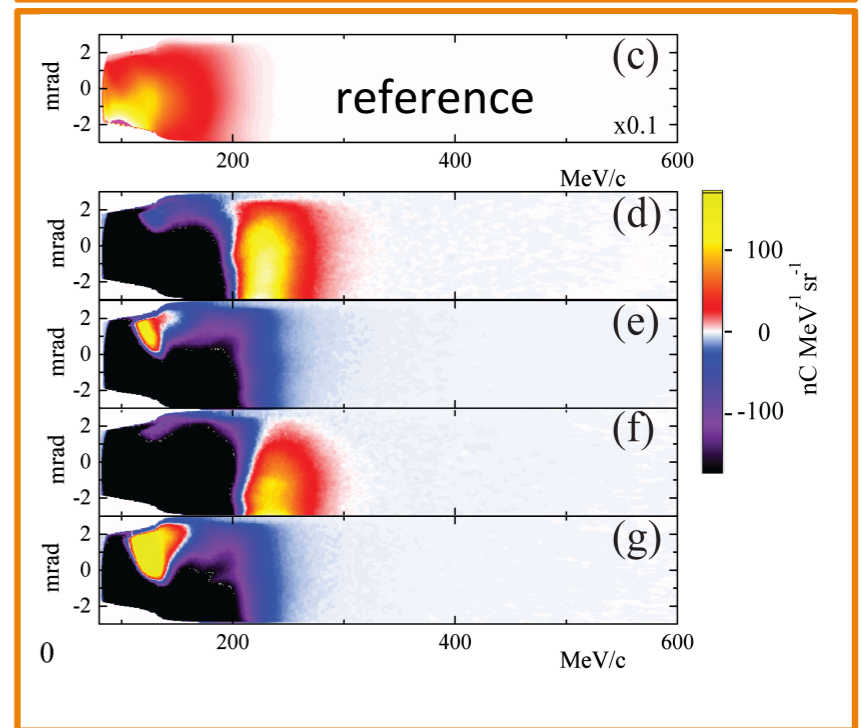


Cap 1: active plasma lens



lanex screen

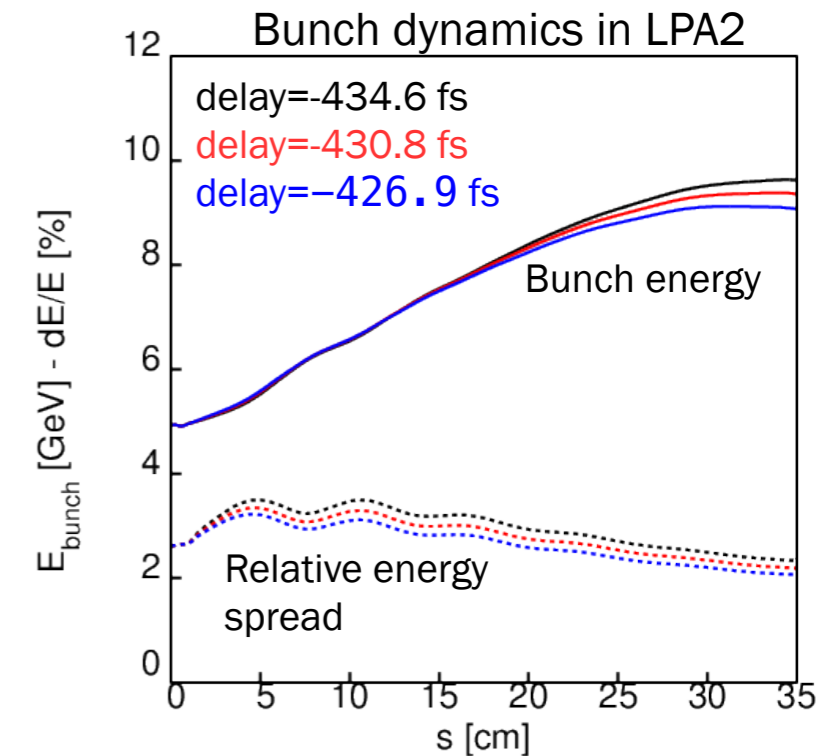
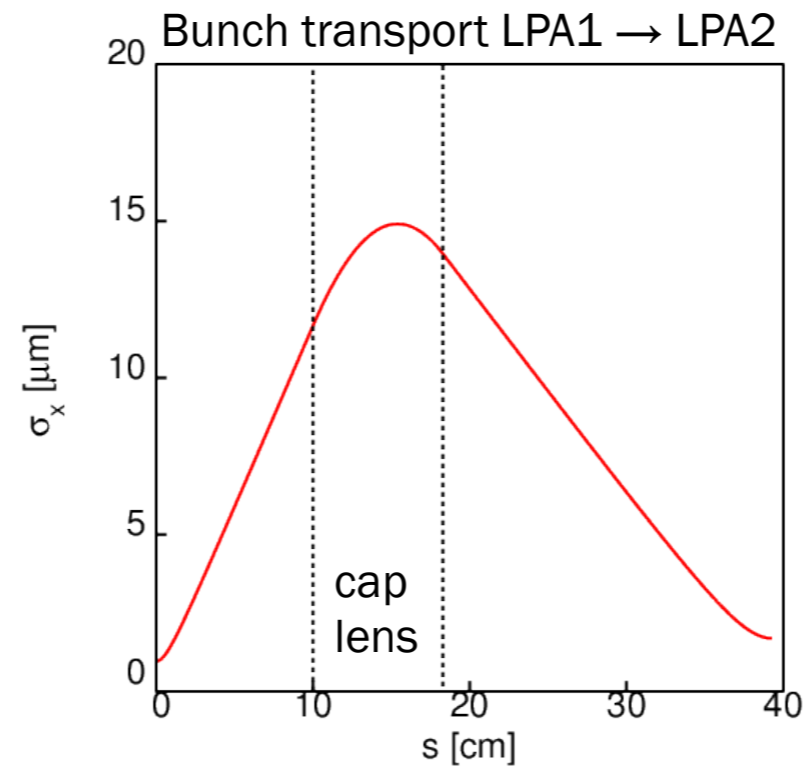
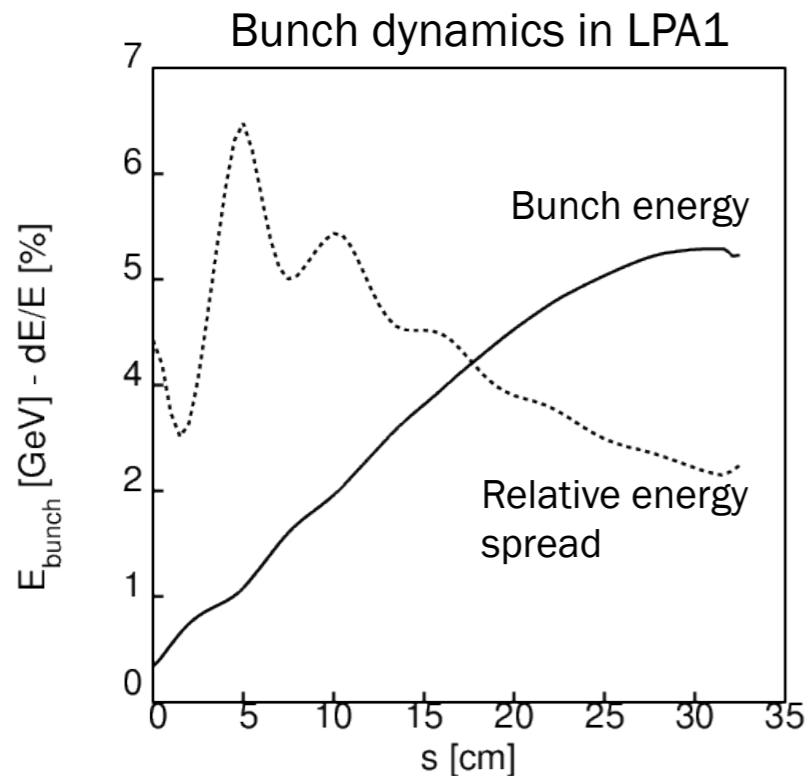
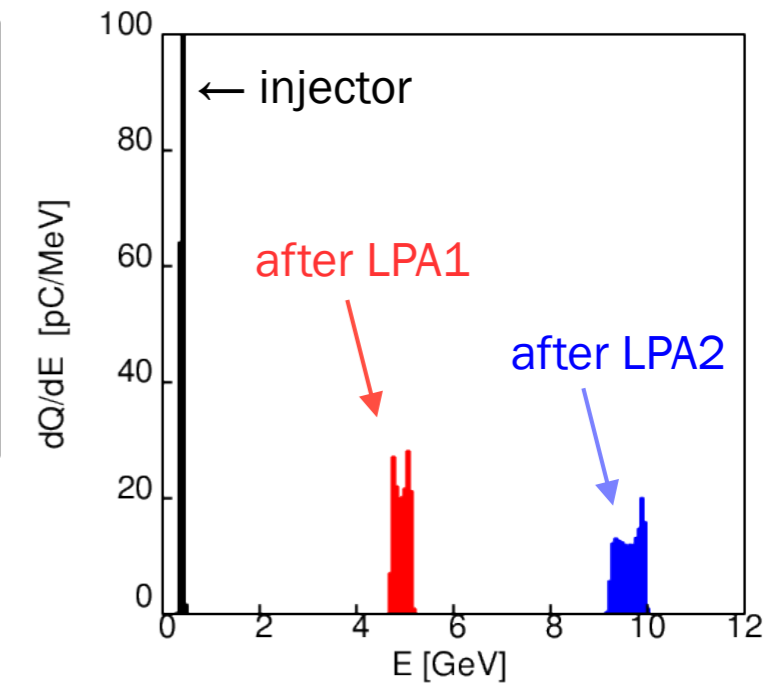
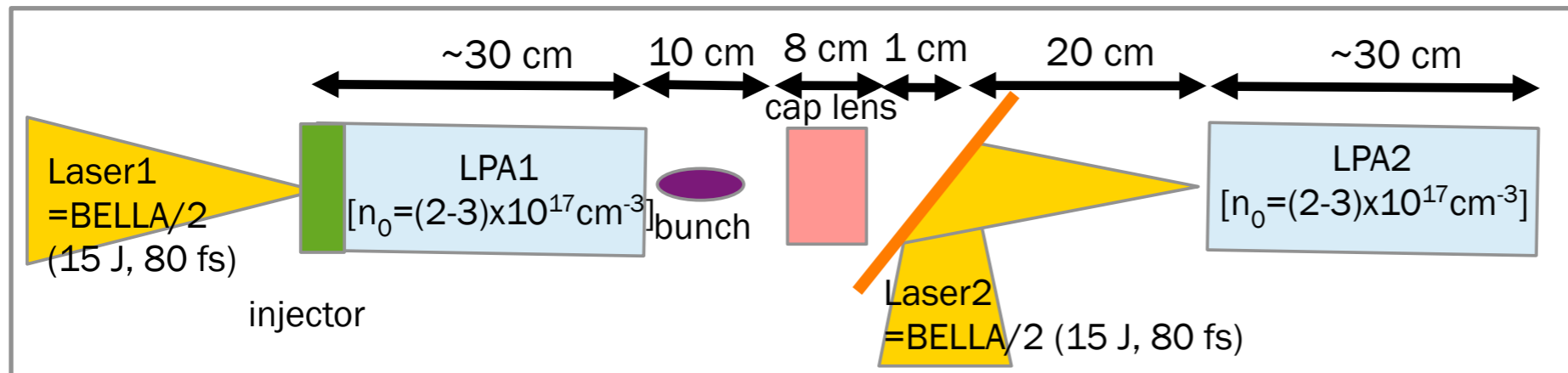
Stage I + II





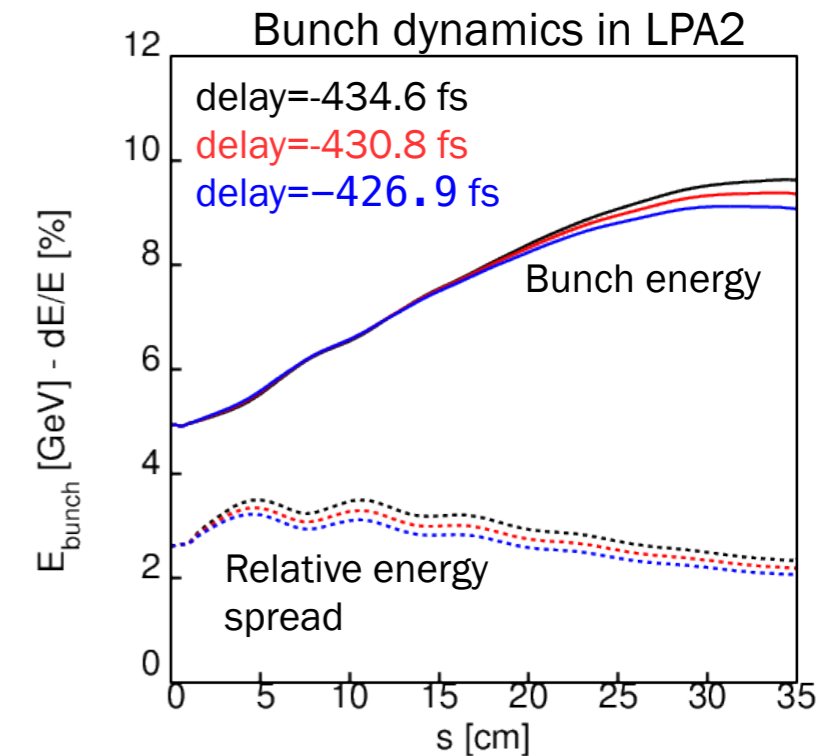
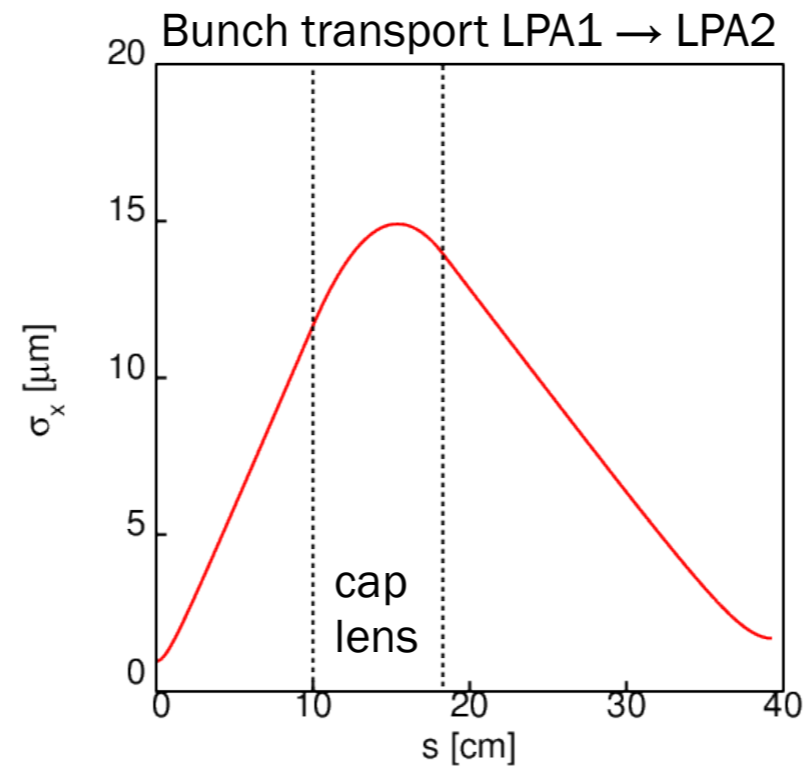
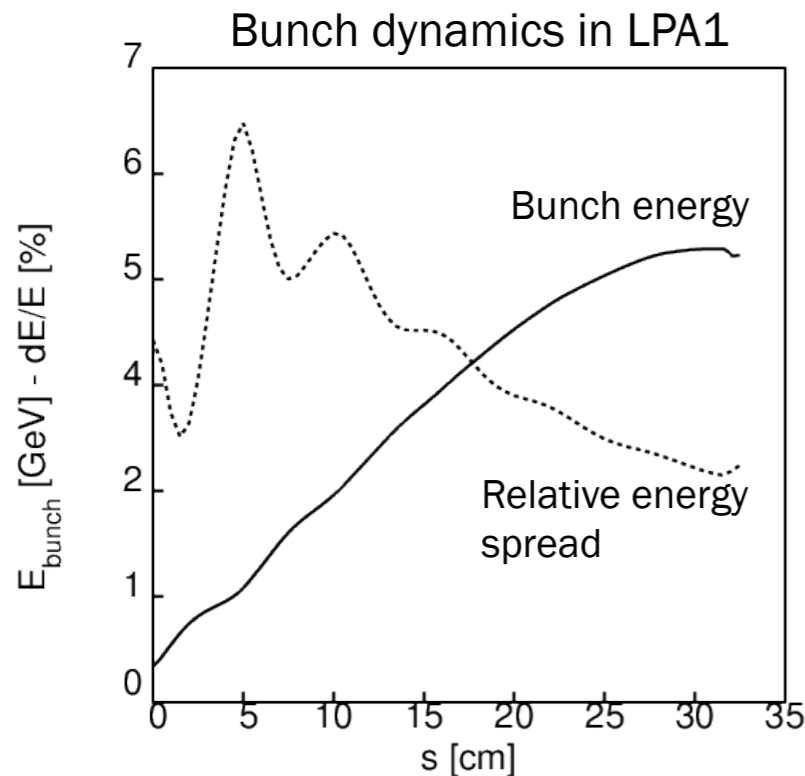
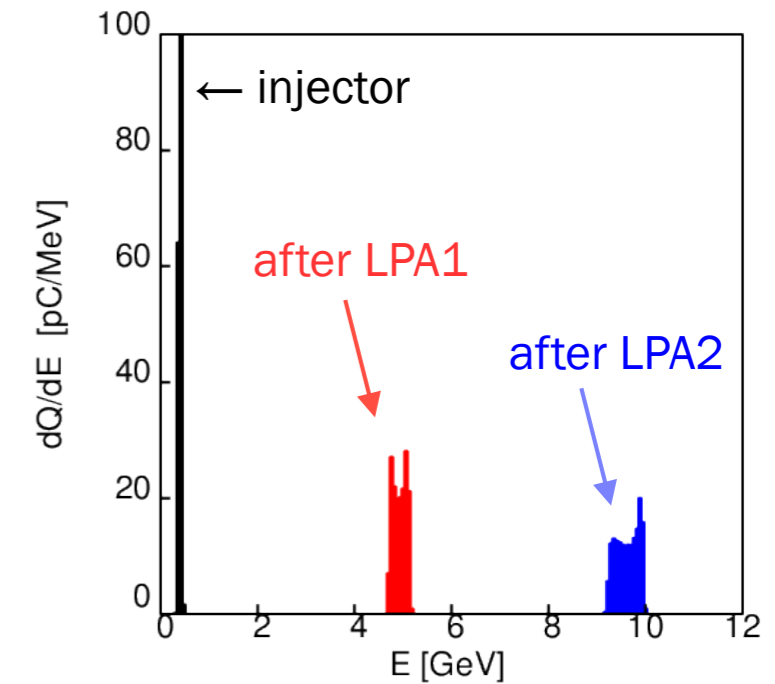
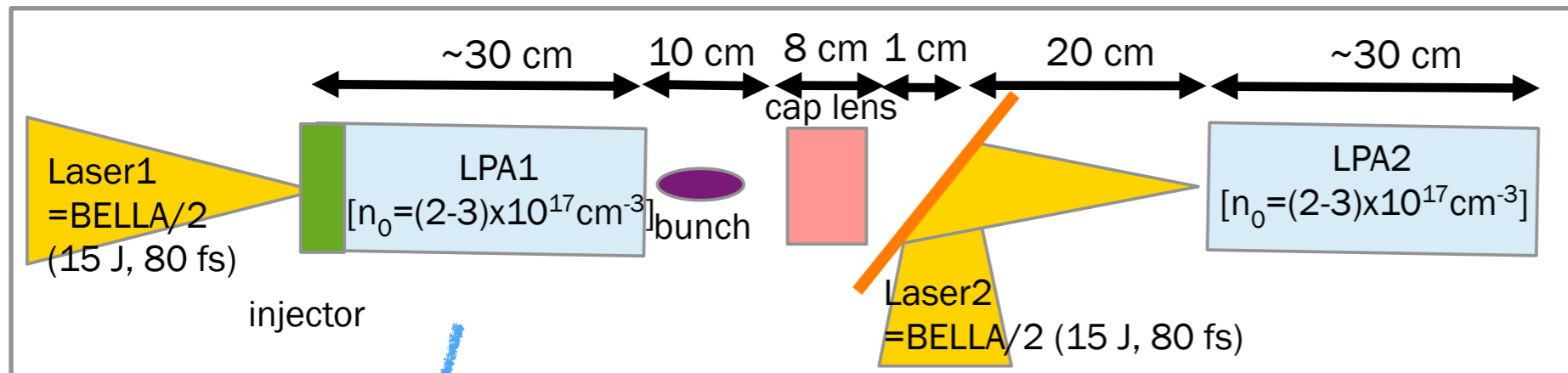
# Future LPA staging experiments: 5 GeV + 5 GeV

Steinke et al., Phys. Plasmas (2016)



# Future LPA staging experiments: 5 GeV + 5 GeV

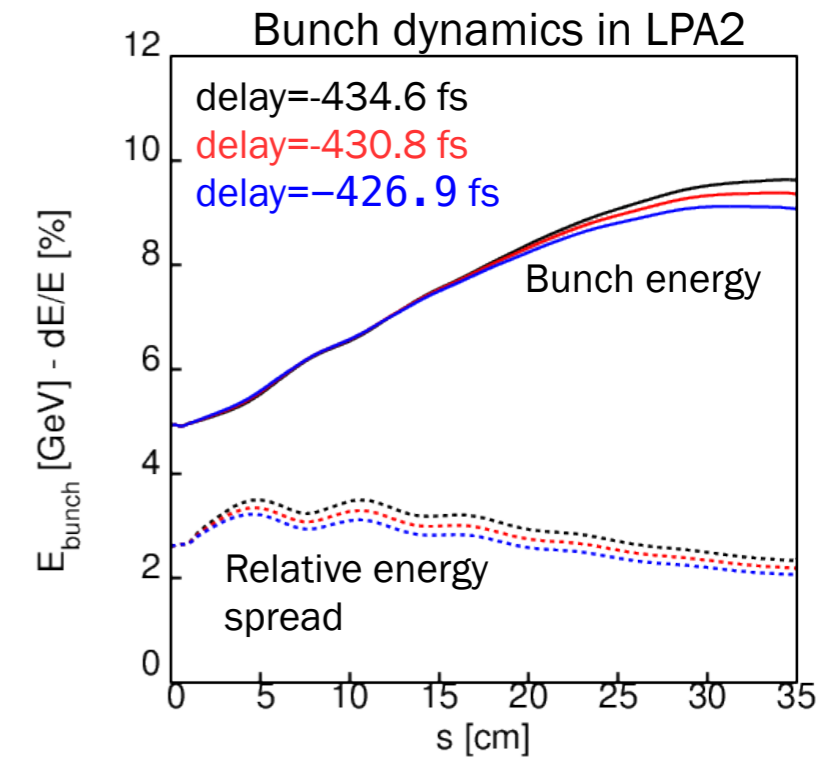
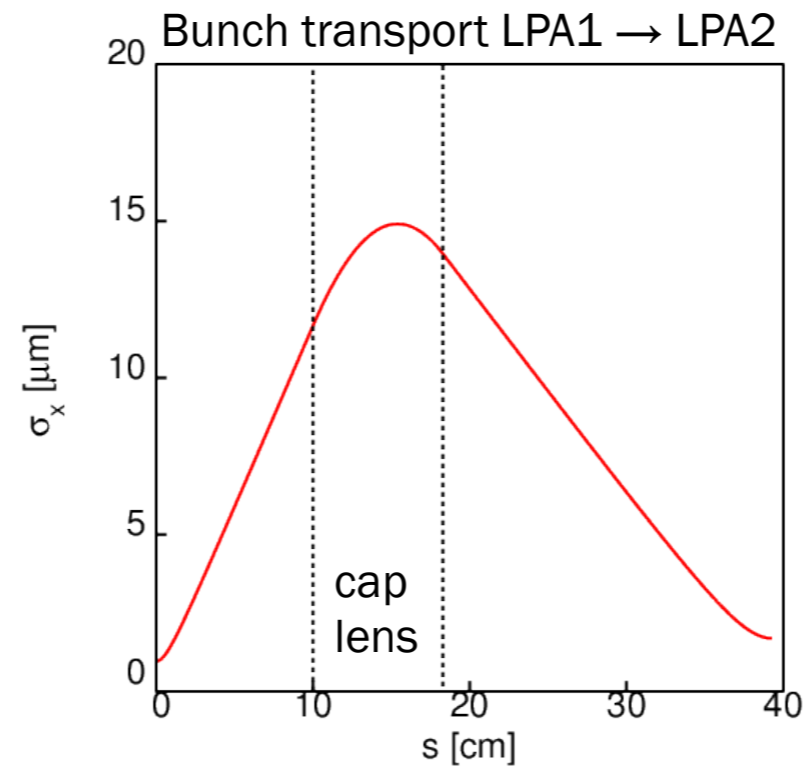
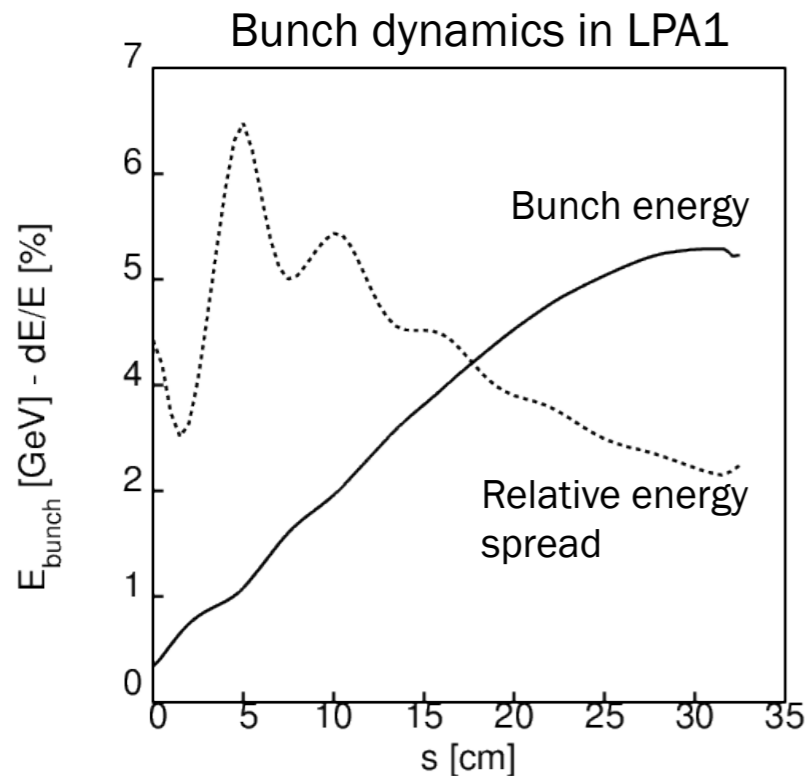
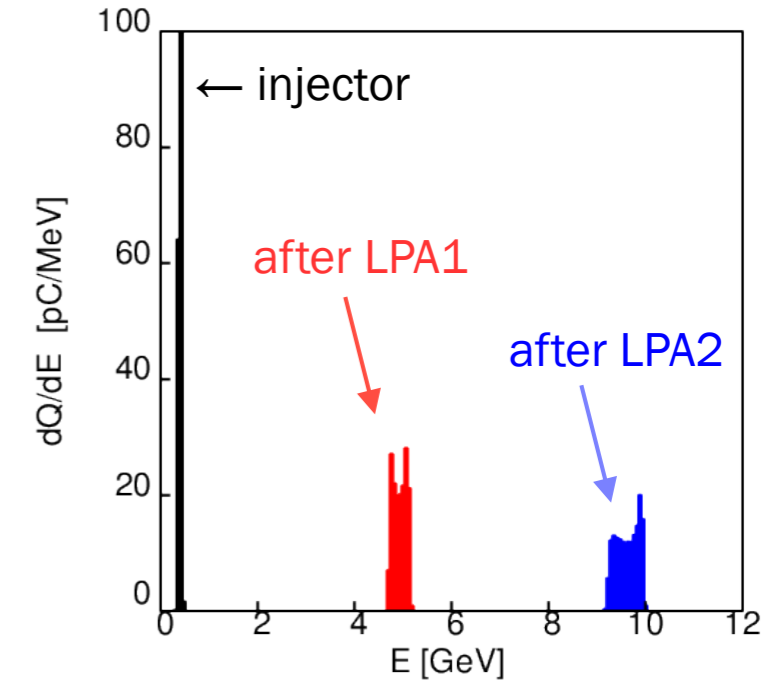
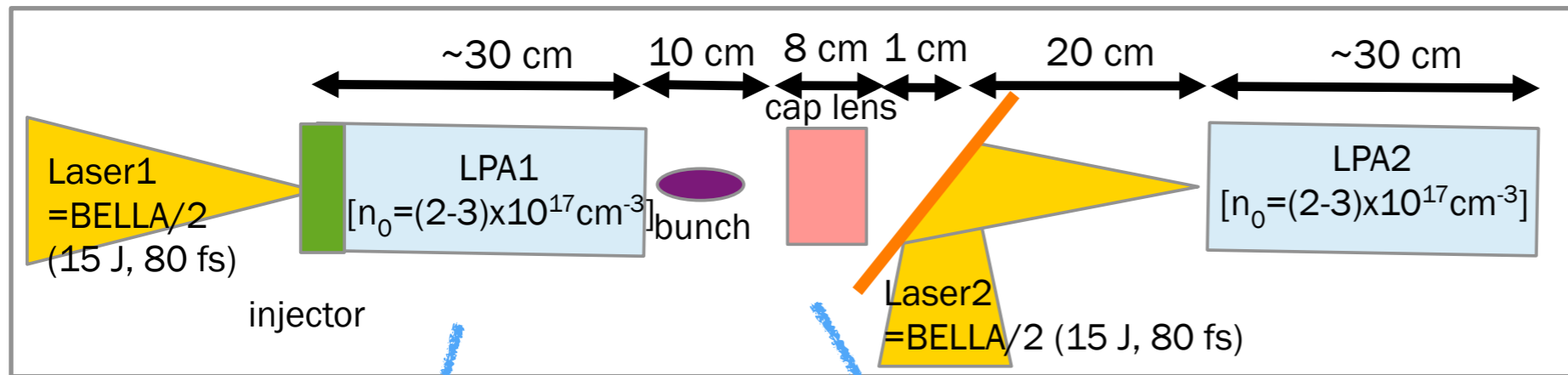
Steinke et al., Phys. Plasmas (2016)





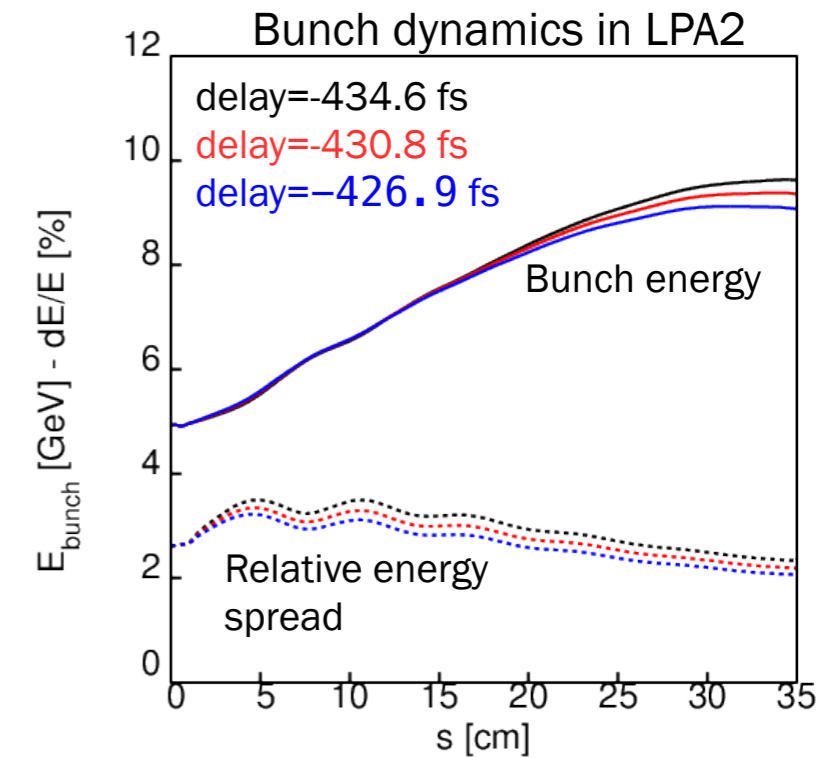
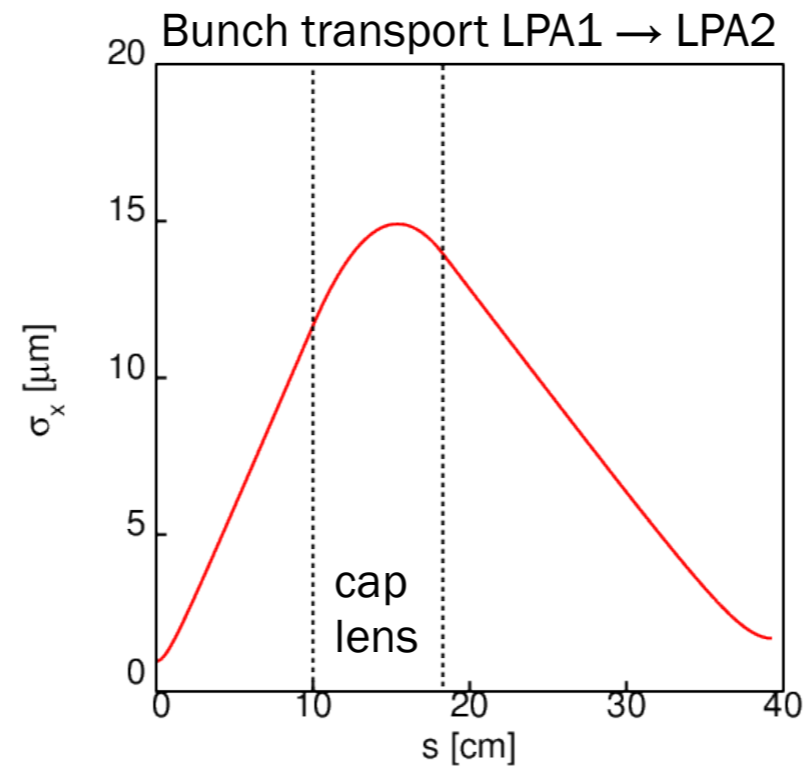
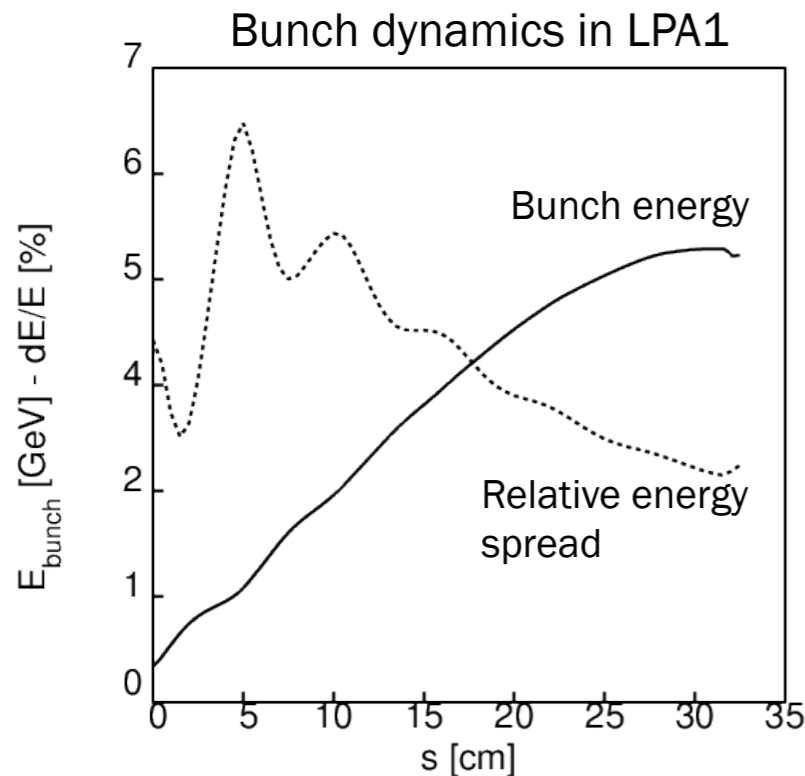
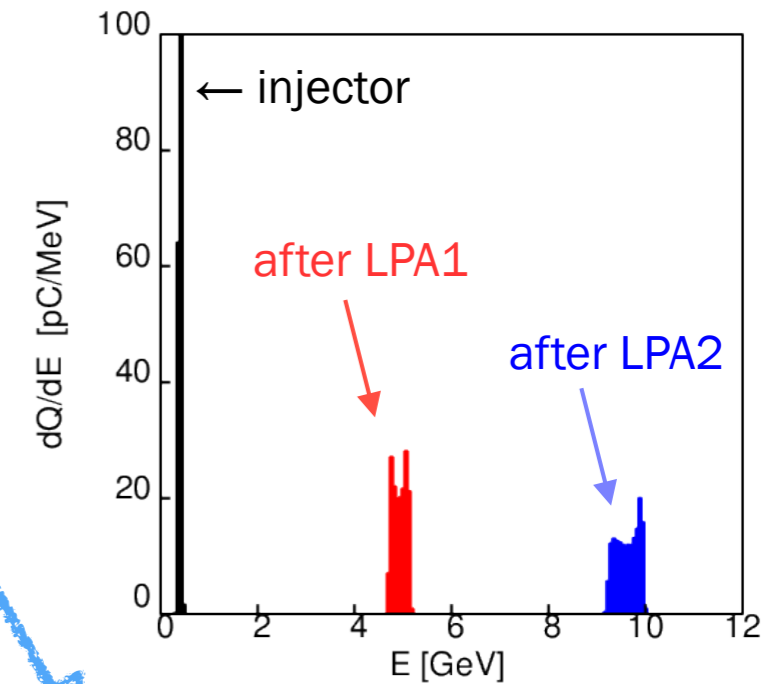
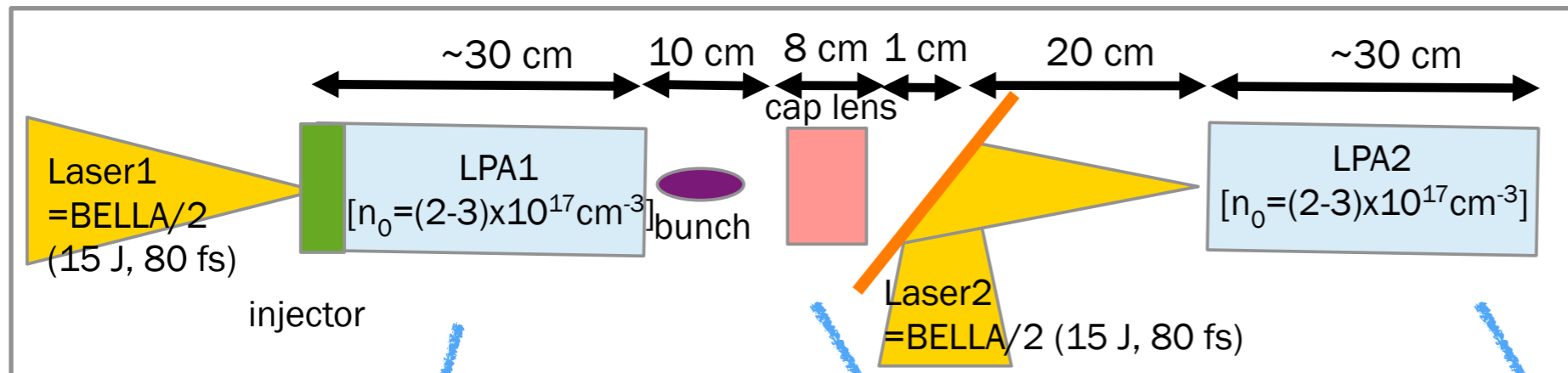
# Future LPA staging experiments: 5 GeV + 5 GeV

Steinke et al., Phys. Plasmas (2016)



# Future LPA staging experiments: 5 GeV + 5 GeV

Steinke et al., Phys. Plasmas (2016)



# Summary and Conclusions

- ▶ Laser-plasma accelerators provide ultra-high gradients (compact accelerators) generating short (fs) beams (high peak current)
  - 4 GeV beams in 9-cm plasma using LPA at BELLA
    - 10 GeV beams in <1m will be available in next (few) years
- ▶ Beam-beam interaction effects modify the basic scalings and determine the operational plasma density (determines required laser technology, etc.)
  - Basic scalings indicate  $\sim 10^{17}/\text{cc}$  operating density.
  - Operating at low plasma density increases beamstrahlung effects (higher bunch charge and longer beams)
  - Bunch charge constrained by beamstrahlung - requiring multi-bunch format or increase repetition frequency (power).
  - Detailed modeling still to be done (including beam shape, delivery system, etc.)
- ▶ Laser-plasma accelerator-based linear collider has many technical challenges and R&D for collider application ongoing
  - Staging LPA experiments underway at LBNL
  - High peak & average power laser technology development



Many thanks to my colleagues:

Carlo Benedetti

Eric Esarey

Cameron Geddes

Wim Leemans

Jean-Luc Vay

and the members of the *BELLA Center*

Work supported by the U.S. DOE, Office of Science, Office of High Energy Physics,  
under Contract No. DE-AC02-05CH11231



**BERKELEY LAB**

LAWRENCE BERKELEY NATIONAL LABORATORY

ACCELERATOR TECHNOLOGY &  
APPLIED PHYSICS DIVISION



U.S. DEPARTMENT OF  
**ENERGY**

Office of Science