## Liquid Xenon Scintillation Measurements and Pulse Shape Discrimination in the LUX Dark Matter Detector

Brian Lenardo UC Davis / LLNL LIght Detection In Noble Elements September 22, 2017





## The LUX dark matter experiment



**LIDINE 2017** 

122 low-radioactivity photomultiplier tubes for high-efficiency light collection

## Scintillation PSD analysis in LUX

### Goals:

- Measure electron recoil and nuclear recoil scintillation time distributions at DM search energies
- Develop pulse shape discriminant for use in future dark matter searches with LUX

## **ER/NR calibration data**

### Fast neutrons from D-D generator

- Collimated 2.45 MeV beam
- Elastic scatters  $\rightarrow$  nuclear recoils 0 74 keV.

### Beta decays from <sup>3</sup>H and <sup>14</sup>C source

- Methane with <sup>3</sup>H or <sup>14</sup>C dissolved into liquid xenon circulation system
- Source removed by standard purification system
- Populates detector uniformly with electron recoils from beta decay (0 150 keV)



### LUX PRELIMINARY

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## Photon timing algorithm



## Timing calibration using LEDs



Six blue LEDs (440nm) embedded in each PMT array at top and bottom of xenon volume.

### Timing calibration procedure:

- 1. Pulse LEDs with 20 ns FWHM pulses to produce light at known times
- 2. Build distribution of photon arrivals in each channel
- 3. Use rising edge (10% height) as reference time to correct for relative timing offsets
- 4. Repeat with 4 different LEDs to compute uncertainties in calibration



Relative offsets between channels: ~20 ns Uncertainty in calibration:  $\sigma = 2$  ns

## Modeling pulse shapes

## **Detector effects:**

- Timing resolution
  - PMT transit time spread
  - Uncertainties in timing calibrations
  - Uncertainties from template fitting
- Optical transport
  - Modeled using optical simulations in Geant4
  - Empirical analytic model fit to simulation

## Liquid xenon physics:

- Ratio and time constants of singlet and triplet states

$$\begin{cases} P_r = \frac{1}{\sqrt{2\sigma^2}} e^{-x^2/(2\sigma^2)} \\ P_o(t) = A \,\delta(t) + (1-A) \left[ \frac{B_a}{\tau_a} e^{-t/\tau_a} + \frac{B_b}{\tau_b} e^{-t/\tau_a} \right] \\ P(t) = C_1 \, e^{-t/\tau_1} + C_3 \, e^{-t/\tau_3} \end{cases}$$

 $\tau_h$ 

## Fitting to calibration data



Fully convolved pulse shape model P(t) =  $\sum_{i=1,3} \sum_{j=a,b} \frac{C_i A}{2} e^{\frac{\sigma^2}{2\tau_i^2} - \frac{t}{\tau_i}} \left[ 1 + \operatorname{erf}\left(\frac{t - \frac{\sigma^2}{\tau_i}}{\sigma\sqrt{2}}\right) \right] +$   $\frac{C_i (1 - A) B_j}{2(\frac{\tau_j}{\tau_i} - 1)} e^{\frac{\sigma^2}{2\tau_j^2} - \frac{t}{\tau_j}} \left[ 1 + \operatorname{erf}\left(\frac{t - \frac{\sigma^2}{\tau_j}}{\sigma\sqrt{2}}\right) \right] \frac{C_i (1 - A) B_j}{2(\frac{\tau_j}{\tau_i} - 1)} e^{\frac{\sigma^2}{2\tau_i^2} - \frac{t}{\tau_i}} \left[ 1 + \operatorname{erf}\left(\frac{t - \frac{\sigma^2}{\tau_j}}{\sigma\sqrt{2}}\right) \right]$ 

Free Parameter	Expected
$(C_1 \tau_1) / (C_3 \tau_3)$ for ER	~0.1
$(C_1 \tau_1) / (C_3 \tau_3)$ for NR	~1.52 (at ~100 MeV)
τ <sub>1</sub>	2.2 - 4.3 ns *
$\mathbf{\tau}_3^{}$ for ER	21 - 27 ns *
$\mathbf{\tau}_{3}^{}$ for NR	21 - 27 ns *
σ	≥ 3.1 ns

\* Range of measured values from Kubota (1978 & 1979), and Hitachi (1983)

## Fit results

Model is fitted to all histograms for ER and NR at all energies simultaneously

- Allows us to vary parameters common among different energy / particle type bins (i.e.  $\sigma$ ,  $\tau_1$ , etc. )

### LUX PRELIMINARY

Parameter	Expected	Best fit ± stat.	Fit sys. err.	Optical sys. err.
$(C_1 \tau_1) / (C_3 \tau_3)$ for ER	~0.1	$0.042 \pm 0.006$	± 3.1%	+75% / -66%
$(C_1 \mathbf{r}_1) / (C_3 \mathbf{r}_3)$ for NR	1.52 (at ~100 MeV)	0.269 ± 0.034	± 3.1%	+20% / -10%
τ <sub>1</sub>	2.2 - 4.3 ns *	3.27 ± 0.66 ns	± 1%	+11% / -70%
$\mathbf{\tau}_{_{3}}$ for ER	21 - 27 ns *	25.89 ± 0.06 ns	± 1.9%	+0.5% / -0.6%
$\mathbf{\tau}_{_3}$ for NR	21 - 27 ns *	23.97 ± 0.17 ns	± 1.9%	+0.1% / - 1.1%
σ	≥ 3.1 ns	3.59 ± 0.09 ns	± 1.1%	± 1.2%

\* Range of measured values from Kubota (1978 & 1979), and Hitachi (1983)

## Singlet/triplet ratio energy dependence



Three important results:

- NR ratio is much smaller at low energies than published measurements made at high energies
- ER ratio has no significant energy dependence under applied field
- NR may show energy dependence? Not prominent in our data

## **Prompt fraction discriminator**

PSD accomplished through optimized prompt fraction

Also, **built a toy MC** that draws photon times from best-fit timing distribution and

- Adds fluctuations in PMT signal size
- Adds fluctuations in 5% area time from digitization
- Computes prompt fraction discrimination



## **Prompt fraction discrimination**



## ER leakage into NR 50% acceptance region



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September 22, 2017

## Summary

- We are studying pulse shape discrimination for background rejection in dark matter analyses with LUX
- Built a framework for timing photon arrivals
- Used an analytic model to reconstruct singlet/triplet ratios at low energies (NR for the first time)
- Demonstrated prompt fraction discrimination with LUX calibration data
- Constructed a Monte Carlo model that reproduces ER/NR distributions, can be used in LUX simulations and analysis

## Acknowledgements

## The LUX collaboration Sanford Underground Research Facility (SURF)

### LUX PSD subgroup

- Dev Ashish Khaitan (U of Rochester)
- Mongkol Moongweluwan (U of Rochester)
- Daniel Hogan (UC Berkeley)
- Prof. Matthew Szydagis (U Albany, SUNY)
- Dr. Kareem Kazkaz (LLNL)

This work performed under the auspices of the U.S. Department of Energy by Lawrence Livermore National Laboratory under Contract DE-AC52-07NA27344.



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## Back up

## WIMP dark matter

## Weakly Interacting Massive Particles

- New neutral particle, beyond the standard model
- Weak-scale interaction cross-section gives us the right amount of dark matter
- Predicted to produce NUCLEAR RECOILS (no EM interactions)
- Most backgrounds (γ's and β's from radioactive decay) produce ELECTRON RECOILS



γ, β

γ, β

### Assumptions

- Weak scale scattering cross section with nuclei
- Mass density ~ 0.3 GeV/c<sup>2</sup>/cm<sup>3</sup>
- Maxwellian velocity distribution with  $v_0 = 220$  km/s
- Velocity distribution truncated at galactic escape velocity

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# Possible future searches with higher energy nuclear recoils

### Dark matter inelastic scattering:

- $\chi + N \rightarrow \chi^* + N$
- Kinematically suppresses low-energy recoils

These searches would require us to extend our acceptance at higher energies

Larger window could introduce new/more backgrounds



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## Liquid xenon TPC advantages

### Low threshold

 Can detect events that produce 10's of scintillation photons and 1's of ionization electrons

### Low background

- No long-lived radioactive xenon isotopes
- High-Z and high density provides self-shielding in large detectors

### Scalable technology

- ~20 years development experience
- Ton-scale detectors are in operation

### Particle ID (ER/NR) capabilities

- Charge/light ratio (~99.9% rejection)
- PSD???



## LXe scintillation physics

Emission can be modeled as:

$$P(t) = C_1 e^{-t/\tau_1} + C_3 e^{-t/\tau_3}$$

with three free parameters:

- Singlet time  $\tau_1$
- Triplet time  $\tau_3$
- Singlet/triplet ratio  $(C_1 \tau_1)/(C_3 \tau_3)$

**Recombination** of electrons and ions can contribute to timing

- Only observed in electron recoils (ER)
- Suppressed by applied electric field and high LET at low energies
  - NEST model (Mock et al.) predicts
    ~1ns effect in LUX data
- We treat it as a different  $\tau_3$  for ER and NR



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## **Optical transport**

Photons in LUX typically scatter before arriving in PMTs.

Studied using ray-tracing simulations in LUXSim (LUX Geant4 simulation package)

- Direct-path transit time subtracted
- Short-time behavior driven by geometric efficiency of bottom PMTs

Sims are fit to an analytic model for easy convolution and simulation.

A = direct-hit fraction $B_a = \text{weights exponential terms}$  $B_b = (1 - B_a)$  $\tau_a = \text{long-time constant (11.2 ns)}$  $\tau_b = \text{short-time constant (varies to fit short-time behavior)}$ 



$$P_o(t) = A \,\delta(t) + (1 - A) \left[ \frac{B_a}{\tau_a} \, e^{-t/\tau_a} + \frac{B_b}{\tau_b} \, e^{-t/\tau_b} \right]$$

International Conference on Applications of Nuclear Techniques (Crete17)

June 15, 2017