Particle Accelerators

Physics 290E, Spring 2025

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Why Accelerators?

- "A better magnifying glass"
	- \Box Need to look deeper into the structure of matter
	- \Box Spatial resolution is limited by wavelength of the probe
		- Microscope: visible light λ ∼ 1 µm → cell structure
		- \mathcal{F} Higher energy particles: $\lambda = 2\pi/p$
		- X-rays: λ ∼ 0.01−10 nm → atomic, crystal structure
	- □ Charged particles:
		- Rutherford experiment: p ~ 10 keV, λ ∼ 10−¹⁰ m
		- Discovery of quarks: p ~ 10 GeV, λ ∼ 0.1 fm
		- LHC: p ~ 1−10 TeV, λ ∼ 10−16−10−¹⁷ m → quark structure, compact extra dimensions,

New Particles

- Most elementary particles are unstable (WIMPs not seen yet \odot) \rightarrow create them in high-energy collisions
	- Energy conservation needs no explanation \odot
- Most recent:
	- c quark (J/Ψ meson): SLAC/SPEAR (3-7 GeV c.m. energy) and BNL (28 GeV protons on Be target)
	- b quark (Υ(1S) meson): FNAL (400 GeV pN fixed target)
	- \Box t quark: FNAL (1.8 GeV c.m. $p\bar{p}$ collisions)
	- τ lepton: SLAC/SPEAR
	- W/Z bosons: CERN/SppS: 540 GeV c.m. pp collisions
	- Higgs boson: CERN/LHC: 12-14 TeV c.m. pp collisions

Fixed Target vs Collider

• Energy reach

Fixed (stationary) target:

$$
E_{cm} \approx \sqrt{2 m_t E_b}
$$

- Other advantages of collider beams
	- \Box Re-circulating beams possible
	- Collisions between fundamental particles
- Advantages of fixed target facilities
	- \Box High event rates (see next)
	- Variety of targets and beams (including beams of unstable particles)
	- Boost for secondaries (this may be a disadvantage too !)

- What is relevant for high-energy collisions is CM energy of the fundamental process
- Colliding beams (assuming $E \gg m$):

$$
E_{cm} \approx \sqrt{4 E_1 E_2}
$$

- R = σ*L* : *L* is *instantaneous* luminosity [cm−2s−¹ or fb-1year−1]
- □ For colliders

$$
\mathcal{L} = f n \frac{N_1 N_2}{4 \pi \sigma_x \sigma_y}
$$

Event Rates

• Colliders:

- **Here f is revolution frequency,** *n* is number of bunches in a ring, N_i is # of particles in the bunch, and σ is the beam size
- Typical luminosities:
	- \Box FNAL: 10³¹–10³² cm⁻²s⁻¹ (10-100 μb⁻¹s⁻¹, 0.1-1 fb⁻¹year⁻¹)
	- \Box LEP,SLC: ~10³¹ cm⁻²s⁻¹ (10 μb⁻¹s⁻¹, 0.1 fb⁻¹year⁻¹)
	- PEP-II, KEKB: 1034 cm−2s−1 (10 nb-1s−1, 100 fb-1year−1)
	- HL-LHC: 5×1034 cm−2s−1 (50 nb-1s−1, 500 fb-1year−1)

Event Rates L V CHIL NAICS

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theoretical uncertainties are indicated as bands.

- Much higher rates achievable with modest beam currents and sizes \Box For example, a 1m liquid hydrogen target and a beam of 10¹³ particles per second are equivalent to \sim 10³⁸ cm⁻²s⁻¹ luminosity
	- Example from E158 at SLAC

Components of an Accelerator

- Particle source
- Damping systems (reduce beam size)
- Particle acceleration
- Beam delivery systems (to fixed-target or collider regions)
- Beam dump

- Electrons
	- Thermal
	- Photo-electric
	- Polarized
- (Heavy) ions including protons **ECR** sources
- Indirect (secondary beams)
	- \Box High energy photons
	- **Positrons**
	- □ Neutrons
	- □ Muons, pions, kaons

Particle sources

Thermionic Gun

' Stanford Linear Accelerator Center, Stanford University, Stanford, Stanford, CA 94309, USA099, USA099, USA09

Fig. 1: Cross-sectional view of the 1-1/2 on-axis cells.

and during several microwave gun cavity fill times.

the barrier is made via a tungsten-wire spring in the form of a

ECR Source

• Electron Cyclotron Resonance

Polarization

- Spin is the basic quantum number
- Electroweak couplings are helicity-specific
- Access to polarized beams and/or targets provides additional degree of freedom
	- \Box Precision electroweak studies of couplings
	- Precision studies of nucleon structure
- Polarized sources:
	- □ Semiconductor photocathodes (electrons)
	- Sokolov-Ternov effect (electrons/positrons)
	- □ Secondary production (positrons, protons)
	- \Box Electron or laser cooling (protons)
- Polarized targets
	- \Box Optical pumping
	- Dynamic nuclear polarization

NOTES:

- 1. High polarization ($>80\%$) and available charge (\sim 10¹¹ electrons/pulse)
- 2. Non-strained cathodes: limit 50% polarization
- 3. Small anisotropy in strain: ~3% analyzing power for residual linear polarization

Reference: T. Maruyama et al., SLAC-PUB-9133, March 2002; (submitted to Nucl. Inst. Meth. A).

Fig. 2.2. Energy levels in strained GaAs.

Polarized Source Photocathode

Polarized Laser System

Particle Acceleration

- Physics depends on CM Energy
- Use Electric (EM) fields to accelerate charged particles
	- $\Delta E = V e$
- Non-relativistic
	- \Box Static (or slowly-varying) E fields
	- \Box <100 kV/m, limited by corona
- Relativistic
	- Dynamic fields
	- \Box >10 MV/m, limited by breakdown of conductor

• Static E-field Vandergraaf □ Crockroft-Walton

Low-energy Accelerators

• Cyclotron Static B field, variable E-field

Cyclotrons

Lawrence's 4.5" cyclotron

184" cyclotron

RF Acceleration

- Electrons are accelerated in alternating electric field form in a series of cavities
	- \Box SLAC cells are copper with length of ~1/3 λ at frequency 2856 MHz
	- \Box NLC/CLIC design used X band cavities at f=11.4 GHz to get accelerating gradients up to 75 MeV/m
	- Accelerating cavities at JLab (and ILC) are superconducting $f \sim 1$ GHz
- "Auto-phasing" effect produces bunch structure
- RF power is generated by series of high-power klystrons with output of 60-70 MW

Laser-driven

Future Accelerator Technology

• Beam-driven RF CLIC

• Plasma Wakefield

Synchrotron Radiation \sim 1 motion; motion; motion; motion; more easily is much mo erated with a magnetic field compared to an electric field compared to an electric field of 1 Testima produces the 1 Testima produ same force as an electric field of 300 MV/m. α small - only 7 μ at 100 GeV. A more directly meaningful a parameter is the total energy loss per turns to the theory

- All accelerated charges radiate. For relativistic particles in a magnetic field, radiation is called synchrotron *diate* For relativistic particles in a mag $\frac{1}{2}$
- Radiated power *P* = 2 3 e^2c $4\pi\varepsilon_0$ $\beta^4 \gamma^4$ $\frac{\rho I}{\rho^2}$
- Energy loss per particle per turn - Elicigy loss per particle per turn In particular parti

*U*⁰ =

(8)

$$
U_0
$$
 [eV] = 8.85 10⁴ $\frac{E^4$ [GeV]}{\rho [m]} = 2.65 10⁴ E³ [GeV] B [T]

- Grows as 4th power of energy highest energies prohibitively expensive U U \mathbf{r} ^β*^c* ∫ *ds* ⁼ *v* — *c* **ghest** er $f^{\mu\nu}$ power of energy – ingrest energies promotively ex-
- Need to increase bend radius for higher energies (but reduction is only linear) $\frac{1}{2}$ For an isomagnetic lattice (uniform bending radius in the bending radius in the bending magnetic magnetic system of $r_{\rm max}$) simple to: o increase bend radius for higher energies (but redu
	- Examples: LEP: 100 GeV beams, 27 km circumference *e* <u>2022</u>
1922 - Στρατικό τρόπος
1932 - Στρατικό τρόπος

Accelerator Types

- Linear accelerators (electrons, protons)
	- \Box Avoids synchrotron radiation losses for electrons, compact for protons
	- High gradients now available (up to 50 MeV/m)
	- But beams used only once \rightarrow luminosity at a premium
	- Example: SLAC, future Linear Colliders (ILC)
- Circular accelerators: synchrotrons and storage rings
	- Highest energies possible with protons, muons
	- Synch. radiation losses for electrons
	- Multiple passes more tolerant to errors in steering

Circular Accelerators

 $|p=0.3BR|$ (*B* in T, *R* in m, *p* in GeV)

- Cyclotrons: constant *B*, variable *R*; accelerating RF electric field Lawrence, 1932
- Betatrons: a little variation
	- *E* field generated by a variable *B* field (a la transformer)
- Synchrotrons: constant *R*, variable *B*, accelerating *EM* field \Box Most high-energy machines are synchrotrons

Beam Optics

- Beamlines made of a series of discrete magnets (lattice)
- Two main types of magnets
	- Dipoles: bending.
	- Quadrupoles: focusing
- Strong focusing
	- □ Net focusing using 1 "cell"
	- Many benefits:
		- ^s stability
		- \mathcal{F} small beams (= compact beam pipe)
		- \mathcal{F} small beamspots (= more luminosity)

Accelerators as Secondary Particle Sources

- Fixed target facilities can produce secondary particles
	- **Most common: positrons, photons, pions, kaons (including neutral** K_L **), muons,** neutrinos
- Secondaries can be captured and accelerated
	- Anti-protons and positrons
	- Muons
	- Other species, including radioactive isotopes
- Particle decays in accelerators
	- High-brightness photon sources
	- Neutrino factories

- Solenoid
- Liquid-metal lens
- Particle cooling
	- □ 'Phase-space' cooling
	- \Box Synchrotron cooling (damping rings)
	- Matter-based cooling
	- □ Electron cooling
	- Dynamic cooling

Particle capture and cooling

- Vacuum
	- □ colliders are also 'fixed target'!
- Beam stability (higher-order correctors)
- Final Focus
	- nano-beams for SuperB, XLC
- Spin rotation (for polarized beams)
- Asymmetric (or PP) colliders
	- Finite crossing angle, bunch rotation, …
- Beam energy measurement

Other Issues

Accelerator Parameters (PDG)

https://pdg.lbl.gov/2024/reviews/rpp2024-rev-hep-collider-params.pdf

01/29/2024 YGK, Phys290E

Table 32.3: Updated in March 2022 with numbers received from representatives of the colliders (contact E. Pianori, LBNL). The table shows the parameter values achieved. Design parameters for SuperKEKEB may be found in our 2018 edition (Phys. Rev. **D98**, 030001 (2018)) Quantities are, where appropriate, r.m.s.; unless noted otherwise, energies refer to beam energy; *H* and *V* indicate horizontal and vertical directions; s.c. stands for superconducting.

> *^e*⁺ : +0*.*78*/* [≠] ⁰*.*⁷³ ${\rm ad\,\, cone}$

interaction point (m)

 $\mathcal{B}(\mathcal{B})$ beam tune shifted by the shifted shifted by

p: 2.45 (H), 0.18 (V)

e+: 1.2 (0.23*†*) (*H*), 0.0059 (*V*)

e+: 0.50 (*H*), 0.012 (*V*)

^e+: 0.080 (*H*), ¹ ◊ ¹⁰≠³ (*^V*)

e:190 (H), 450 (V)

 $\sqrt{22}$ (*V*) *e*+: 17.9 (*H*), 0.22 (*V*) $\frac{1}{28}$

e≠: 1020 (*H*), 900 (*V*)

e≠: 703 (*H*), 498 (*V*)

e≠: 15 (*H*), 278 (*V*)

SLAC Accelerator Complex

Past glory…

- 50 GeV electron and positron beams
- Collider (SLD) and fixed target facilities
- Injector for PEP-II storage ring
- SPEAR: separate ring

Fermilab Complex

- Main Tevatron ring: 1 TeV protons and antiprotons (RIP)
- Main injector: 120 GeV booster
	- Now used for neutrino beams
- 8 GeV complex
- Injector for the MI/Tevatron: now used for muon/neutrino beams
- Superconducting magnet technology

- LEP: 100 GeV
	- electrons and positrons
- LHC: 7 TeV protons
- Same tunnel, new magnets, RF

CERN

Electron Ion Collider Eicher Overeign Over

Future Accelerators

- Will certainly want to extend the energy frontier beyond the current generation of facilities
	- \Box Probe deeper into the mysteries of the microworld
- Most straightforward: proton machines \Box FCC etc: E_{CM} > 50 TeV (+possible intermediate stages in e+e- mode)
- Complementarity: lepton machines
	- Electron linear colliders: ILC (R&D), E_{CM} > 500 GeV, CLIC: E_{CM} = 3 TeV, FCC-ee
	- Muon colliders: $E_{CM} > 3$ TeV (R&D, possible future)

Supplemental material

Polarization

- Spin is the basic quantum number
- Electroweak couplings are helicity-specific
- Access to polarized beams and/or targets provides additional degree of freedom
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Example: Polarized Source Photocathode

Eq-1.43 eV

Polarized Laser System

Target Technology

E154 3He Target and **Spectrometers**

E155: NH_3 , ND_3 , ⁶LiD

5 Tesla Superconducting Microwave Target

CERN: Cryogenic Targets

Dedicated, Low **luminosity**

CERN: 5*107 of polarized u per SPS cycle (10+ sec) (80% polarization from π decay)

Polarized Targets

 $NH₃(80\%) ND₃(50\%)$

7 June 1994
Peter Berglund
Jukka Kyynäräine