

---

# Deep Inelastic Scattering at HERA

---

Sam Kohn  
Physics 290E Seminar  
20 April 2016

---

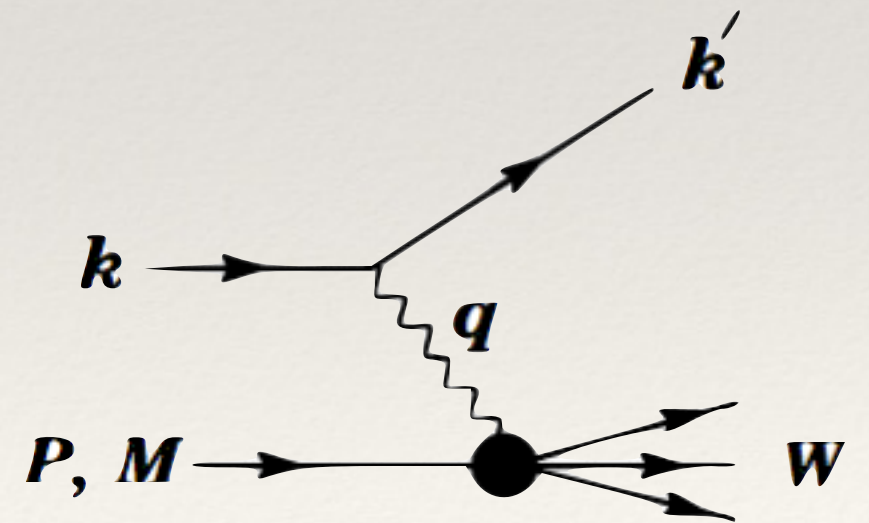
# Outline

---

- ❖ DIS theory review
- ❖ HERA = ZEUS + H1 detectors
- ❖ Results

# Deep Inelastic Scattering ⚡ Review

- ❖ High-energy lepton-nucleon collision
  - ❖ High = relative to QCD scale  $\sim 200$  MeV
- ❖ Interpret as lepton interacting directly with quark or gluon
- ❖ Long list of “invariants” formed from two basic quantities: incoming and outgoing lepton momentum  $k$  and  $k'$



Classic DIS “Feynman” diagram [1]



# Invariants...here they are

$\nu = \frac{q \cdot P}{M} = E - E'$  is the lepton's energy loss in the nucleon rest frame (in earlier literature sometimes  $\nu = q \cdot P$ ). Here,  $E$  and  $E'$  are the initial and final lepton energies in the nucleon rest frame.

$Q^2 = -q^2 = 2(EE' - \vec{k} \cdot \vec{k}') - m_\ell^2 - m_{\ell'}^2$ , where  $m_\ell(m_{\ell'})$  is the initial (final) lepton mass.  
If  $EE' \sin^2(\theta/2) \gg m_\ell^2, m_{\ell'}^2$ , then

$\approx 4EE' \sin^2(\theta/2)$ , where  $\theta$  is the lepton's scattering angle with respect to the lepton beam direction.

$x = \frac{Q^2}{2M\nu}$  where, in the parton model,  $x$  is the fraction of the nucleon's momentum carried by the struck quark.

$y = \frac{q \cdot P}{k \cdot P} = \frac{\nu}{E}$  is the fraction of the lepton's energy lost in the nucleon rest frame.

$W^2 = (P + q)^2 = M^2 + 2M\nu - Q^2$  is the mass squared of the system  $X$  recoiling against the scattered lepton.

$s = (k + P)^2 = \frac{Q^2}{xy} + M^2 + m_\ell^2$  is the center-of-mass energy squared of the lepton-nucleon system.

# Invariants...here they are

$$\nu = \frac{q \cdot P}{M} = E - E'$$

$$\diamond x = -q^2 / (2q \cdot P)$$

$$Q^2 = -q^2 = 2M\nu$$

$\diamond$  In what universe is that the "fraction of nucleon momentum carried by the quark?"

$$\approx 4EE' \sin^2(\theta/2)$$

beam direction.

$$x = \frac{Q^2}{2M\nu}$$

where, in the parton model,  $x$  is the fraction of the nucleon's momentum carried by the struck quark.

$$y = \frac{q \cdot P}{k \cdot P} = \frac{\nu}{E}$$

is the fraction of the lepton's energy lost in the nucleon rest frame.

$$W^2 = (P + q)^2 = M^2 + 2M\nu - Q^2$$

is the mass squared of the system  $X$  recoiling against the scattered lepton.

$$s = (k + P)^2 = \frac{Q^2}{xy} + M^2 + m_\ell^2$$

is the center-of-mass energy squared of the lepton-nucleon system.



# Invariants...here they are

$$\nu = \frac{q \cdot P}{M} = E - E'$$

$$\diamond x = -q^2 / (2q \cdot P)$$

$$Q^2 = -q^2 = 2$$

$\diamond$  In what universe is that the “fraction of nucleon momentum carried by the quark?”

$$\approx 4EE' \sin^2 \theta$$

beam direction.

$$x = \frac{Q^2}{2M\nu}$$

where, in the parton model,  $x$  is the fraction of the nucleon's momentum carried by the struck quark.

- $\diamond$  A: The infinite-momentum frame, where the struck quark's final momentum is  $0 = (xP + q)^2 = 0 + 2xq \cdot P + q^2$
- $\diamond$  This “intuition” brought to you by Bjorken and Feynman

system.

---

# Why DIS?

---

- ❖ Allows us to probe fundamental QCD physics
- ❖ Measure Parton Distribution Functions (PDFs) which conveniently (and predictably) depend on  $x$  and  $Q^2$
- ❖ Measure strong coupling constant  $\alpha_s$
- ❖ Measure cross section ratios
  - ❖ predictable from QCD given PDF and  $\alpha_s$

---

# Quantity to measure

---

- ❖ No surprise here: cross sections at colliders
- ❖ Work backwards to get to structure functions, PDFs, and  $\alpha_s$

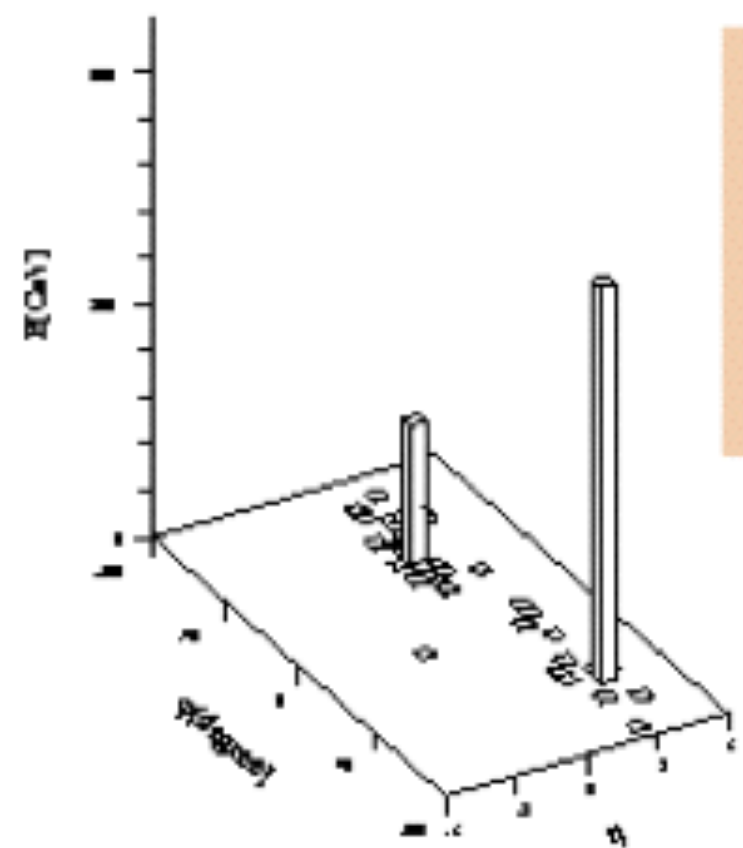
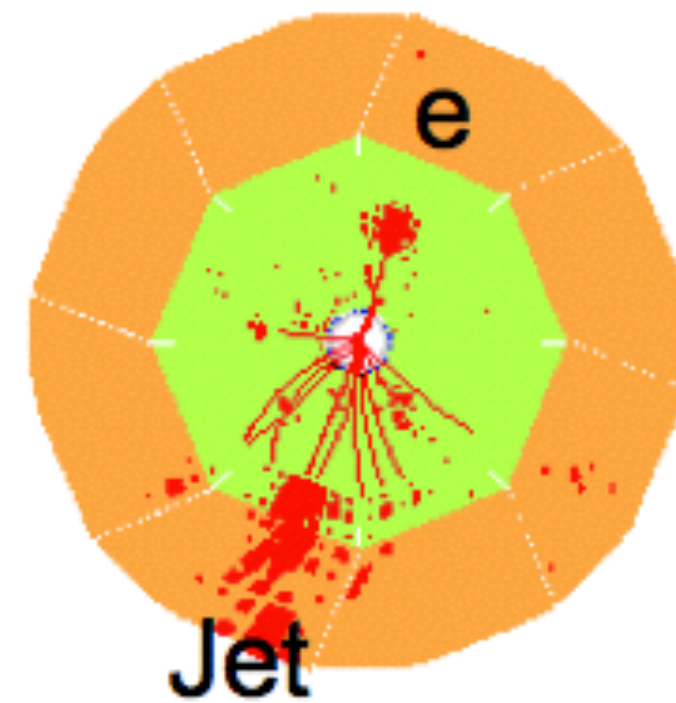
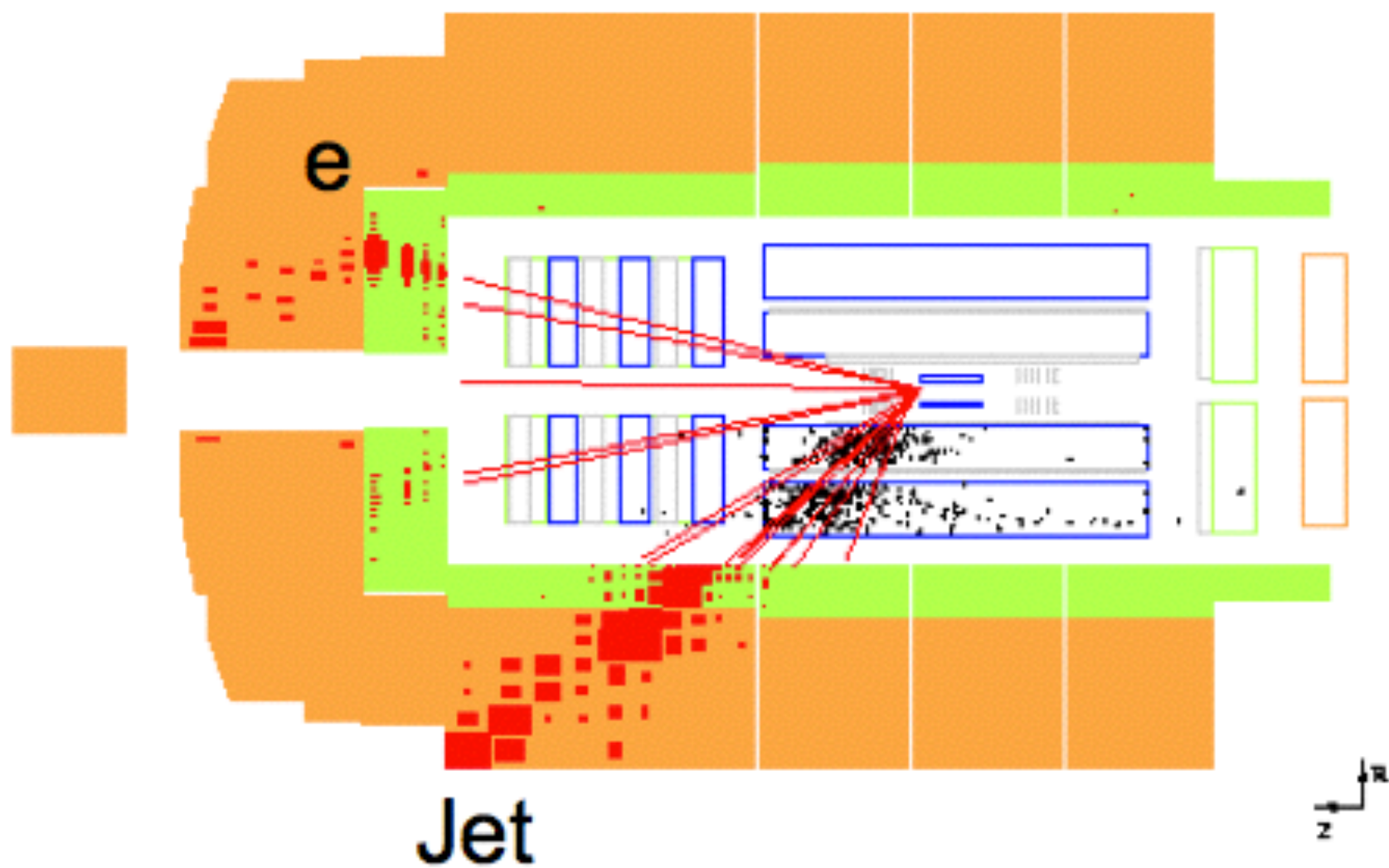




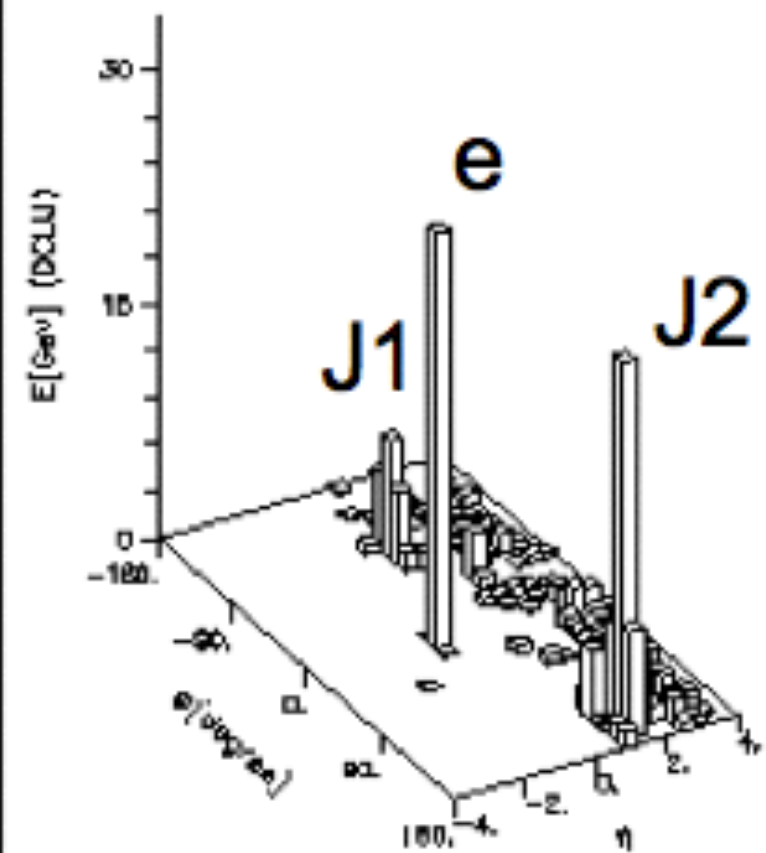
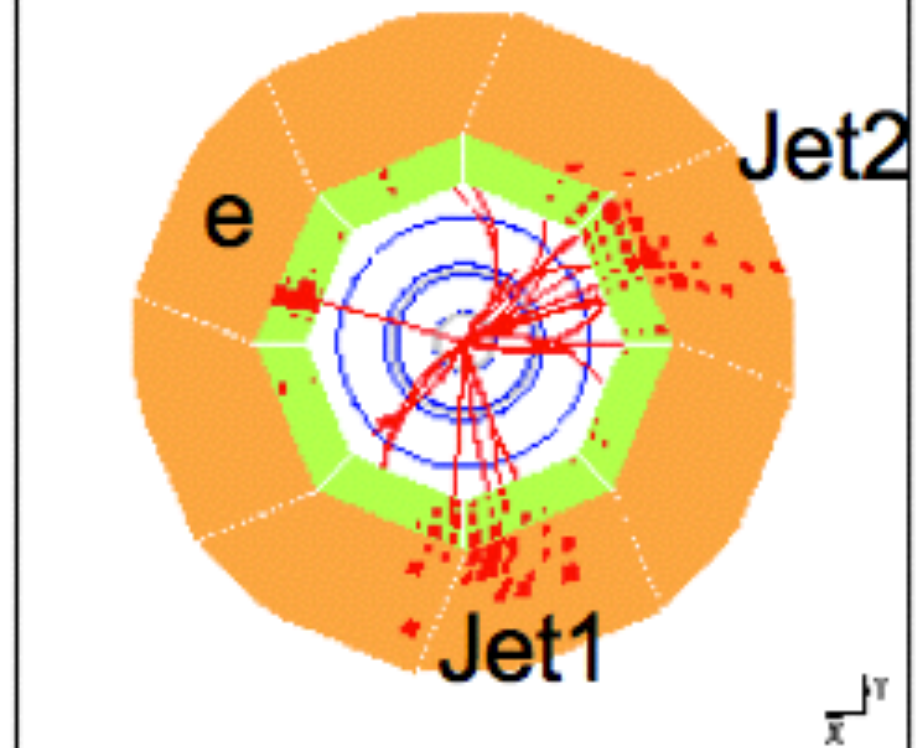
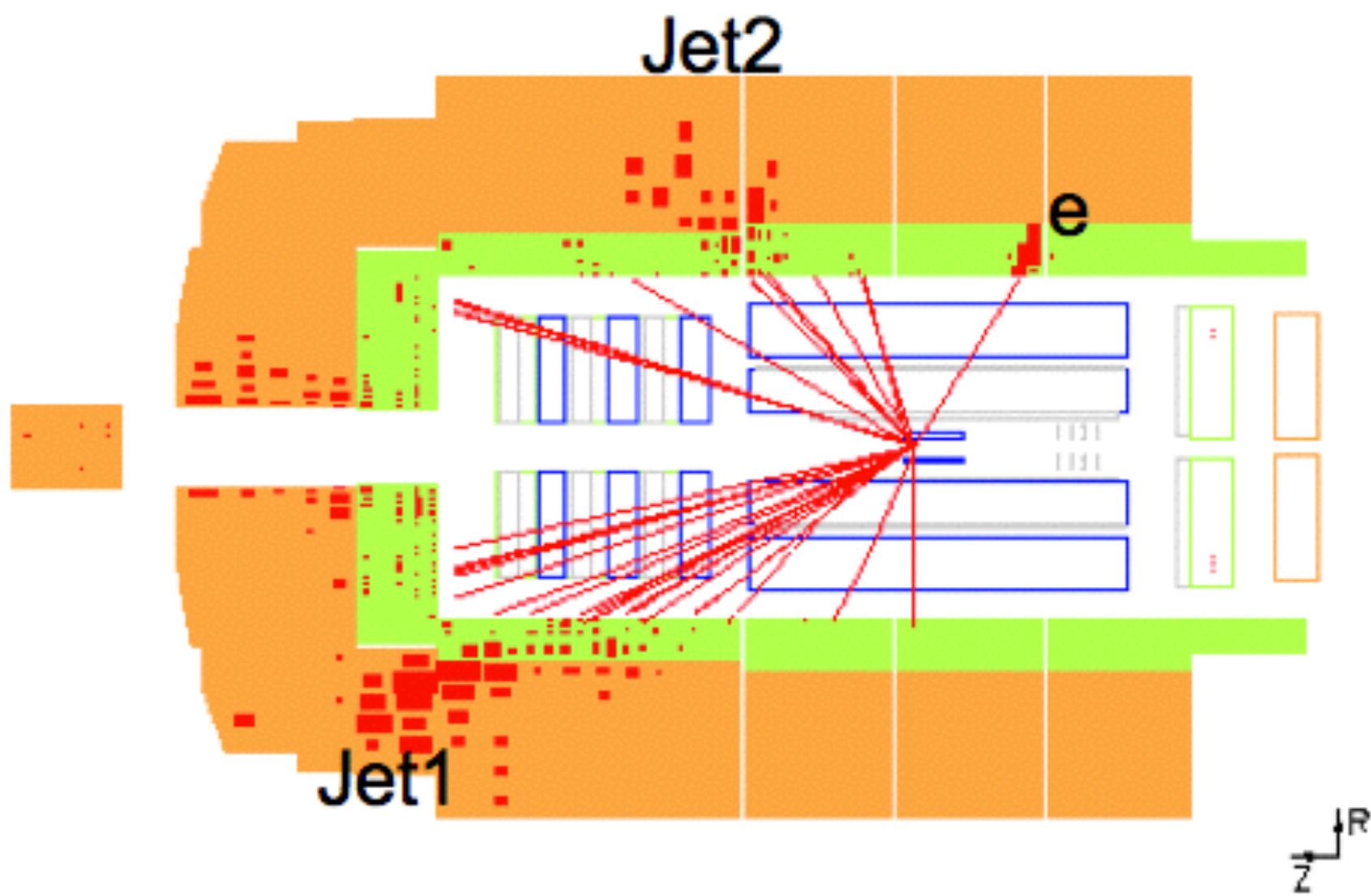




NC :  $Q^{*2} = 49308 \text{ GeV}^2$  ;  $P_{T_e} = 109 \text{ GeV}$



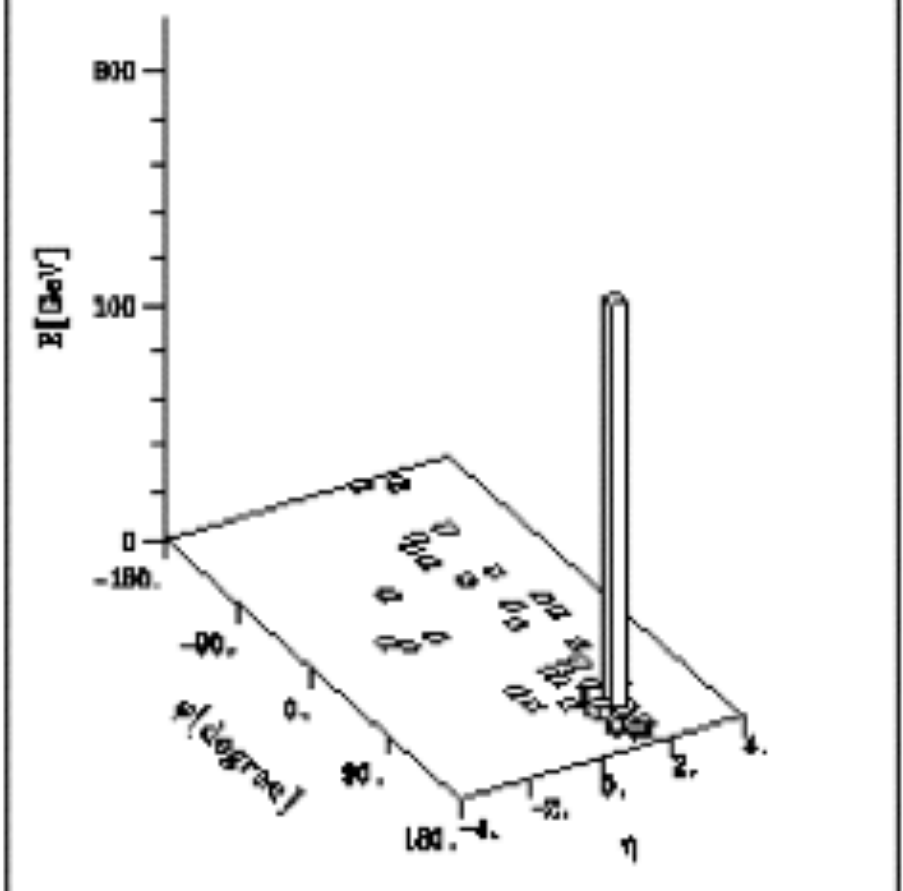
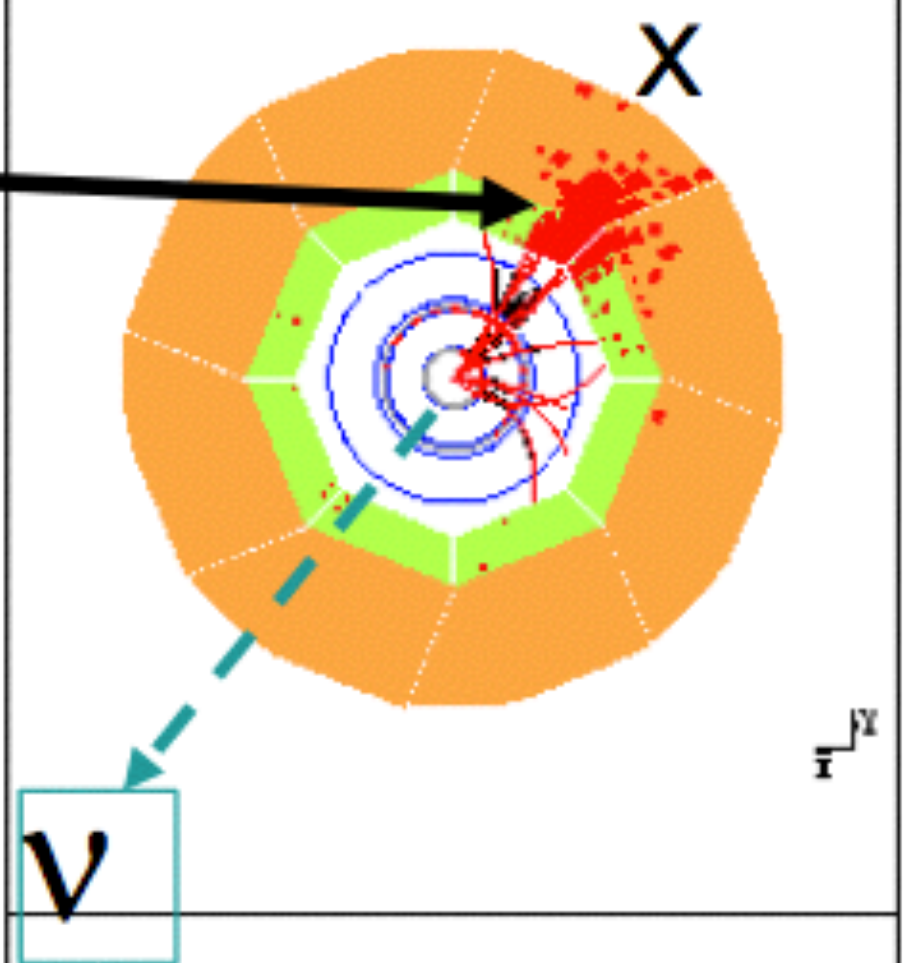
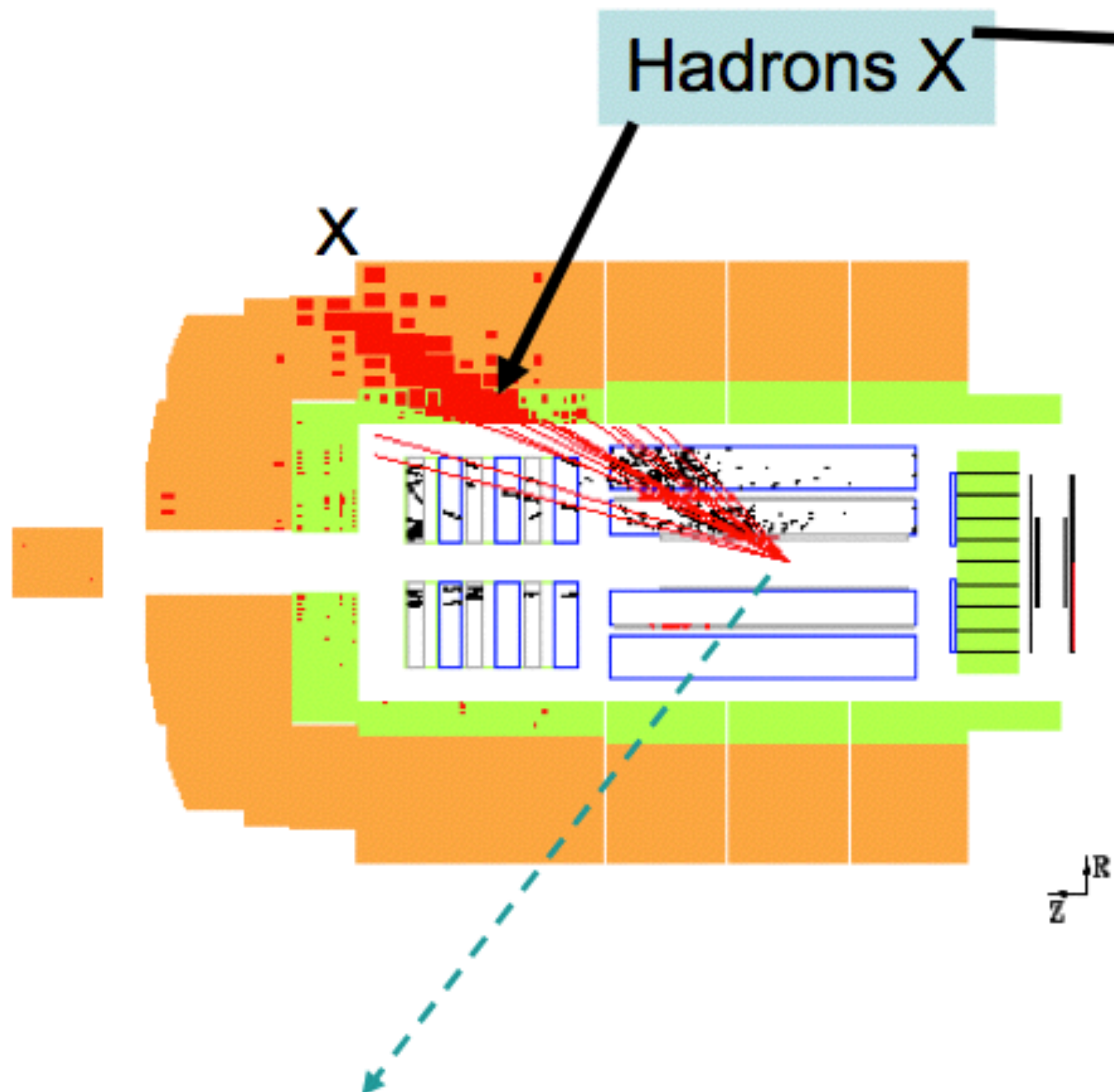




H1 event display [2]

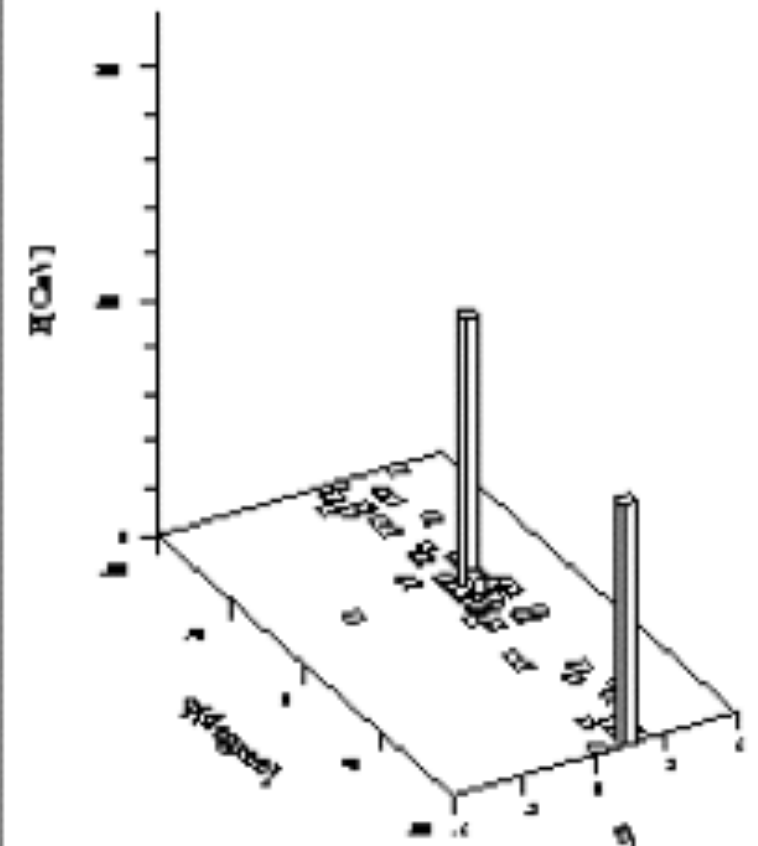
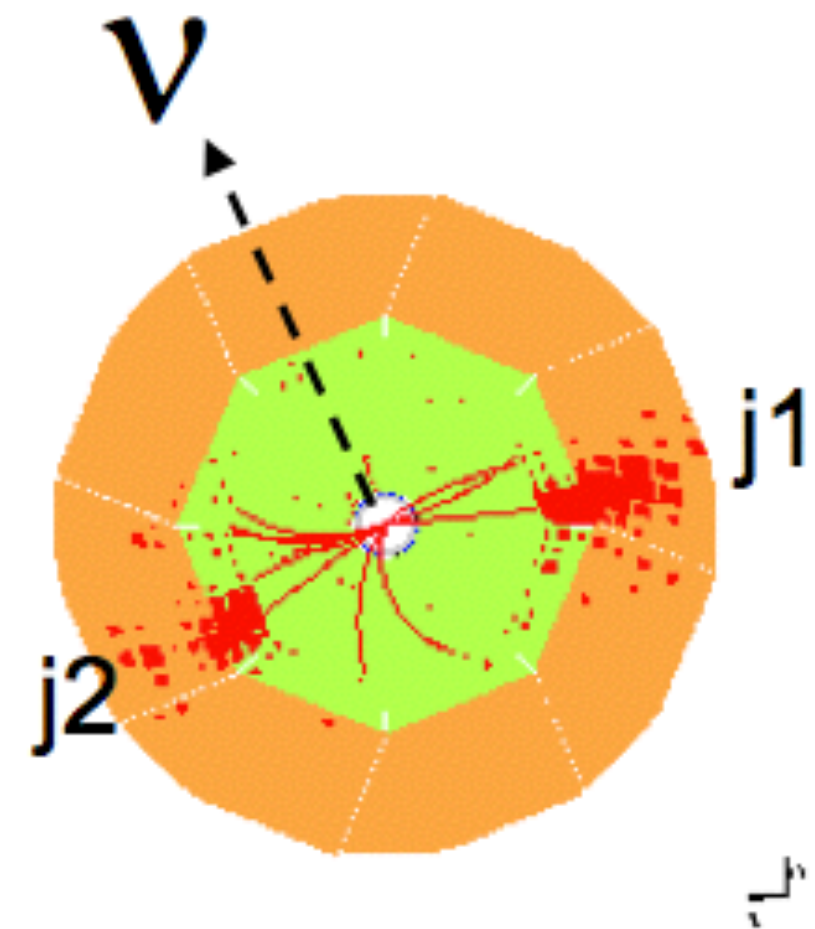
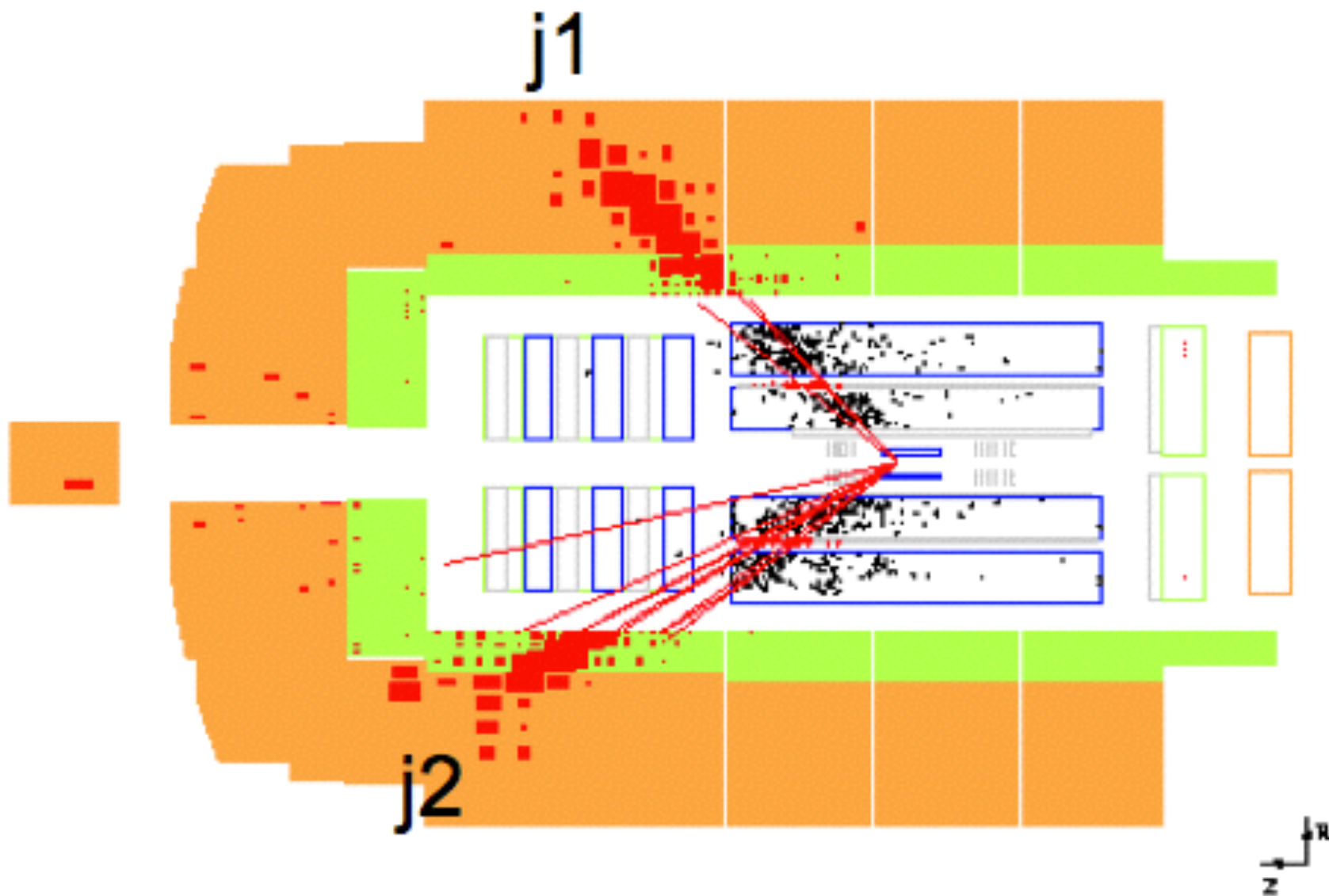
CC event  $Q^2 = 20000 \text{ GeV}^2$   $y=0.49$   $x=0.44$

Hadrons X



v

CC ? :  $M_{jj} = 148 \text{ GeV}$  ;  $P_{T\_miss} = 35 \text{ GeV}$



H1 event display [2]

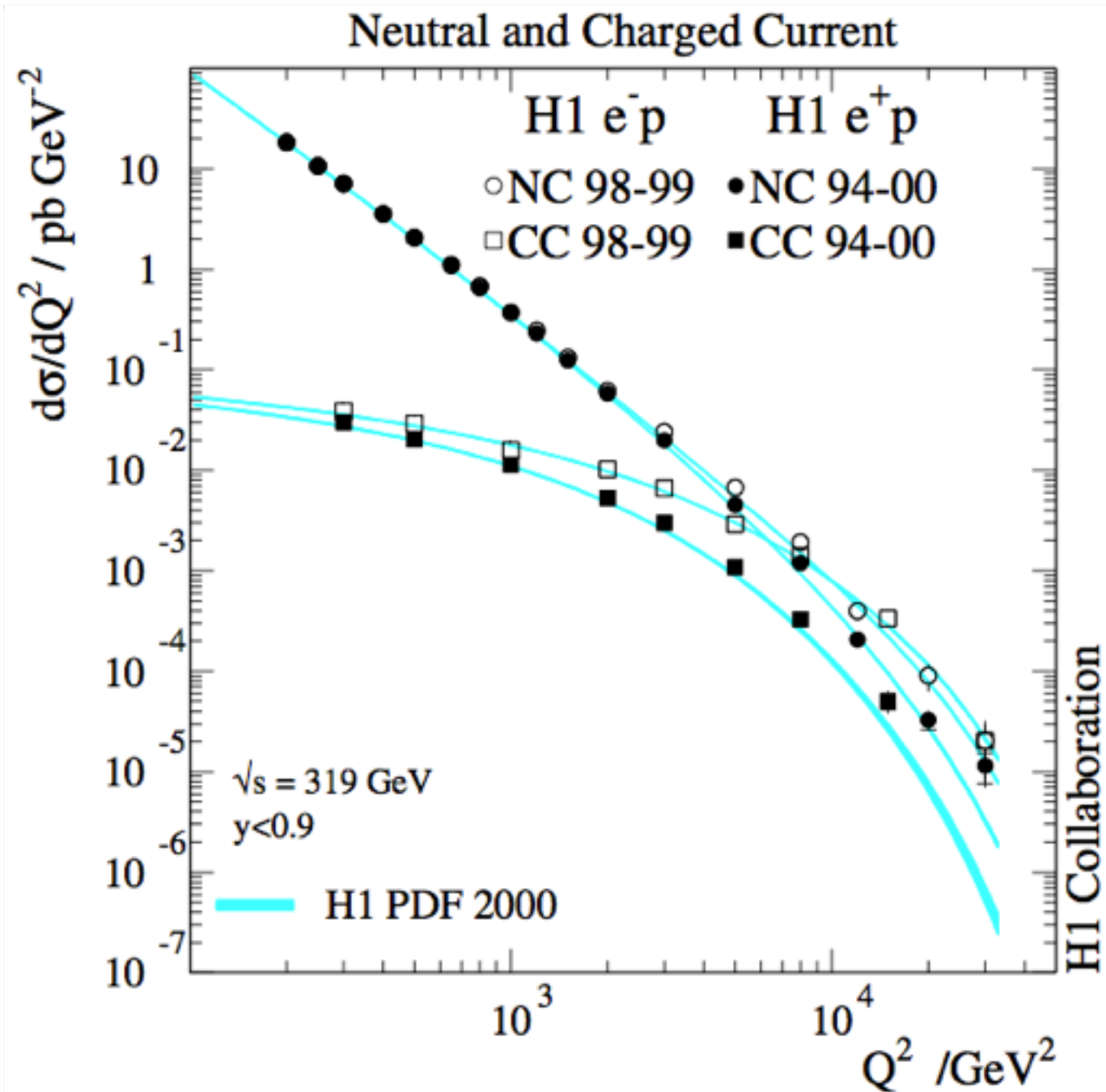


---

# Background and other Considerations

---

- ❖ Photoproduction background
  - ❖ (CC) electron escapes down beam pipe & fluctuation in calorimeter generates “missing”  $E_T$
  - ❖ (NC) electron escapes down beam pip & hadron mimics electron signal
  - ❖ Minimize by requiring energy conservation (NC) and analyzing the energy flow relative to  $P_{hadronic}$  (CC)
- ❖ Initial state bremsstrahlung (radiative corrections)
  - ❖ (NC) Photons escape down beam pipe. Minimize by enforcing energy conservation



*H1 data & fit [4]*

# Measured data

The measured cross sections for CC and NC. The bands show the results of the fit, which are accurate over 6 orders of magnitude in  $\sigma$ .



---

# Performing the $\chi^2$ fit

---

- ❖ Fit to gluons, up-type, anti-up-type, down-type, and anti-down-type. Try to separate out  $s, c, b$  later
- ❖ Start with trial PDFs then “predict” cross section
- ❖ Use smart parametrization that relies on physics
  - ❖  $U, D$  must account for valence quarks while  $\bar{U}, \bar{D}$  are sea quarks only  $\rightarrow$  more parameters for  $U, D$

- ❖  $B_q$ 's must all be equal

Fitter parametrization [4]

$$xg(x) = A_g x^{B_g} (1-x)^{C_g} \cdot [1 + D_g x]$$

$$xU(x) = A_U x^{B_U} (1-x)^{C_U} \cdot [1 + D_U x + F_U x^3]$$

$$xD(x) = A_D x^{B_D} (1-x)^{C_D} \cdot [1 + D_D x]$$

$$x\bar{U}(x) = A_{\bar{U}} x^{B_{\bar{U}}} (1-x)^{C_{\bar{U}}}$$

$$x\bar{D}(x) = A_{\bar{D}} x^{B_{\bar{D}}} (1-x)^{C_{\bar{D}}},$$



# Performing the $\chi^2$ fit (2)

- ❖ Get from PDF to cross section
- ❖ Tons of formulas. Here are parts of the CC & NC computations

$$[xF_3^{\gamma Z}, xF_3^Z] = 2x \sum_q [e_q a_q, v_q a_q] \{q - \bar{q}\}$$

FYI not a commutator

$$x\tilde{F}_3 \equiv -a_e \frac{\kappa Q^2}{(Q^2 + M_Z^2)} xF_3^{\gamma Z} + (2v_e a_e) \left( \frac{\kappa Q^2}{Q^2 + M_Z^2} \right)^2 xF_3^Z$$

$$\frac{d^2\sigma_{NC}^\pm}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} \phi_{NC}^\pm (1 + \Delta_{NC}^{\pm,weak}),$$

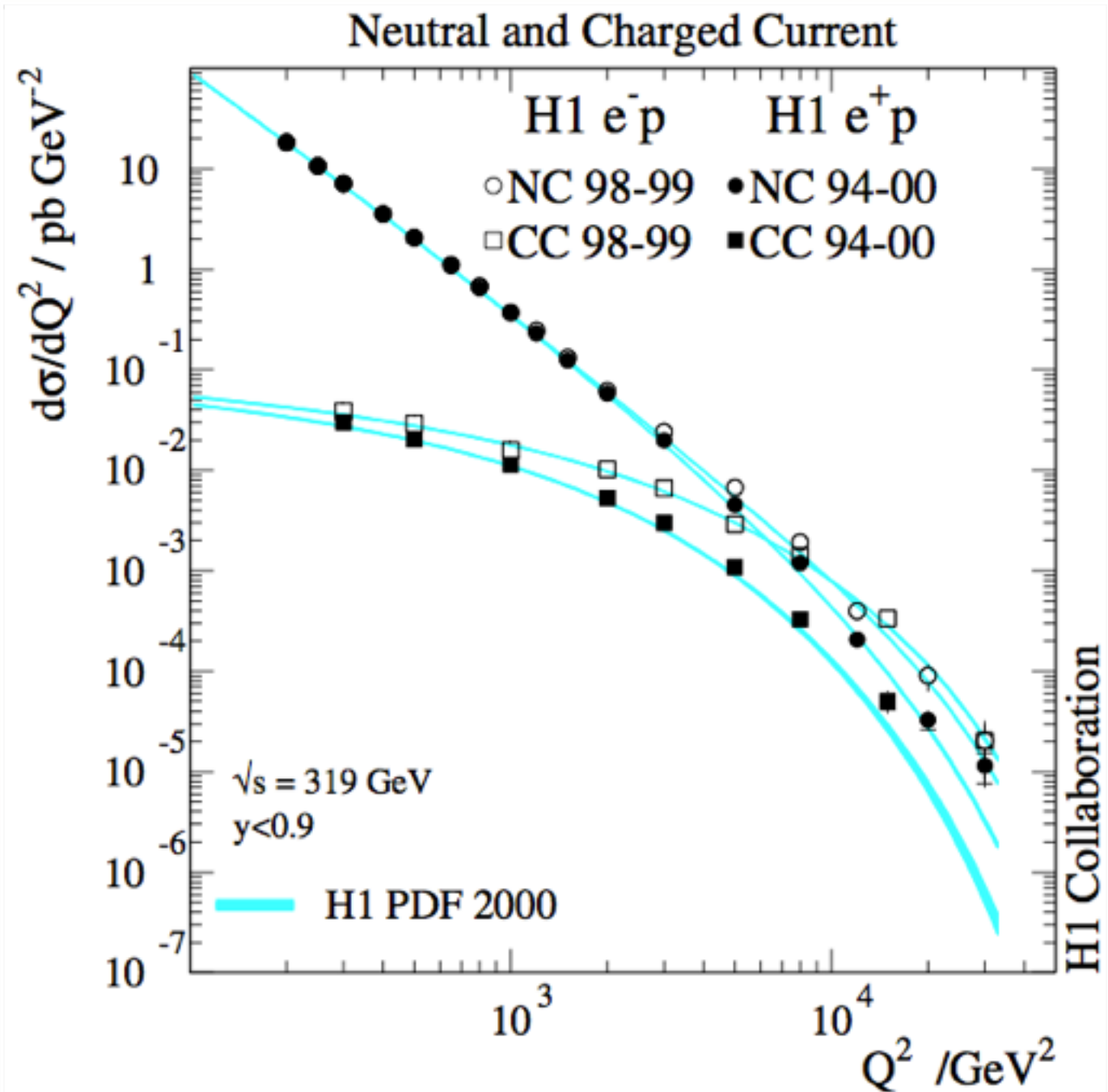
with  $\phi_{NC}^\pm = Y_+ \tilde{F}_2 \mp Y_- x\tilde{F}_3 - y^2 \tilde{F}_L,$

$$W_2^+ = x(\bar{U} + D), \quad xW_3^+ = x(D - \bar{U})$$

$$W_2^- = x(U + \bar{D}), \quad xW_3^- = x(U - \bar{D})$$

$$\phi_{CC}^\pm = \frac{1}{2} (Y_+ W_2^\pm \mp Y_- xW_3^\pm - y^2 W_L^\pm),$$

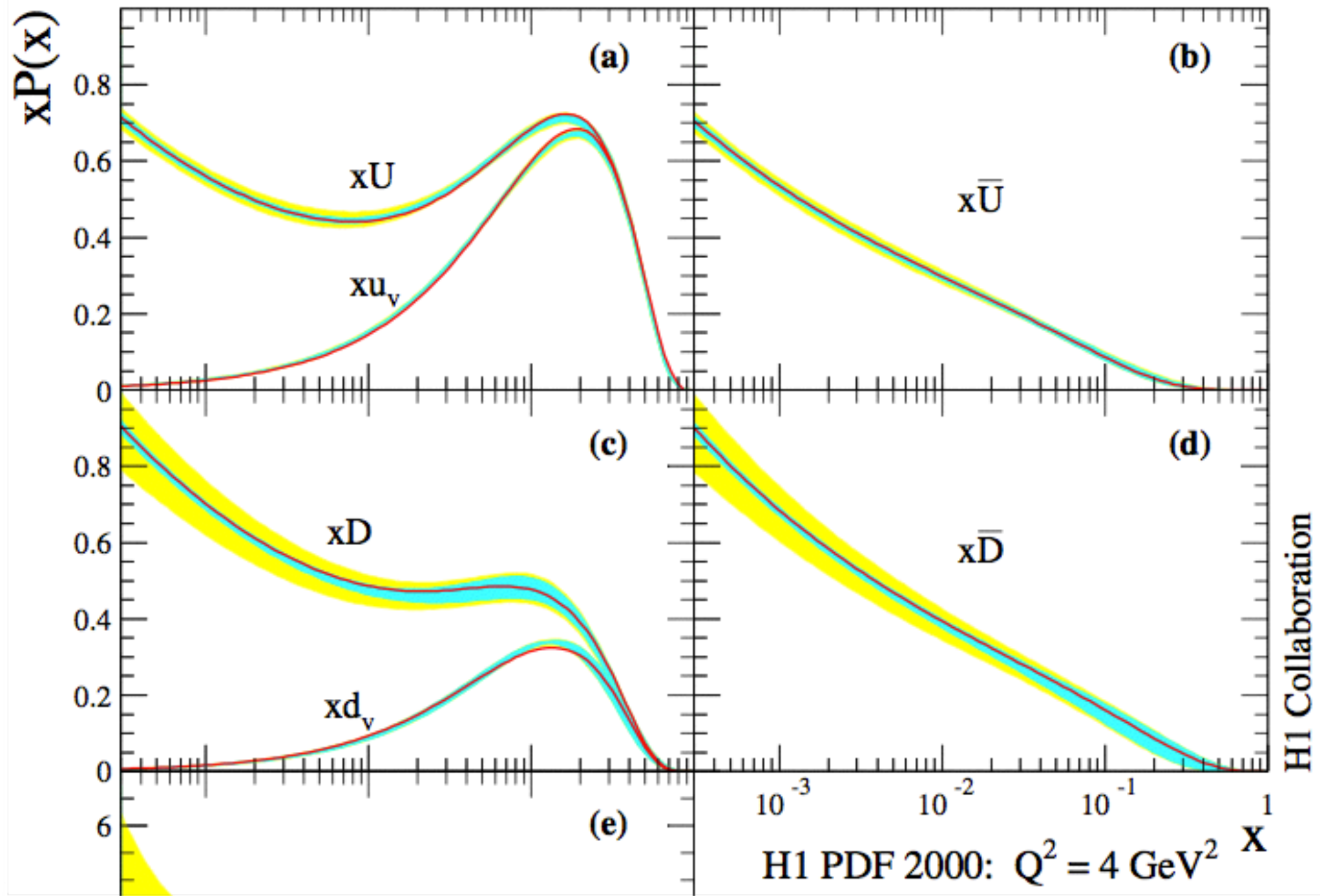
$$\frac{d^2\sigma_{CC}^\pm}{dx dQ^2} = \frac{G_F^2}{2\pi x} \left[ \frac{M_W^2}{Q^2 + M_W^2} \right]^2 \times \phi_{CC}^\pm (1 + \Delta_{CC}^{\pm,weak})$$



H1 data & fit [4]

# Measured data

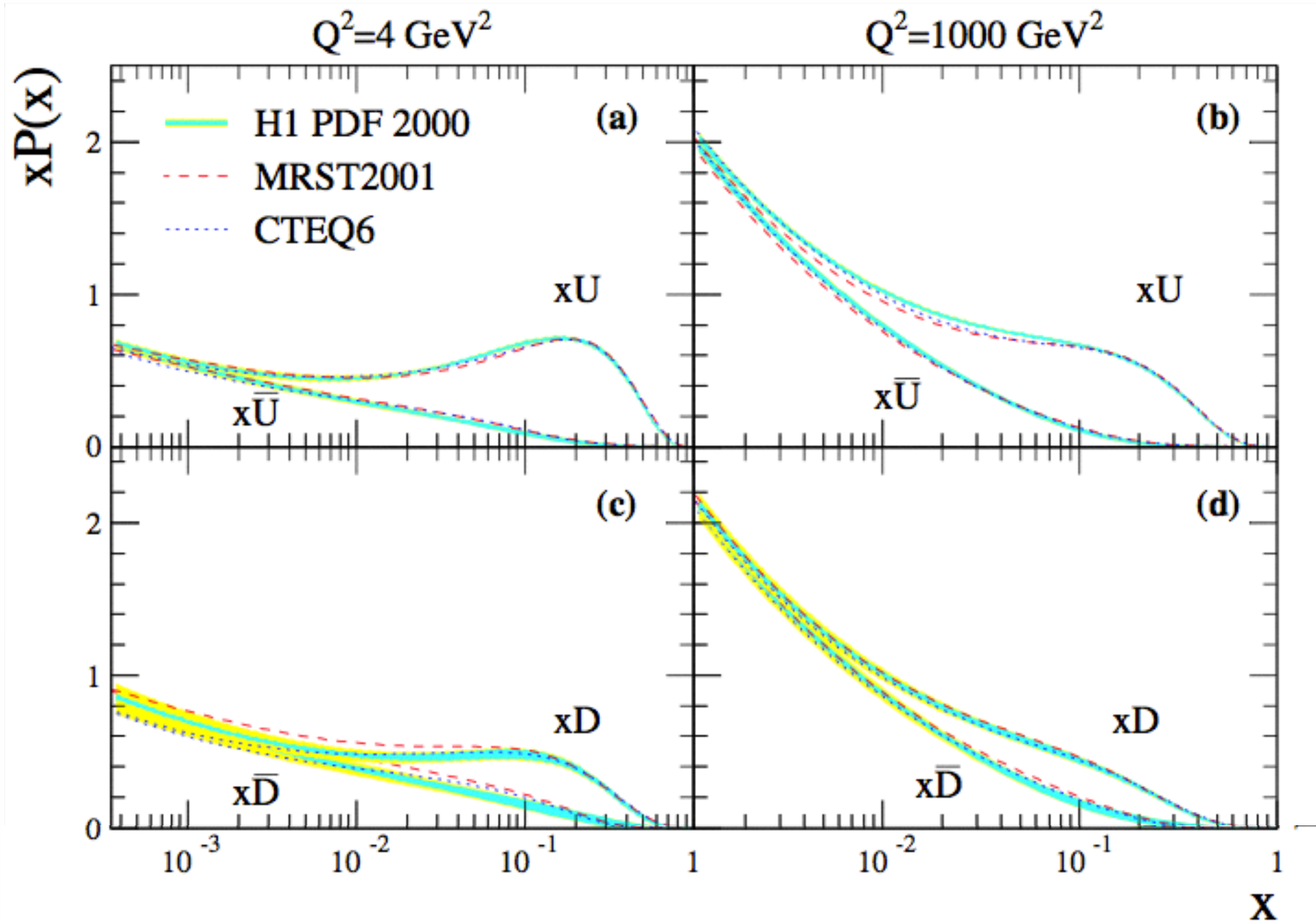
The measured cross sections for CC and NC. The bands show the results of the fit, which are accurate over 6 orders of magnitude in  $\sigma$ .



# Results

Final fitted models of PDFs at  $4 \text{ GeV}^2$ . Note  $q_v = Q - \bar{Q}$  are the valence quark PDFs.





# Results

PDFs varied by energy using the DGLAP evolution equations.

---

# Conclusions

---

- ❖ DIS is a fantastic probe of QCD dynamics
- ❖ There is a lot of complex structure hidden in what look like simple cross section plots
- ❖ Sure enough, the proton is made of  $uud$
- ❖ One final plot...

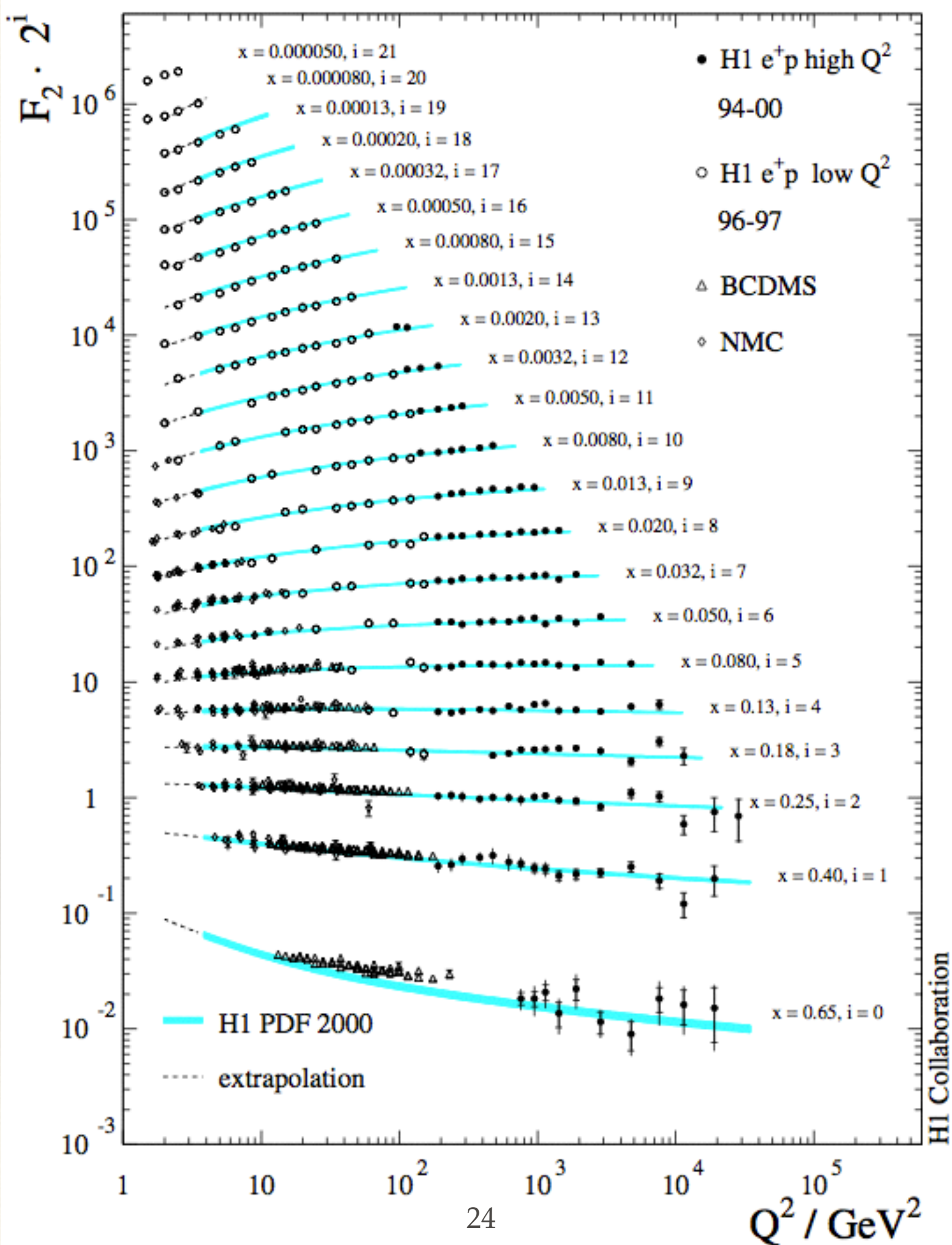
---

# References

---

- ❖ [1] PDG Structure Functions Review. K.A. Olive et al. (Particle Data Group), Chin. Phys. C, 38, 090001 (2014) and 2015 update.
- ❖ [2] H1 Events Tutorial <http://www-h1.desy.de/pictures/H1-short.event.tutorial2.pdf>.
- ❖ [3] ZEUS 1993 Status Report <http://www-zeus.desy.de/bluebook/>.
- ❖ [4] The H1 Collaboration, Eur. Phys. J. C **30**, 1 (2003).
- ❖ (to learn about DGLAP) [http://www.scholarpedia.org/article/QCD\\_evolution\\_equations\\_for\\_parton\\_densities](http://www.scholarpedia.org/article/QCD_evolution_equations_for_parton_densities)





Thank You