

Non-perturbative techniques for constraining the cosmological collider



Oliver Philcox



J. Colin Hill



Lam Hui

Sam Goldstein



*SG, Philcox, Hill, Hui,
(in prep)*

[2209.06228](#)

*SG, Esposito, Philcox, Hui, Hill, Scoccimarro,
Abitbol*

[2310.12959](#)

SG, Philcox, Hill, Esposito, Hui

Introduction

- Spec-S5 surveys will ~~improve constraints on~~ help unveil the physics of inflation
 - Constrain **primordial non-Gaussianity (PNG)** beyond $f_{\text{NL}}^{\text{loc.}}$, $f_{\text{NL}}^{\text{eq.}}$, $f_{\text{NL}}^{\text{ort.}}$,
- **Cosmological collider:** primordial squeezed bispectrum from massive particles during inflation

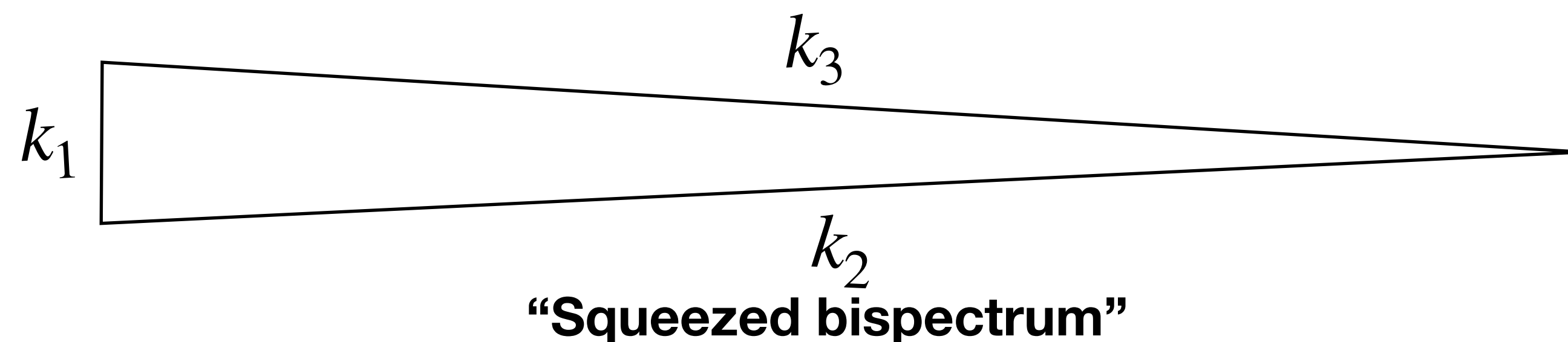
$$\lim_{k_1 \ll k_2 \approx k_3} B_{\Phi}(k_1, k_2, k_3) = 4f_{\text{NL}}^{\Delta} \left(\frac{k_1}{k_2}\right)^{\Delta} P_{\Phi}(k_1)P_{\Phi}(k_2), \quad \Delta = 3/2 - \sqrt{9/4 - m^2/H^2}$$

(Arkani-Hamed & Maldacena, 2015)

Case 1 – Massless ($m \ll H$; $\Delta \approx 0$): local non-Gaussianity

Case 2 – Intermediate mass ($0 < m/H \leq 3/2$; $0 < \Delta \leq 3/2$): between $f_{\text{NL}}^{\text{loc.}}$ and $f_{\text{NL}}^{\text{eq.}}$ (quasi-single field)

Case 3 – Massive ($m/H \geq 3/2$): oscillatory bispectrum



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How can we constrain this scenario using large-scale structure data?

Soft limits of LSS correlation functions

- Late-time LSS correlators contain primordial information
 - **Challenge:** non-linear structure formation
- Exploit two key properties of **soft limits** in LSS

1. Protected by symmetries

$$\lim_{q \rightarrow 0} \left[\frac{B_m(q, k, k')}{P_m(q)} \right] \text{ has no } q^{-\alpha} \text{ poles } (\alpha > 0)$$

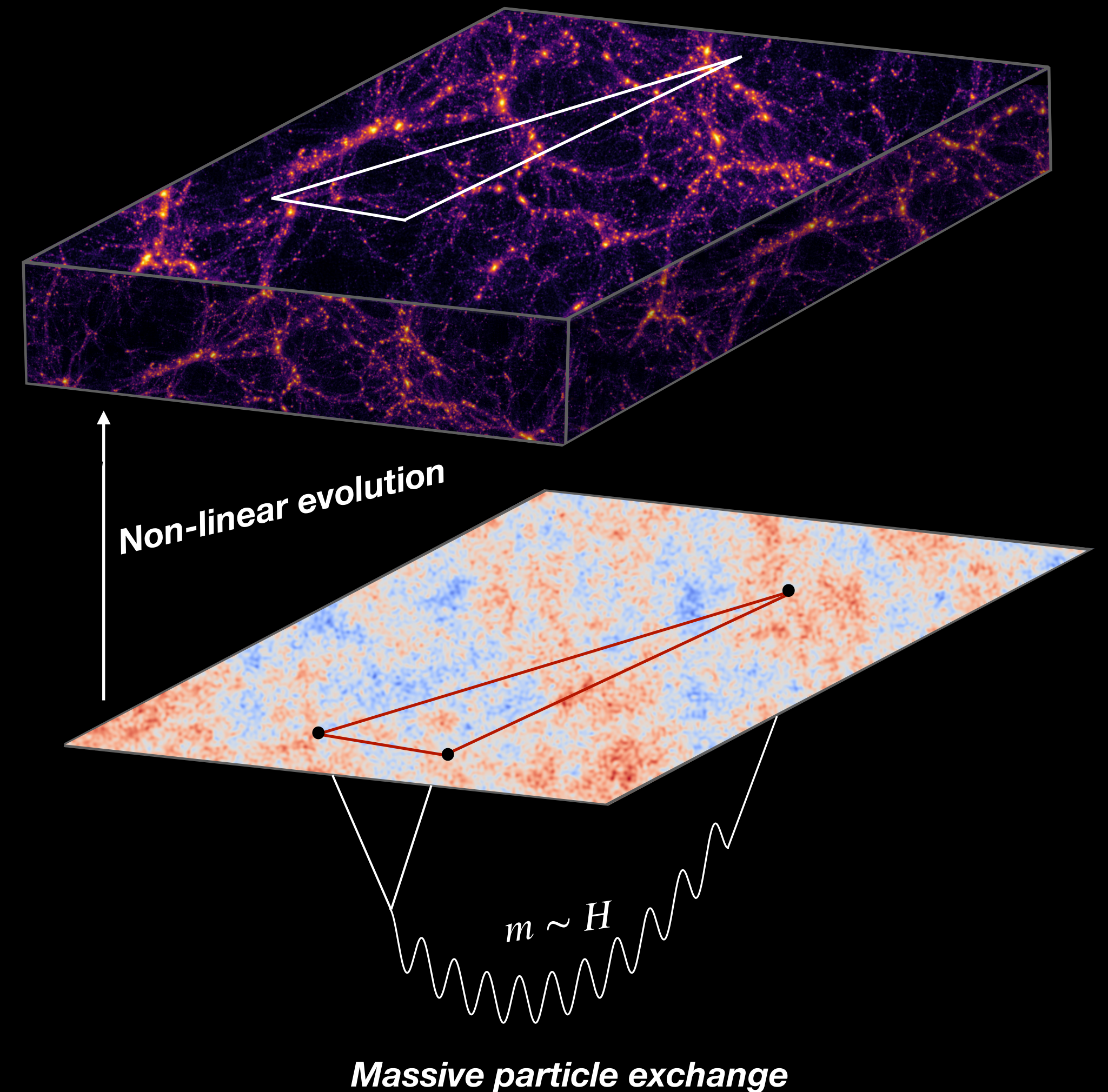
Newtonian LSS consistency relation

Kehagias & Riotto 2012;
Peloso & Pietroni, 2013

2. Tractable non-perturbative calculations using separate universe

- **(1)+(2)** used to derive non-perturbative estimators for $f_{\text{NL}}^{\text{loc.}}$ (**SG+22/Giri+23**)

- How do we generalize these to cosmo. collider?



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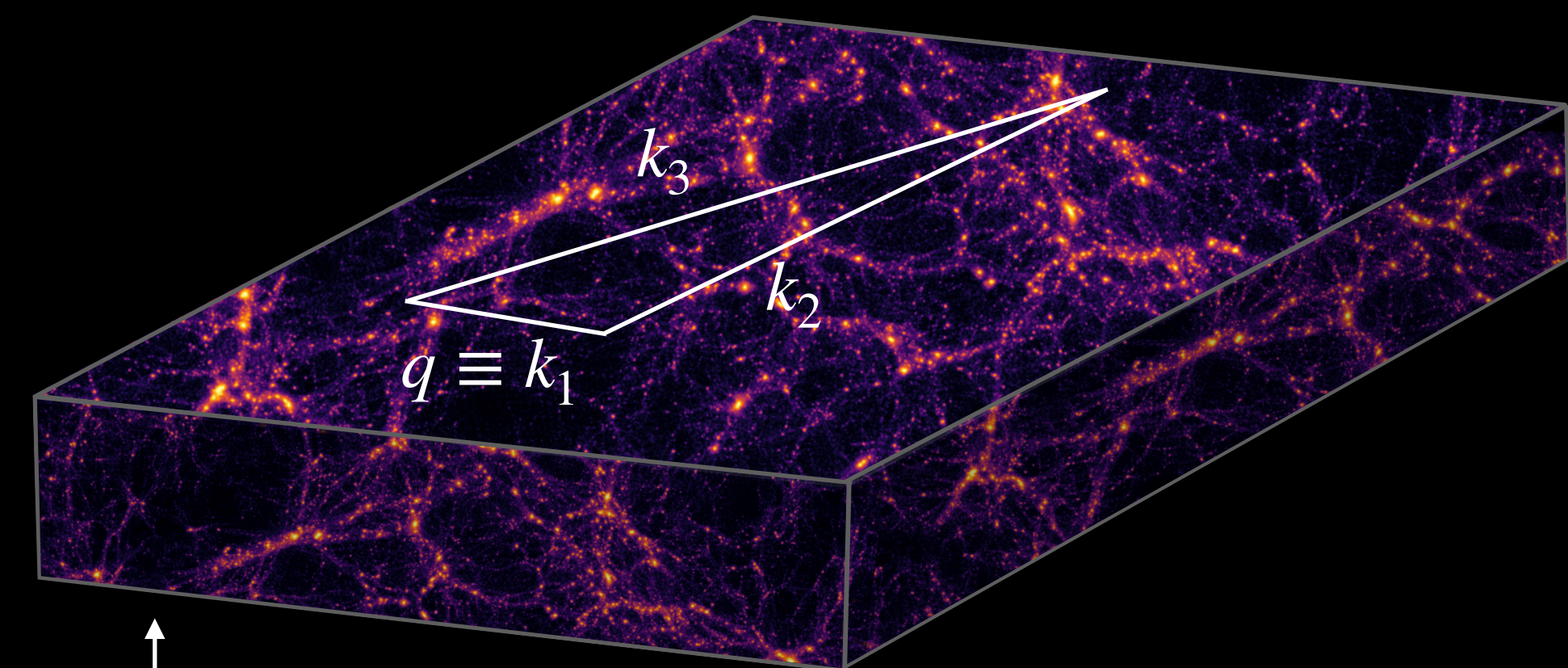
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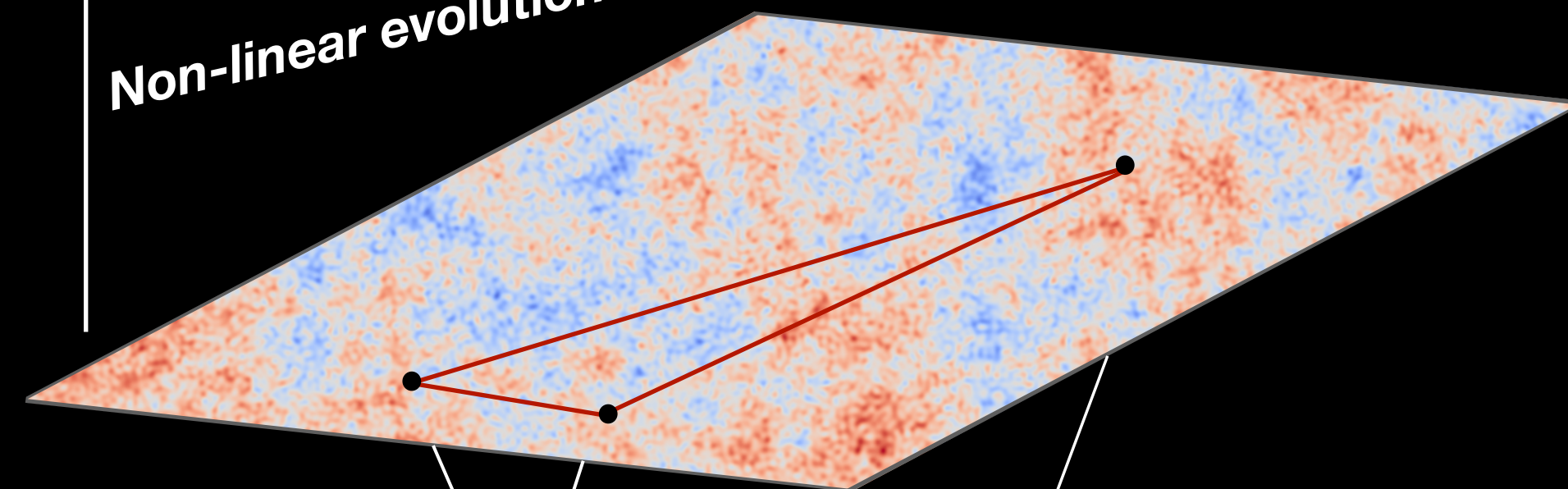
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Non-linear evolution



$$q \equiv k_1 \ll k_2 \approx k_3$$

“Soft limit”

Massive particle exchange

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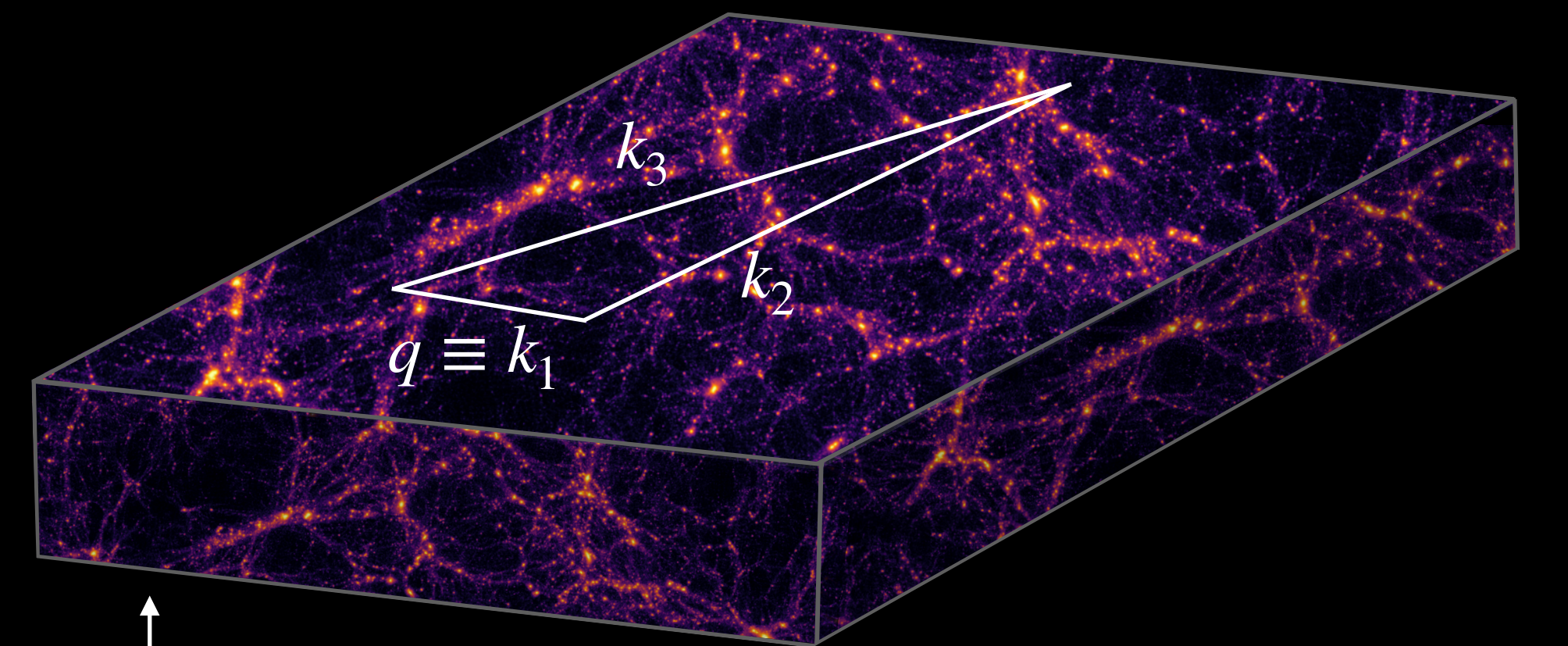
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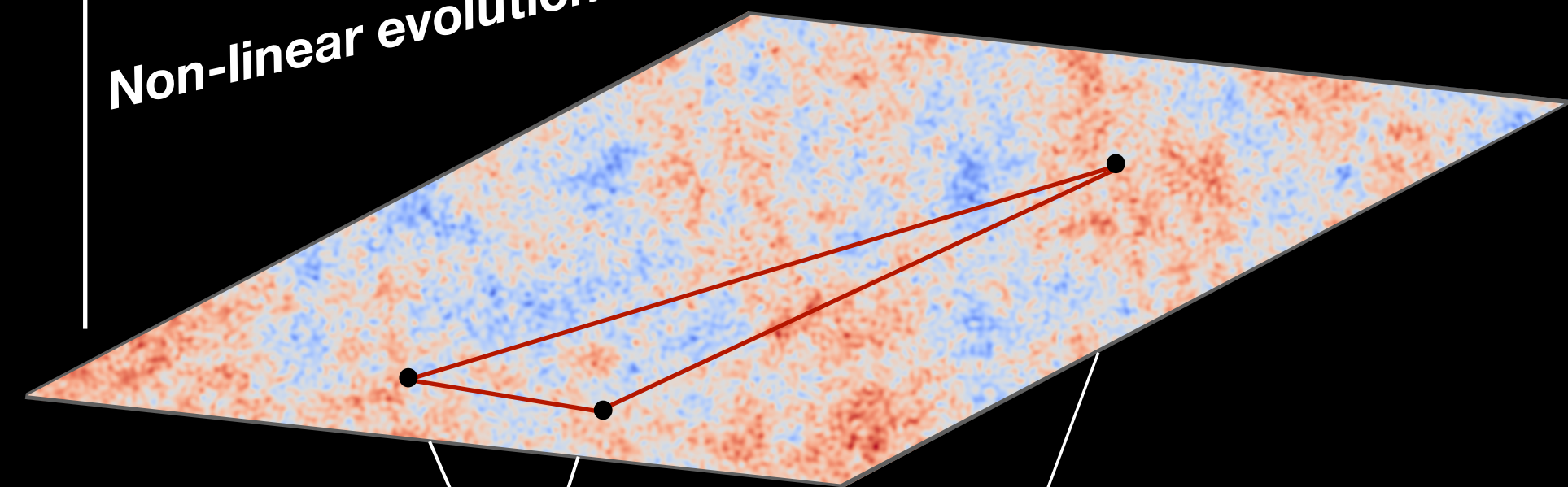
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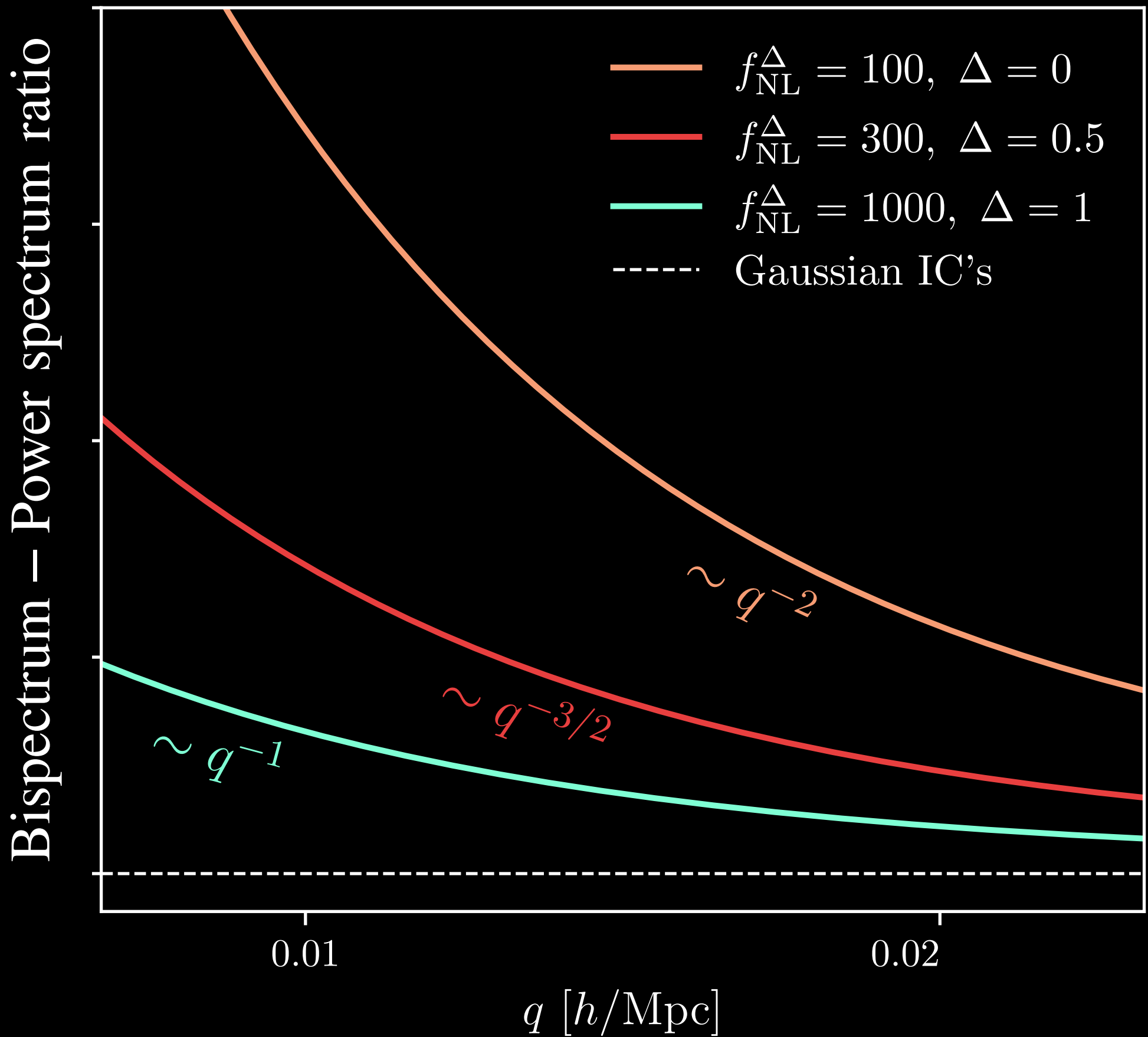
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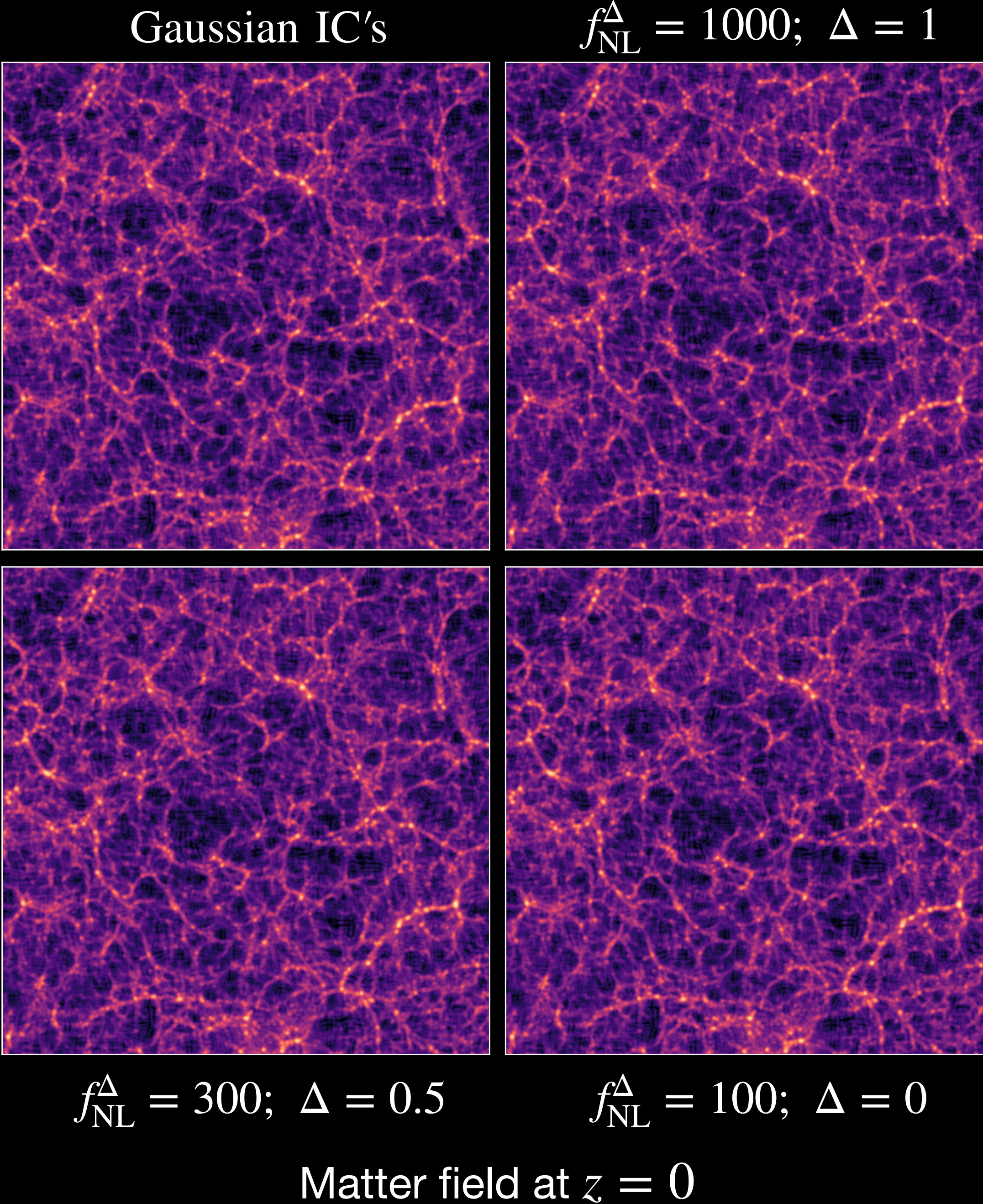
$$m \sim H$$

Massive particle exchange

N-body simulations with collider bispectrum

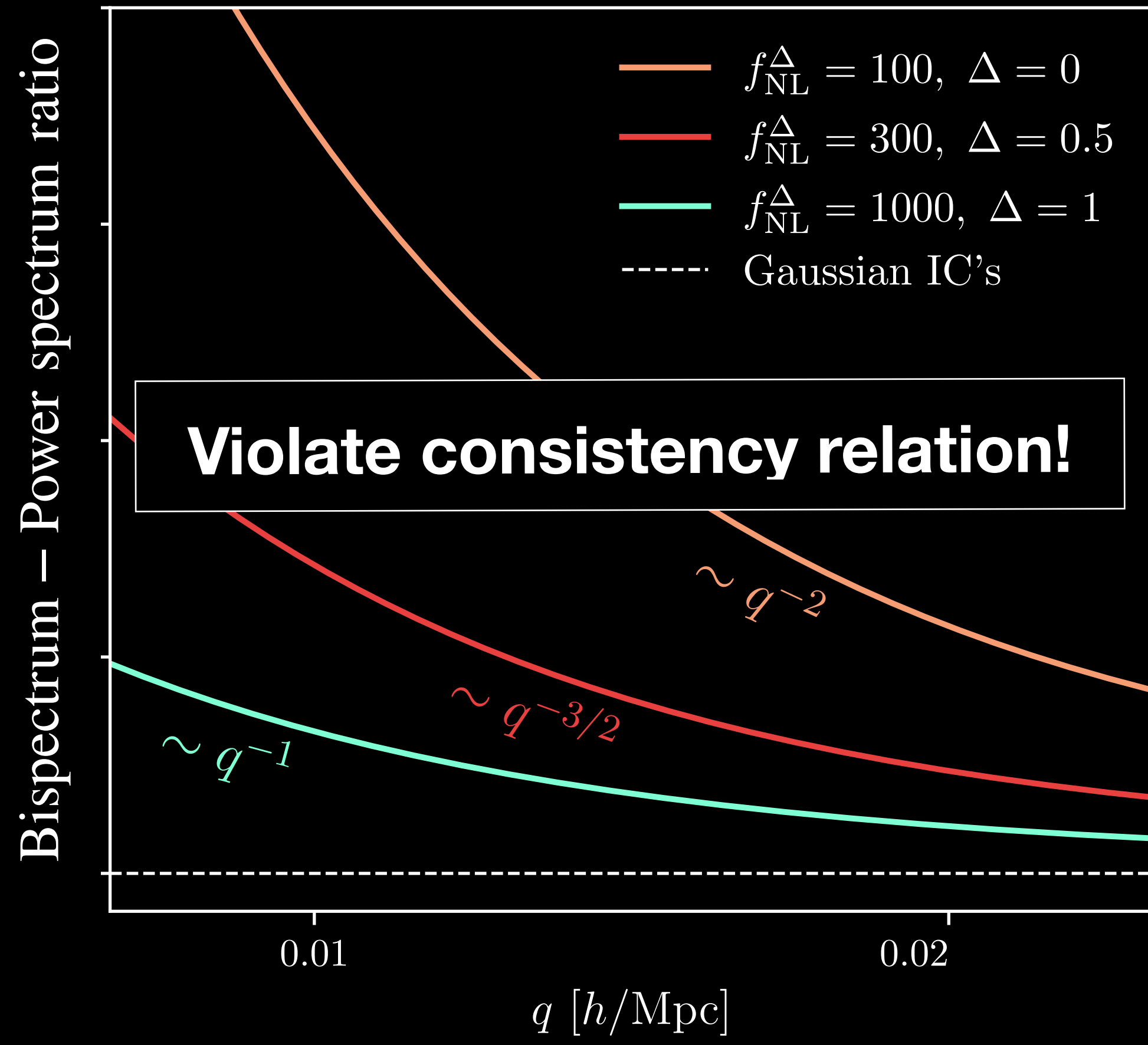


→
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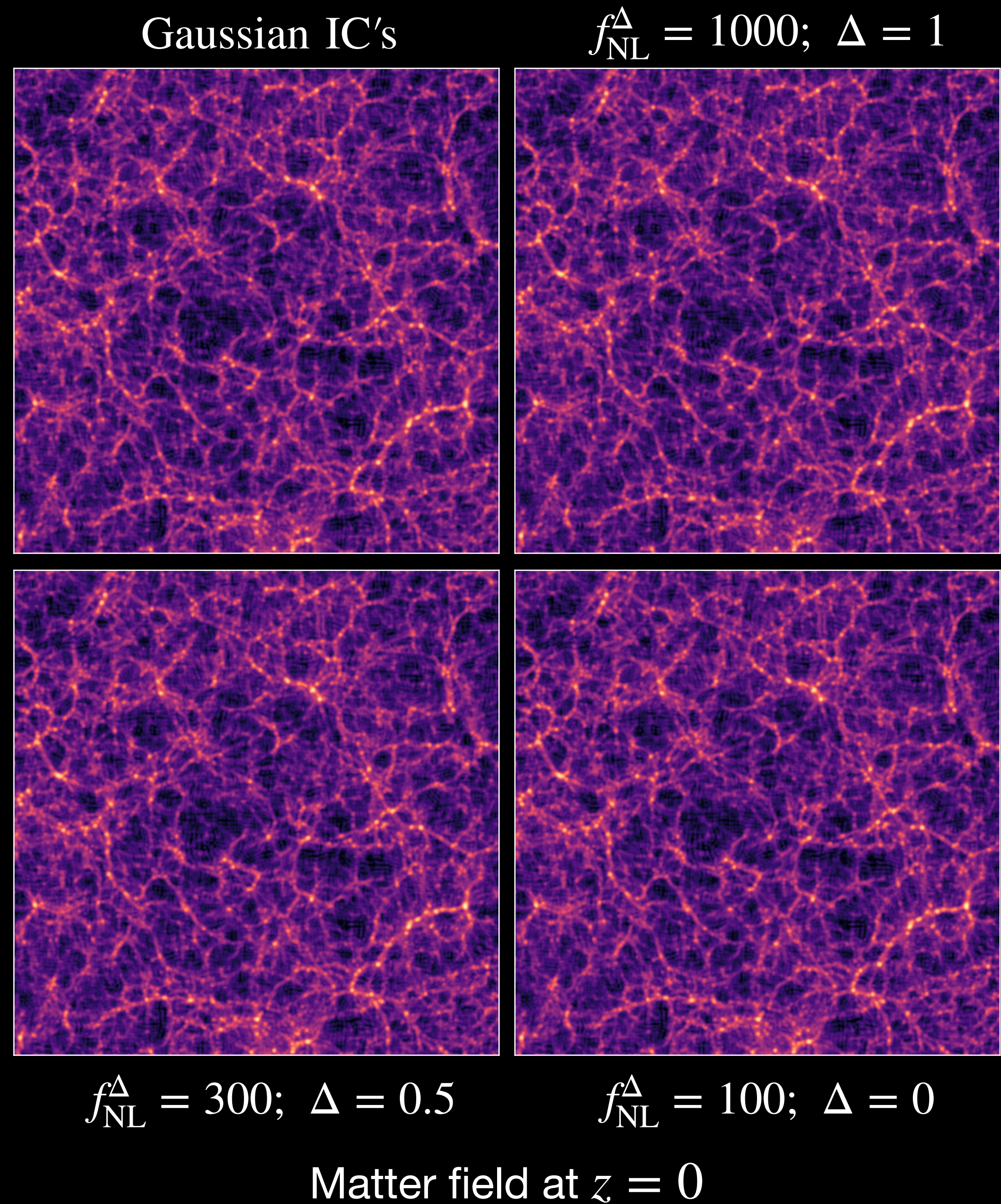


- Added squeezed collider templates to **2LPTPNG**
- Ran suite of simulations with same settings as **QuijotePNG**, but collider primordial bispectrum

N-body simulations with collider bispectrum



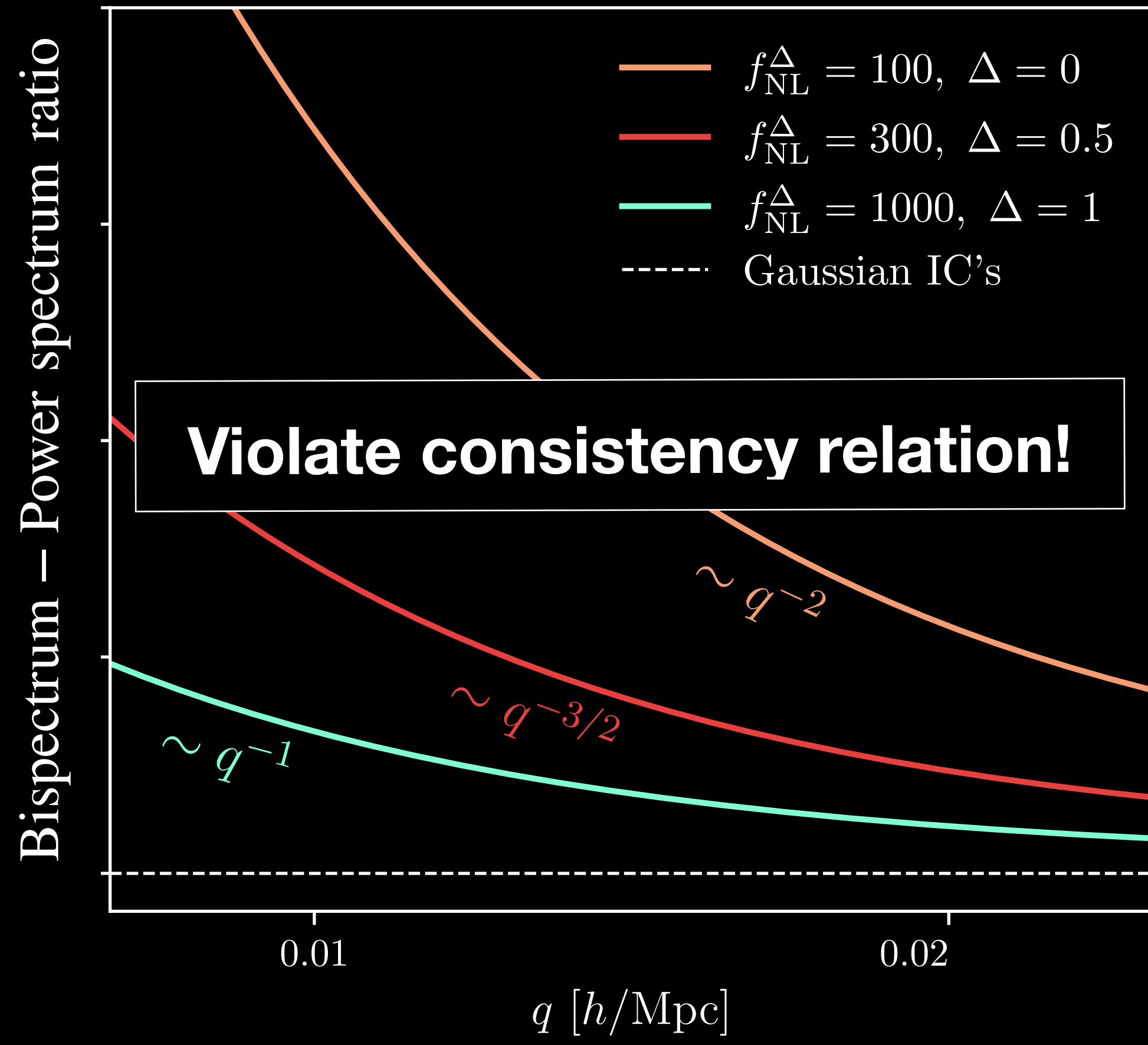
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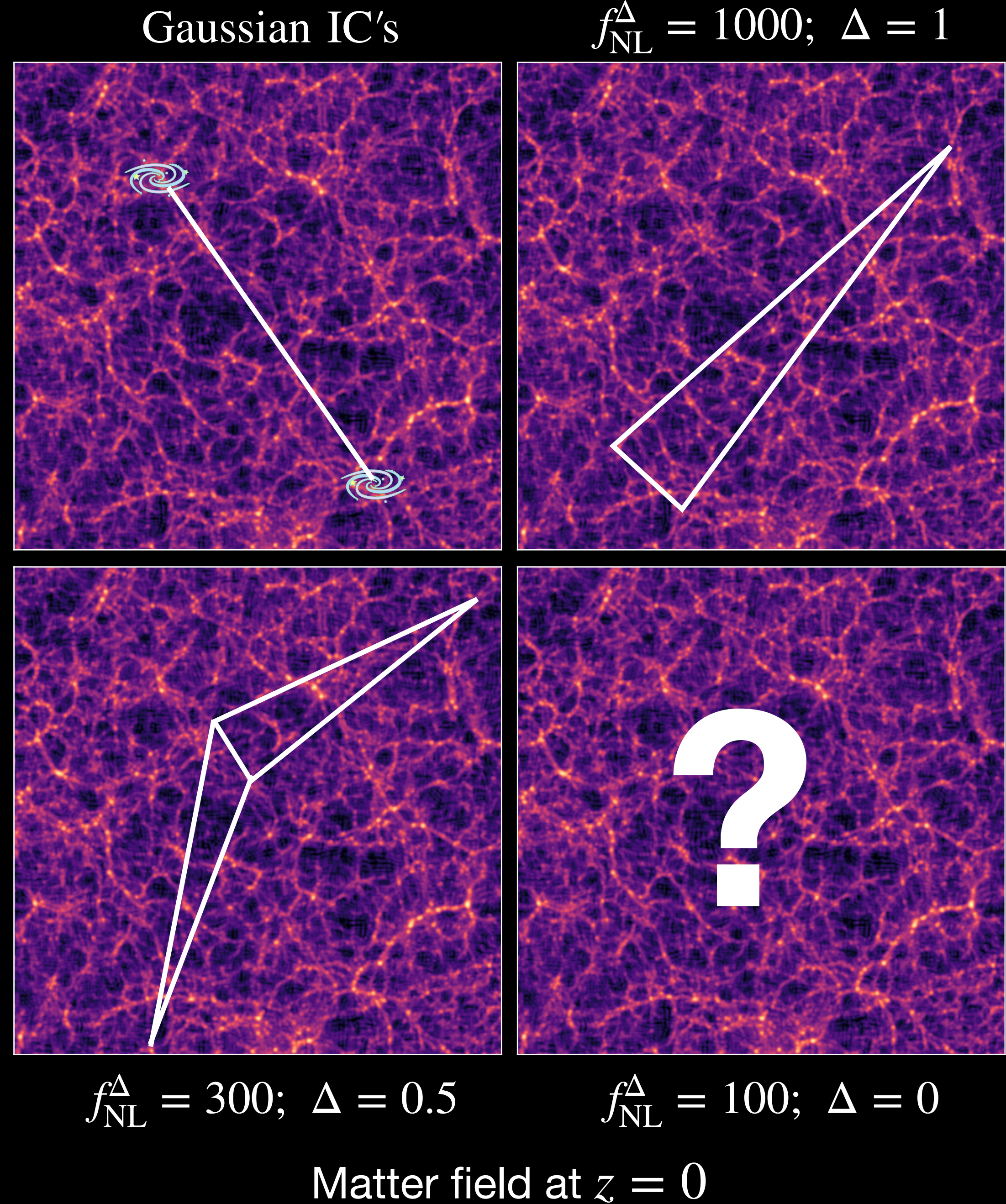
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Matter field at $z = 0$

N-body simulations with collider bispectrum

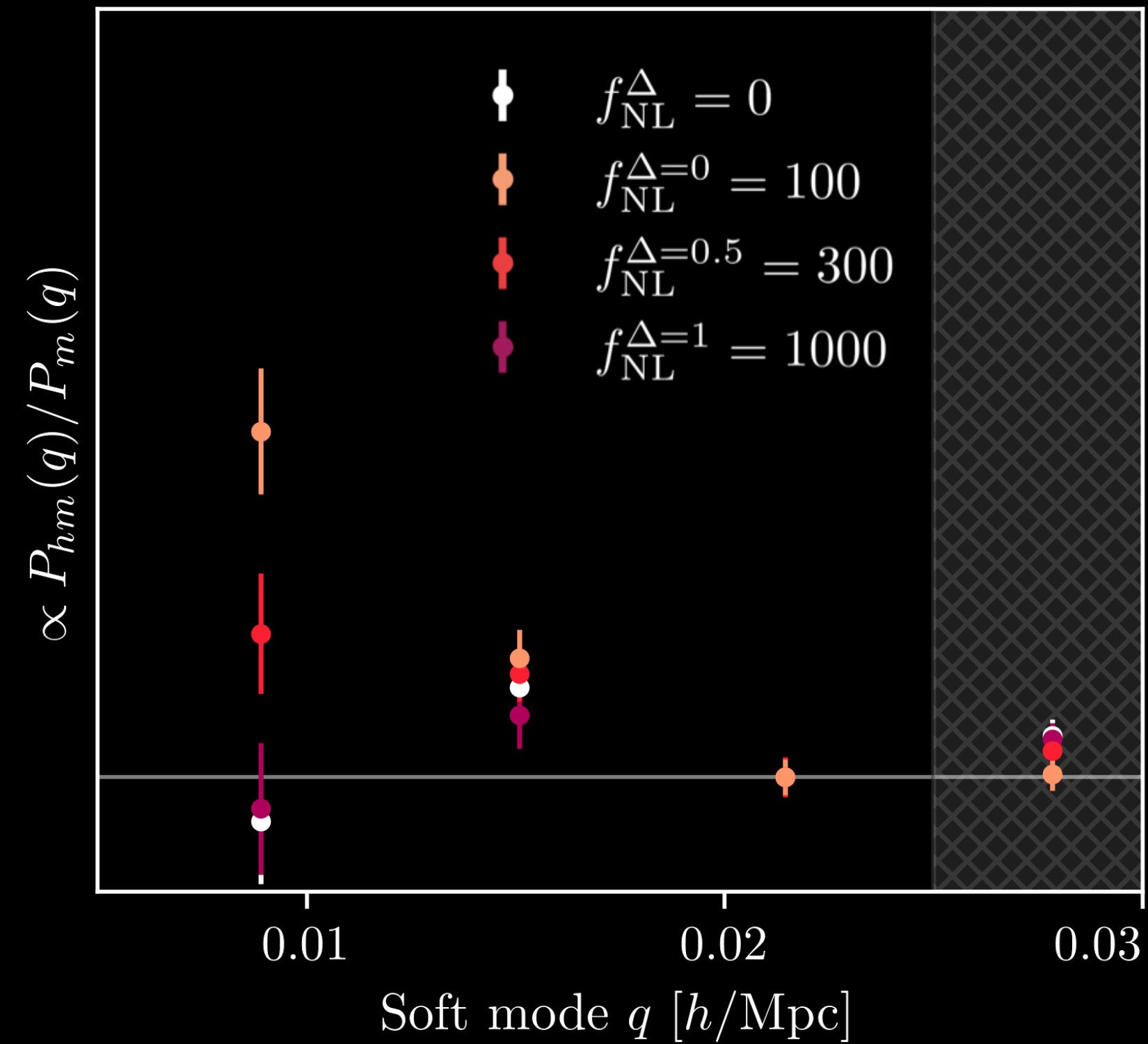


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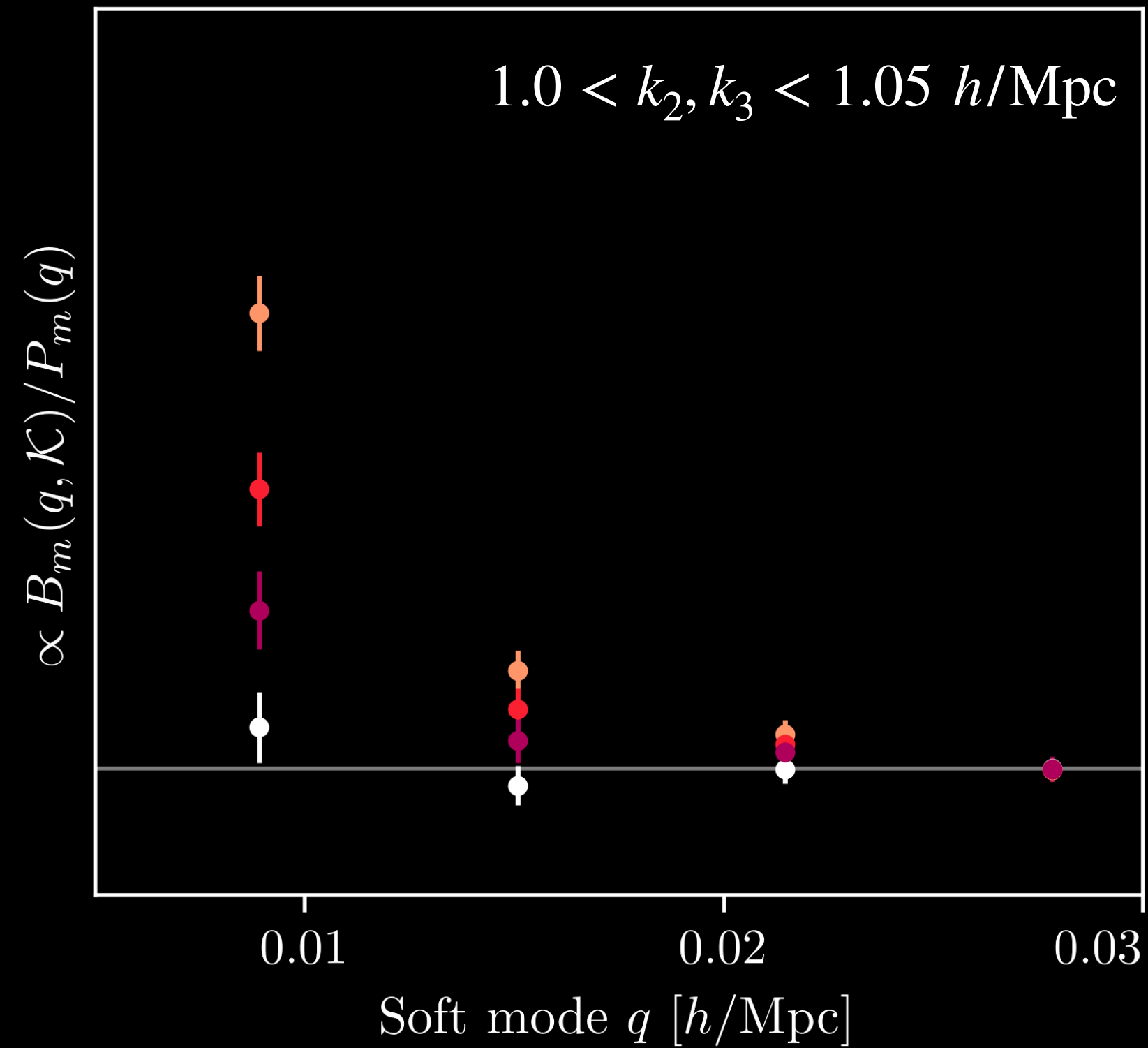
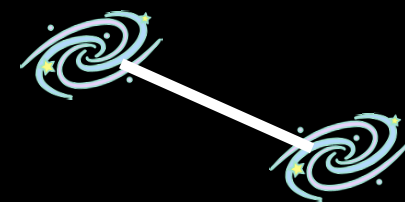


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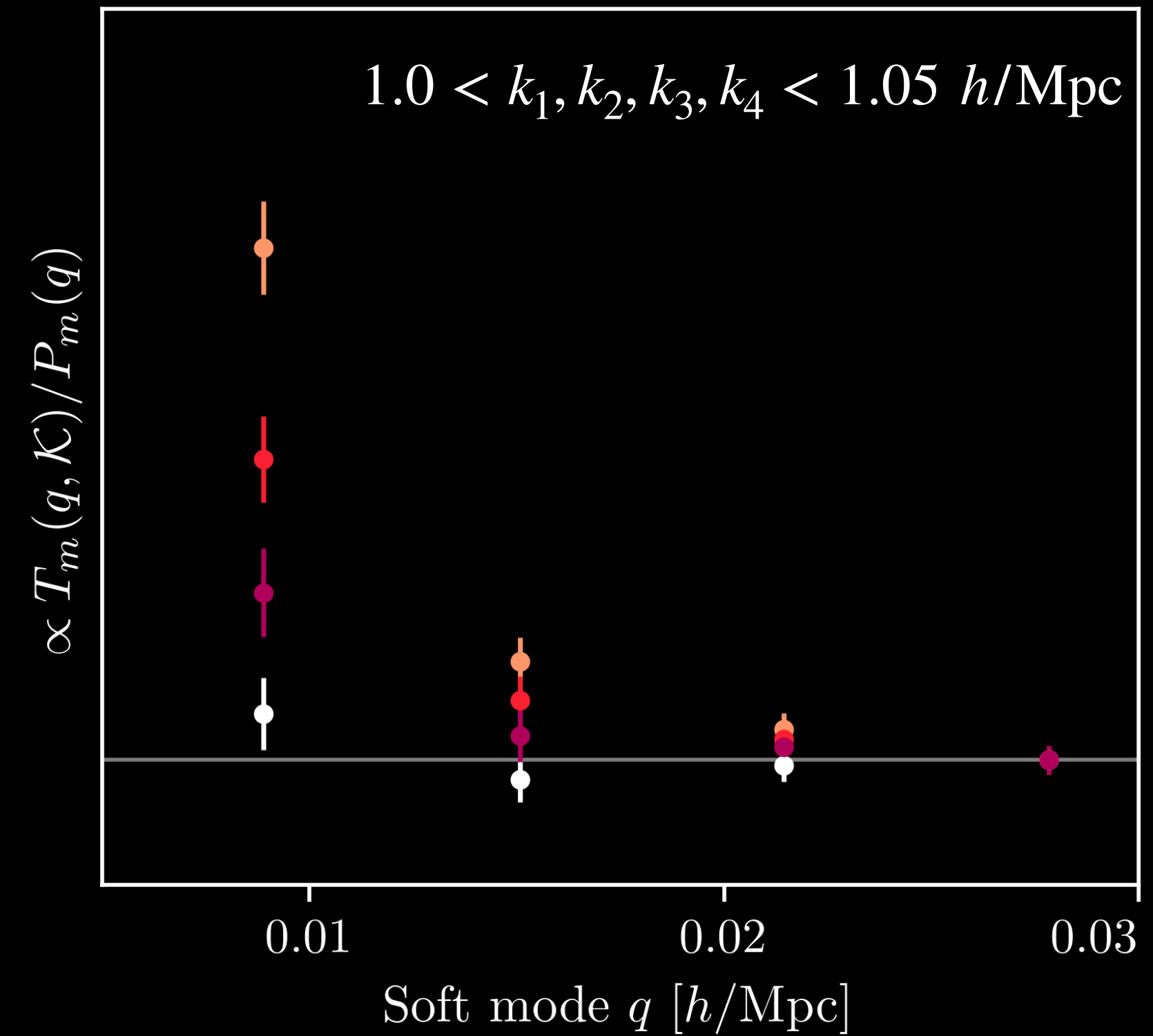
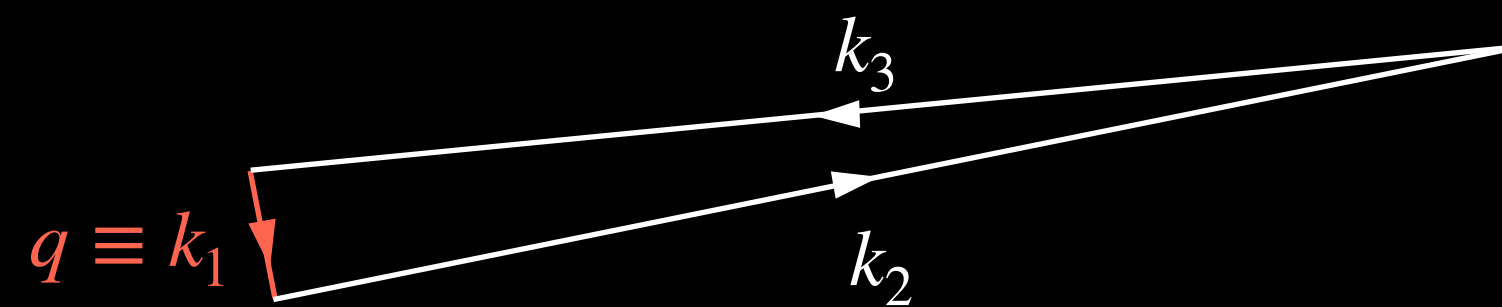
Imprints of the cosmological collider on LSS correlators



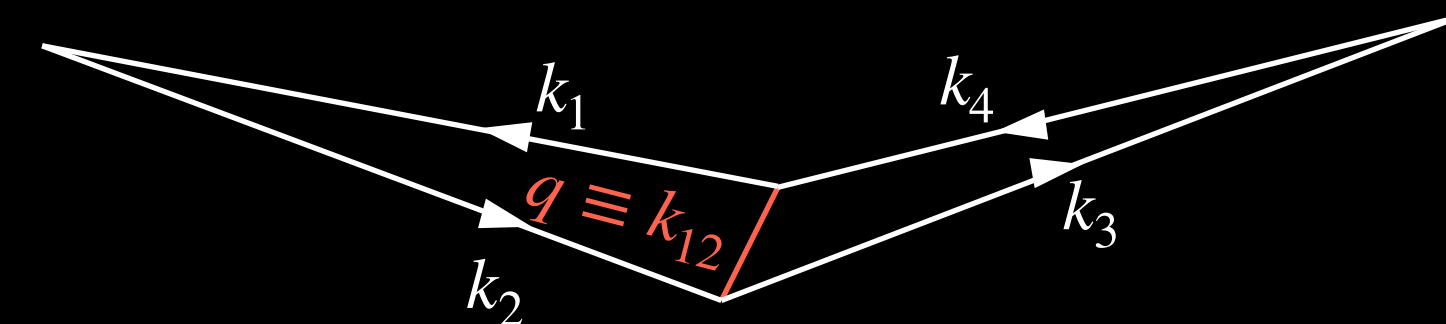
Scale-dependent halo bias



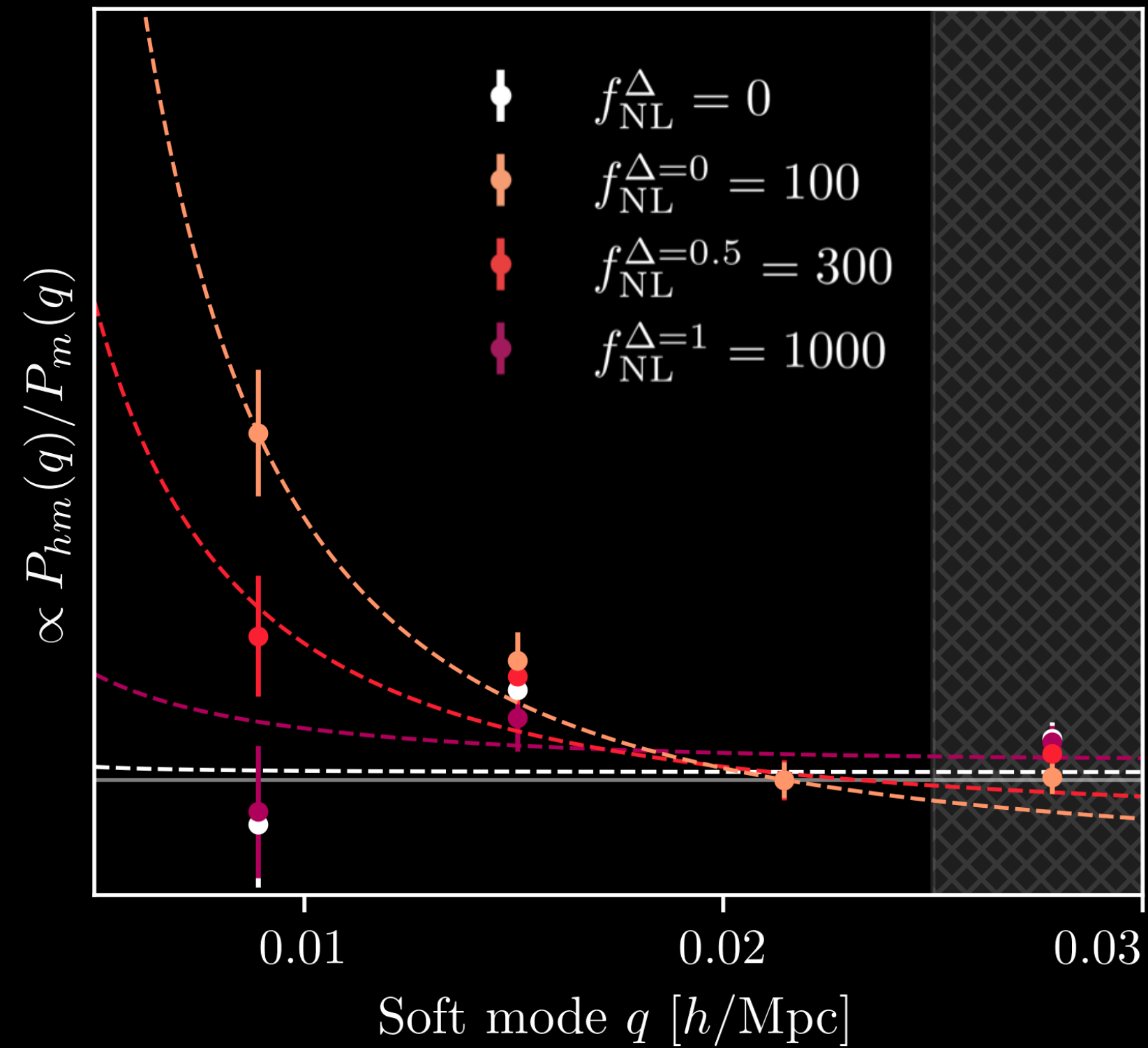
Squeezed matter bispectrum



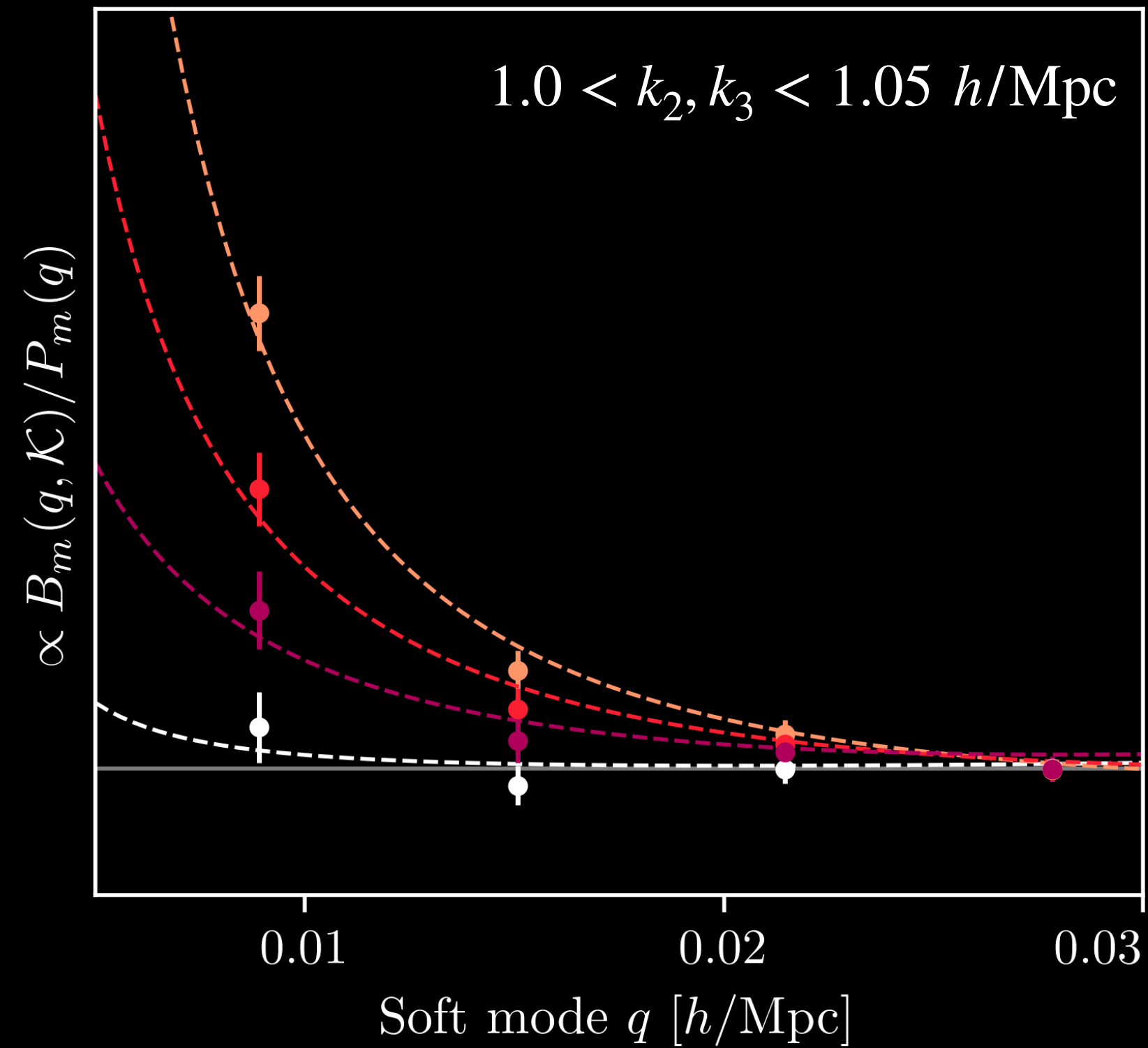
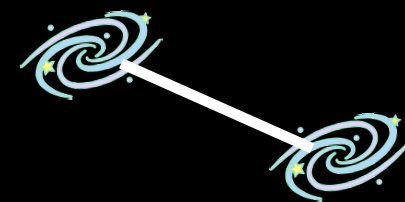
Collapsed matter trispectrum



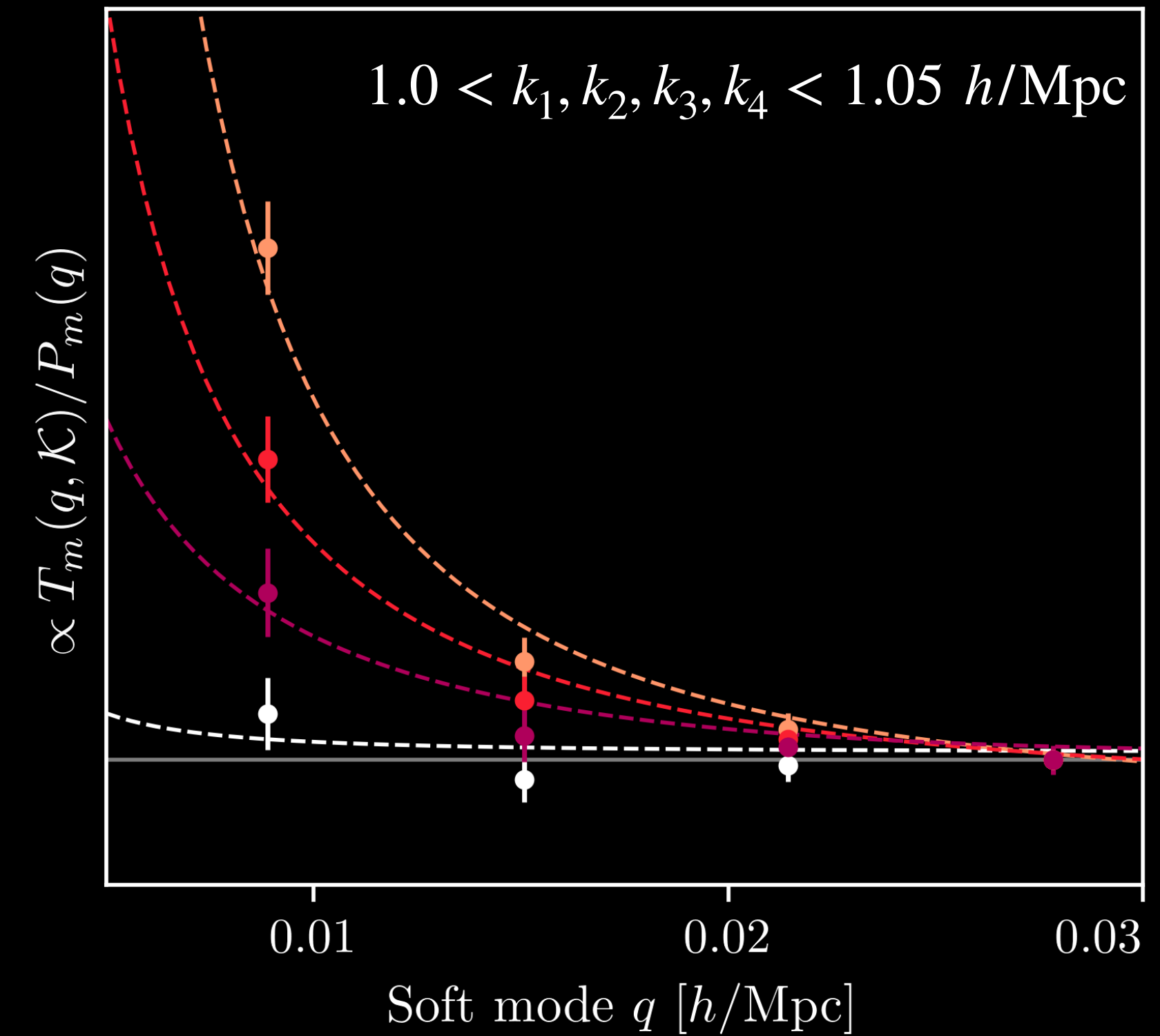
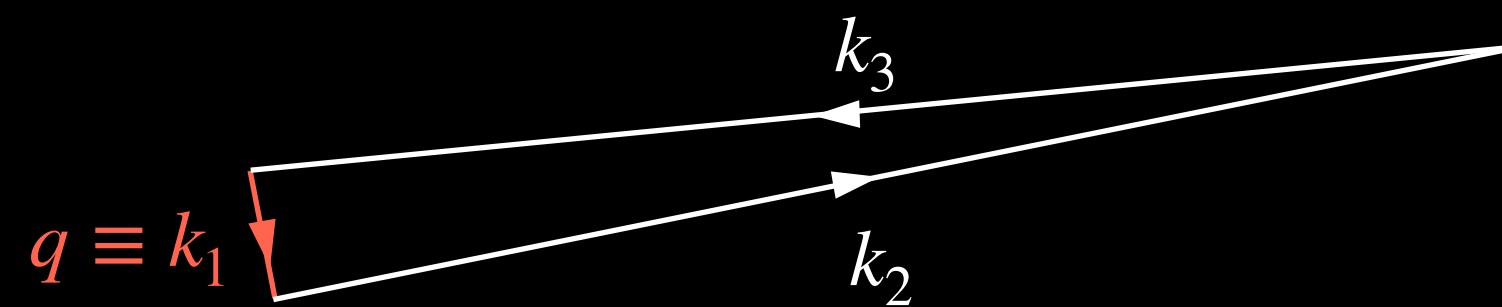
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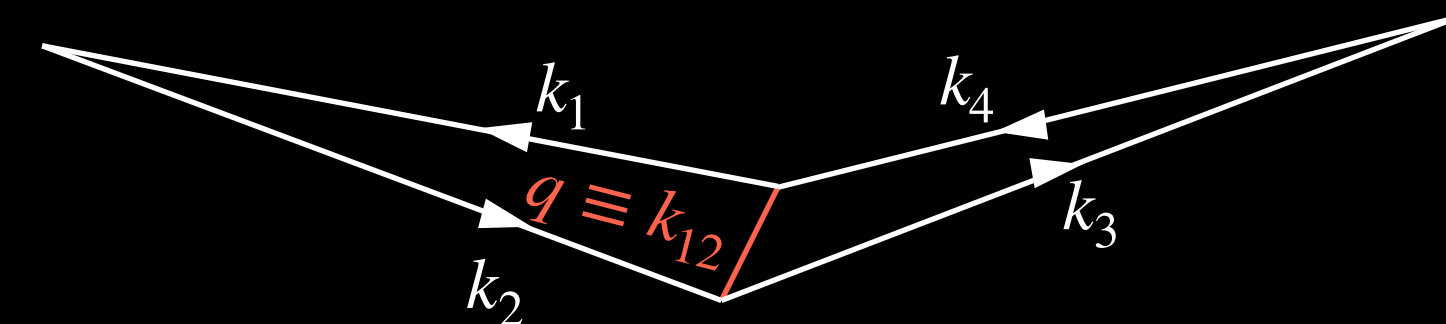
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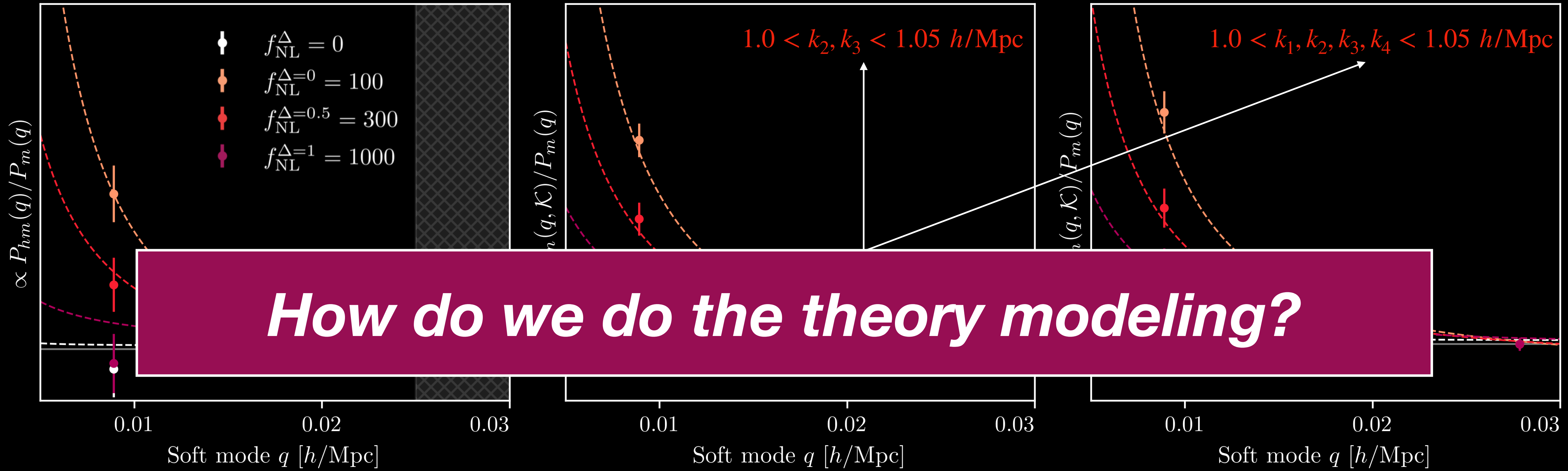
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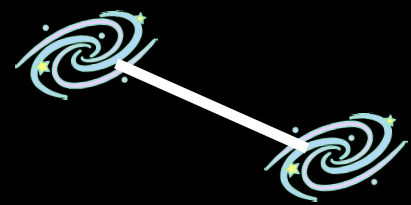
Collapsed matter trispectrum



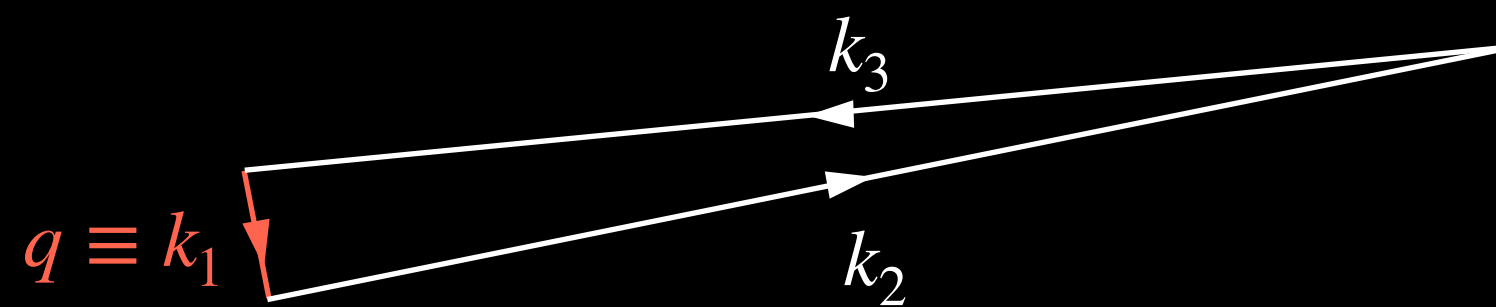
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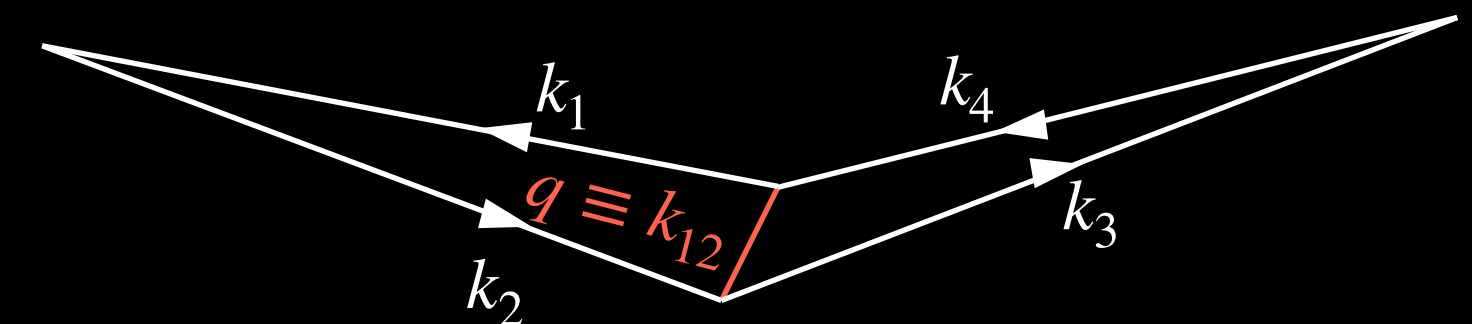
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Squeezed matter bispectrum



Collapsed matter trispectrum



Non-perturbative bispectrum model

- Squeezed bispectrum ($q \ll k$) is **modulation** of small scale power spectrum by a long-wavelength gravitational potential $\Phi_L(\mathbf{x})$
- Contributions from **PNG** are associated with $\partial P_m(k)/\partial\Phi_L(q)$
- Contributions from **gravitational non-Gaussianity** modeled using response approach

$$\lim_{q \ll k_{\text{NL}}, k} B_m(q, k, k') = \frac{3\Omega_{m0}H_0^2}{2q^2 T(q) D_{\text{md}}(z_q)} P_m(q) \left(\frac{\partial P_m(k)}{\partial \Phi_L(q)} \right) + \bar{a}_0 P_m(q) P_m(k) + \bar{a}_2 \frac{q^2}{k^2} P_m(q) P_m(k)$$

(e.g., Valageas 13; Chiang+17; Esposito+19; Biagetti+22)

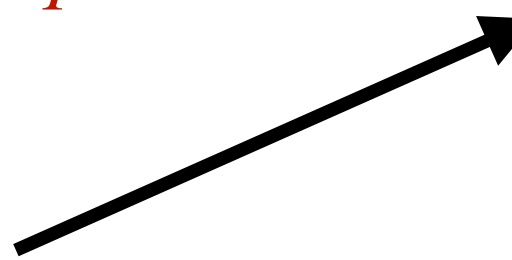
- How to estimate this *non-linear* potential derivative?

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- How to estimate this *non-linear* potential derivative?

Run separate universe simulations

Separate universe and the cosmological collider

- Consider small-scale modes $\delta_m(\mathbf{k}_1)$ and $\delta_m(\mathbf{k}_2)$ with **fixed amplitude**, but $k_1 < k_2$
- Add in **background** ($\sim \text{const.}$) potential fluctuation $\Phi_L(\mathbf{q})$
- PNG induces mode coupling

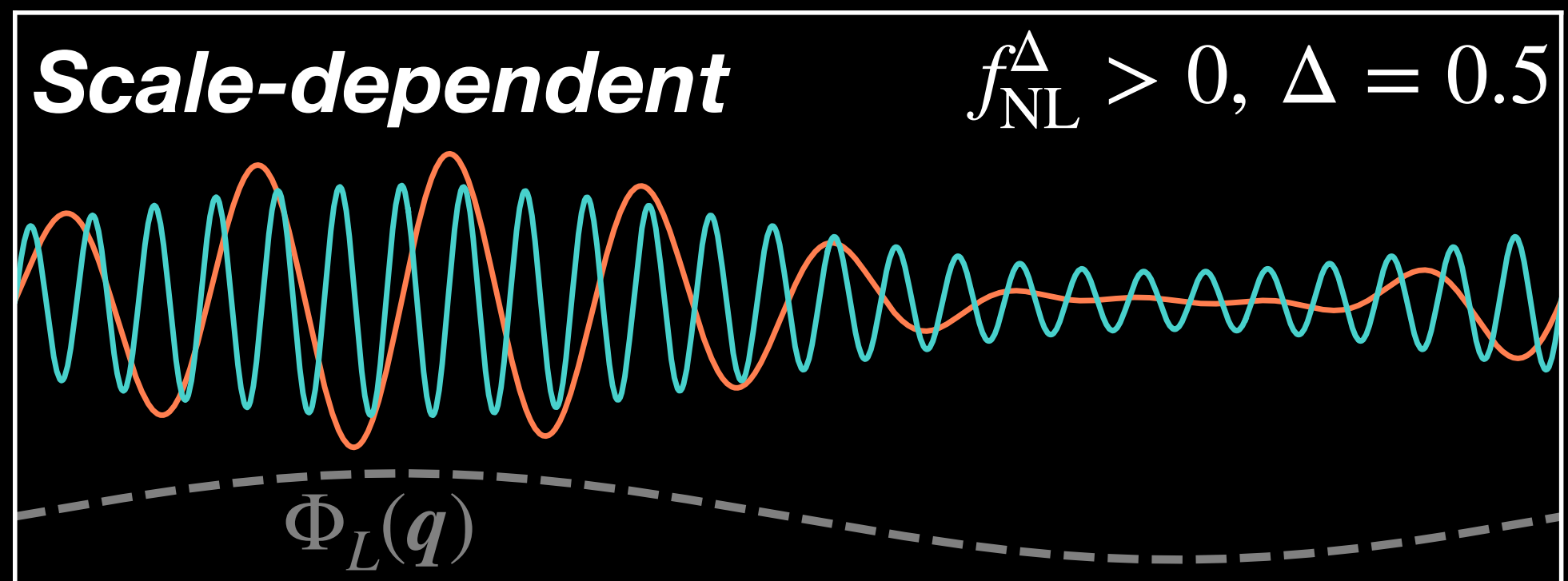
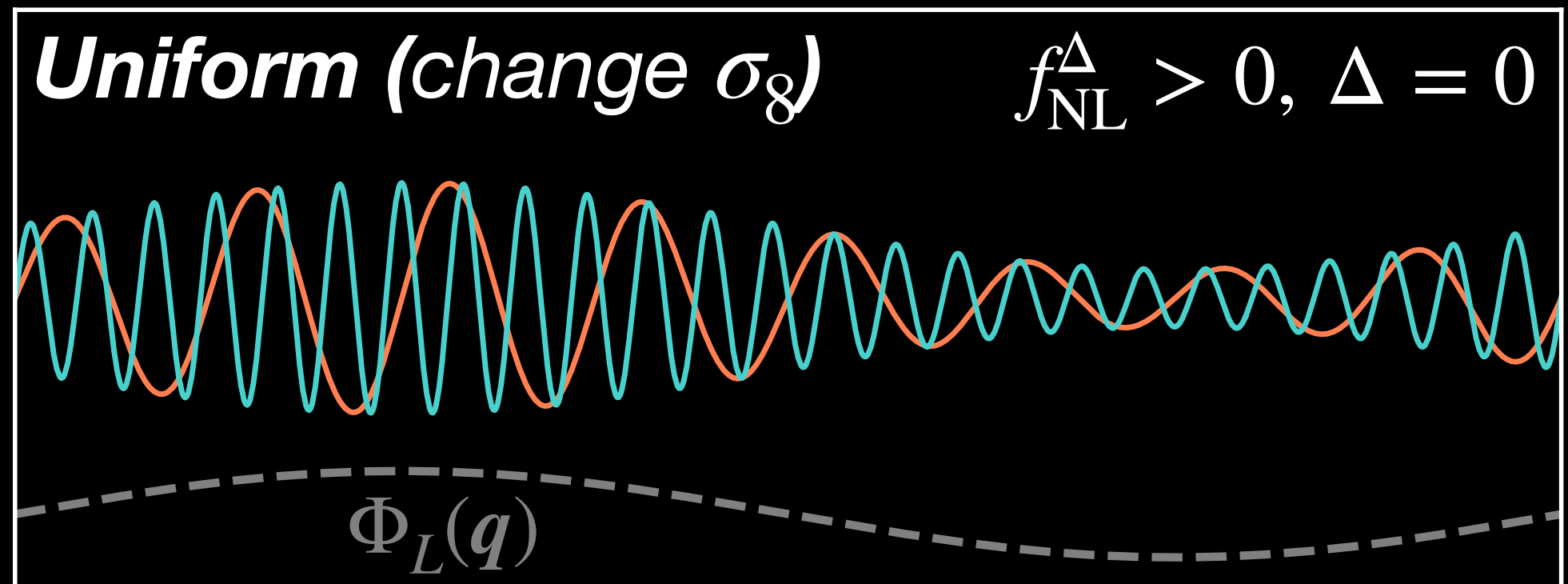
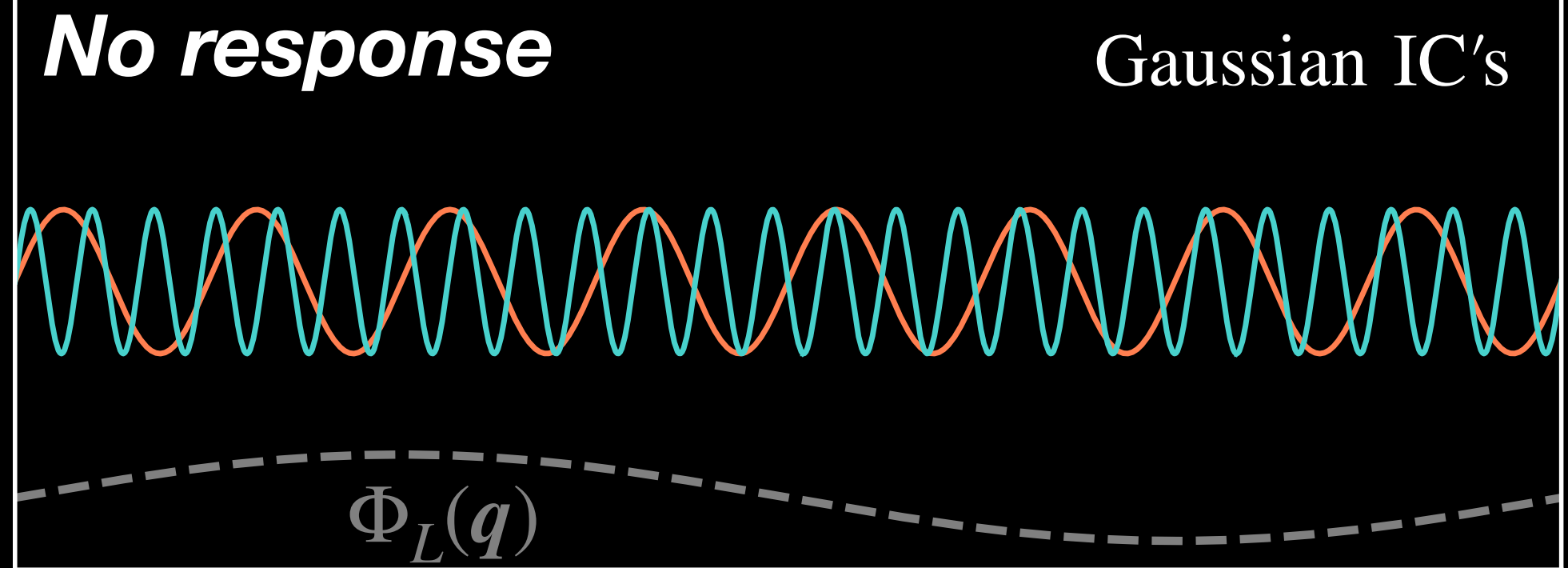
$$\delta_m(\mathbf{k} | \Phi_L(\mathbf{q})) = \left(1 + 2f_{\text{NL}}^\Delta \left(\frac{q}{k} \right)^\Delta \Phi_L(\mathbf{q}) \right) \delta_m(\mathbf{k})$$

Schmidt, Jeong, Desjacques, 2012

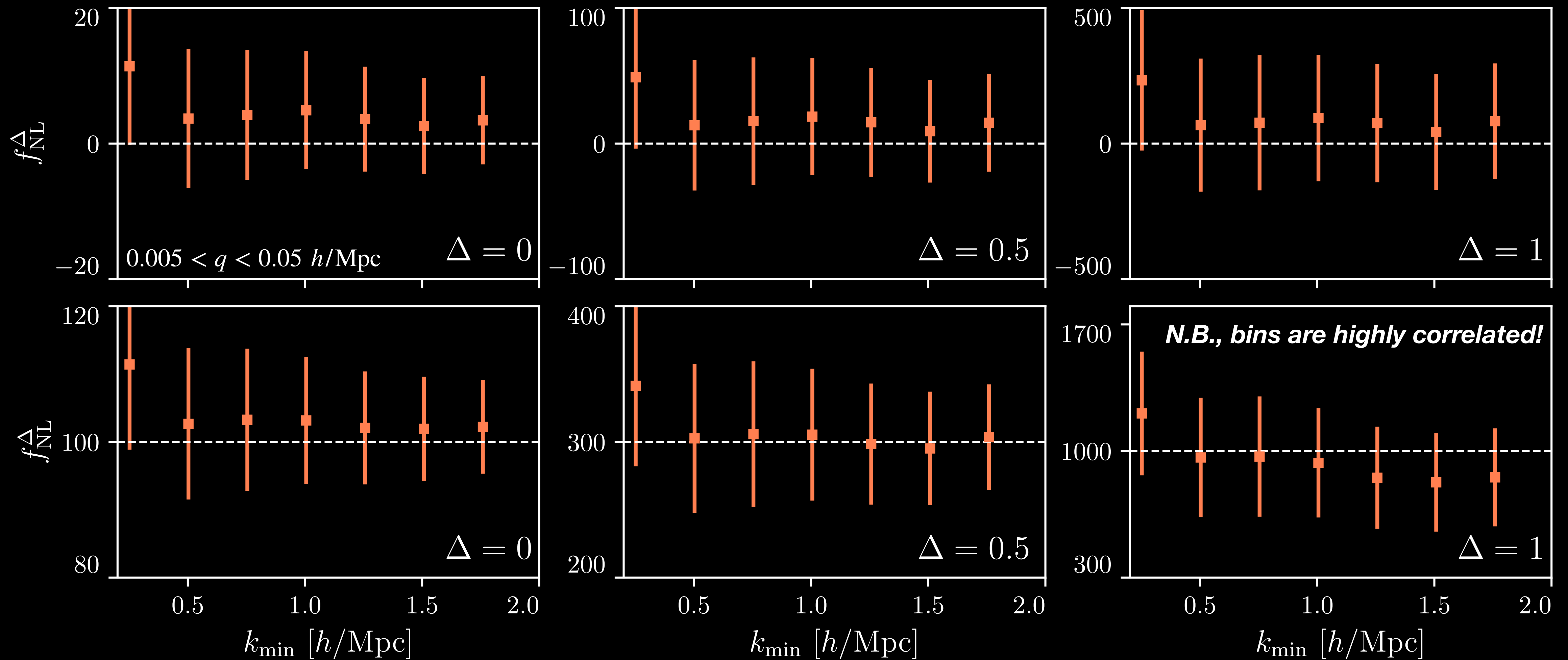
- Can compute $\partial P(k)/\partial \Phi_L(q)$ from sims with modified $P_m^{\text{lin.}}$

$$P_m^{\text{lin.}}(k | \epsilon, \Delta) \equiv (1 + 2\epsilon k^{-\Delta}) P_m^{\text{lin.}}(k).$$

$$\frac{\partial P_m(k)}{\partial \Phi_L(\mathbf{q})} = 2f_{\text{NL}}^\Delta q^\Delta \frac{\partial P_m(k | \epsilon, \Delta)}{\partial \epsilon} \Bigg|_{\epsilon=0}.$$

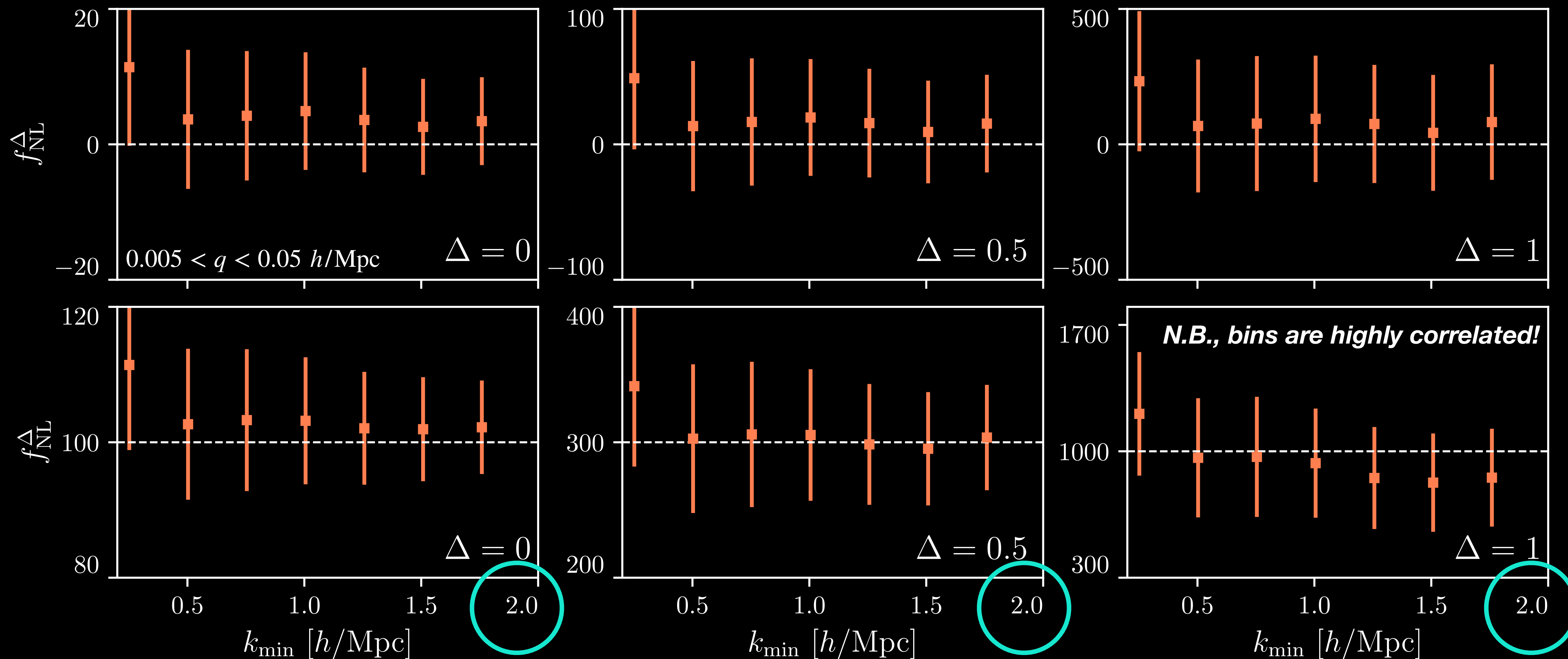


Constraints from squeezed bispectrum at $z = 0$



Unbiased constraints on f_{NL}^{Δ} using ***non-linear*** squeezed matter bispectrum!

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Scale-dependent halo bias

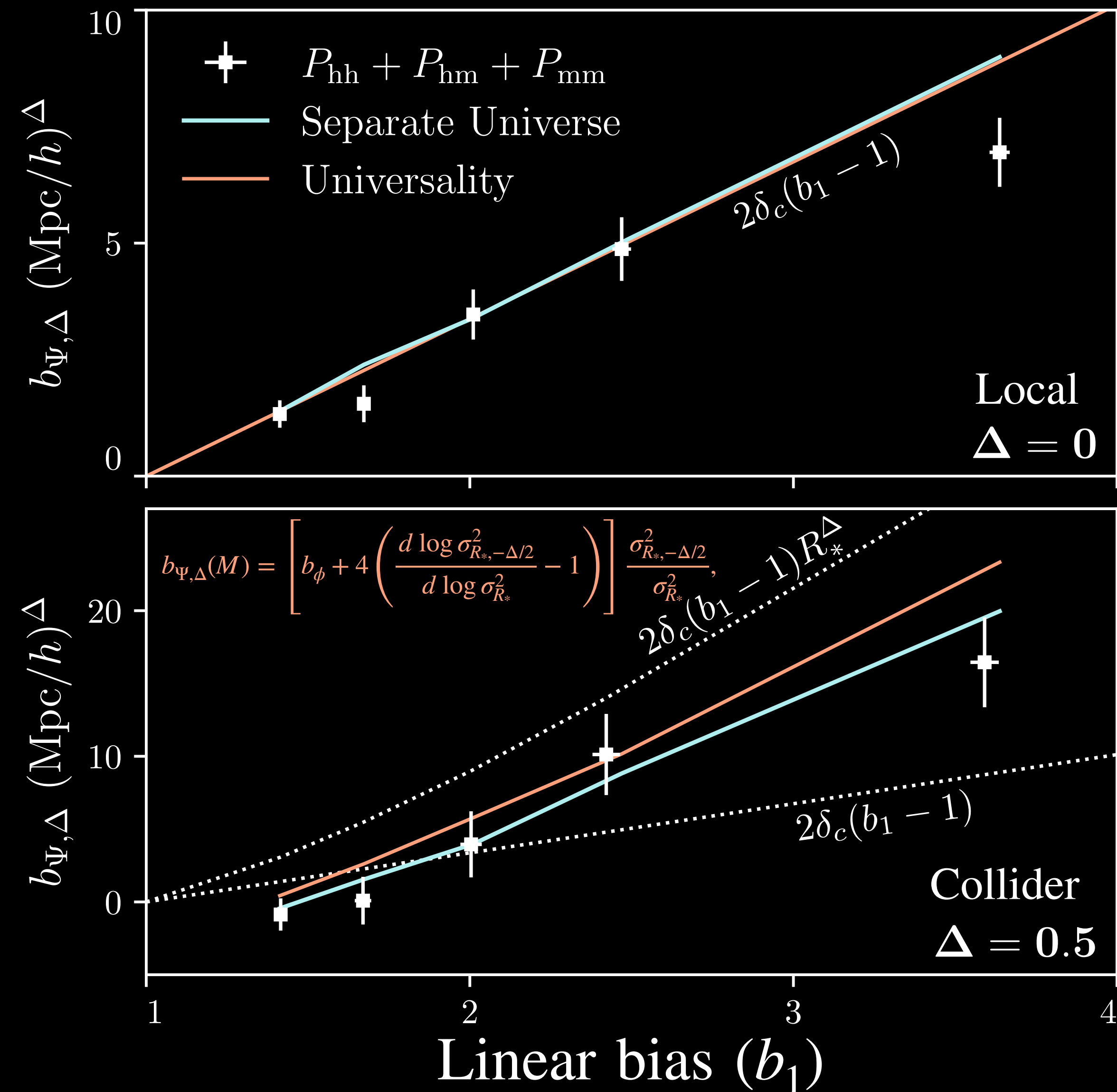
- Squeezed bispectrum leads to **scale-dependent bias** (*Dalal+07, Slosar+08, Desjacques+08*)

f_{NL}^Δ degenerate with non-Gaussian bias

$$P_{hm}(q) = \left[b_1 + \frac{3\Omega_{m0}H_0^2}{2D_{\text{md}}(z)} \frac{b_{\Psi,\Delta} f_{\text{NL}}^\Delta}{q^{2-\Delta}} \right] P_m(q)$$

Power depends on Δ

- Fit for $b_{\Psi,\Delta}$ at fixed f_{NL}^Δ and Δ
- Test predictions for galaxy biasing in non-local PNG *e.g., Shandera+2010, Schmidt & Kamionkowski, 2010, Schmidt+2012*
- **Galaxies will be more challenging...**



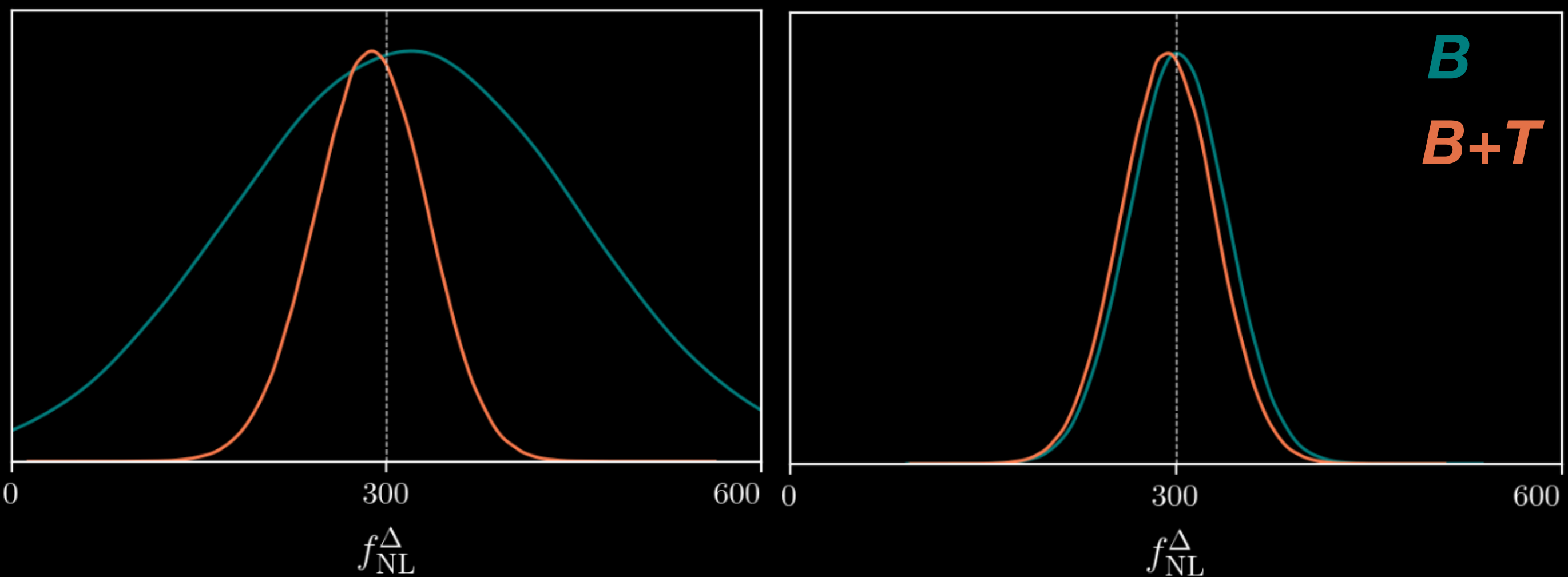
Fits to **halo** catalogues from sims agree with separate universe and universality.

Collapsed trispectrum

- What can we learn from the trispectrum?
- Higher-order PNG \longrightarrow
- Sample variance cancellation \downarrow

Unknown $P_{mm}(k)$

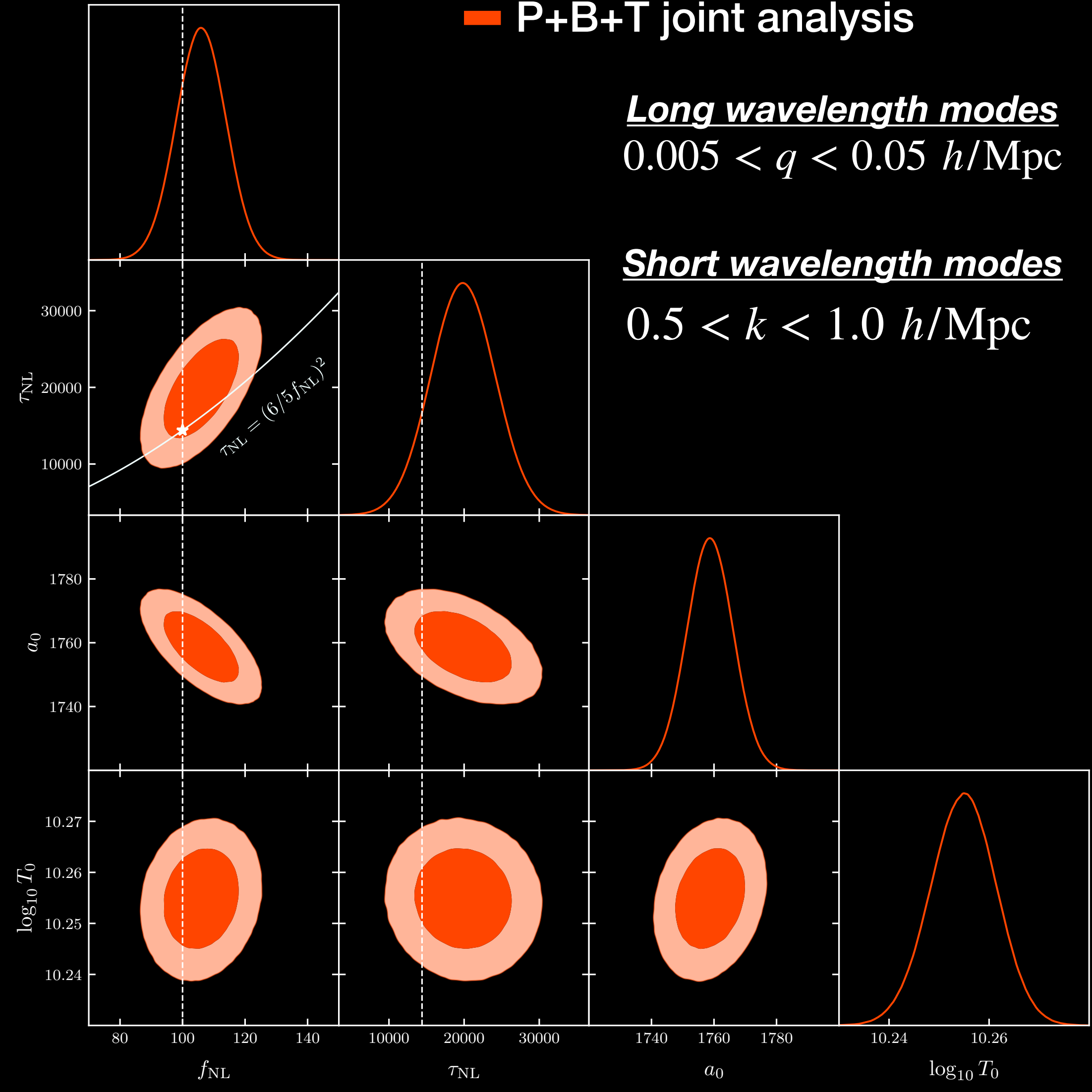
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— P+B+T joint analysis

Long wavelength modes
 $0.005 < q < 0.05 \text{ h/Mpc}$

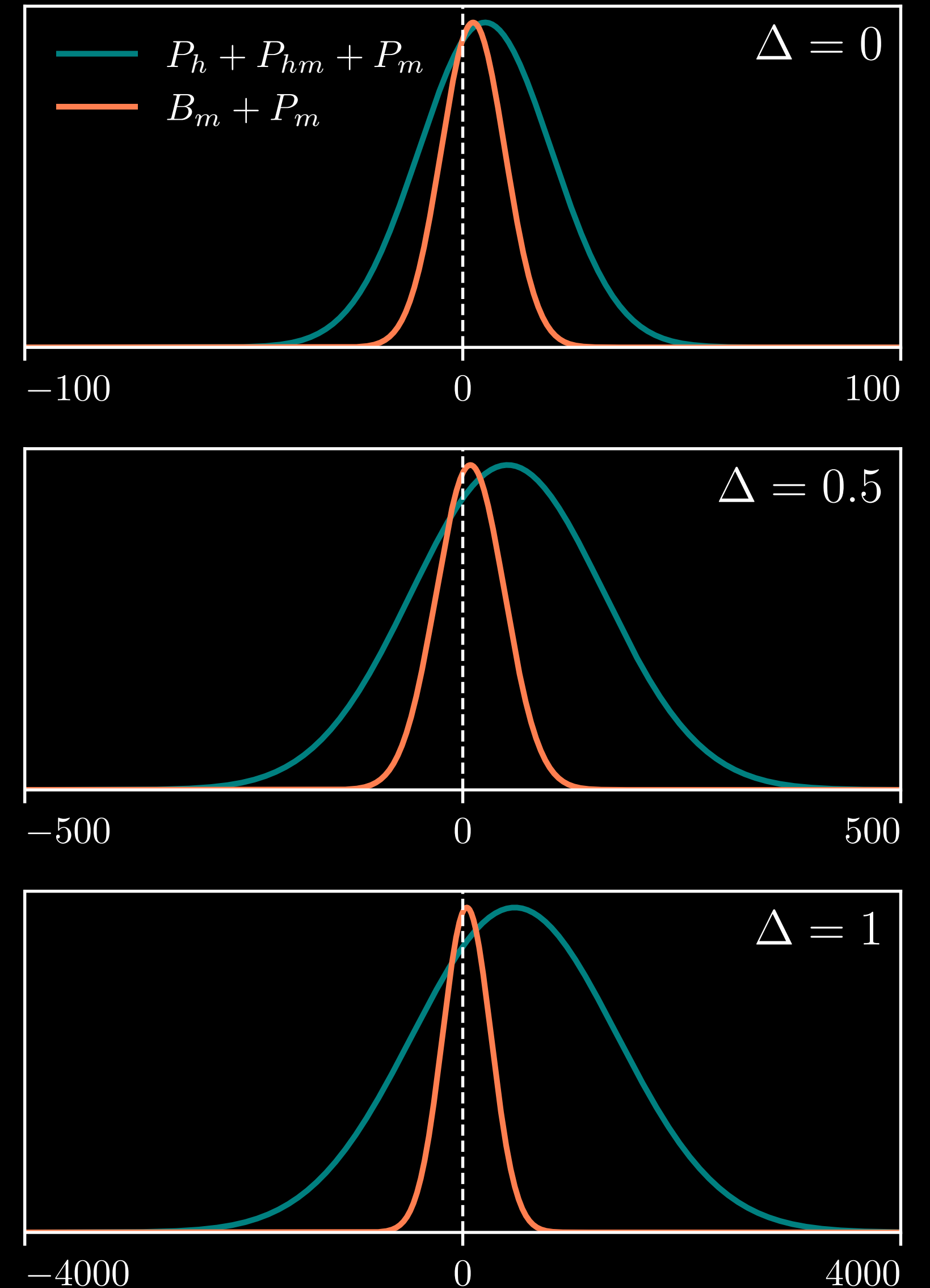
Short wavelength modes
 $0.5 < k < 1.0 \text{ h/Mpc}$



Constraints on (f_{NL}, τ_{NL}) from squeezed bispectrum and collapsed trispectrum using Quijote.

Conclusions

- Particles with $m \sim H$ during inflation leave distinct imprint on LSS correlators (“cosmo. collider”)
- Methods to run **N-body simulations with cosmo.collider** squeezed bispectrum (for $0 < m/H \leq 3/2$)
- **Non-perturbative** models for squeezed B_m and collapsed T_m
 - Validated up to $k_{\max} = 2 h/\text{Mpc}$ (useful for deep surveys?)
- Simulations can be used to test analysis methods based on **scale-dependent bias** (e.g., [Green, Guo, Han, Wallisch, 2023](#))
 - Test separate universe halo bias for non-local PNG
- More to do
 - E.g., **oscillatory bispectra**, B_g and T_g , **multi-tracer**/ $\mathbf{b}_{\Psi, \Delta}$, CMBxLSS, optimal estimators/field-level inference/alternative summary statistics, **can we constrain Δ ?**



*An unfair, unfinished, but not uninteresting comparison

Backup

Initial conditions for N-body simulations

- How to generate IC's with following **squeezed** bispectrum (🌲-level)

$$\lim_{k_1 \ll k_2 \approx k_3} B_{\Phi}(k_1, k_2, k_3) = 4f_{\text{NL}}^{\Delta} \left(\frac{k_1}{k_2} \right)^{\Delta} P_{\Phi}(k_1)P_{\Phi}(k_2),$$

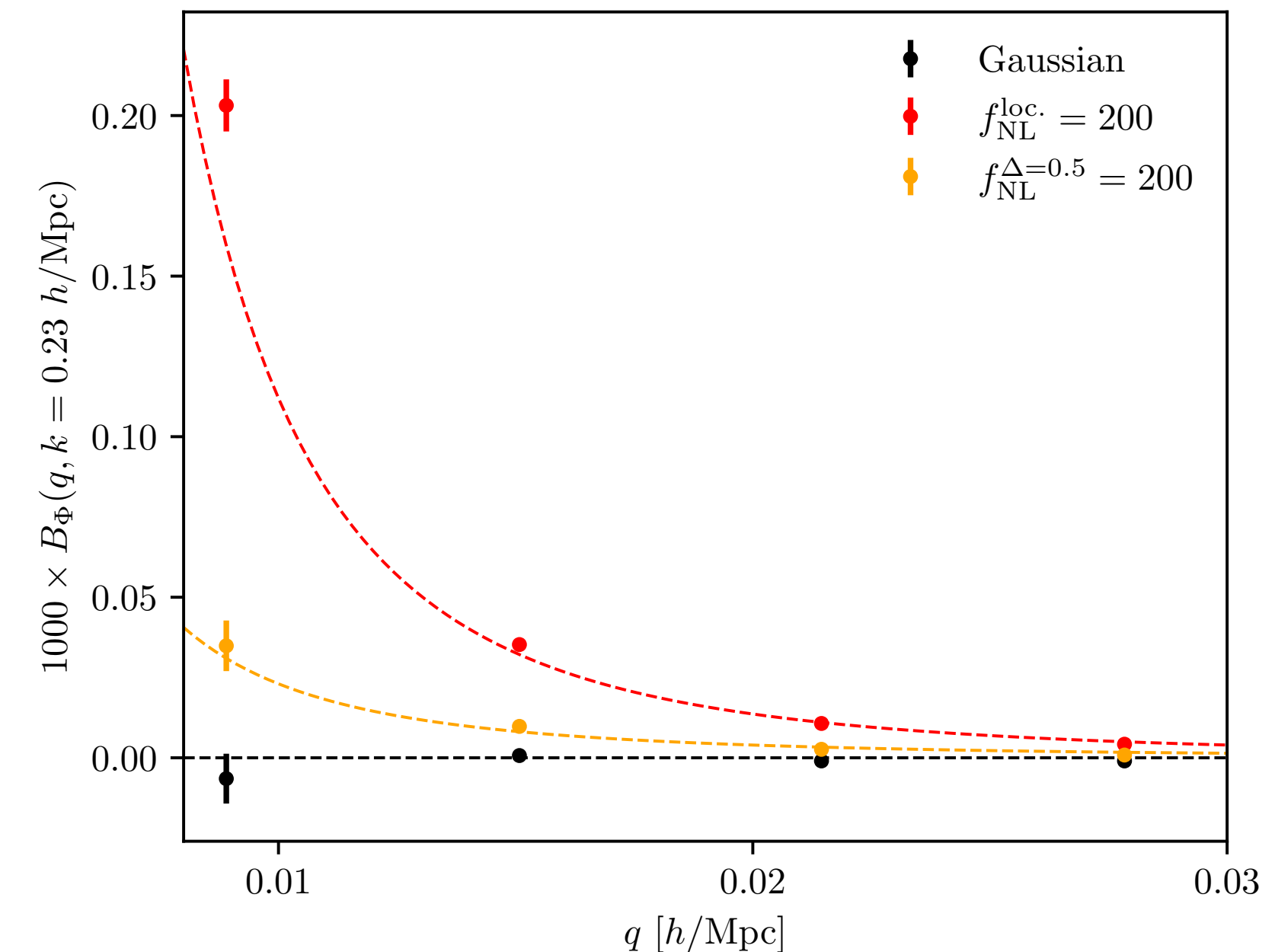
- Construct Φ quadratically from ϕ_G (e.g., [Scoccimarro, Hui, Manera, 2012](#))

$$\Phi(\mathbf{k}) = \phi_G(\mathbf{k}) + f_{\text{NL}}^{\Delta} [\Psi_{\Delta}(\mathbf{k}) - \langle \Psi_{\Delta}(\mathbf{k}) \rangle],$$

$$\Psi_{\Delta}(\mathbf{k}) = \int_{\mathbf{k}_1, \mathbf{k}_2} (2\pi)^3 \delta_D(\mathbf{k} - \mathbf{k}_{12}) K_{\Delta}(\mathbf{k}_1, \mathbf{k}_2) \phi_G(\mathbf{k}_1) \phi_G(\mathbf{k}_2),$$

- Choose $K_{\Delta}(\mathbf{k}_1, \mathbf{k}_2)$ that gives target bispectrum and doesn't "break" P_{Φ} with loop corrections

- No fun if I write it out for you!
- Convolution can be computed with FFTs!



Check IC's with 🌲-level bispectrum

Bispectrum Estimator Validation

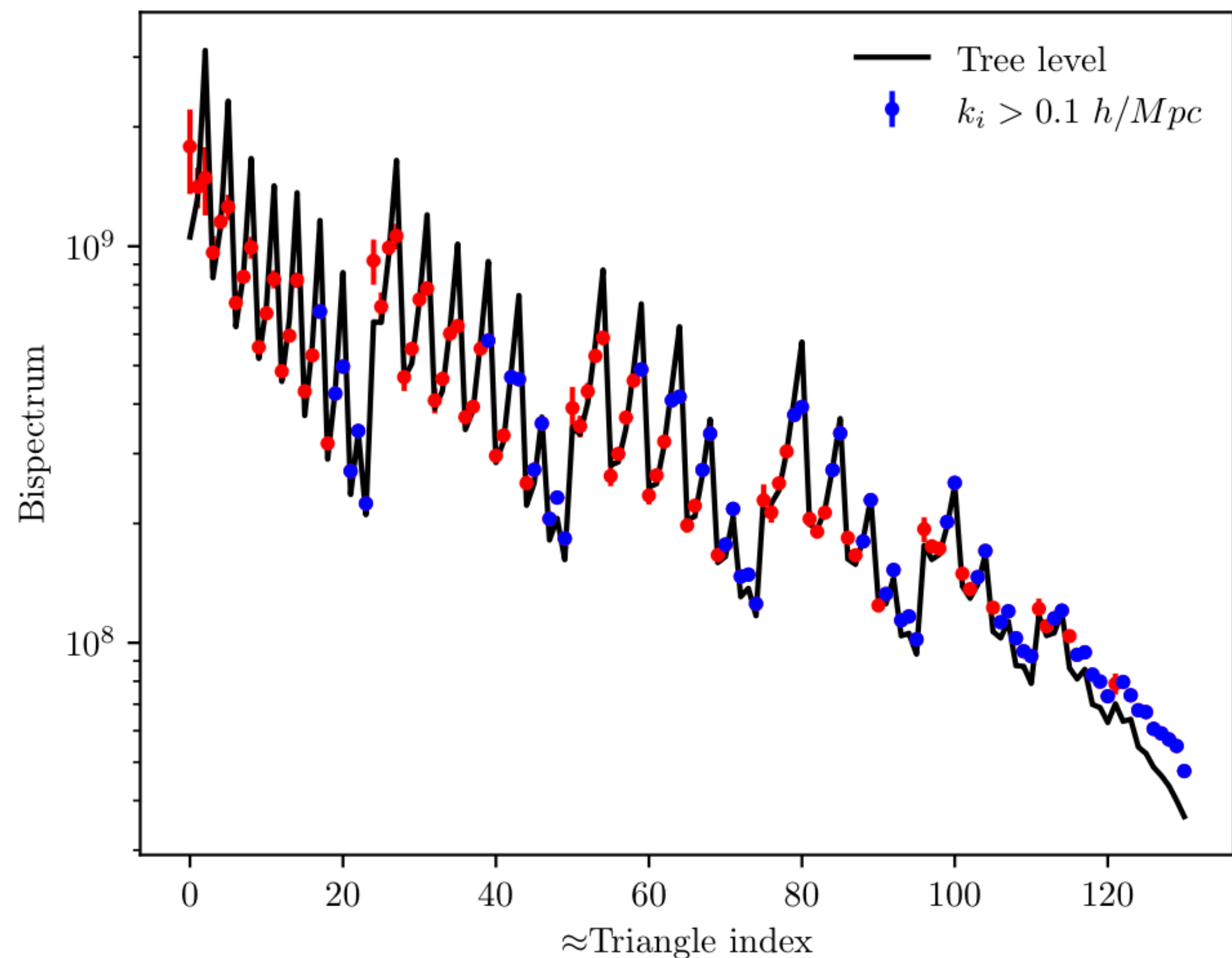


Figure 1: Validation on measurements from Quijote

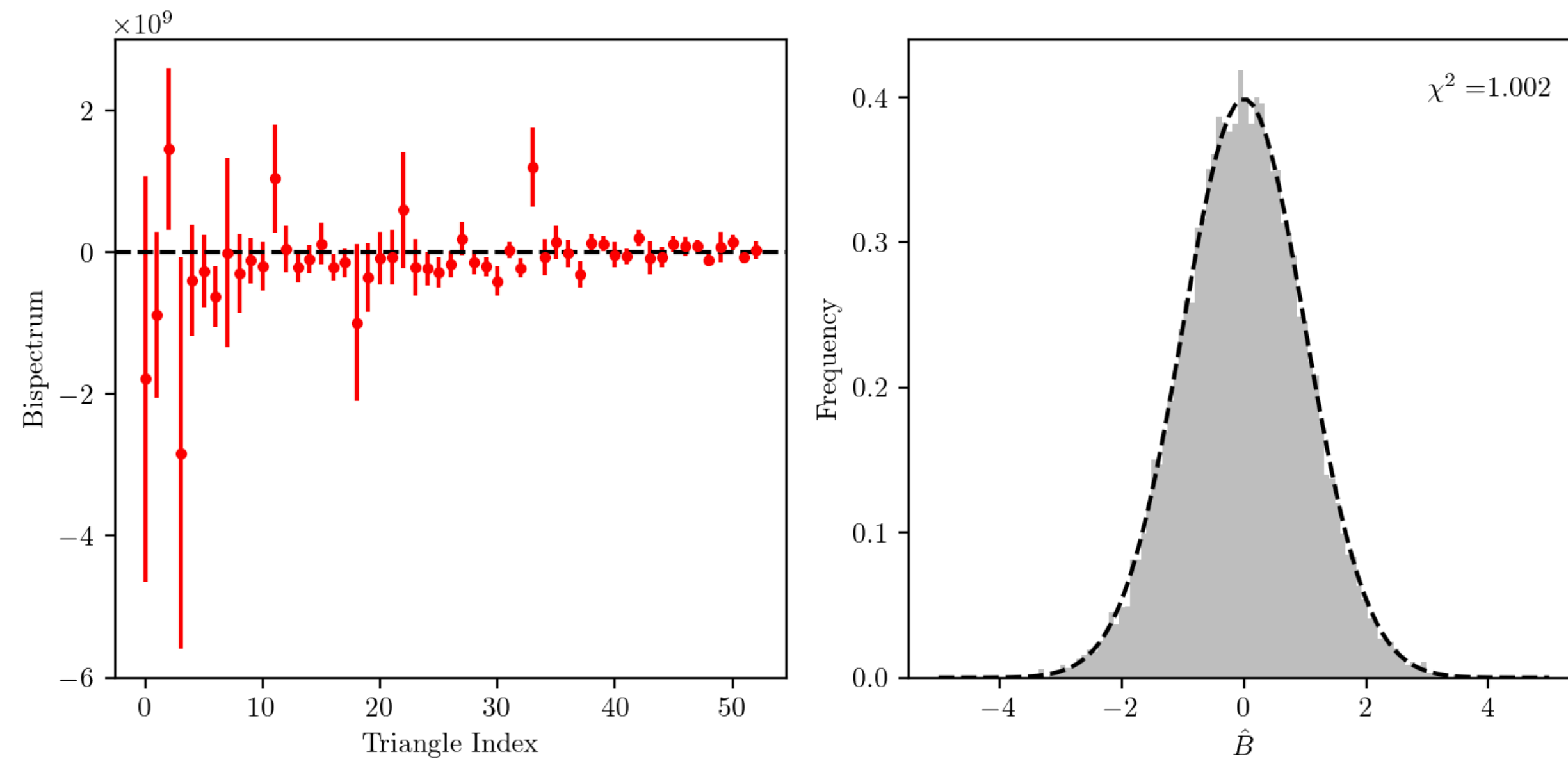


Figure 2: Validation on Gaussian mocks

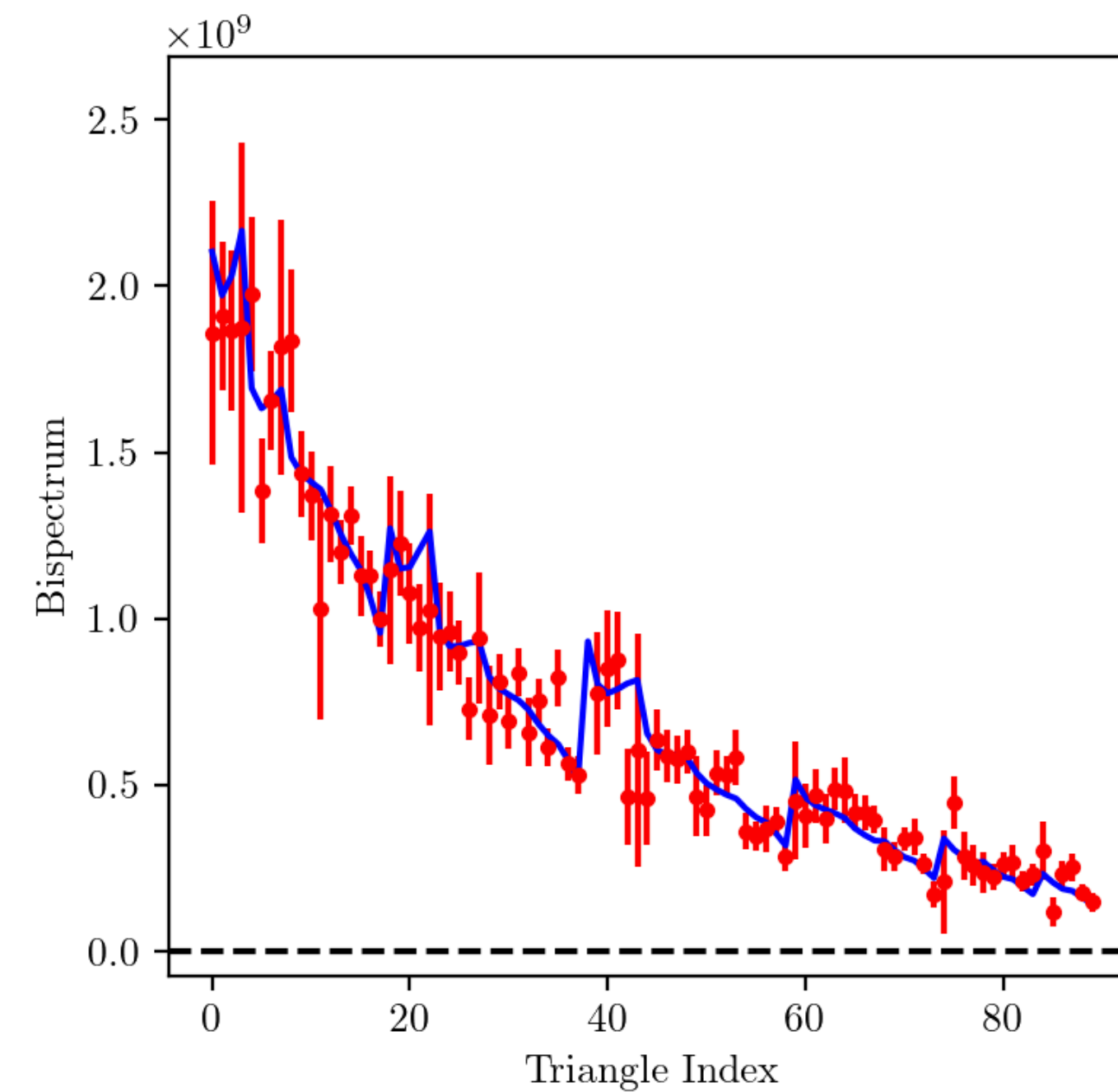


Figure 3: Validation on mocks with f_{NL}^{loc}

Trispectrum Estimator

- Trispectrum estimator in terms of external legs k_a, \dots, k_d and diagonals k_{ab}, k_{bc} is **not separable**
- Use estimator integrated over q_{23} (see Appendix A of [2306.11782](#))

$$\left\langle \hat{T}_{\text{tot}}(k_a, k_b, k_c, k_d, k_E) \right\rangle = \frac{1}{N_{a,b,c,d,E}} \int_{\mathbf{Q}} W_E(\mathbf{Q}) \int_{\mathbf{k}_1, \dots, \mathbf{k}_4} \left\{ W_a(k_1) W_b(k_2) W_c(k_3) W_d(k_4) \left\langle \delta_{\mathbf{k}_1} \delta_{\mathbf{k}_2} \delta_{\mathbf{k}_3} \delta_{\mathbf{k}_4} \right\rangle \right. \\ \left. (2\pi)^3 \delta_D^{(3)}(\mathbf{k}_{12} - \mathbf{Q}) (2\pi)^3 \delta_D^{(3)}(\mathbf{k}_{34} + \mathbf{Q}) \right\},$$

- Can be efficiently evaluated by computing FFT's of product fields

$$D_{ij}(\mathbf{Q}) \equiv \int d^3x e^{-i\mathbf{Q}\cdot\mathbf{x}} \delta_{W_i}(\mathbf{x}) \delta_{W_j}(\mathbf{x}).$$

$$\hat{T}_{\text{tot}}(k_a, k_b, k_c, k_d, k_E) = \frac{1}{N_{a,b,c,d,E}} \int_{\mathbf{Q}} W_E(\mathbf{Q}) D_{ab}(\mathbf{Q}) D_{cd}^*(\mathbf{Q}),$$

Biased estimator because it includes disconnected terms!!!

Trispectrum Estimator (continued)

- Estimate and subtract disconnected terms
 - Simple estimator: $\delta^4 - 3\langle\delta^2\rangle^2$
 - More optimal estimator: $\delta^4 - 6\delta^2\langle\delta^2\rangle^2 + 3\langle\delta^2\rangle^2$ (See Appendix F. of [Shen, Schaan, Ferraro, 24](#))
- Can be efficiently implemented using FFT's

Depends on fiducial $P(\mathbf{k})$

$$F_{ij}^P(\mathbf{x}) \equiv \int_{\mathbf{k}} W_i(k)W_j(k)P(k)e^{i\mathbf{k}\cdot\mathbf{x}}$$

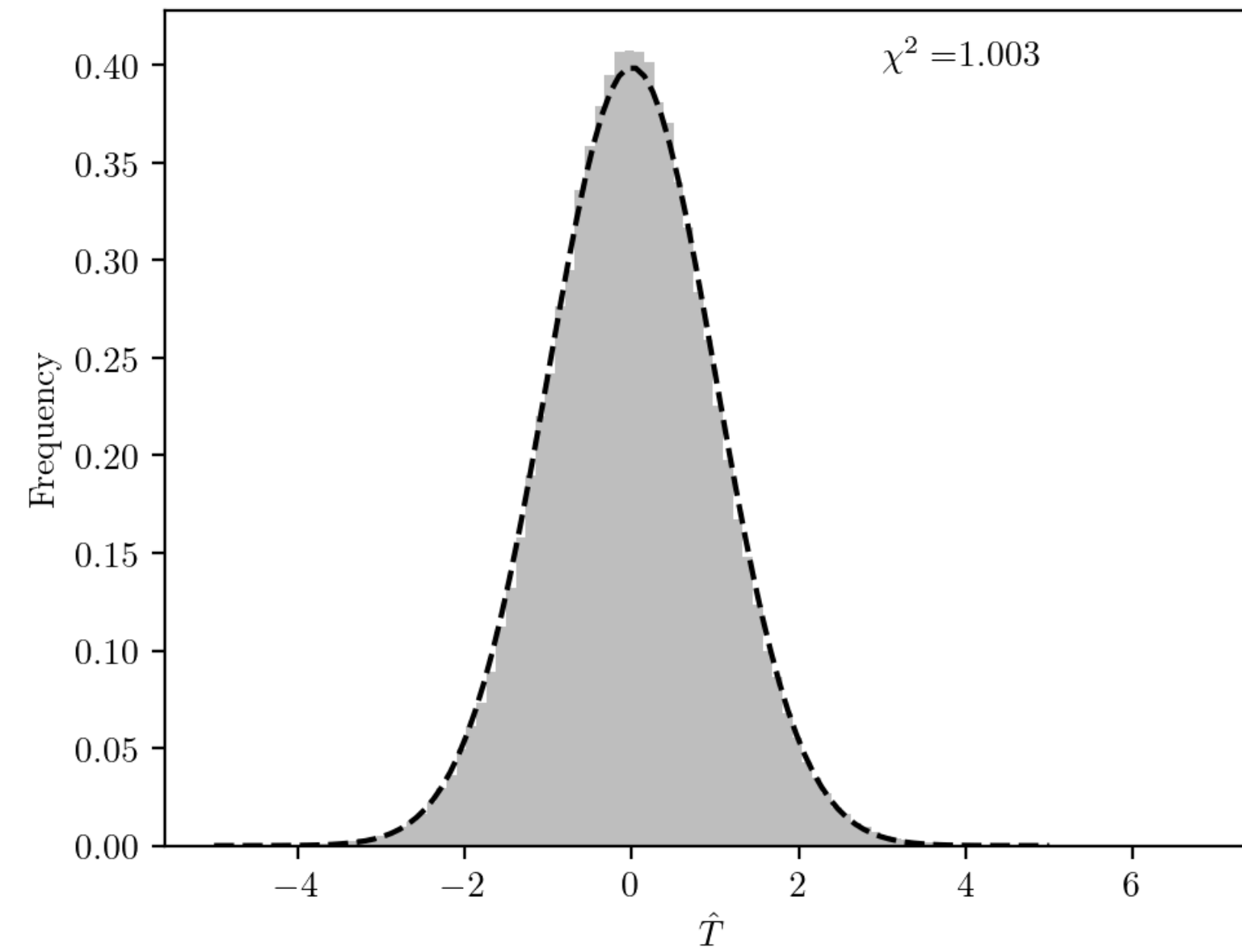
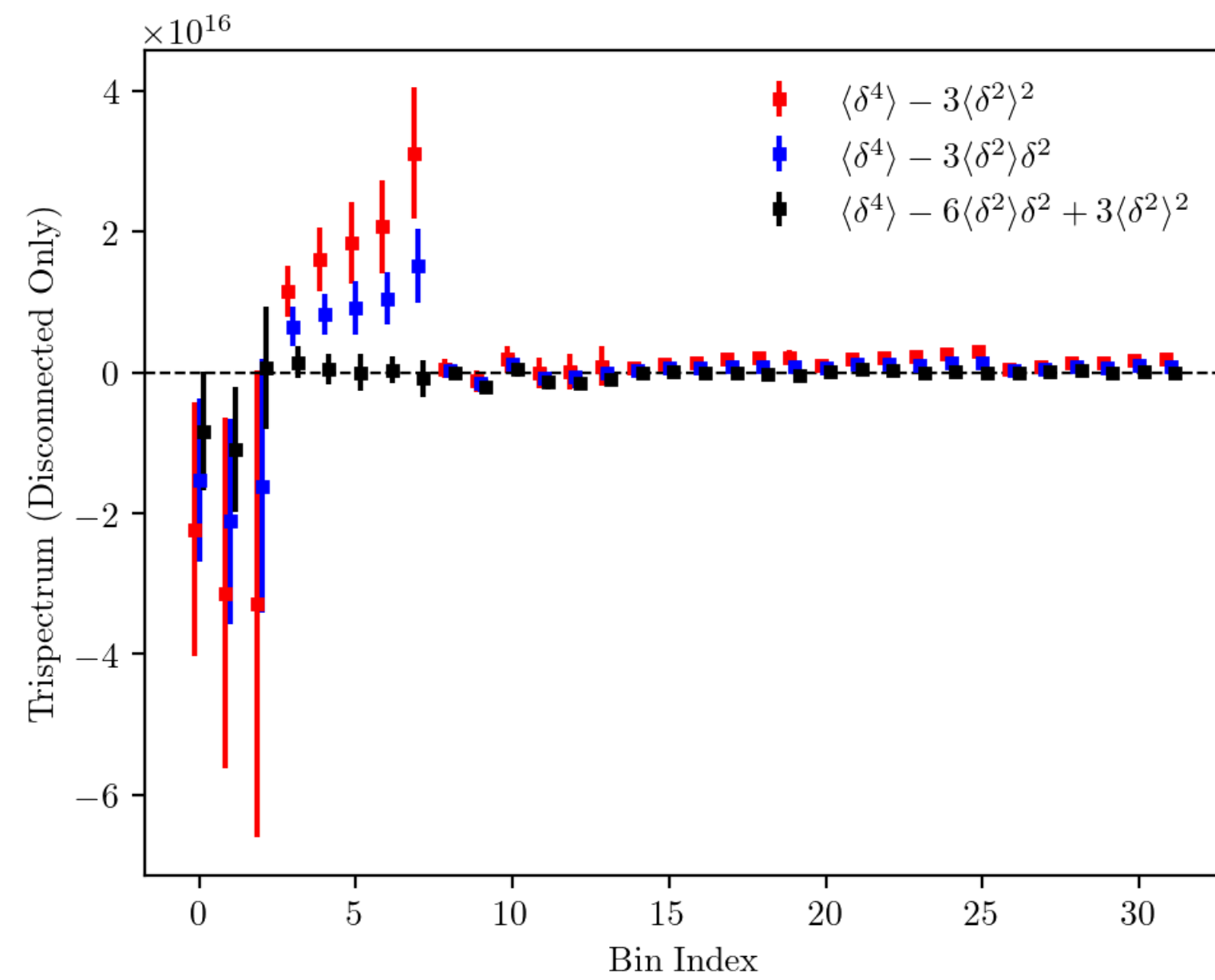
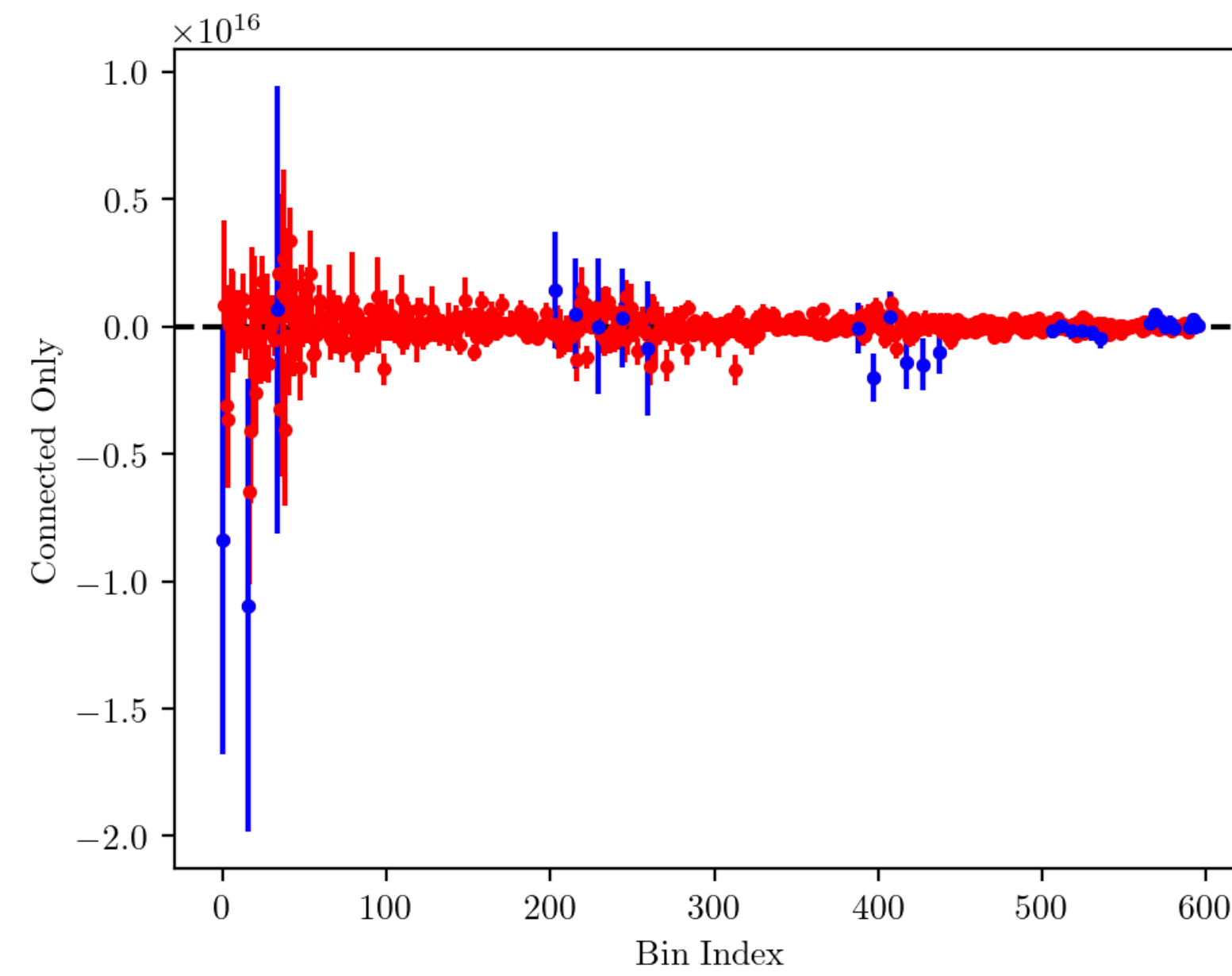
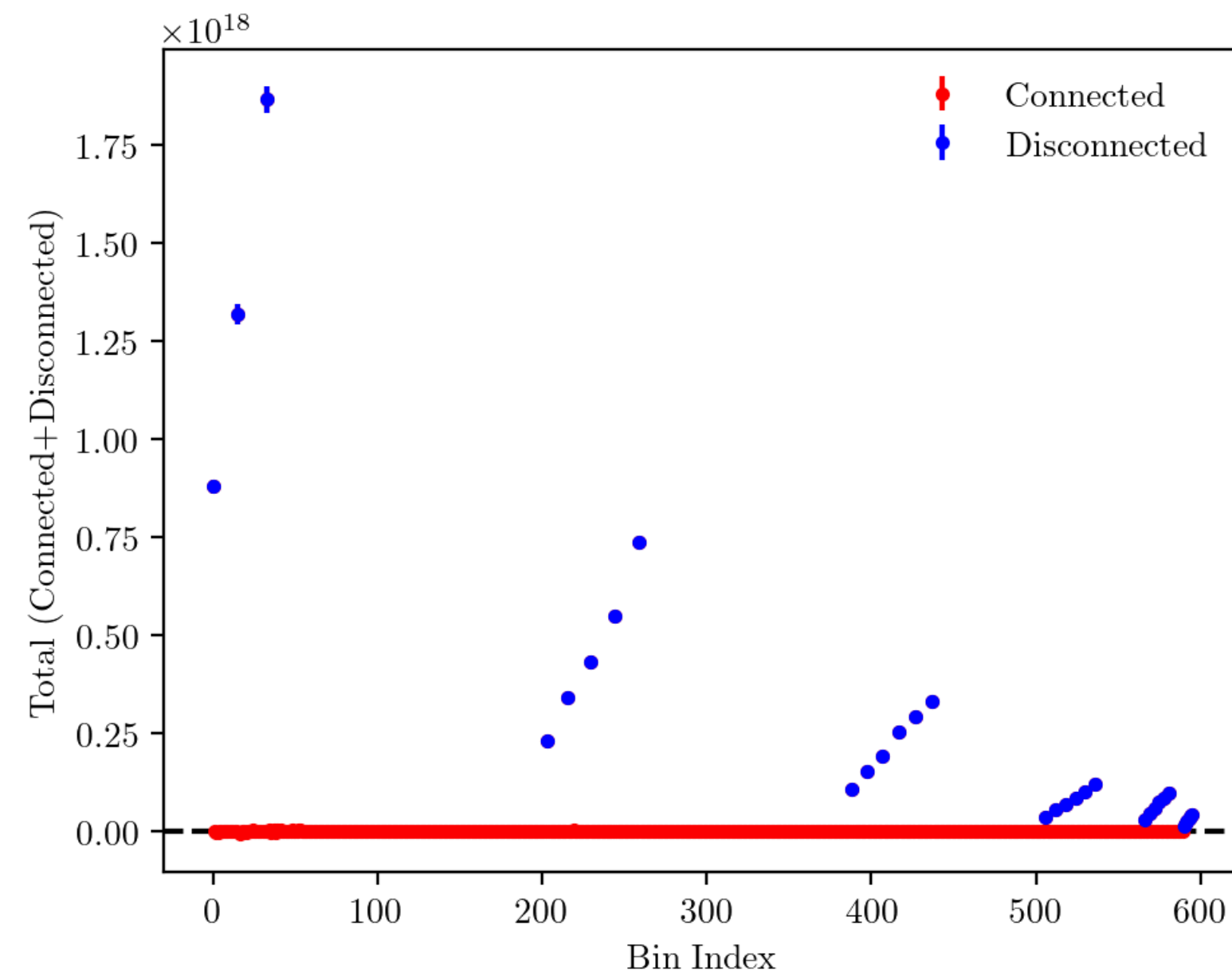
Depends on realization of δ

$$F_{ij}^\delta(\mathbf{x}) \equiv \int_{\mathbf{k}} W_i(k)W_j(k)|\delta_{\mathbf{k}}|^2 e^{i\mathbf{k}\cdot\mathbf{x}}$$

$$\hat{T}_{\text{disc}}^{\langle\delta^2\rangle^2} = \frac{V}{N_{a,b,c,d,E}} \int_{\mathbf{Q}} W_E(Q) \int d^3x e^{-i\mathbf{Q}\cdot\mathbf{x}} [F_{ac}^P(\mathbf{x})F_{bd}^P(\mathbf{x}) + F_{ad}^P(\mathbf{x})F_{bc}^P(\mathbf{x})]$$

$$\hat{T}_{\text{disc}}^{\delta^2\langle\delta^2\rangle} = \frac{V}{N_{a,b,c,d,E}} \int_{\mathbf{Q}} W_E(Q) \int d^3x e^{-i\mathbf{Q}\cdot\mathbf{x}} [F_{ac}^P(\mathbf{x})F_{bd}^\delta(\mathbf{x}) + F_{ac}^\delta(\mathbf{x})F_{bd}^P(\mathbf{x}) + F_{ad}^P(\mathbf{x})F_{bc}^\delta(\mathbf{x}) + F_{ad}^\delta(\mathbf{x})F_{bc}^P(\mathbf{x})].$$

Trispectrum Estimator Validation



Potential derivatives

- Compute logarithmic derivative from separate universe sims

$$P_m^{\text{lin.}}(k | \epsilon, \Delta) \equiv (1 + 2\epsilon k^{-\Delta}) P_m^{\text{lin.}}(k).$$

$$\frac{\partial P_m(k)}{\partial \Phi_L(\mathbf{q})} = 2f_{\text{NL}}^\Delta q^\Delta \frac{\partial P_m(k | \epsilon, \Delta)}{\partial \epsilon} \Big|_{\epsilon=0}.$$

Get from finite differencing SU sims with $\pm\epsilon$

- Strong dependence on Δ
- Same sims used to estimate non-Gaussian bias

$$b_{\Psi, \Delta}(M, z) = \frac{\partial \log \bar{n}_h(M, z)}{\partial \epsilon} \Big|_{\epsilon=0}$$

