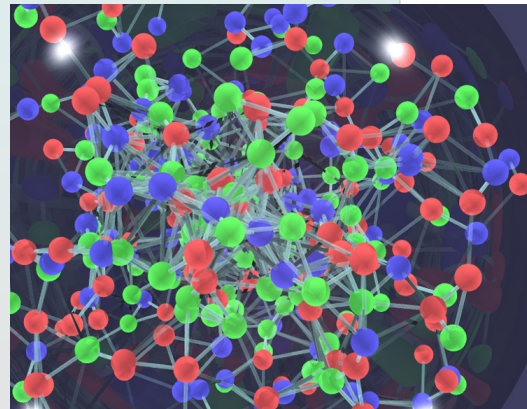
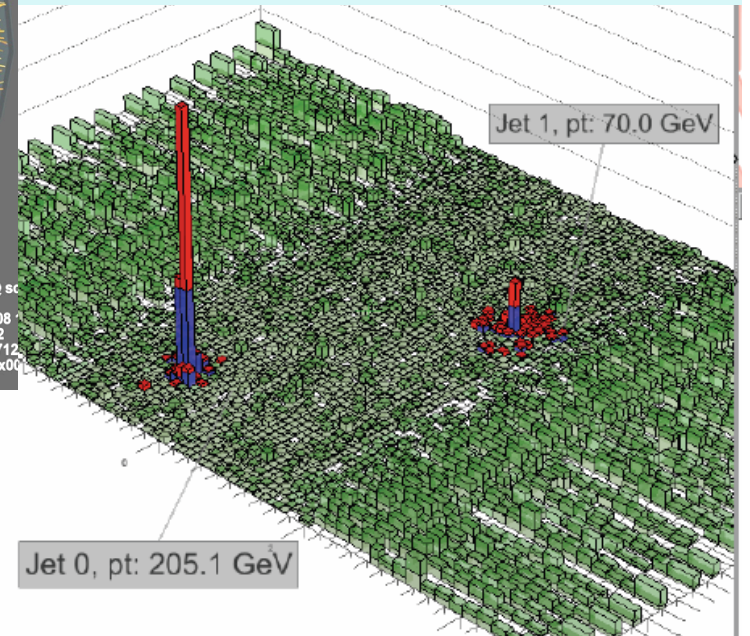
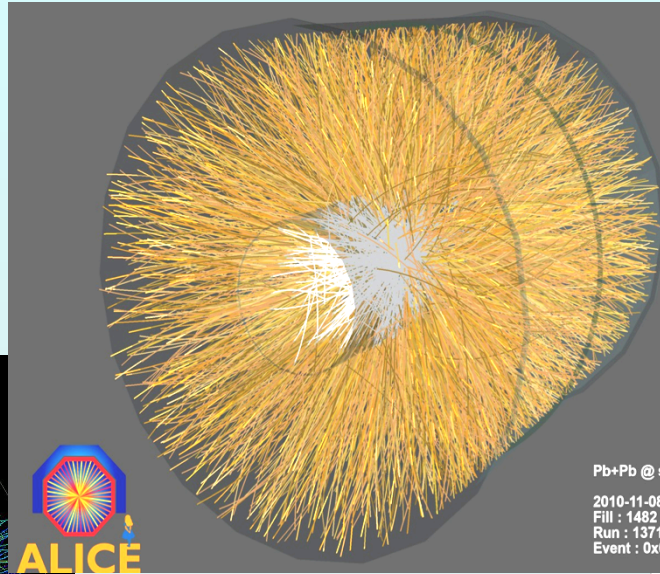
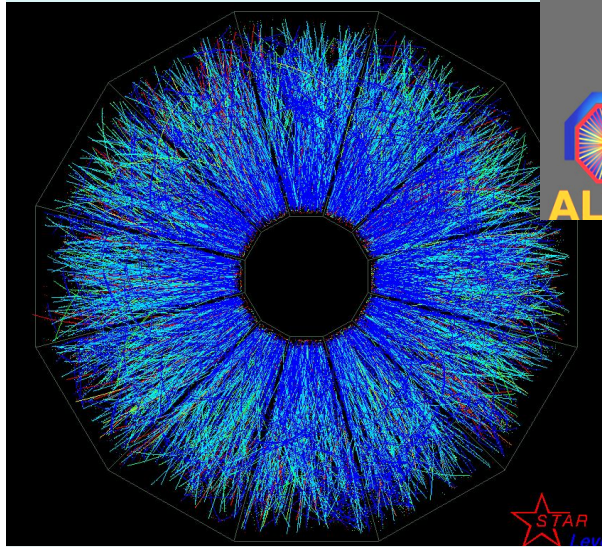


Hot QCD

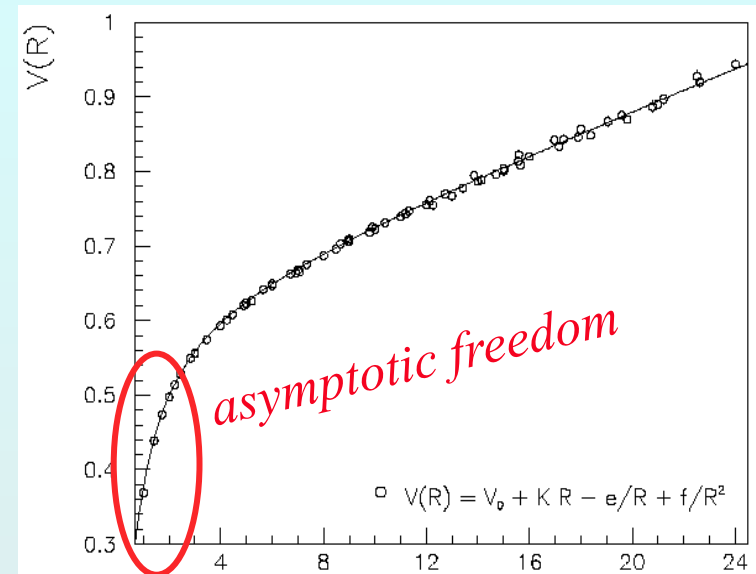
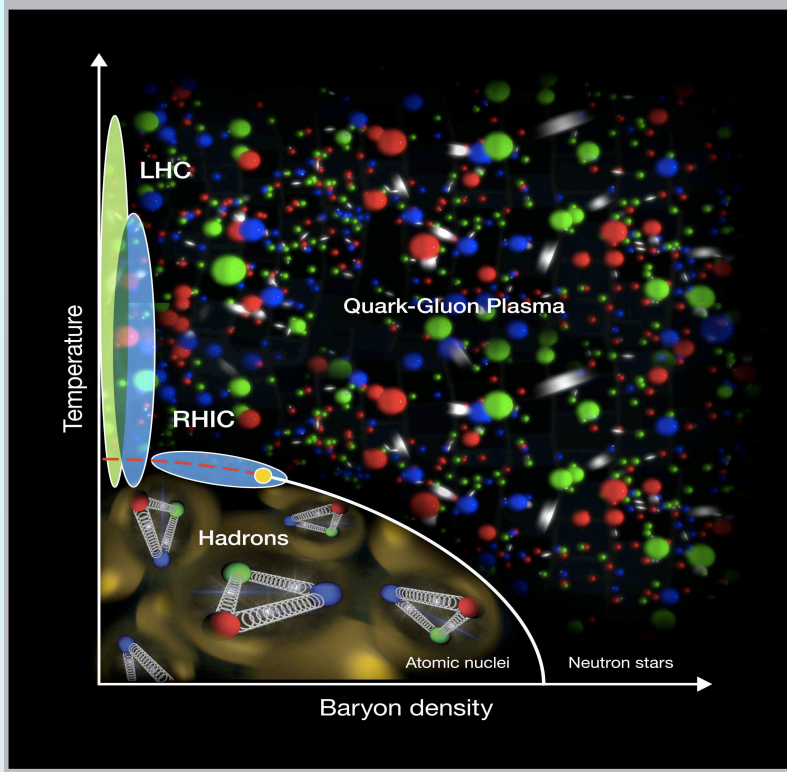


Barbara Jacak
UC Berkeley
Feb. 17, 2016

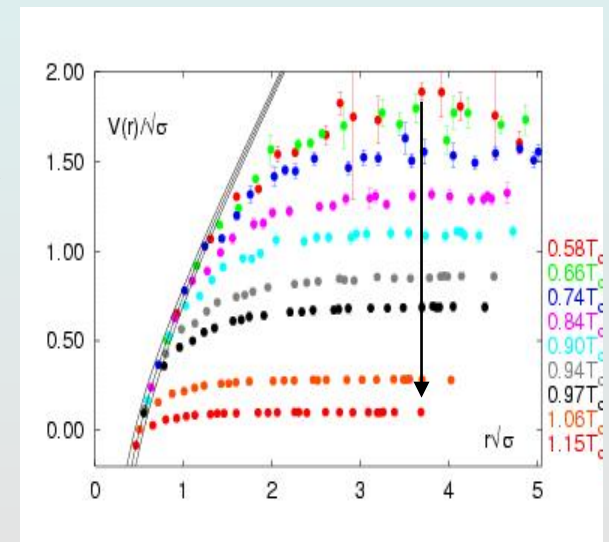
outline

- **Why hot QCD**
- **How – and what - we measure**
- **Controlling the collision**
- **Collective behavior in QCD**
- **Parton interactions with quark gluon plasma**
- **Color screening in quark gluon plasm**

Hot, dense QCD matter



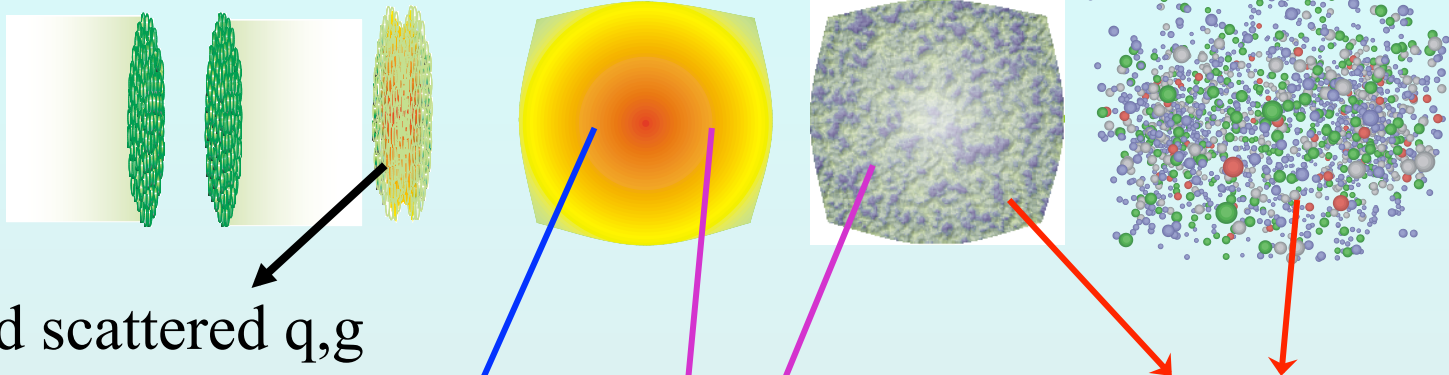
At high temperature/density
screening by produced colored particles
Expect phase transition to
deconfined quark gluon plasma
Lattice QCD $\rightarrow T_c \sim 150$ MeV



heavy ion collision & diagnostics

time \longrightarrow

ϵ , pressure builds up



Hard scattered q, g
(short wavelength)
probes of plasma
formed

Hot QGP

thermal radiation
($\gamma, \gamma^* \rightarrow e^+e^-, \mu^+\mu^-$)

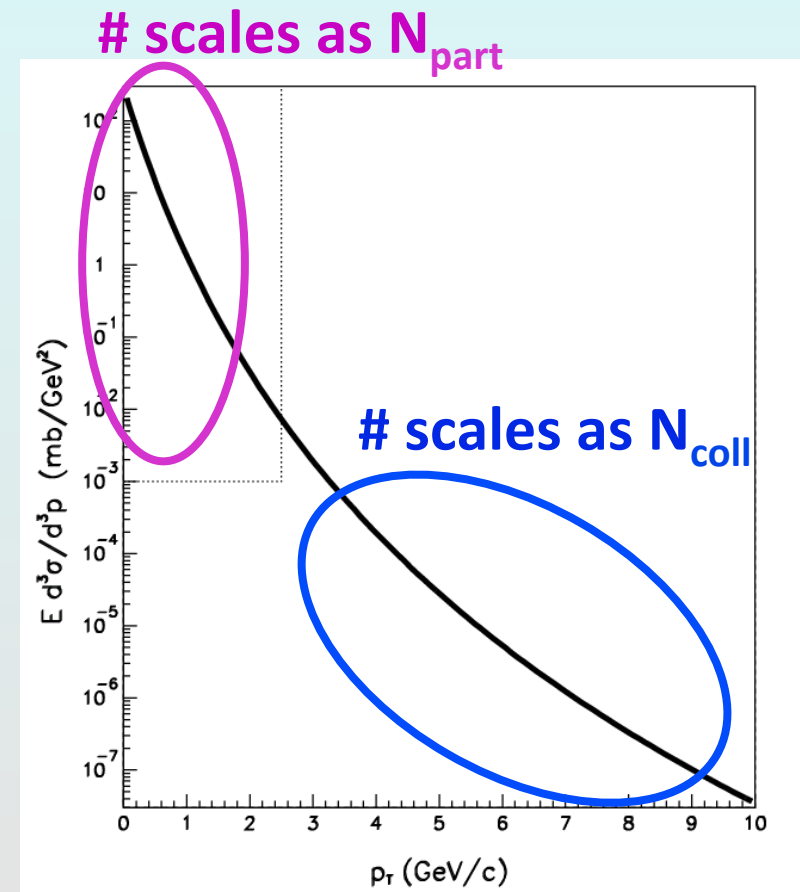
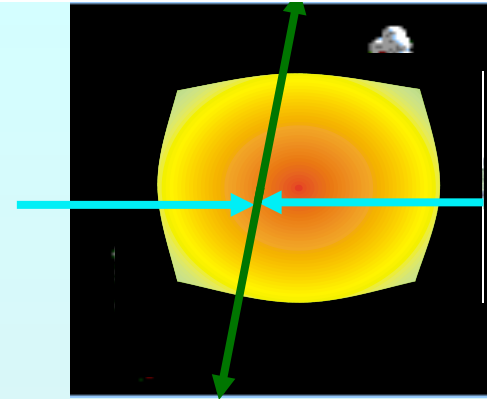
$\pi, K, p, n, \phi, \Lambda, \Delta, \Xi, \Omega, d,$

Hadrons qqq baryons & $q\bar{q}$ mesons
reflect (thermal) properties when
inelastic collisions stop

*not possible to measure as a function of time
nature integrates over the entire collision history*

study plasma with radiated & “probe” particles

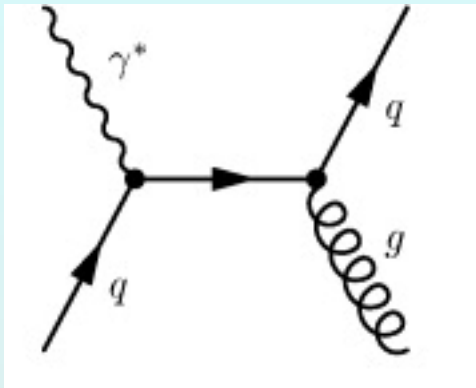
- as a function of transverse momentum
90° is where the action is (max T, ρ)
 p_L between the two beams: midrapidity
- $p_T < 1.5$ GeV/c
“thermal” particles π, K, p , etc.
radiated from bulk medium
“internal” plasma probes
- $p_T > 3$ GeV/c
“large E_{tot} ” (p_T or $M \gg T$)
set scale other than T(plasma)
autogenerated “external” probe
describe by perturbative QCD
- control probe: photons
EM, not strong interaction
produced in Au+Au by QCD
Compton scattering




Some cool observables

- Compare p+A to p+p for gluon distribution in A

Use QCD Compton scattering

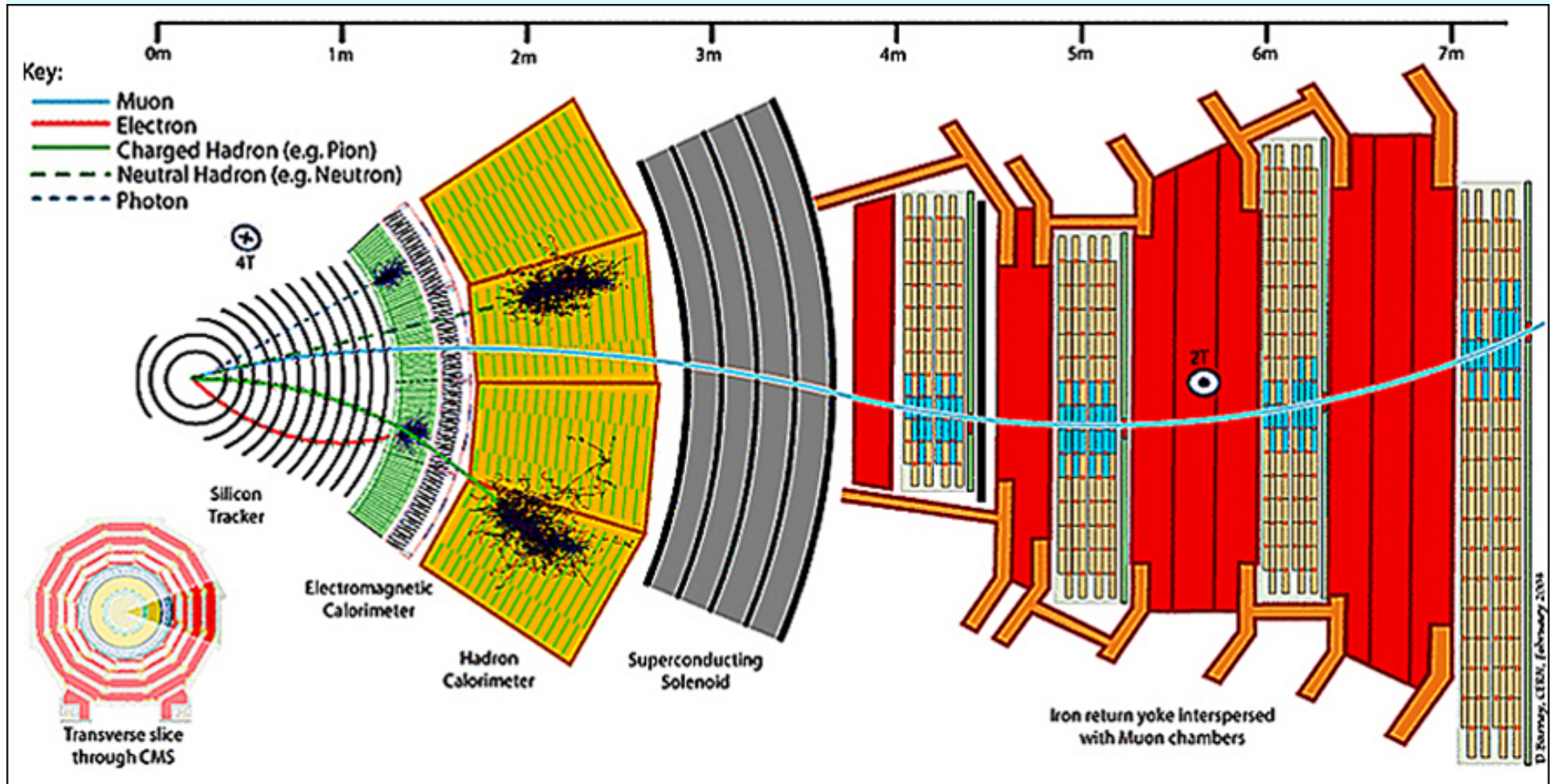


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What to measure to study thermalization (pre-equilib)?

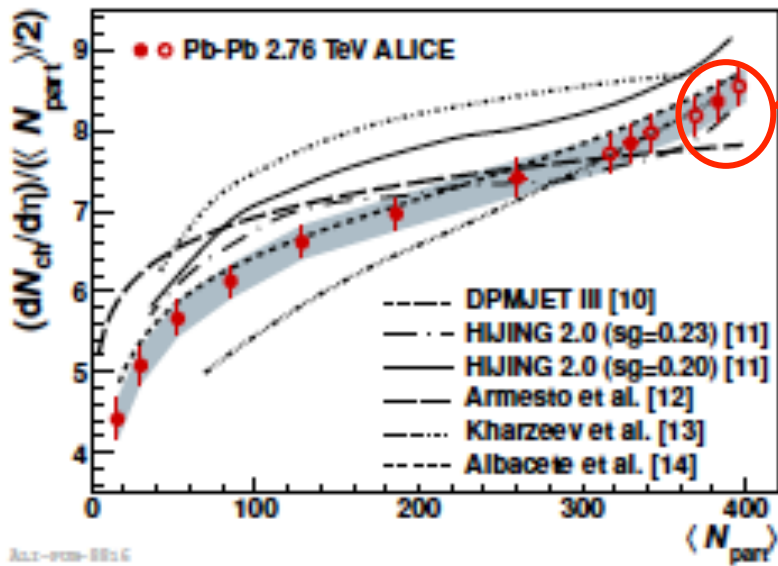
- Penetrating probes
 - γ or e^+e^- emitted during thermalization?
- Lasting probes (which may survive later stages)
 - Correlations and fluctuations among groups of particles
 - Pre-equilibrium flow?

How to measure all this stuff?



+ dE/dx & Time-of-flight: hadron ID + Cherenkov & TRD: electron ID
High granularity to handle high multiplicity

Handling high multiplicities



1700 particles per unit η

Challenge for Si: # of space points

Challenge for TPC: space charge

Trick: reconstruct collision vertex & use it in track reconstruction

Luminosity = frequency * n_1 * n_2 / $(4\pi\sigma_x\sigma_y)$

higher Z -> lower n

@RHIC $\langle \mathcal{L} \rangle$ in Au+Au / pp $\sim 1/600$

Just how tough is this?

For p+p collisions, value is about 4 particles per unit η

But, luminosity is much higher!

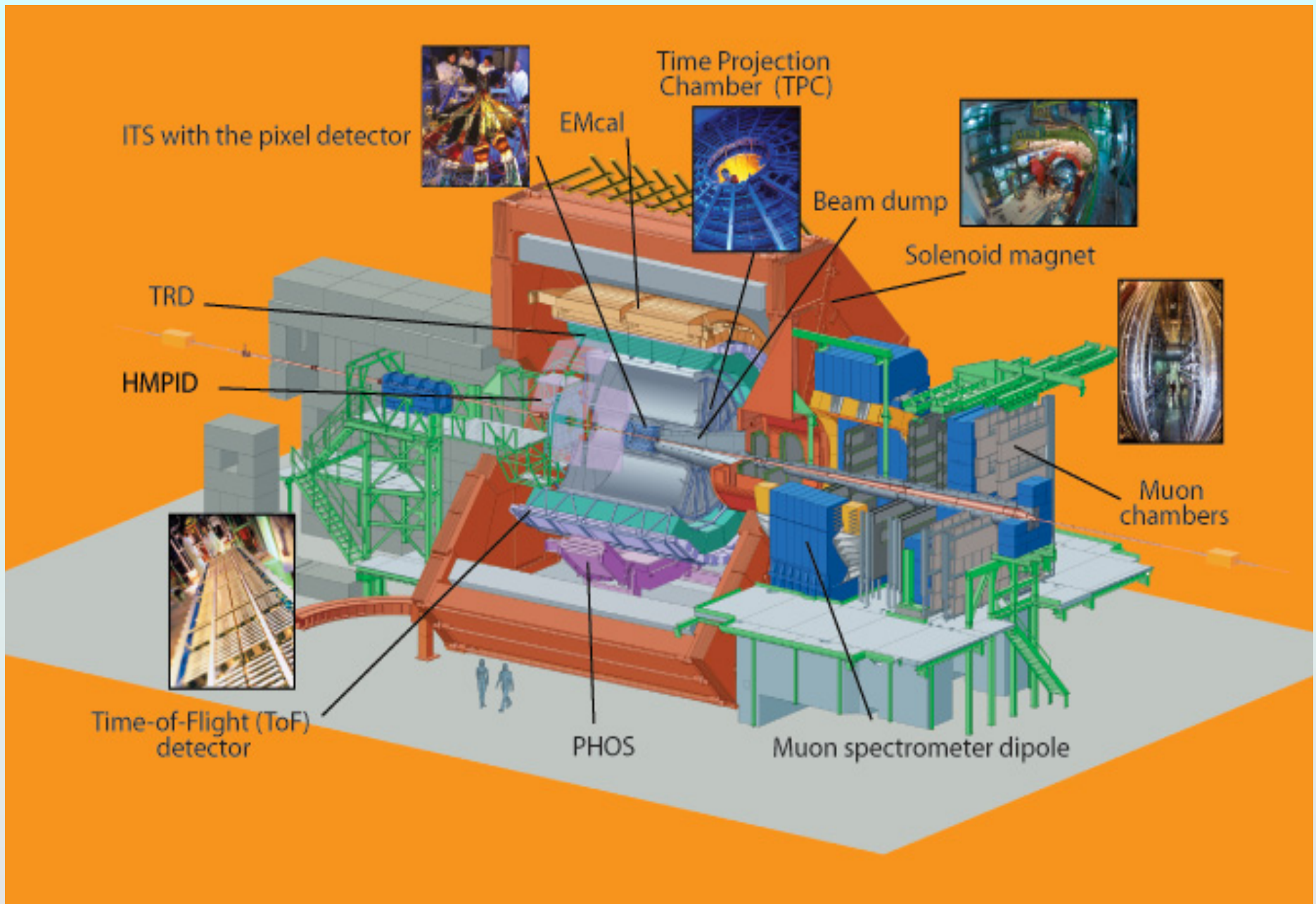
Pileup in high luminosity p+p = 150 events/crossing ($dN/d\eta \sim 600$)

but these are from different vertices!!

Pile-up rate in ATLAS in 2015?

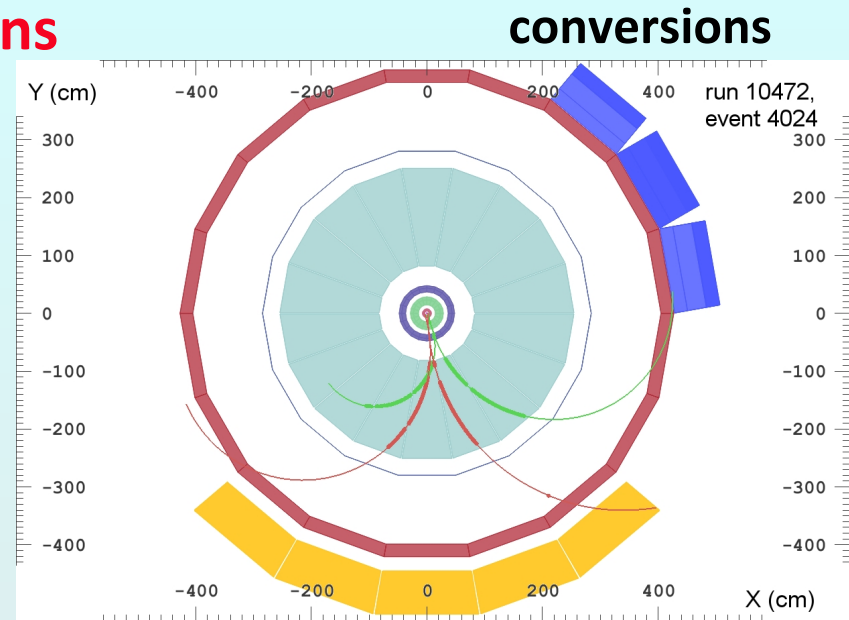
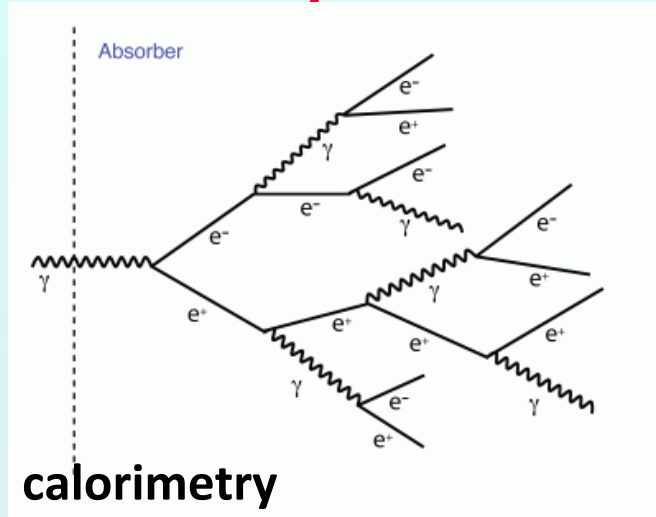
ATLAS handles a similar tracking problem (it can do HI collisions!)

PID in ALICE

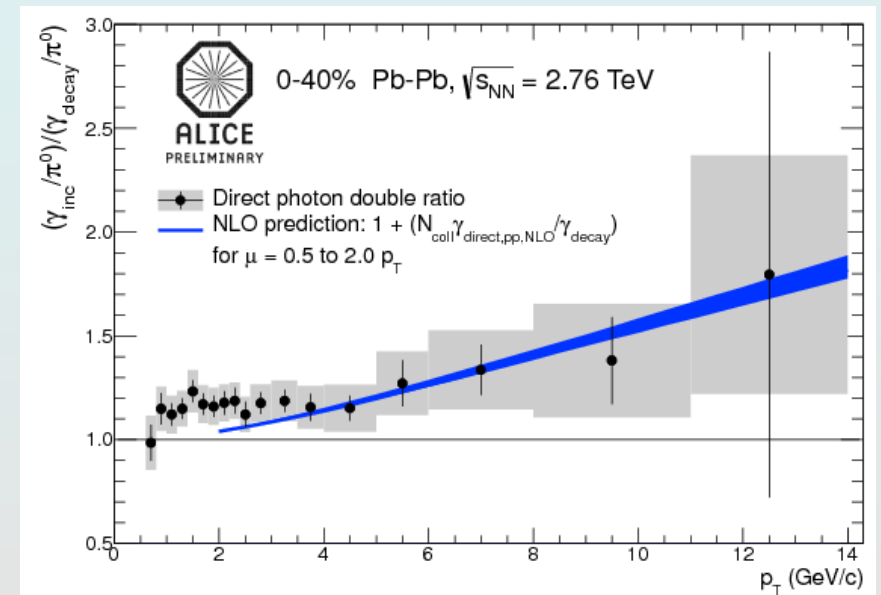
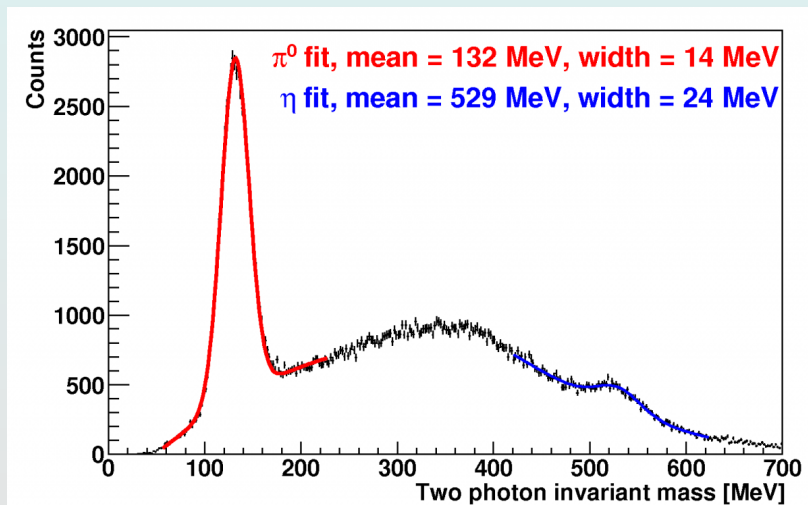


Measuring photons

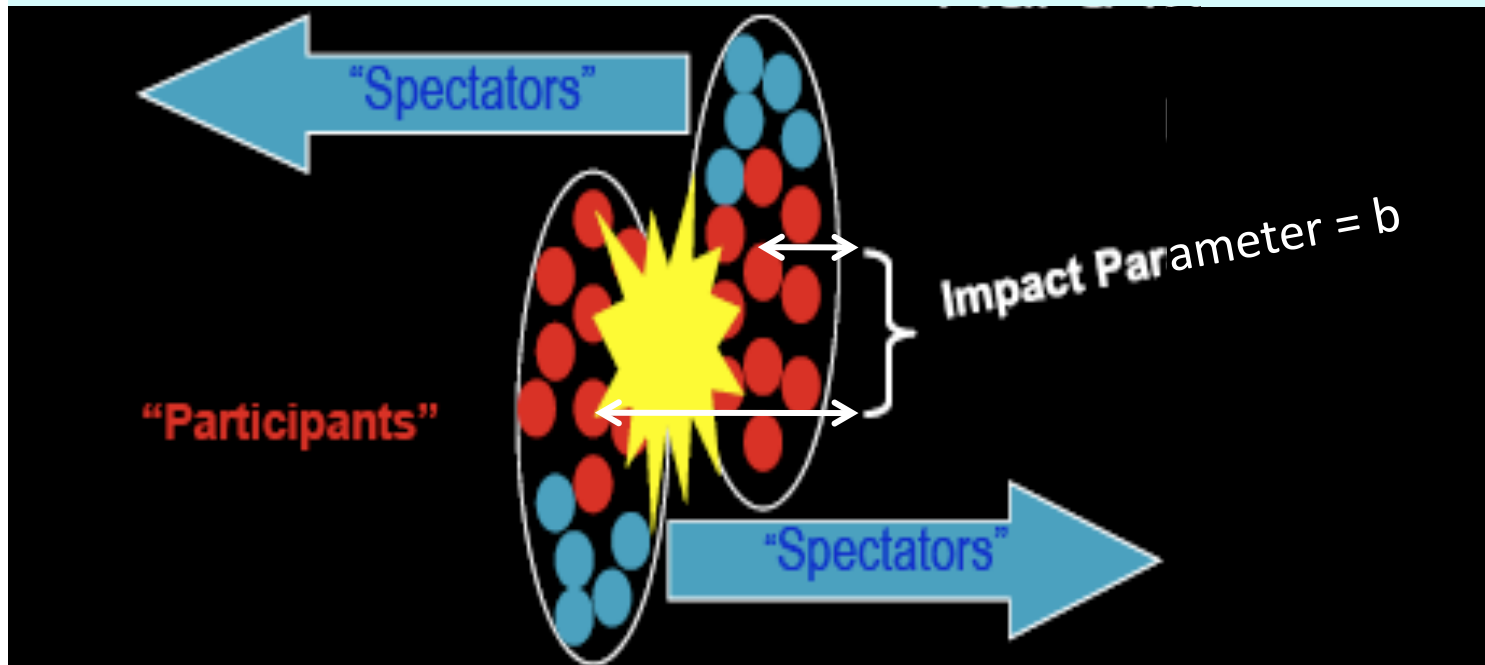
Two ways to measure photons



Big backgrounds

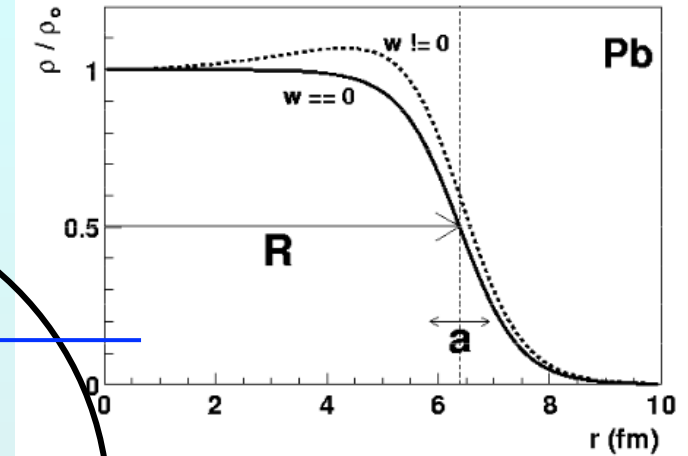
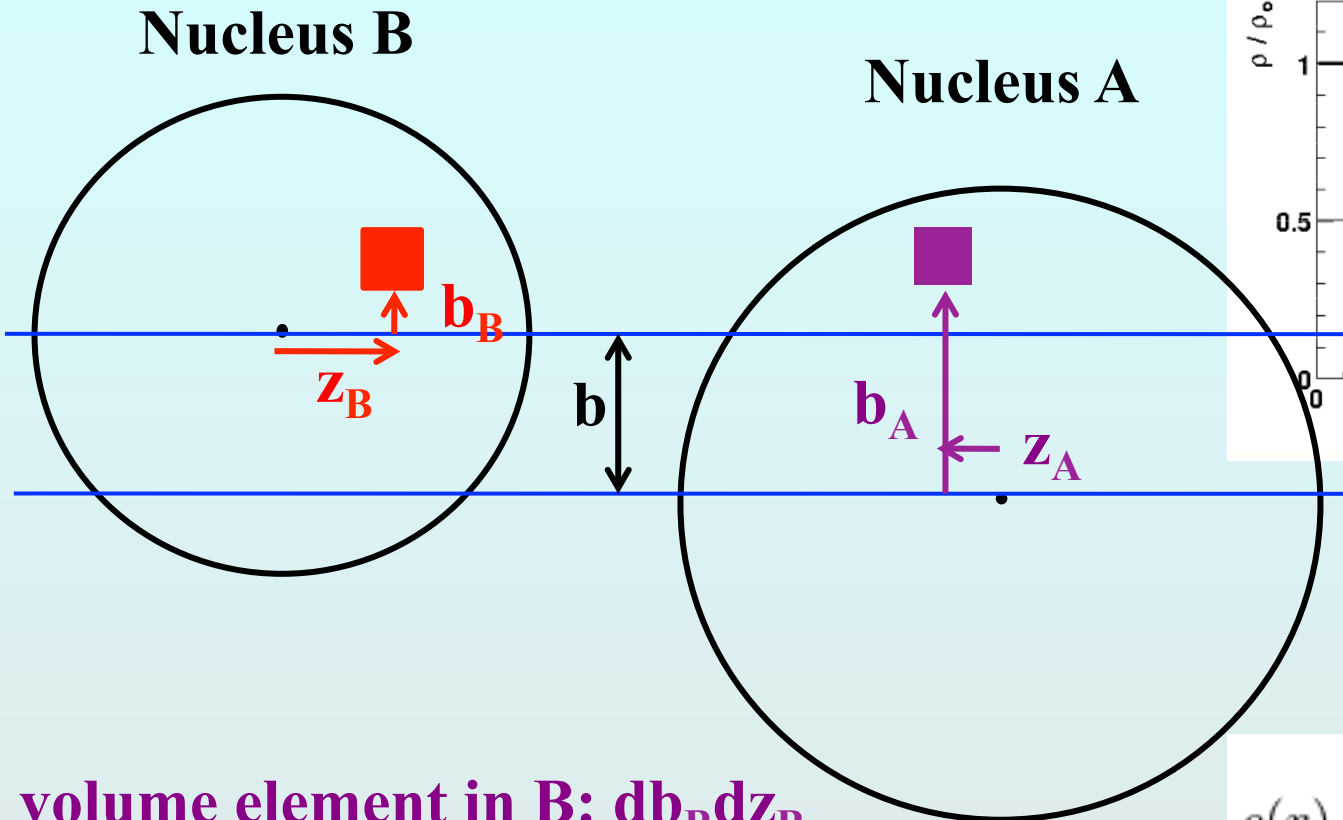


Geometry of heavy ion collisions



Use Glauber model of nucleons in nuclei
calculate # of participant nucleons N_{part}
of binary NN collisions N_{coll}

Glauber model: calculate probabilities



Woods-Saxon nuclear density profile:

$$\rho(r) = \frac{\rho_0 (1 + wr^2/R^2)}{1 + \exp((r - R)/a)}$$

volume element in B: $db_B dz_B$

volume element in A: $db_A dz_A$

Probability of finding a nucleon in volume element A =

$$\rho_A(b_A, z_A) db_A dz_A \quad (\rho_A \text{ is nuclear density * nucleons in A})$$

Probability of n interactions

Thickness function in nucleus A: $\hat{T}_A(\mathbf{s}) = \int \hat{\rho}_A(\mathbf{s}, z_A) dz_A$
s measures transverse overlap – given by b and nucleus R

Thickness function for A+B: $\hat{T}_{AB}(\mathbf{b}) = \int \hat{T}_A(\mathbf{s}) \hat{T}_B(\mathbf{s} - \mathbf{b}) d^2s$

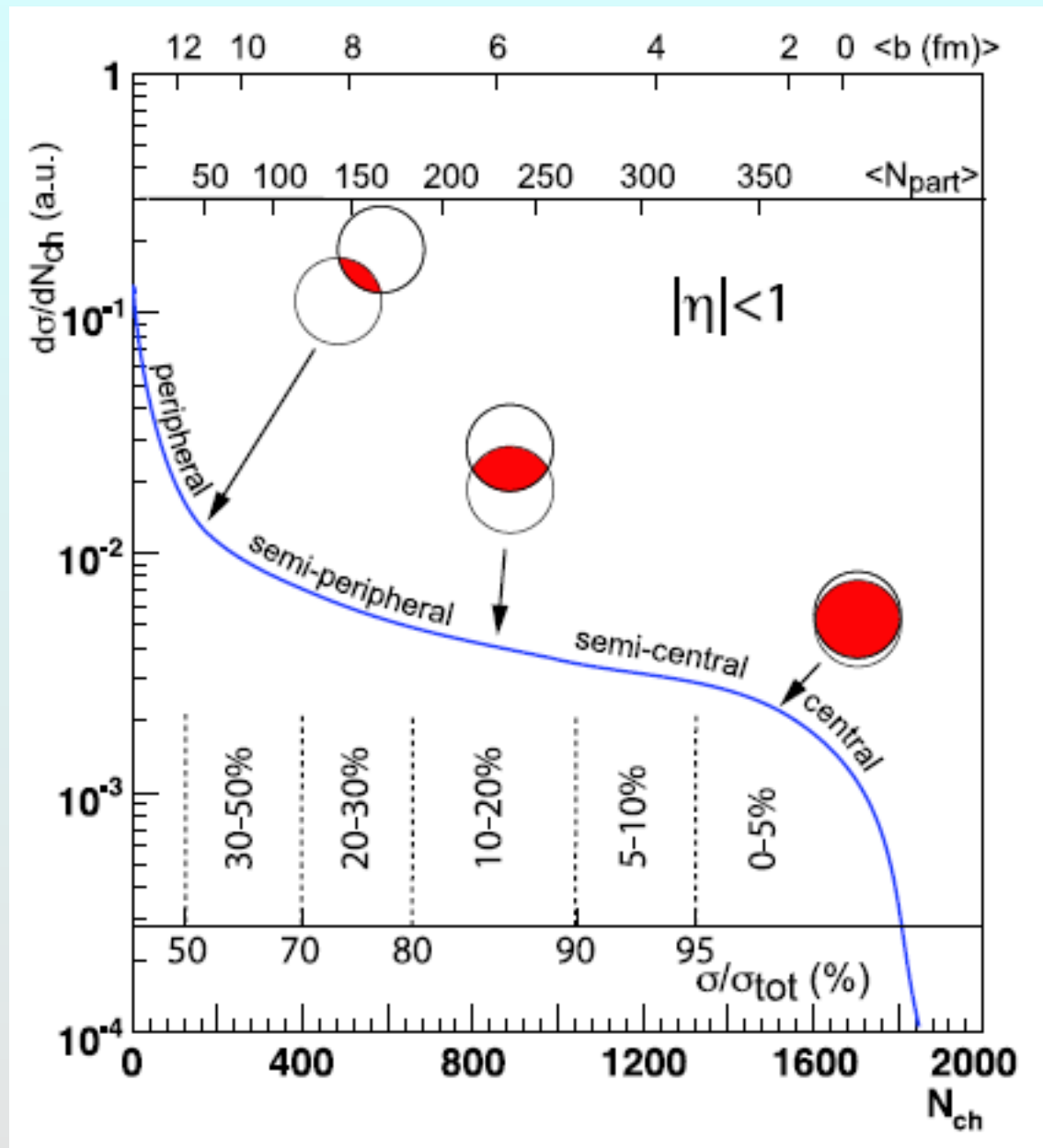
Probability of n interactions is given by binomial distribution:

$$P(n, \mathbf{b}) = \binom{AB}{n} [\hat{T}_{AB}(\mathbf{b}) \sigma_{\text{inel}}^{\text{NN}}]^n [1 - \hat{T}_{AB}(\mathbf{b}) \sigma_{\text{inel}}^{\text{NN}}]^{AB-n}$$

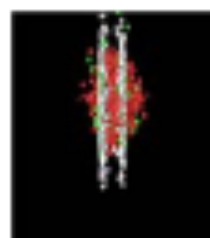
Number of nucleon-nucleon collisions:

$$N_{\text{coll}}(b) = \sum_{n=1}^{AB} n P(n, b) = AB \hat{T}_{AB}(b) \sigma_{\text{inel}}^{\text{NN}}$$

Collision centrality



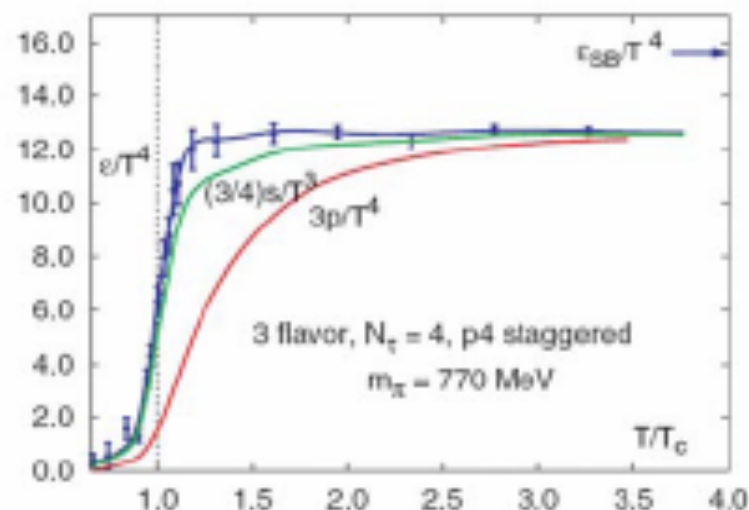
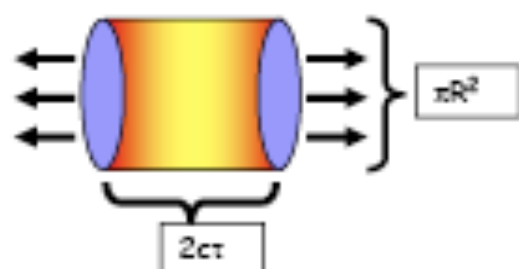
Energy Density



Energy density far above transition value predicted by lattice.

$$\varepsilon_{Bj} = \frac{1}{\pi R^2} \frac{1}{2c\tau} \left(2 \frac{dE_T}{dy} \right)$$

$R \sim 7 \text{ fm}$



Lattice: $T_C \sim 190 \text{ MeV}$ ($\varepsilon_C \sim 1 \text{ GeV}/\text{fm}^3$)

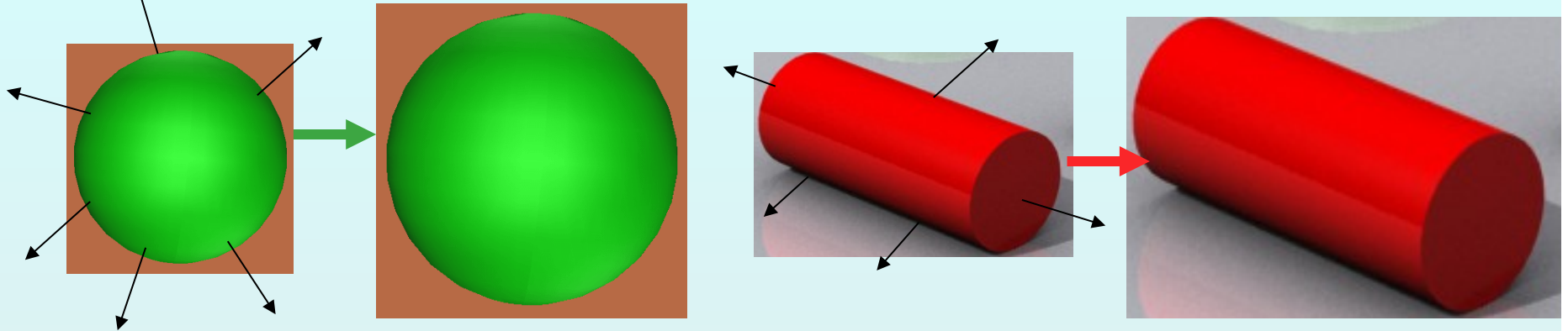
PHENIX: Central Au-Au yields

$$\left\langle \frac{dE_T}{d\eta} \right\rangle_{\eta=0} = 503 \pm 2 \text{ GeV}$$

$$\varepsilon \sim 15 \frac{\text{GeV}}{\text{fm}^3} @ \tau = 0.6 \text{ fm} / c \text{ (thermalization)}$$

Does the matter exhibit collectivity?

Look for collective flow via velocity boosts



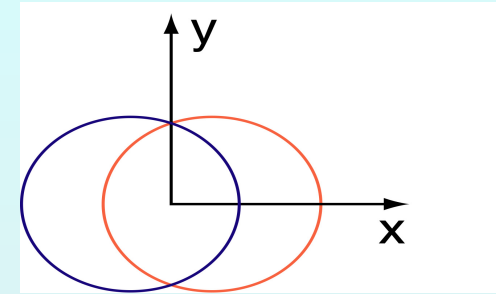
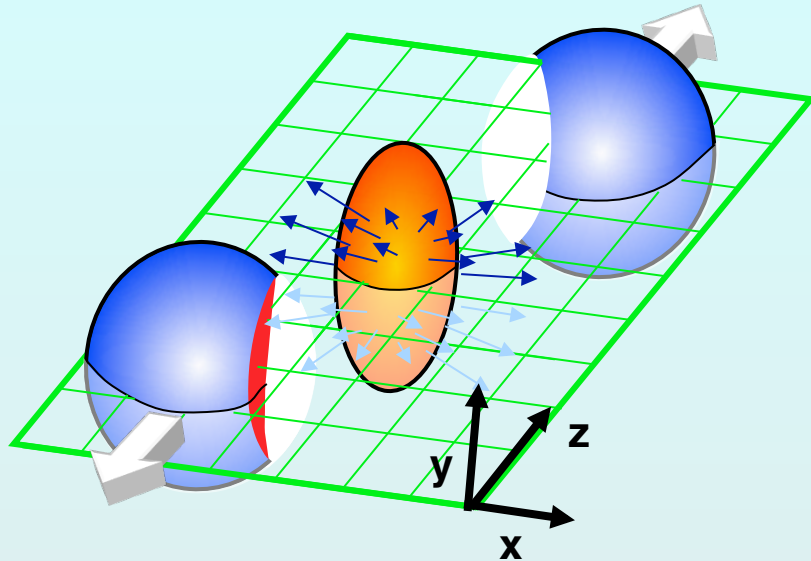
Model expansion of the system with fluid dynamics

$$\partial_t \begin{pmatrix} \rho \\ \rho u \\ \rho v \\ e \end{pmatrix} + \partial_x \begin{pmatrix} \rho u \\ \rho u^2 + p \\ \rho uv \\ u(e + p) \end{pmatrix} + \partial_y \begin{pmatrix} \rho v \\ \rho uv \\ \rho v^2 + p \\ v(e + p) \end{pmatrix} - \partial_x \begin{pmatrix} 0 \\ \tau_{11} \\ \tau_{12} \\ \tau_{11}u + \tau_{12}v + k\partial_x\Theta \end{pmatrix} - \partial_y \begin{pmatrix} 0 \\ \tau_{21} \\ \tau_{22} \\ \tau_{21}u + \tau_{22}v + k\partial_y\Theta \end{pmatrix} = 0,$$

where u and v are the components of the velocity, ρ the density, p the pressure, e total energy density, τ_{ij} the components of the viscous part of the stress tensor, Θ the absolute temperature and k is the heat conductivity.

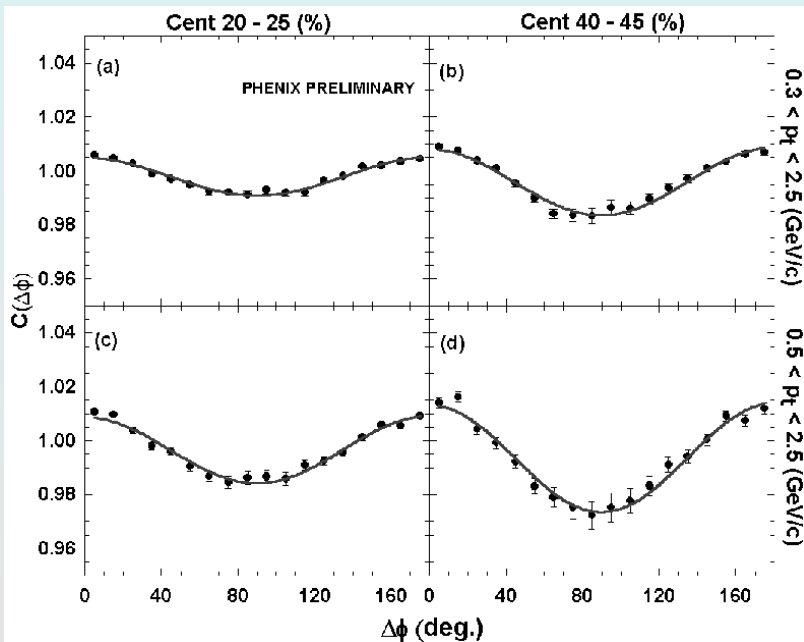
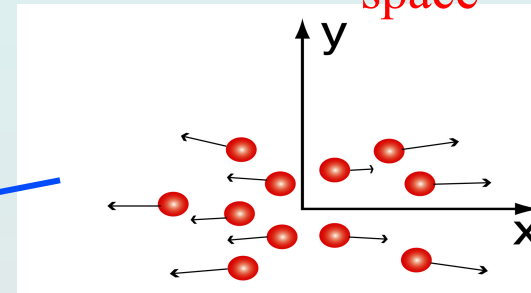
Nuclear theorists have found solutions for 3-D viscous hydro!

Collective motion & elliptic flow (v_2)



Almond shape
overlap region
in **coordinate**
space

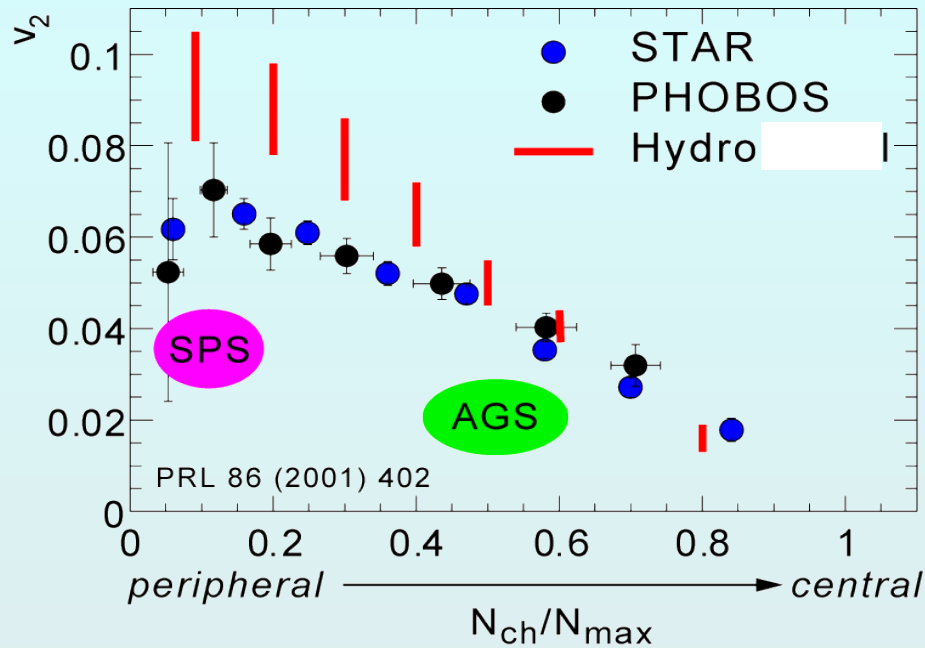
momentum
space



$$dN/d\phi \sim 1 + 2 v_2(p_T) \cos(2\phi) + \dots$$

hydrodynamics works!

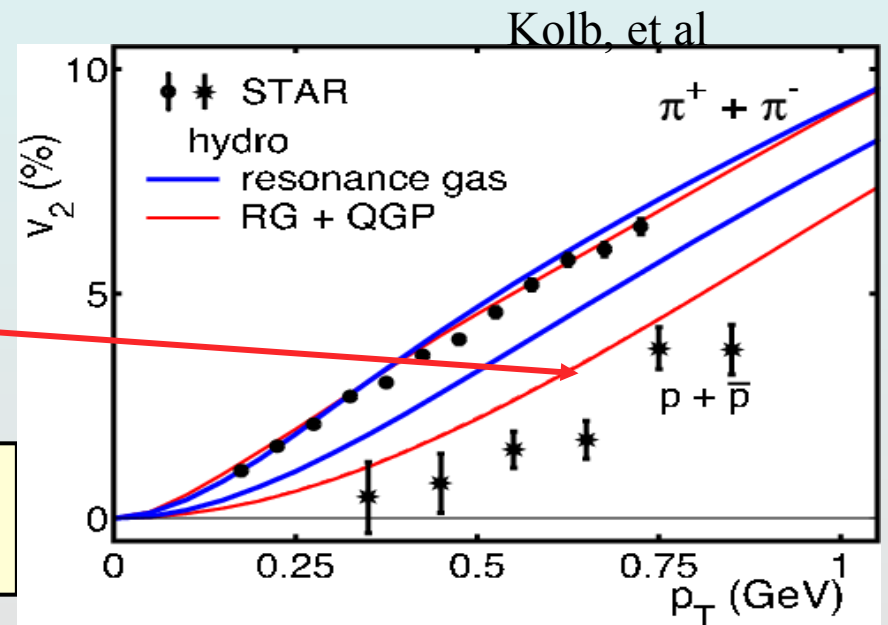
Surprise: matter flows like a liquid



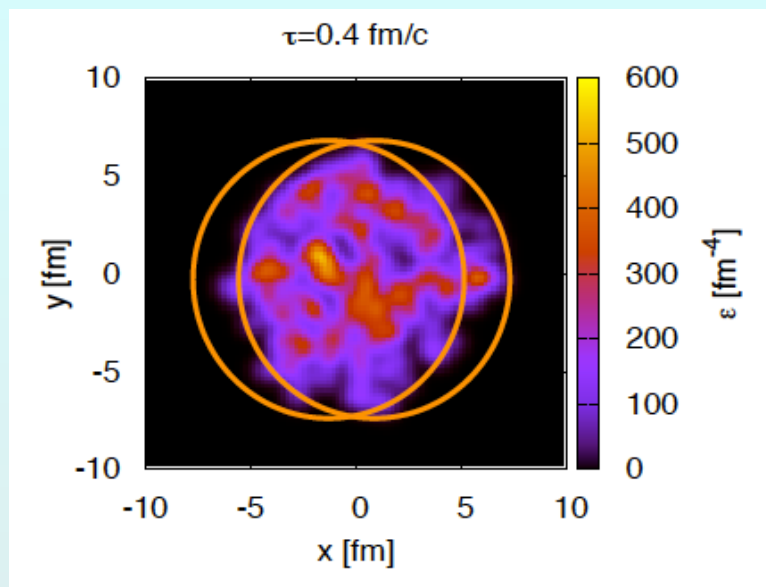
- huge pressure buildup
- large anisotropy \rightarrow it all happens fast
- efficient equilibration mechanism??

Hydrodynamics reproduces elliptic flow of q - \bar{q} and $3q$ states
 Mass dependence requires *soft EOS, NOT gas of hadrons*

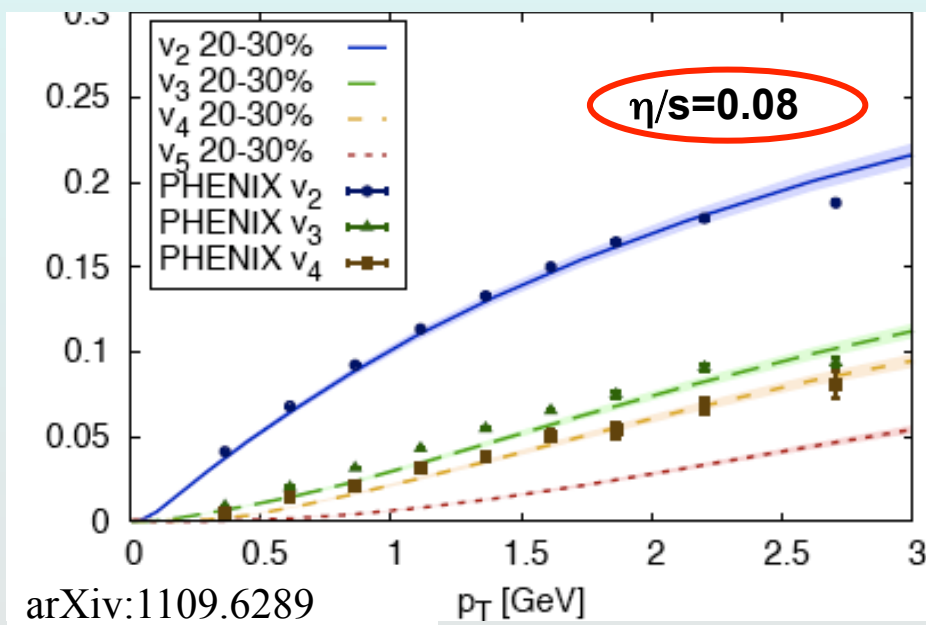
only works with low viscosity/entropy
 “perfect” liquid (D. Teaney, PRC68, 2003)



Fluctuations matter!



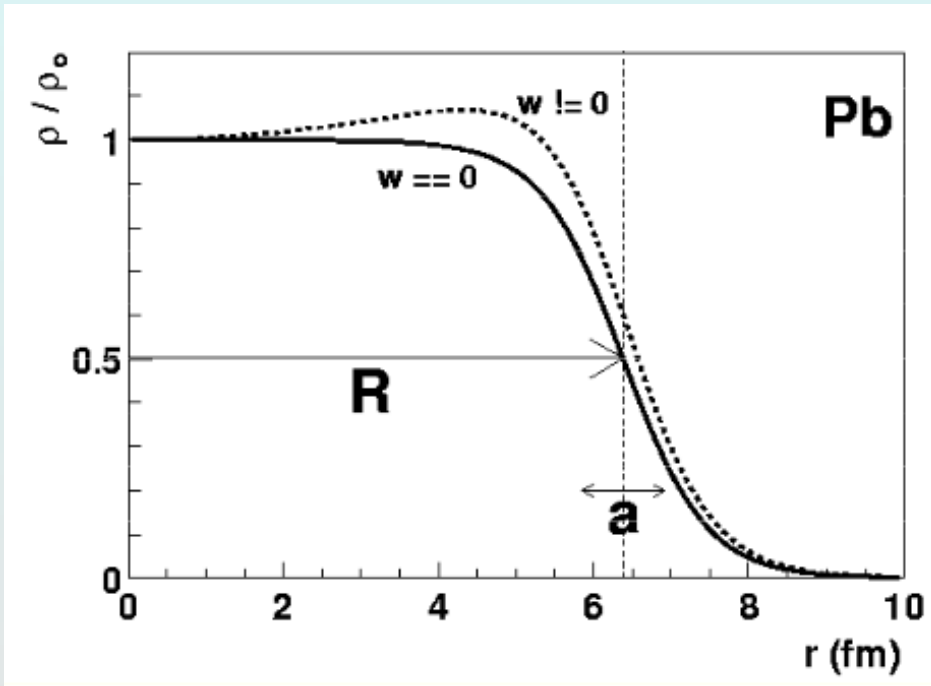
- Nucleons move around inside the nucleus
- > locations of NN scattering fluctuate
- > apparent symmetry effects yielding only even harmonics not realistic



- Reproduce with hydro
- IF include fluctuating initial conditions
- Provides a tool to better pin down the viscosity/entropy ratio

Hydro including event-by-event fluctuations

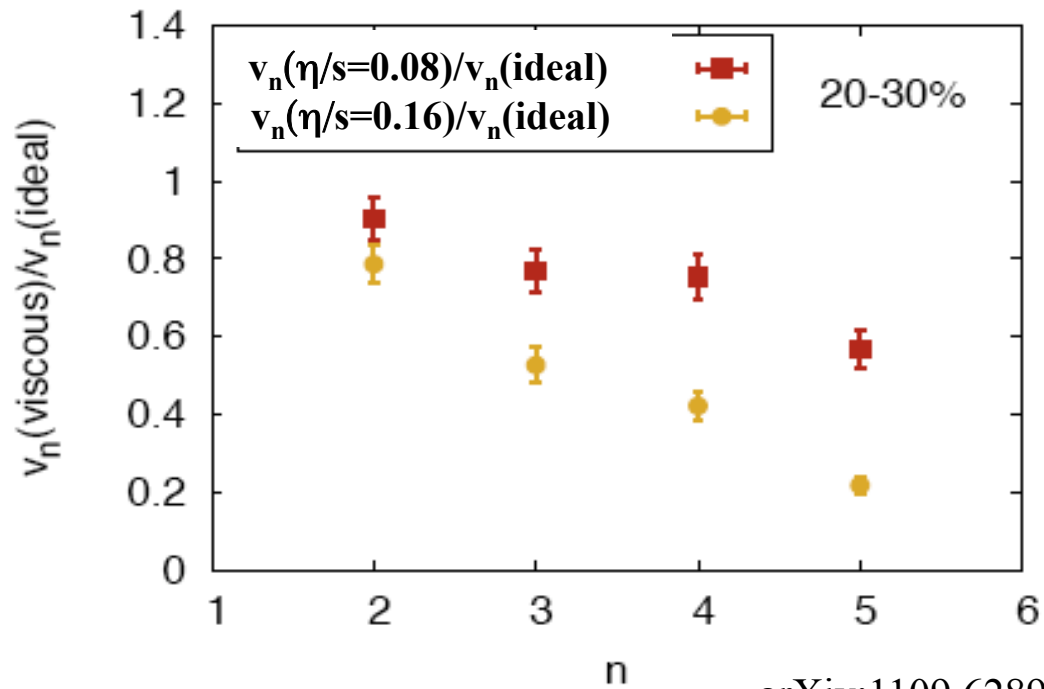
- Use the same Glauber model
 - MC to allow nucleon positions to fluctuate
- Generate an ensemble of hydro calculations with different initial states



Woods-Saxon nuclear density profile:

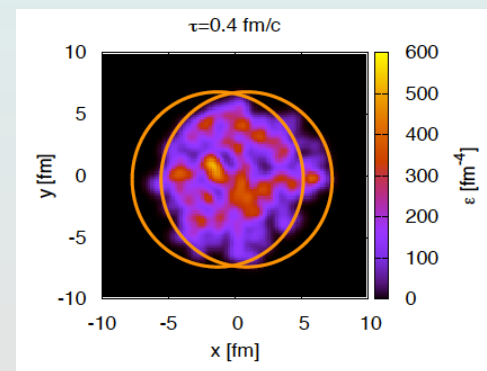
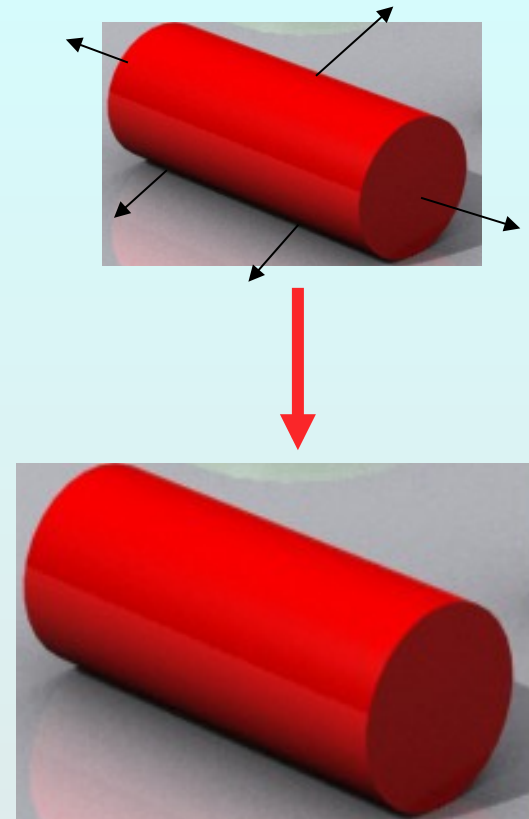
$$\rho(r) = \frac{\rho_0 \left(1 + wr^2/R^2\right)}{1 + \exp\left((r - R)/a\right)}$$

Higher moments more sensitive to viscosity

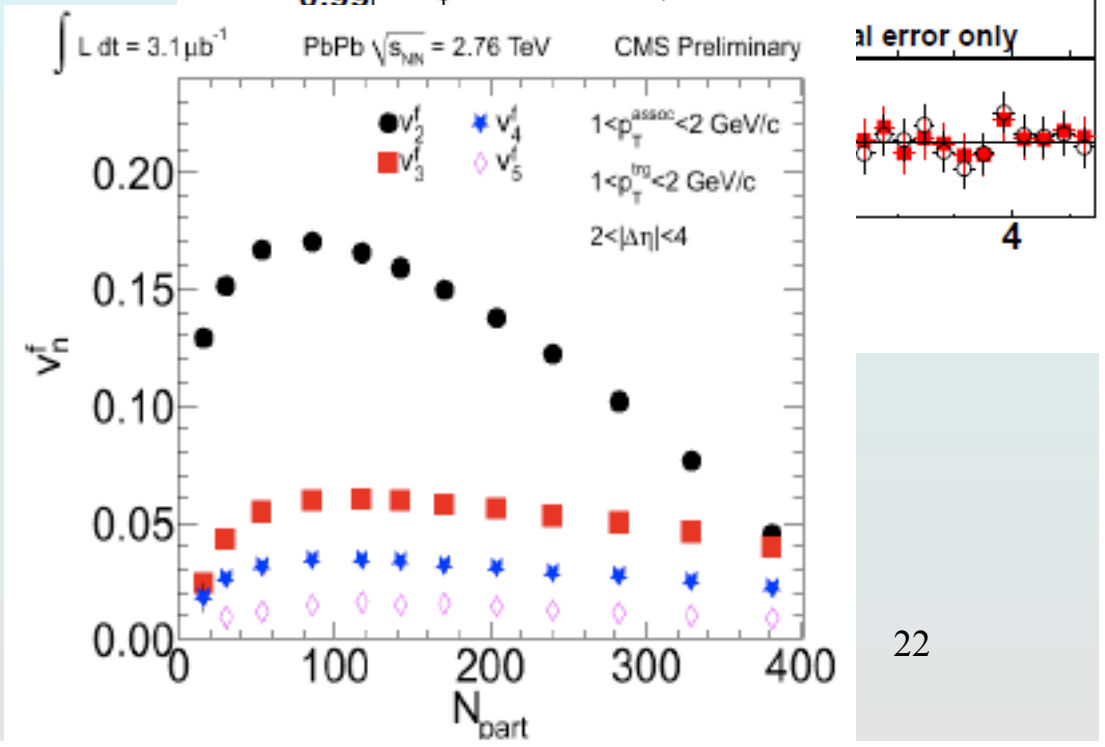
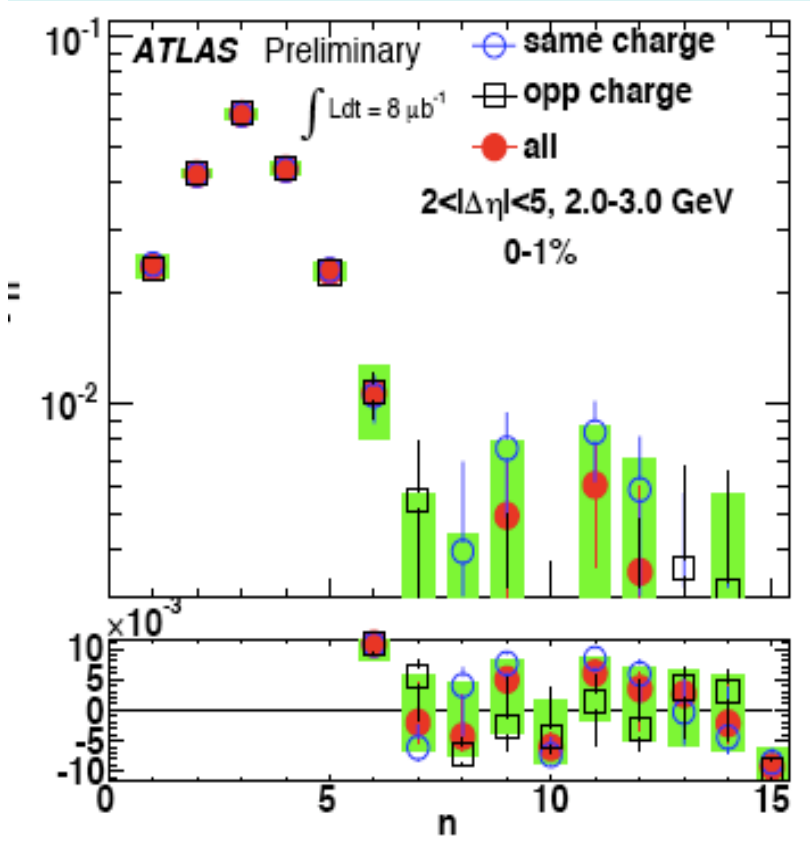
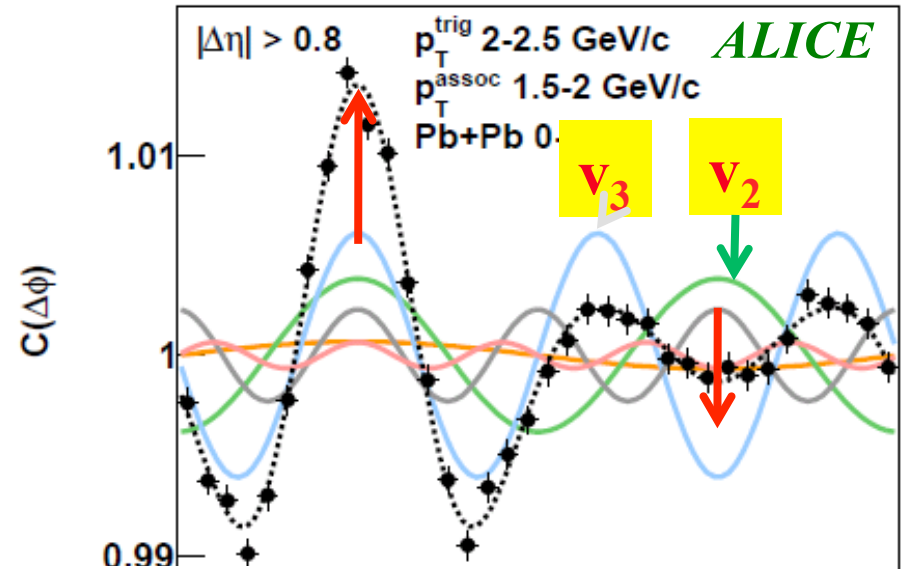
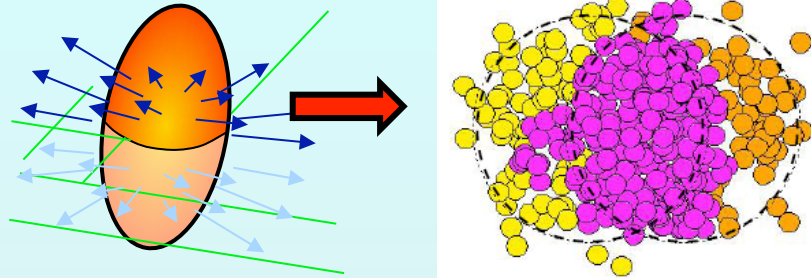


arXiv:1109.6289

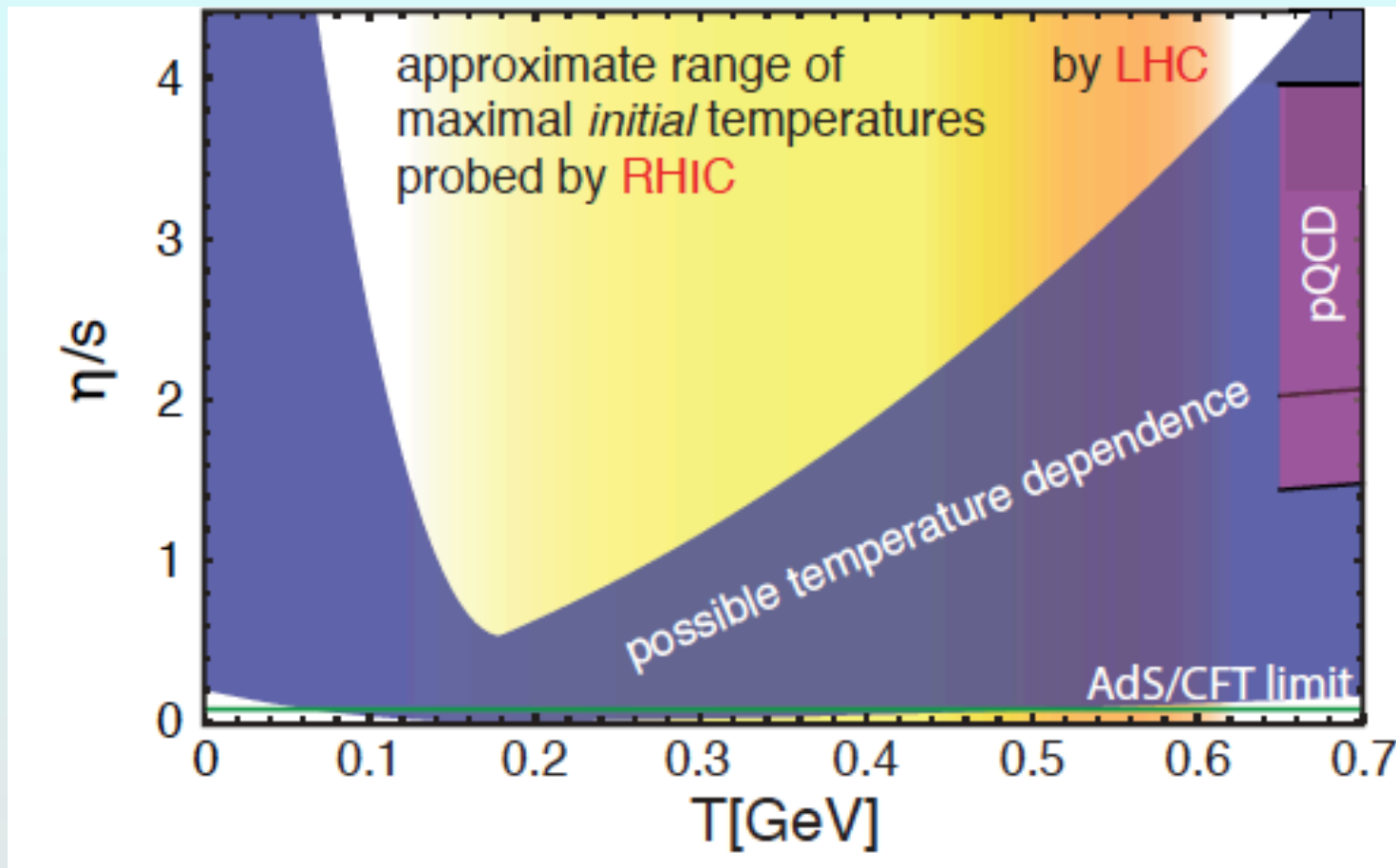
- Longitudinal expansion at $v \sim c$
- “freezes in” small shape perturbations
e.g. triangular fluctuations (v_3)
- Viscosity is like friction!



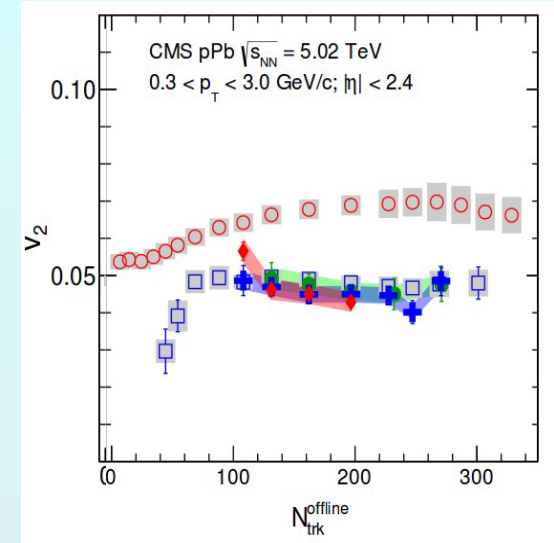
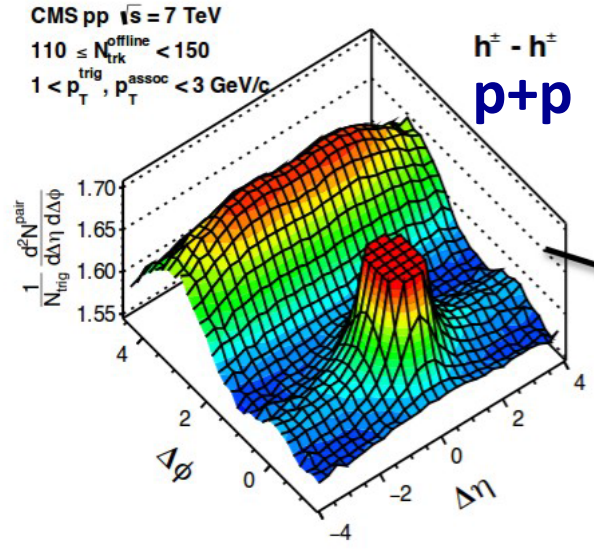
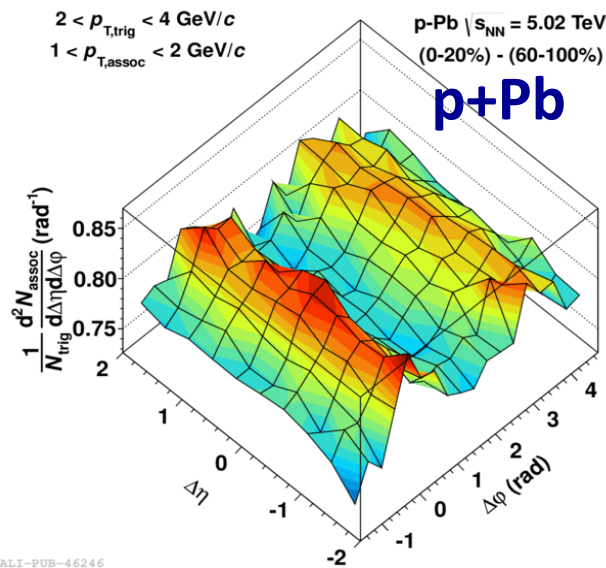
LHC data



Shear viscosity depends on temperature



Collective flow in small systems?



Significant v_N ($n=2$ to 5) with “familiar” ordering + shape in p_T

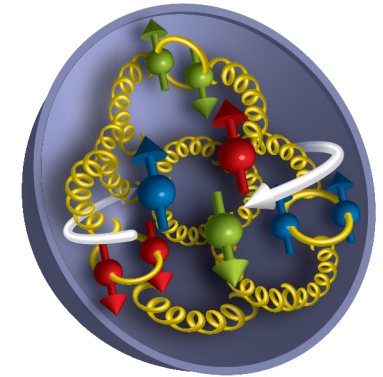
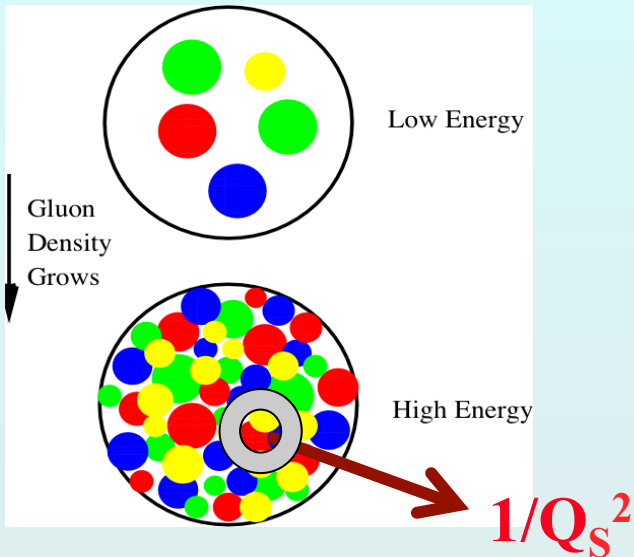
Multi-particle correlations ($v_2\{4\} \approx v_2\{6\} \approx v_2\{8\} \approx v_2\{\text{LYZ}\}$)

● What is going on?

Hydrodynamics in small(ish) systems? Correlations between particles produced in initial state? Radial flow an artefact of constant temperature freezeout surface?

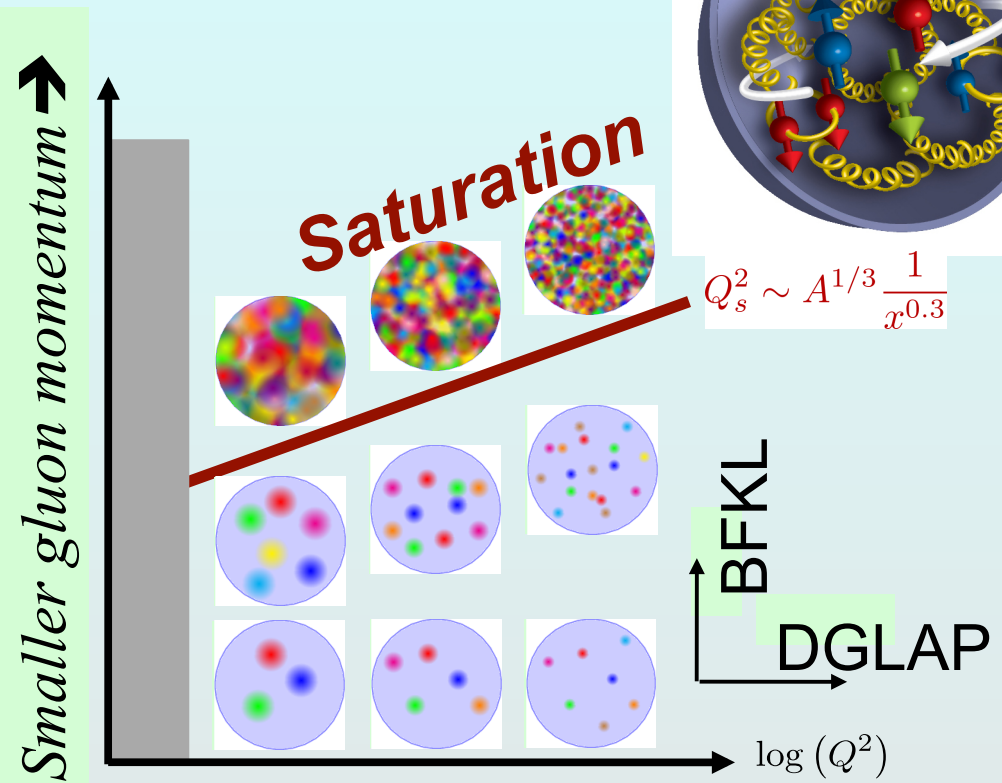
Hot, dense gluonic matter is surprising Are cold dense gluons wierd too?

- Look deep in a nucleus: gluons are numerous

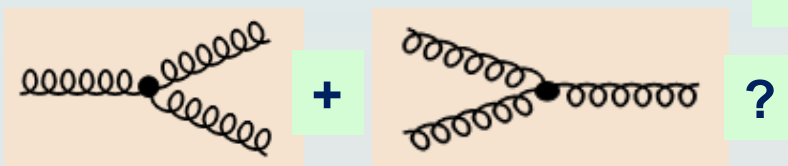


$$Q_s^2 \sim A^{1/3} \frac{1}{x^{0.3}}$$

- At high density what then?



Increasing probe energy \rightarrow



- This is our initial state in heavy ion collisions!

Probe with $e+A$

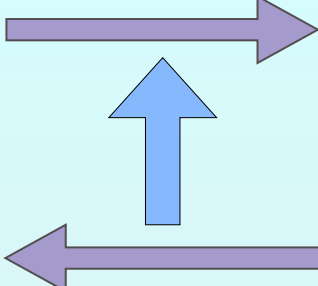
Surprise: viscosity/entropy is small

Viscosity: inability to transport momentum & sustain a wave

low $\eta \rightarrow$ large σ , transports disturbances

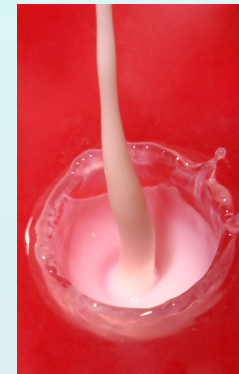
Viscosity/entropy near $1/4\pi$ limit from quantum mechanics!

\therefore liquid at RHIC is “perfect”



The diagram consists of three arrows: a top purple arrow pointing right, a middle blue arrow pointing up, and a bottom purple arrow pointing left. To the right of these arrows is the equation $\eta \approx \frac{1}{3} n \bar{p} \lambda_f$ and below it, $\bar{p} \sim T$.

$$\eta \approx \frac{1}{3} n \bar{p} \lambda_f$$
$$\bar{p} \sim T$$



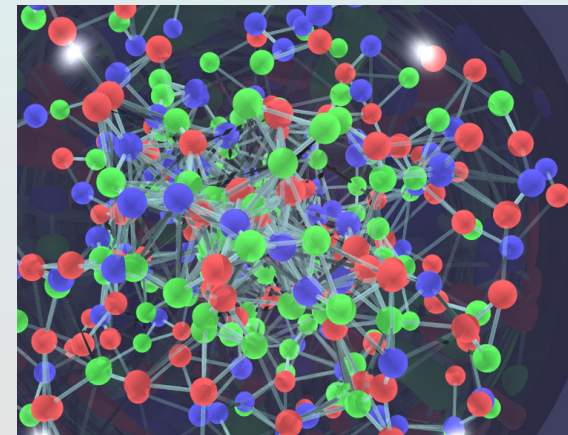
Example: milk.
Liquids with higher viscosities will not splash as high when poured at the same velocity.

Good momentum transport: neighboring fluid elements “talk” to each other

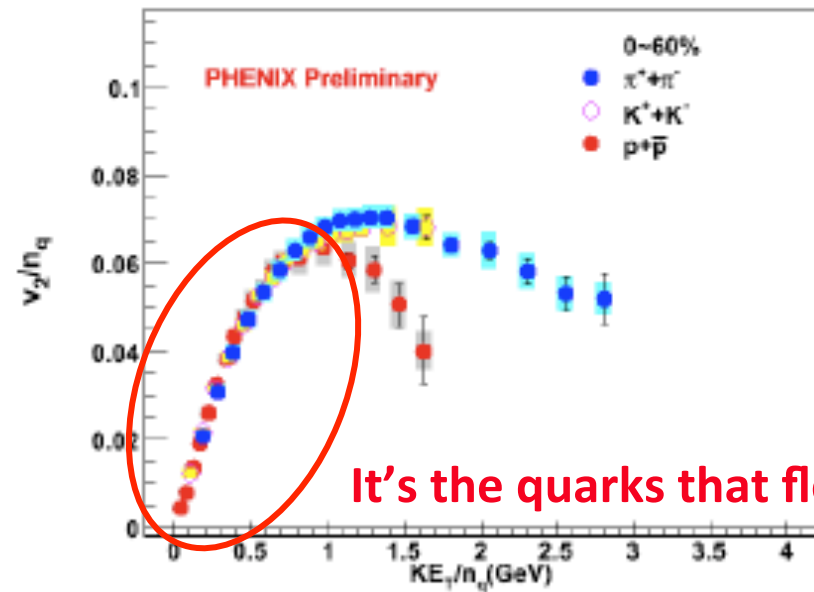
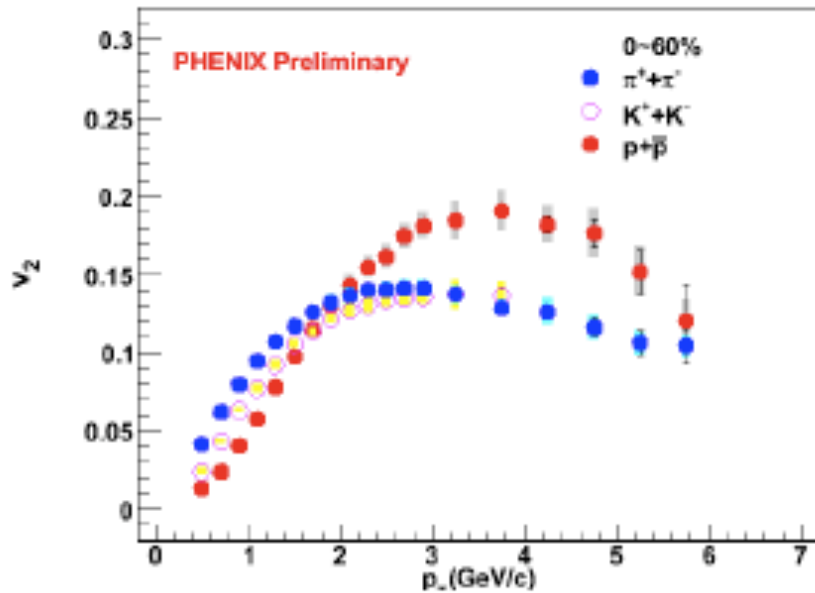
\rightarrow QGP is strongly coupled

Should be opaque:

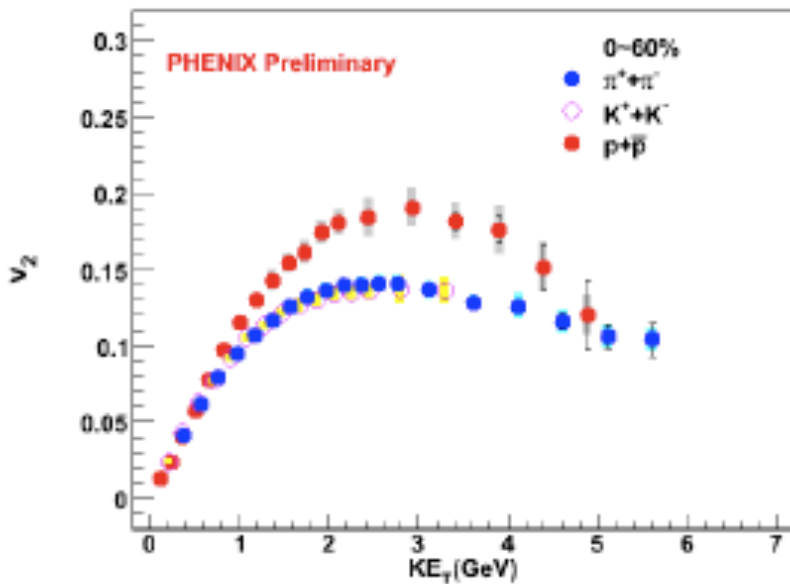
e.g. q,g collide with “clumps” of gluons, not individuals



p_T dependence of v_2



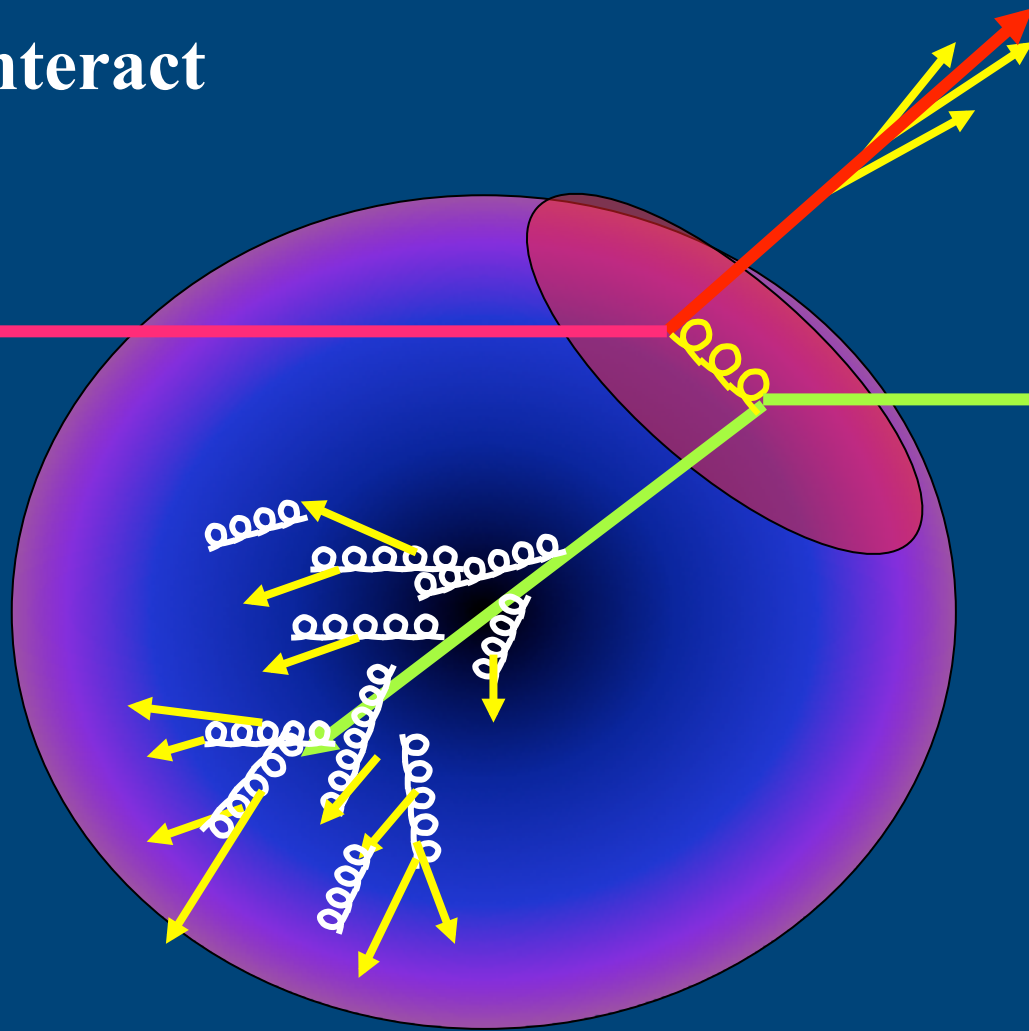
It's the quarks that flow!



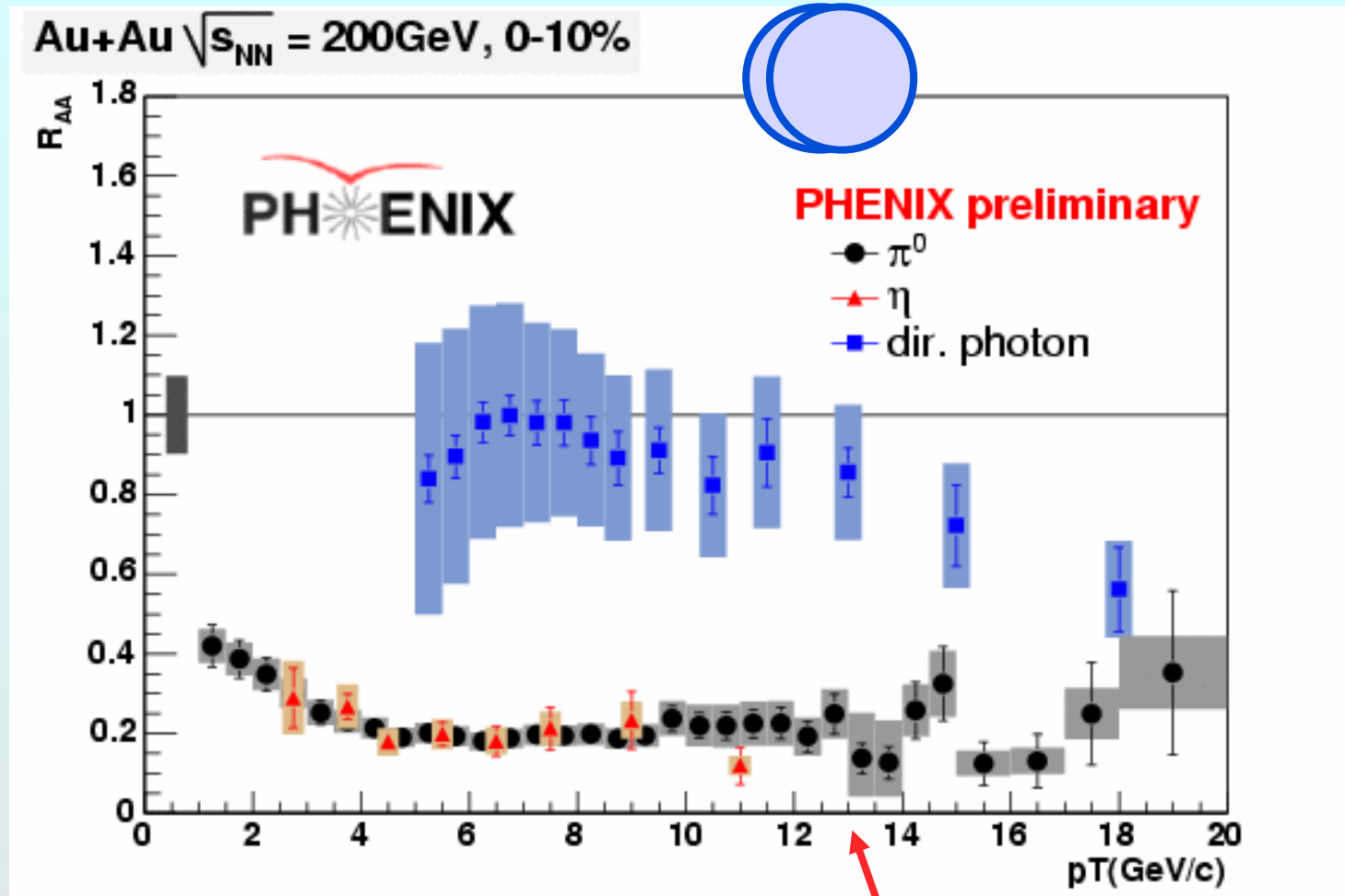
- $KE_T/n_q < 1 \text{ GeV}$ – soft physics
Hydrodynamic flow
- Interplay soft-hard $3.0 < p_T < 5 \text{ GeV/c}$?
- Hard dominates: $p_T > 5 \text{ GeV/c}$

Do fast quarks & gluons escape the plasma?

They feel the strong interaction, so they should interact



colored objects lose energy, photons don't

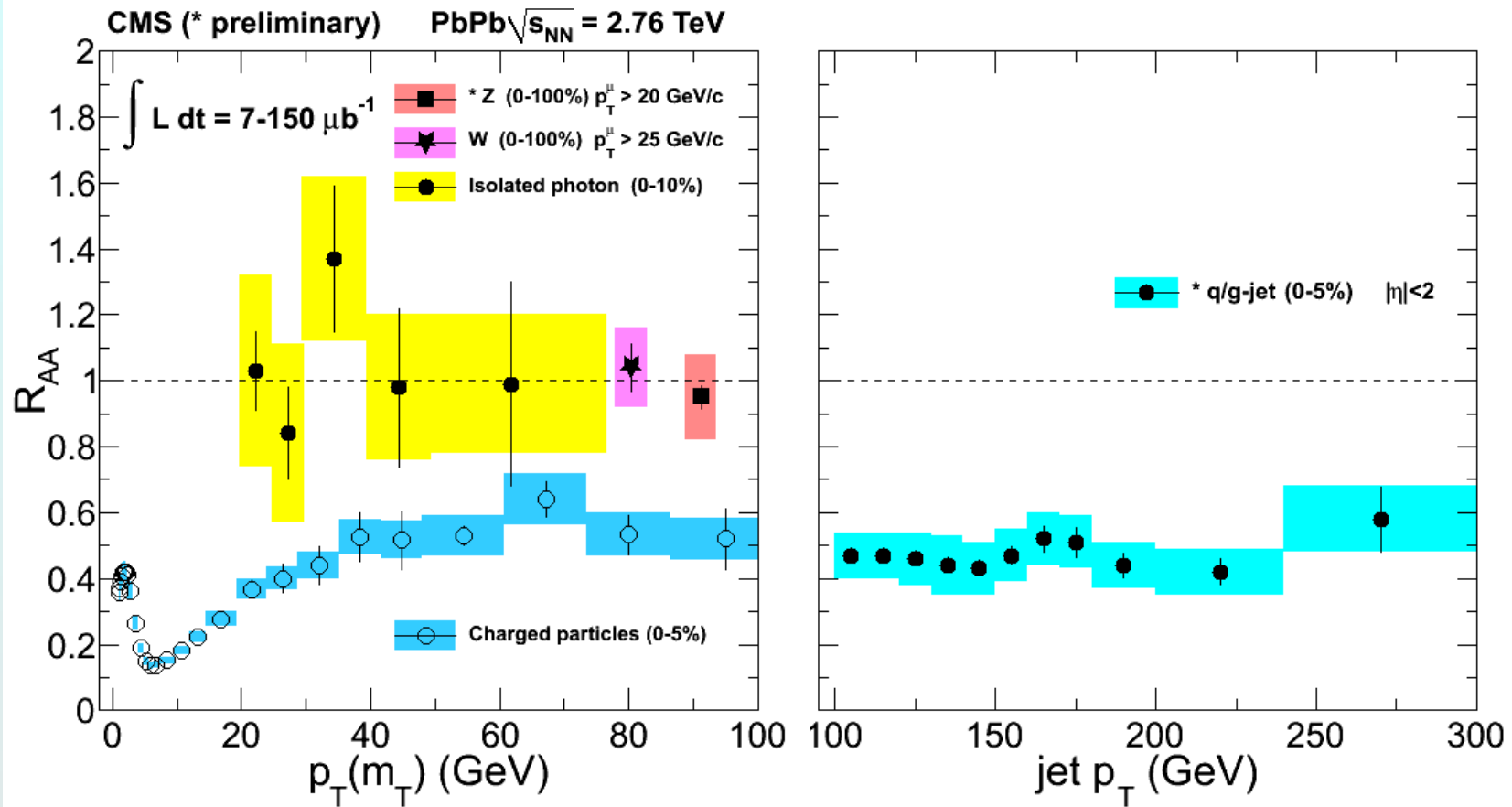


Nuclear modification factor:

$$R_{AA}(p_T) = \frac{d^2 N^{AA} / dp_T d\eta}{T_{AA} d^2 \sigma^{NN} / dp_T d\eta}$$

VERY opaque! Lots of gluon radiation (bremsstrahlung)

Energy loss even by very energetic q & g



● LHC experiments reach to 300 GeV!

Energy loss depends on medium density

- In dilute medium

Independent processes: bremsstrahlung & scattering

Calculate probabilities and add them up

Independent radiations follow Bethe-Heitler

- In dense medium

Mean free path is short: $\lambda = \sigma/\rho$

Formation time of radiated gluon: $\tau = \omega/k_T^2$

Transverse momentum of radiated gluon: $k_T^2 = n\mu^2$

of collisions $n = L/\lambda$, $\mu =$ typical p_T transfer in 1 scattering

λ, μ are properties of the medium, combine to $\hat{q} = \sqrt{\mu^2/\lambda}$

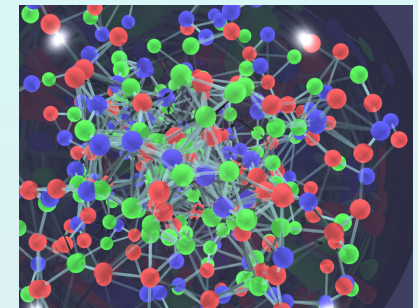
- Coherence in the dense medium!

Next scattering takes place faster than gluon formation

Add amplitudes for all multiple scatterings

In QCD this increases the energy loss!

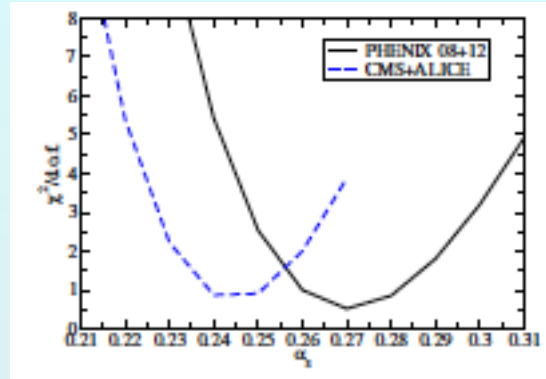
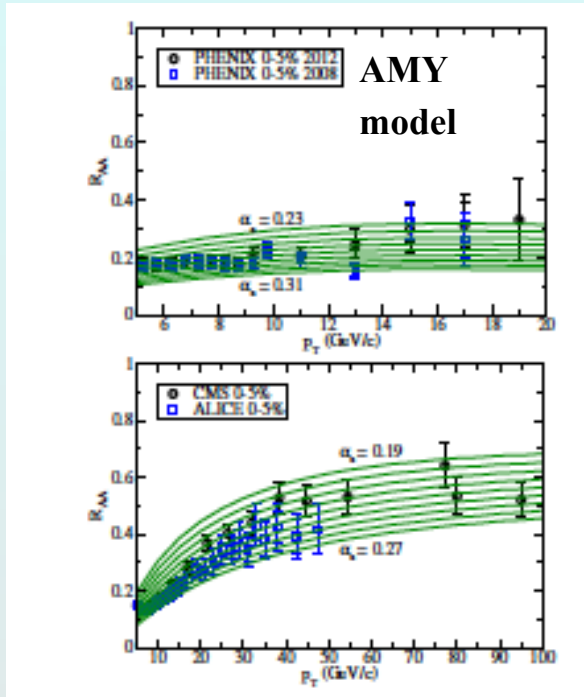
some evidence for coherence already in cold matter! ³¹



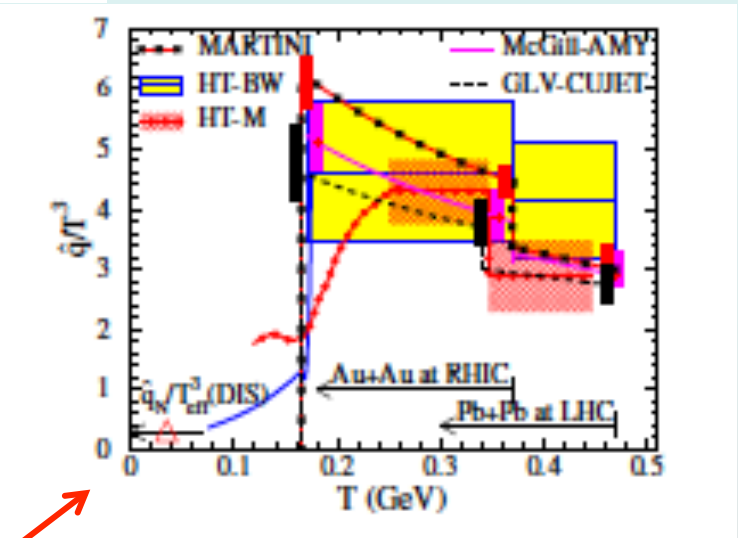
Fit R_{AA} at different \sqrt{s}

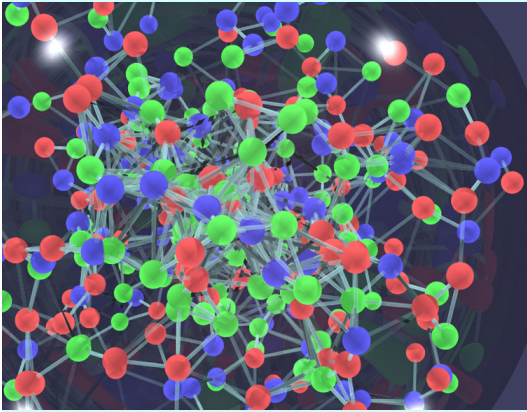
arXiv:1312.5003

JET collaboration fit all data with multiple calculations
 minimize χ^2 for best fit to strong coupling parameter or \hat{q}



Put together all the calculations



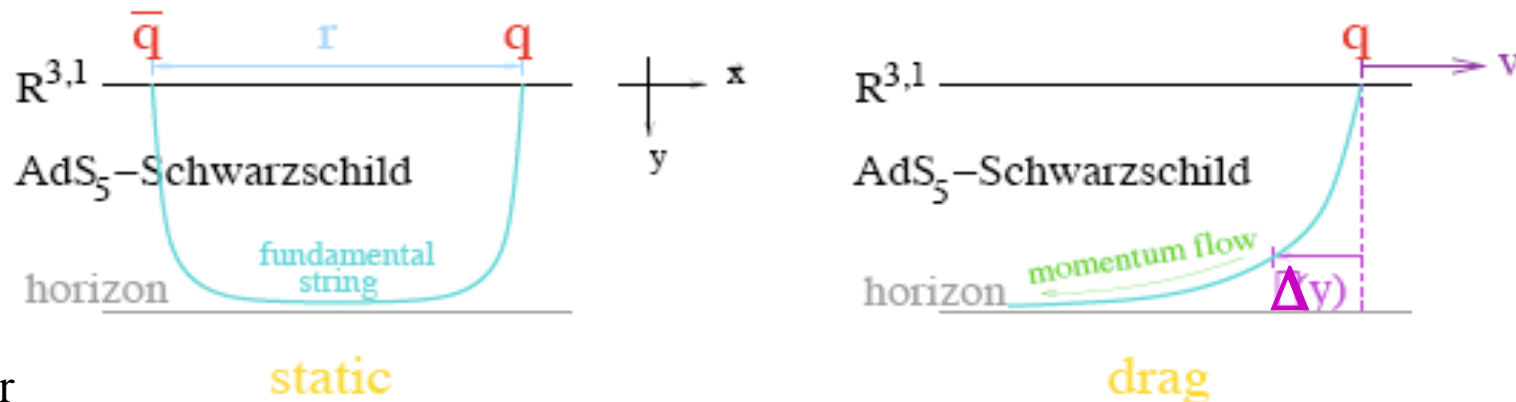


Interactions with plasma?

- radiation (bremsstrahlung)
- collisional energy loss

In plasma: interactions among charges of multiple particles
 charge is spread, screened in characteristic (Debye) length, λ_D
also the case for strong, rather than EM force ... Effect on collisions?

- AdS/CFT says QGP is a strongly interacting field, model as ∞ strong
 Interact with this QGP as with a tiny black hole
 No particles to hit, none can survive inside. Eloss \rightarrow collective excitations



S. Gubser

static

drag

Figure 2: Left: a screened attraction between static quark arises from a string dipping into AdS_5 -Schwarzschild. Right: a drag force arises from a string tailing behind a moving quark.

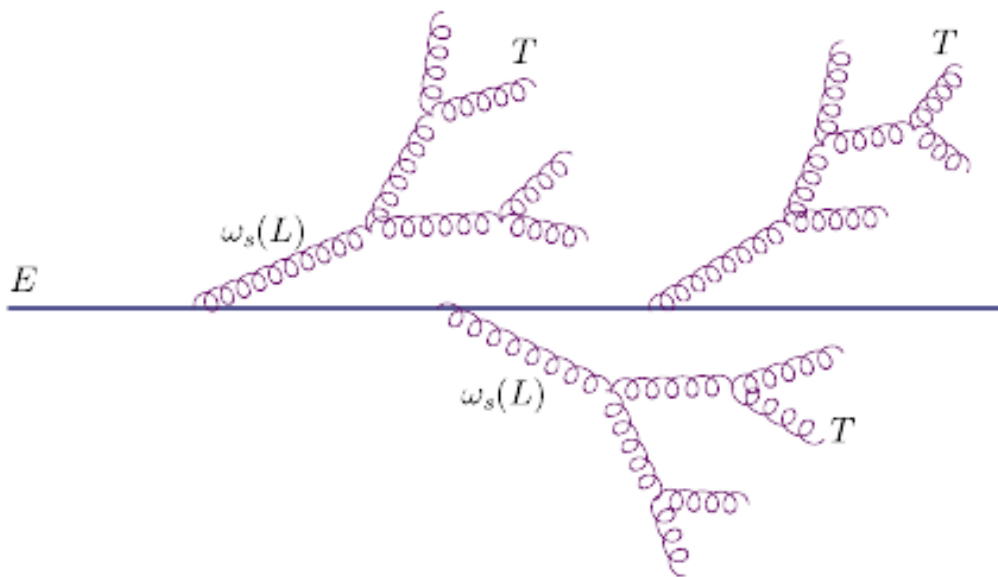
Back to QCD...

Jet versus mini-jets

- At the LHC, the energy E of the leading particle is much higher

$$E \geq 100 \text{ GeV} \gg \omega_{\text{br}} = \alpha_s^2 \hat{q} L^2 \simeq 4 \div 16 \text{ GeV for } L = 3 \div 6 \text{ fm}$$

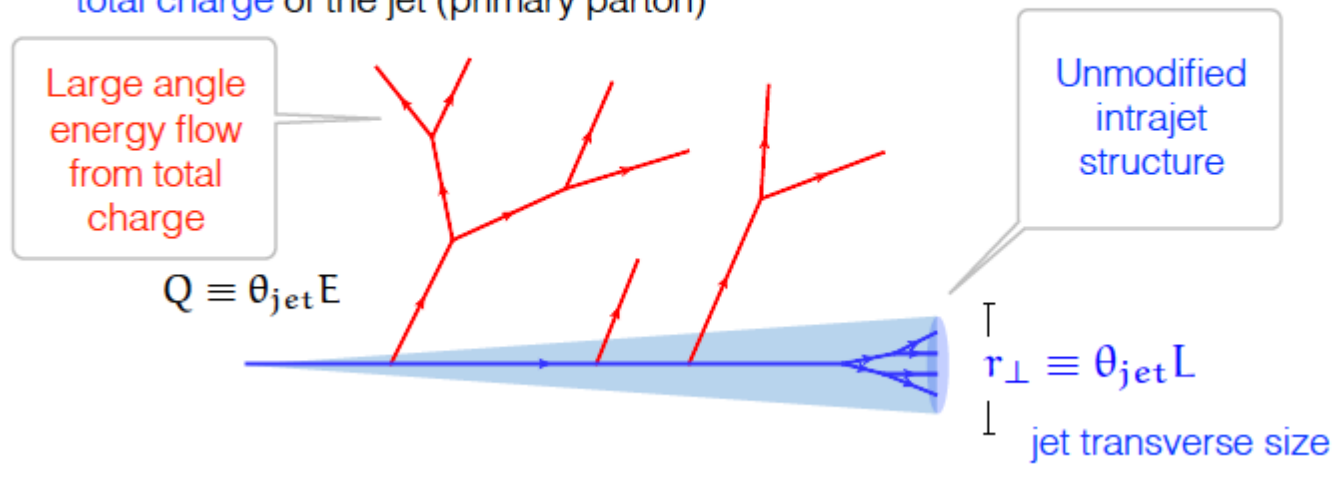
- the leading particle cannot suffer democratic branchings
- it abundantly emits primary gluons with $p \lesssim \omega_{\text{br}}(L)$
- these primary gluons generate 'mini-jets' via democratic branchings
- the energy flows from one parton generation to the next one



$$\Delta E_{\text{flow}} \simeq 2.5 \omega_{\text{br}}(L)$$

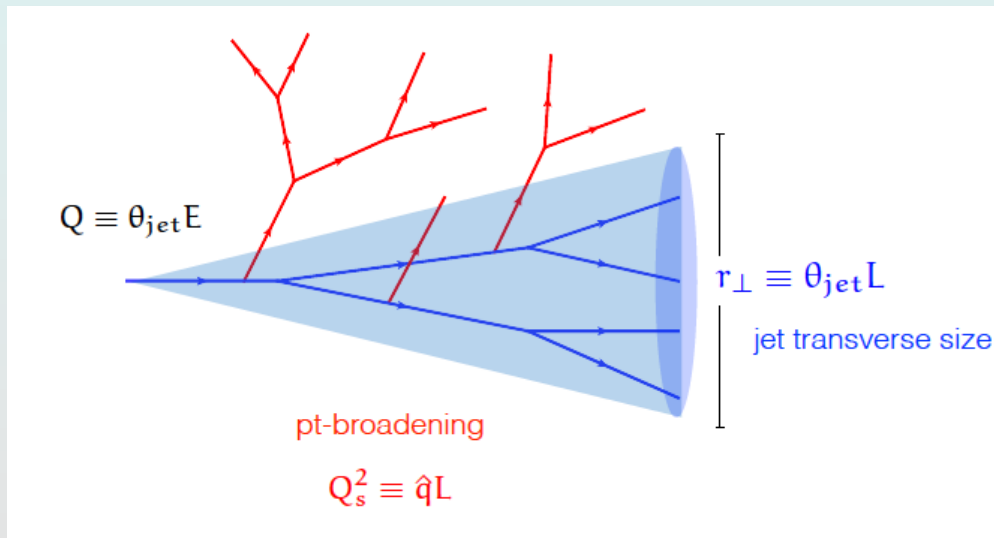
- 'absorbed by the medium'
- is the medium a 'perfect sink' ?
- is the branching process affected by thermalization ?

- When the transverse size r_{\perp} of the jet is smaller than medium resolution scale Q_s^{-1} the medium interacts “effectively” with the total charge of the jet (primary parton)



coherent limit

Subsequent splittings drive color de-coherence of 1st generation emitted gluons & softer leading partons



How this is calculated

- There is a great paper comparing vacuum DGLAP evolution to an in-medium cascade (BDMPS), which depends on \hat{q} .

- See Blaizot and Mehtar-Tani:

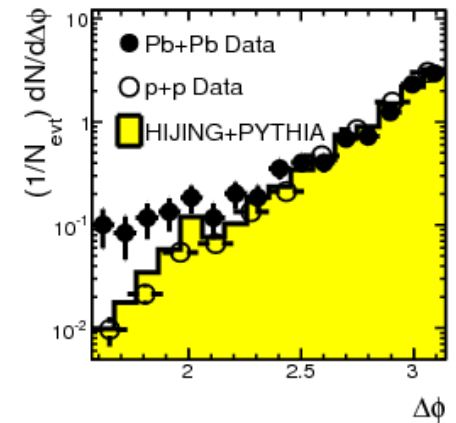
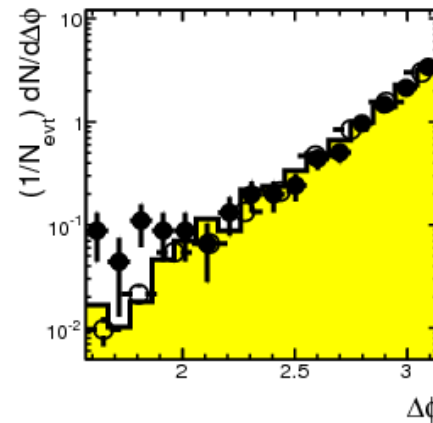
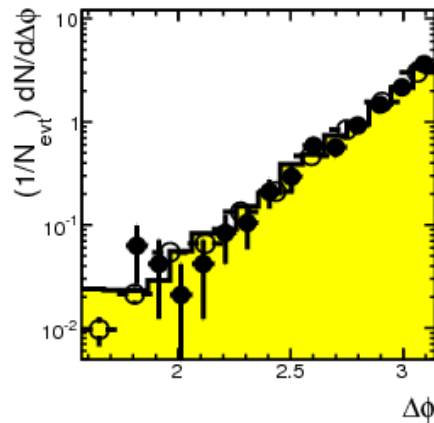
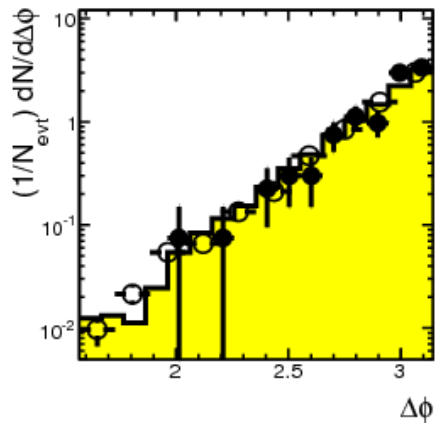
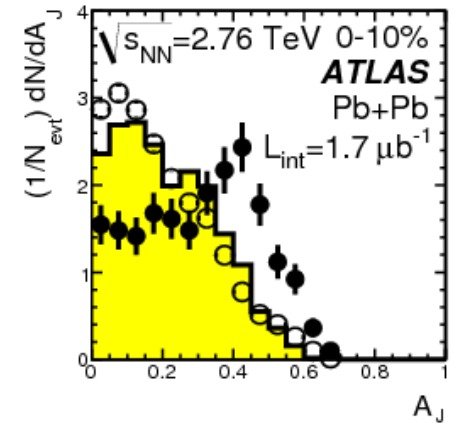
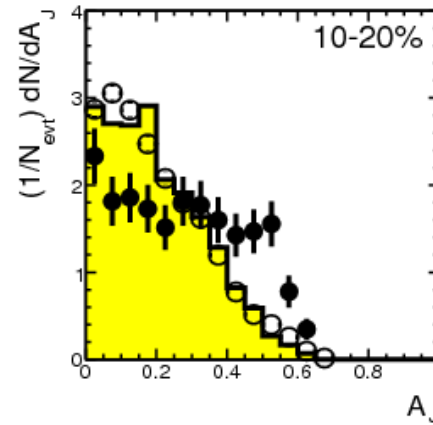
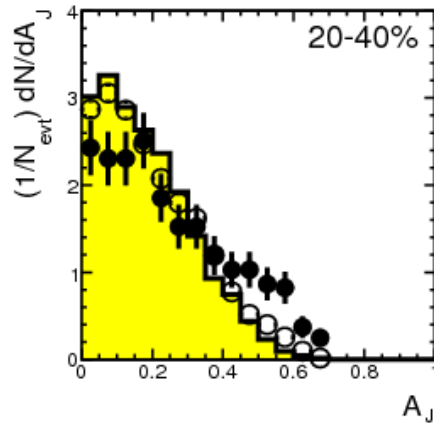
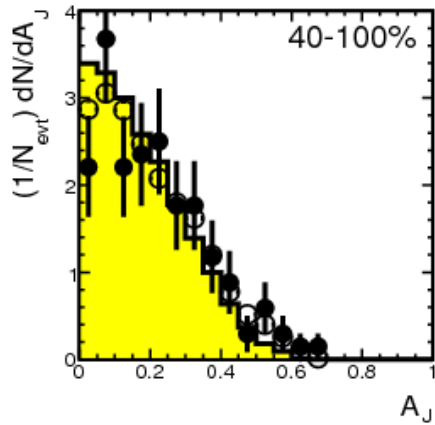
<http://arxiv.org/pdf/1501.03443v1.pdf>

- In medium cascade results in energy flow to $x \sim 0$ due to multiple branchings. DGLAP has an infrared cutoff. DGLAP branching time is constant, while in BDMPS it decreases along the cascade. This allows transport of energy to very soft gluons in a finite amount of time. Result looks a lot like wave turbulence

A good topic for a (somewhat technical) QCD talk!

Jet asymmetry

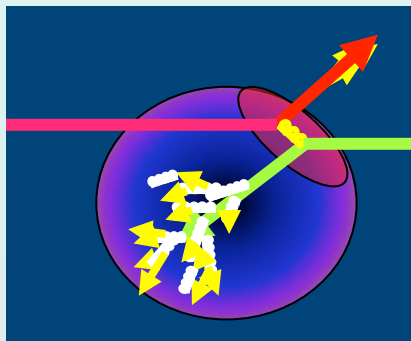
$$A_J = \frac{E_{T1} - E_{T2}}{E_{T1} + E_{T2}}, \Delta\phi > \frac{\pi}{2}$$



Where does the lost energy go?

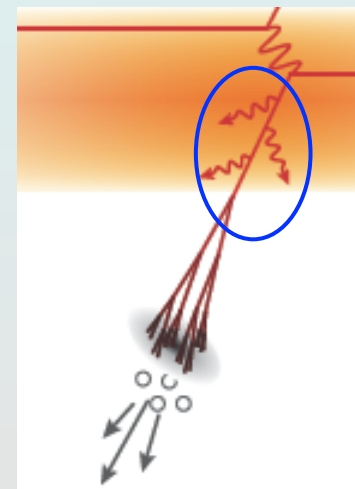
- We don't know yet!
- Medium enhances gluon radiation/splitting:

extra gluons at small angles (in/near jet cone)

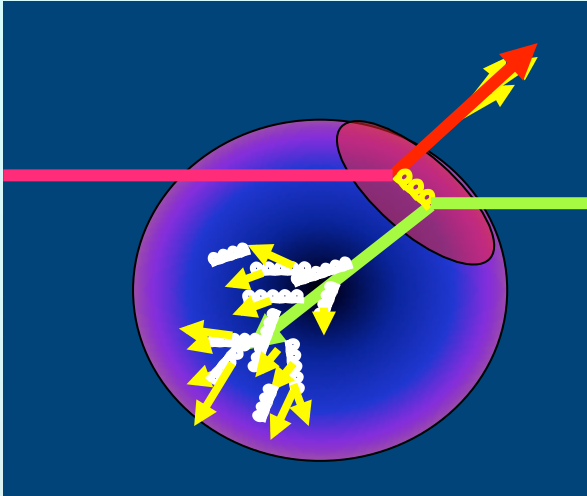


radiated gluons thermalize in medium (i.e. they're gone!)

remain correlated with leading parton, but broaden/change jet

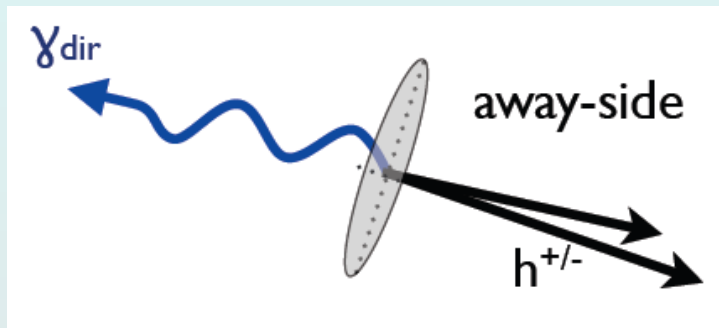


Jet Fragmentation function



$$D(z) = 1/N_{jet} dN(z)/dz; z = p_{had}/p_{jet}$$

Measure: count partners per trigger as fraction of trigger momentum



$$z_T = p_{Ta}/p_{Tt} \sim z \text{ for } \gamma \text{ trigger}$$

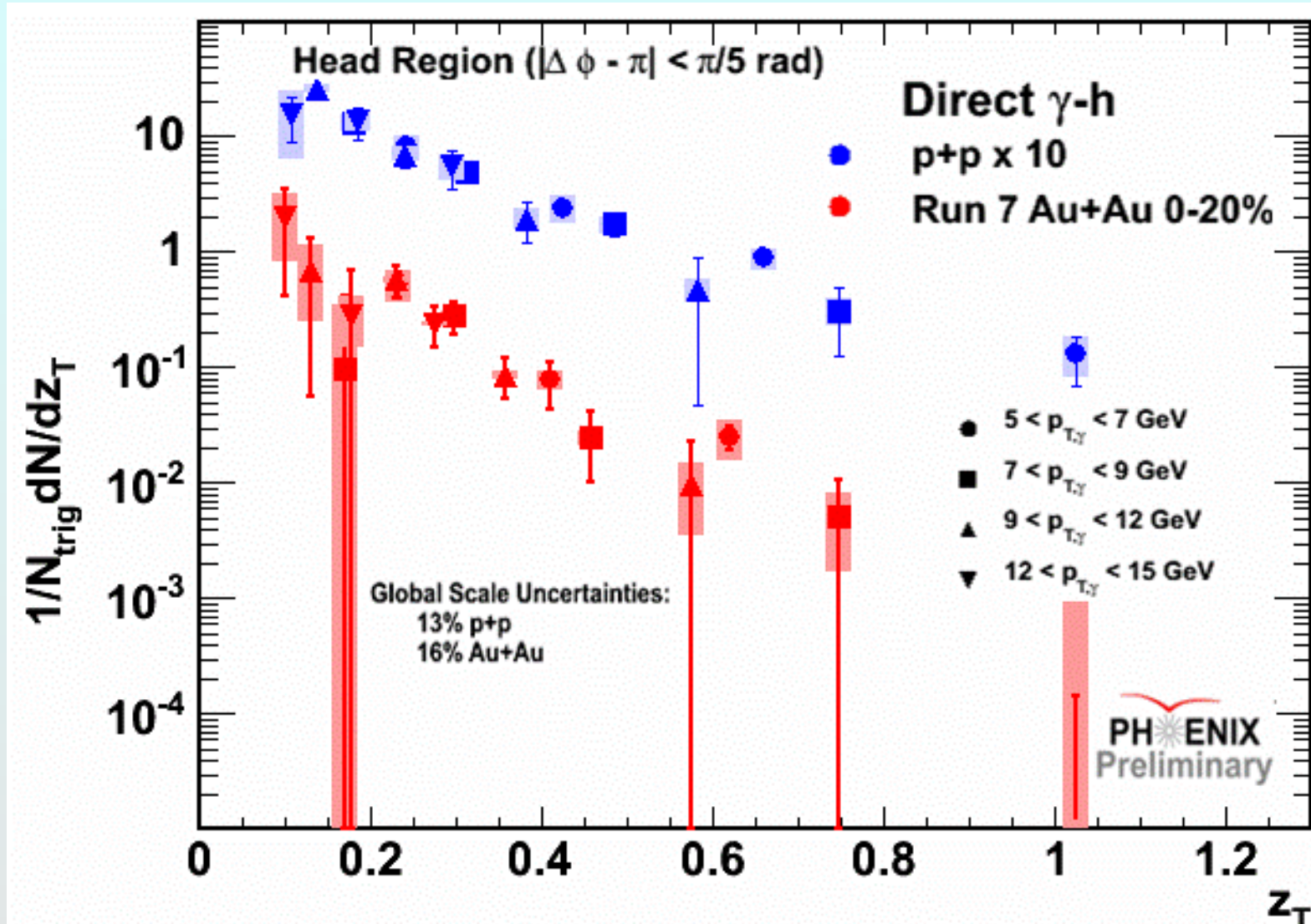
$$\xi = \ln(1/z_T)$$

Modification factor similar to R_{AA} :

**FFn experimental challenge:
measure the parton p
Use trigger γ or jet**

$$I_{AA} \equiv \frac{(1/N_{trig} dN/d\xi)_{AA}}{(1/N_{trig} dN/d\xi)_{pp}}$$

Modified in AA vs. p+p



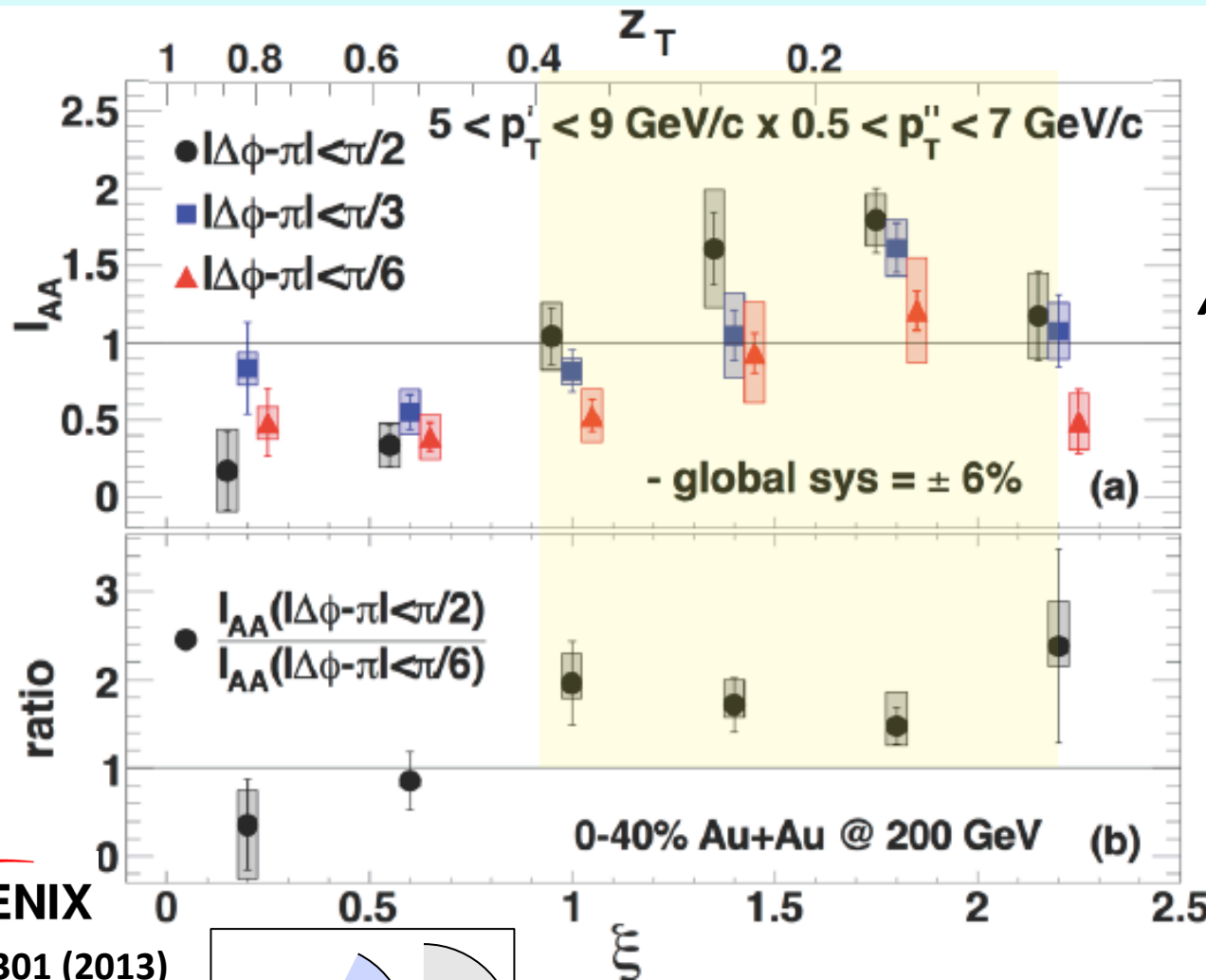
PHENIX: FFn via γ -h correlation

γ : parton energy, h : fragmentation fn.

“Extra” soft particles at larger angles near the away side jet

Provide constraints on gluon splitting

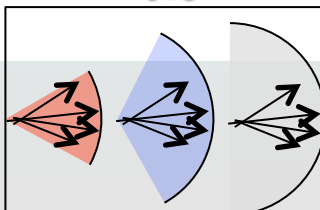
Perturbative?



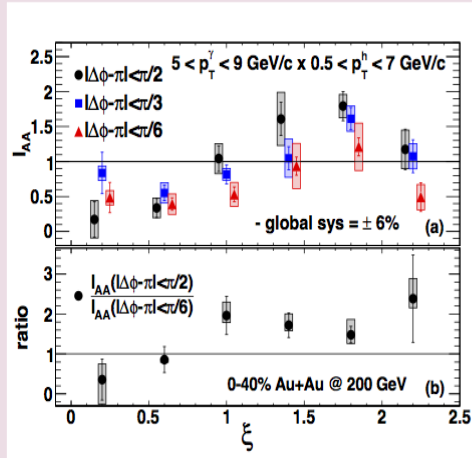
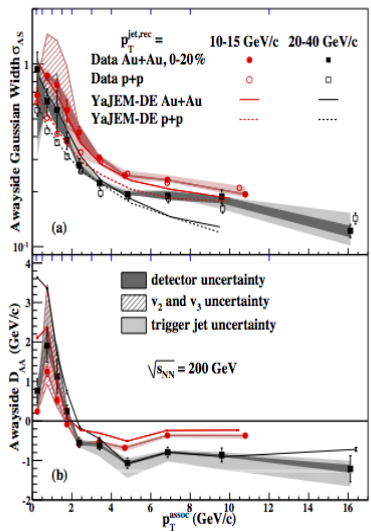
Au+Au/
p+p

PHENIX

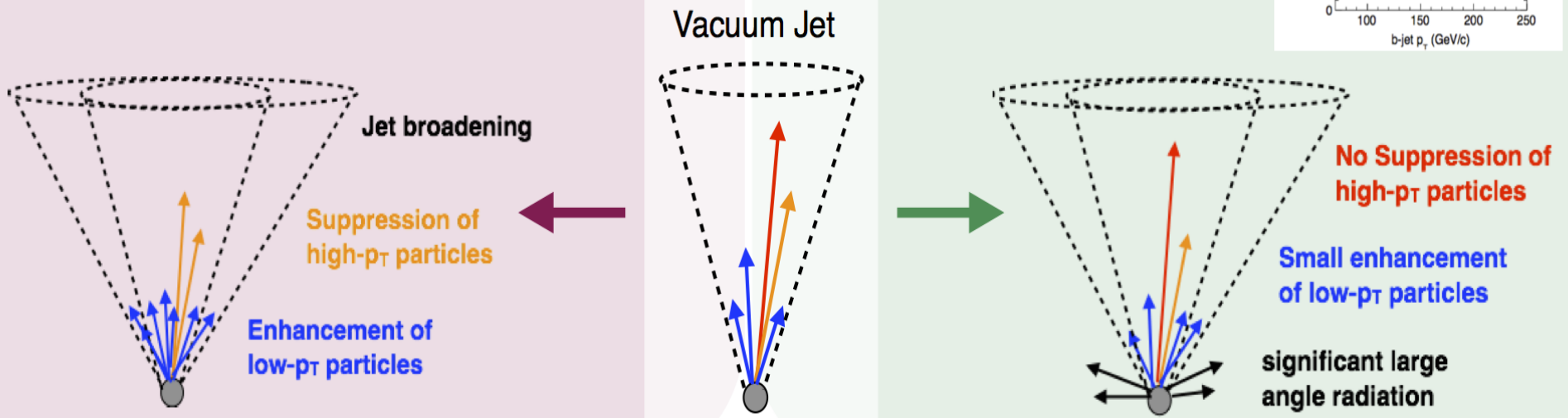
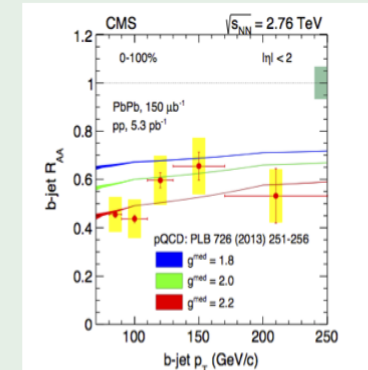
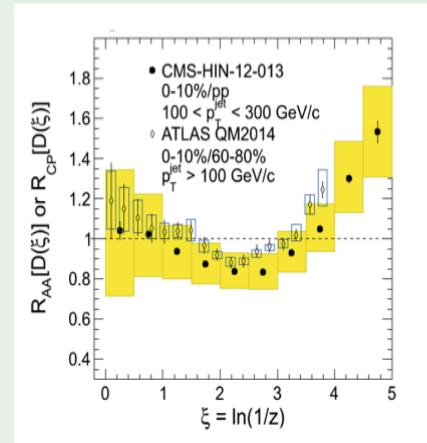
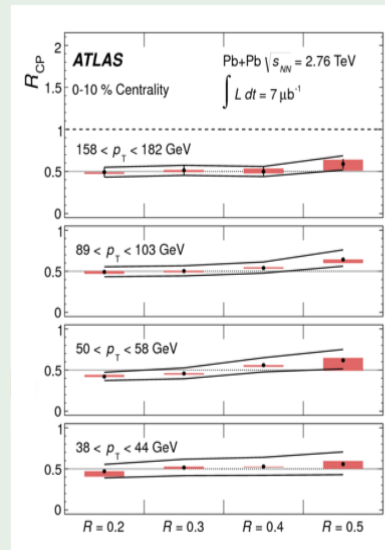
PRL 111, 032301 (2013)



Jets @ RHIC



Jets @ LHC



A jet is a partonic probe

It's not a photon...

- It can lose energy before the hard scattering
- It can lose energy after the hard scattering
- It can experience multiple scattering
- These may not be independent!

These are not just “background”

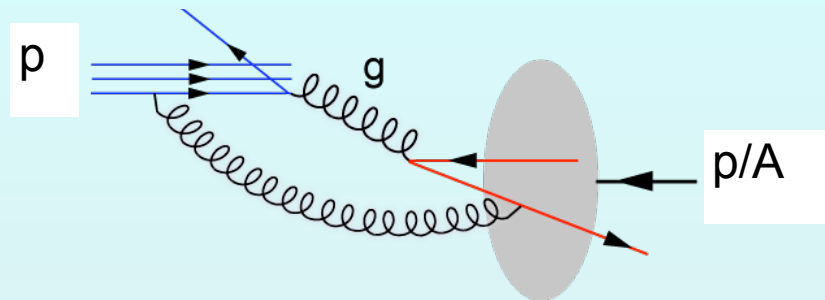
Higher-twist physics is interesting, not well understood,
as well as relevant to parton-plasma interaction !

Should measure γ -h correlations in cold nuclear matter!

this is underway in p+A!

Change the probe! pA vs. eA

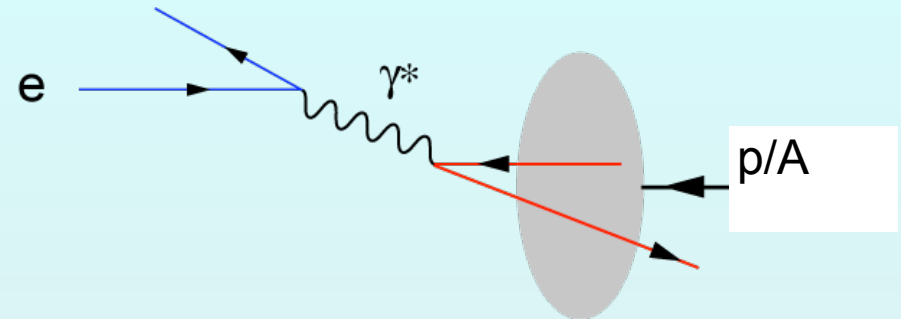
Hadron-Hadron



- More direct information on the response of a nuclear medium to gluon probe
- Probe has structure as complex as the “target”
- Soft color interactions before the collision can alter the nuclear wave function and destroy universality of parton properties (break factorization)

Electron-Hadron (DIS)

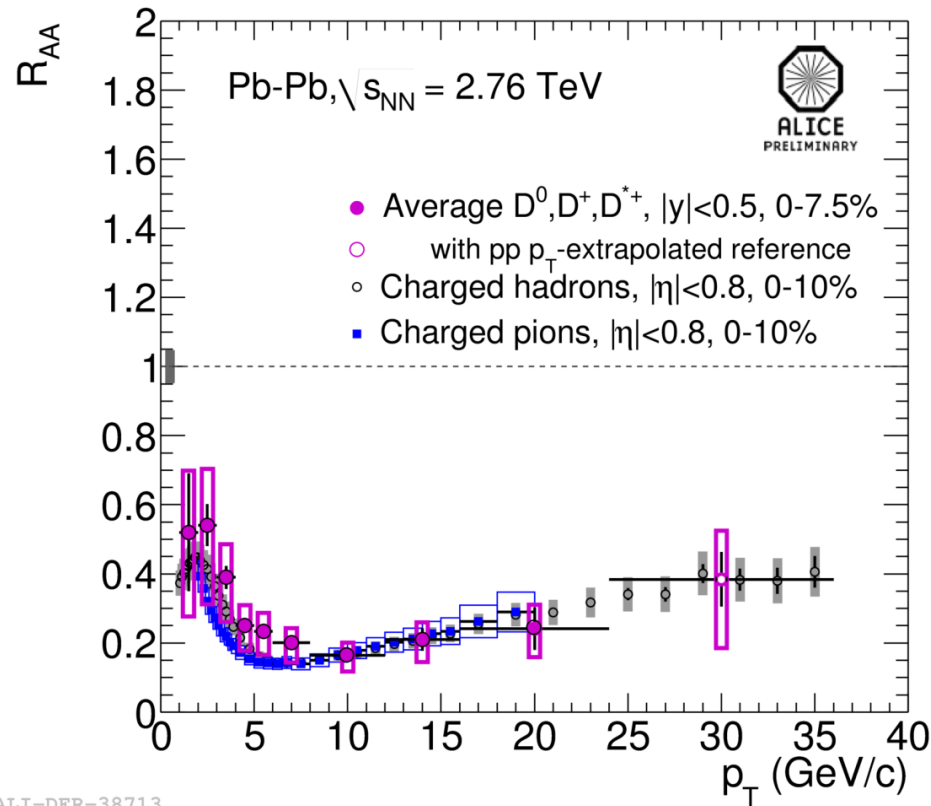
T. Ullrich



- Point-like probe
- Dominated by single photon exchange \Rightarrow no direct color interaction \Rightarrow preserve the properties of partons in the nuclear wave function
- High precision & access to partonic kinematics
- Nuclei always “cold” nuclear matter (CNM)

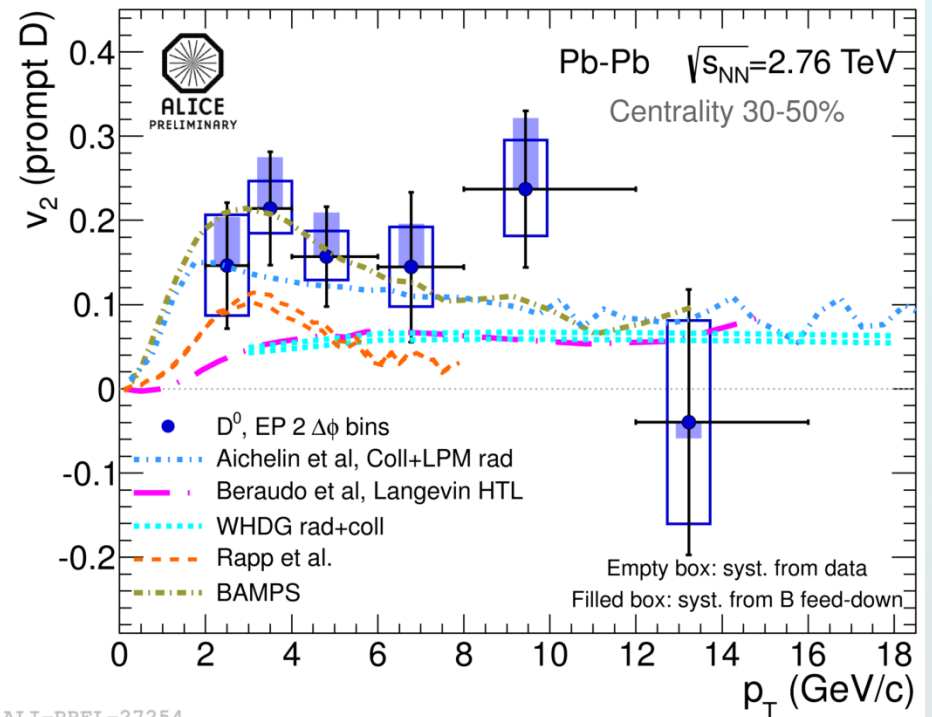
Heavy quark probes of QGP

Charm quarks also lose energy and flow!



ALI-DER-38713

- Charm quarks diffusing through strongly coupled QGP
- Collisions drive energy loss and push charm along with the bulk

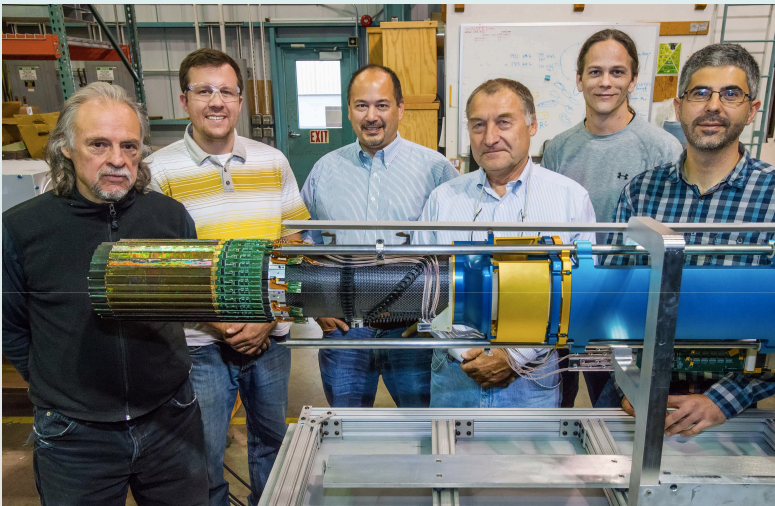
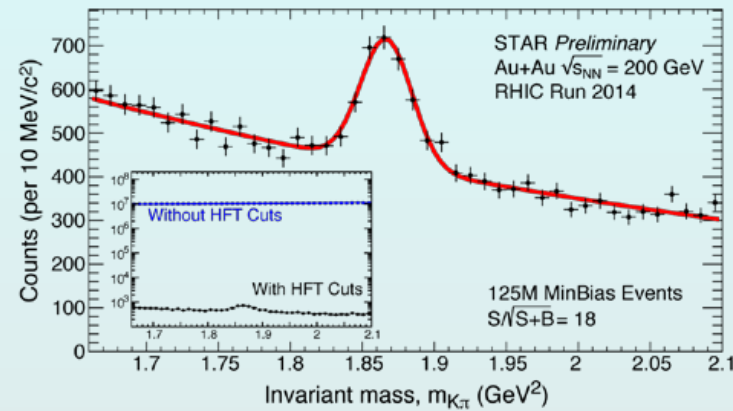
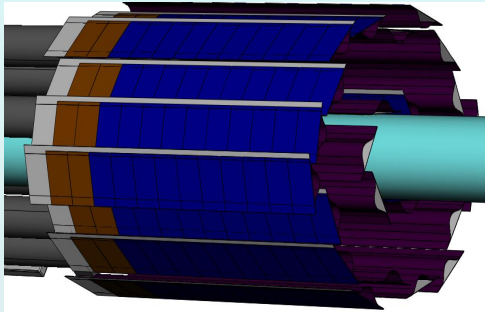


ALI-PREL-27254

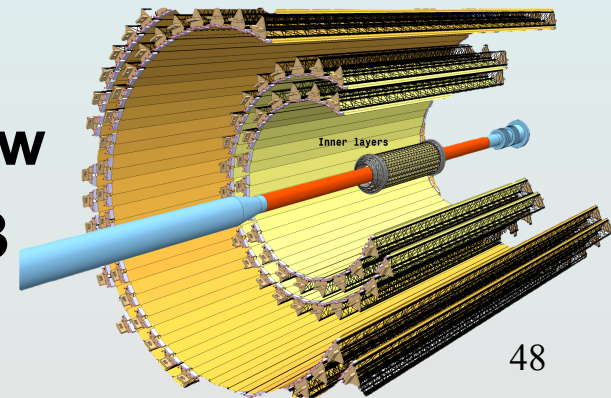
Upgrade ALICE for heavy quark probes in Run3

- Silicon detector arrays around beam pipe
 - Tag displaced vertex to separate c, b
 - Reconstruct D & B mesons from their decay hadrons

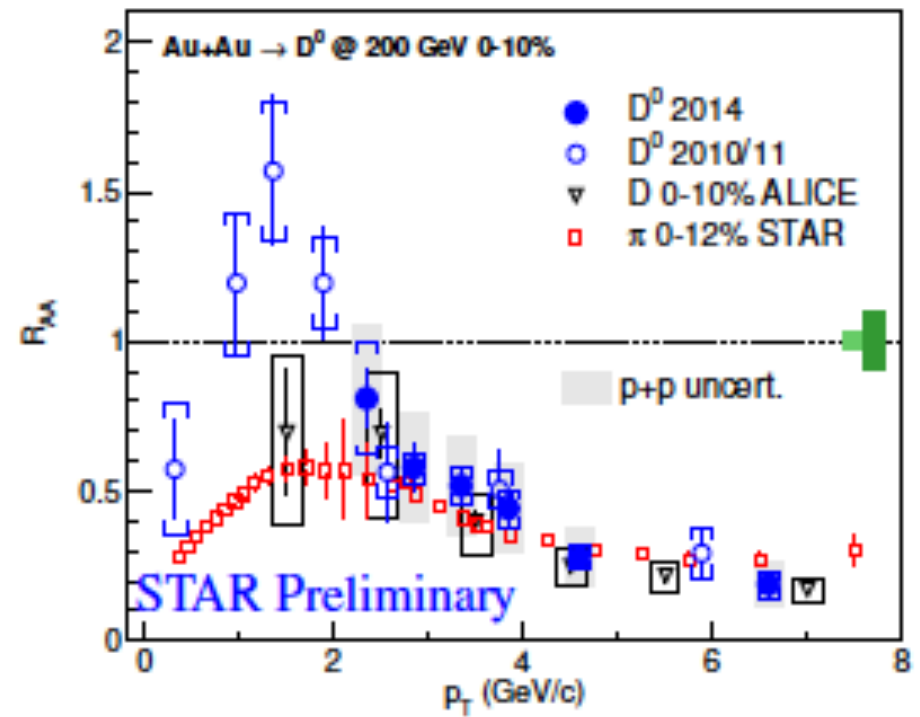
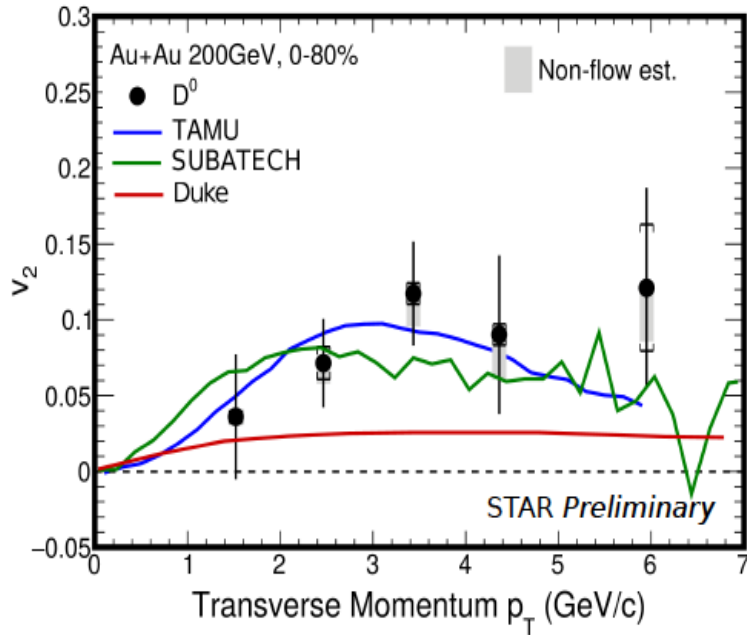
MAPS technology used for STAR at RHIC

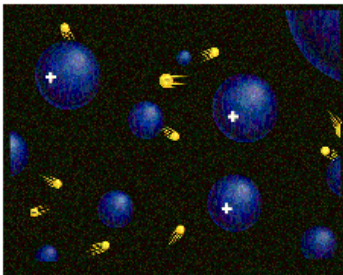


**ALICE
build now
for Run3
(2020)**



Flow and suppression as for light quarks





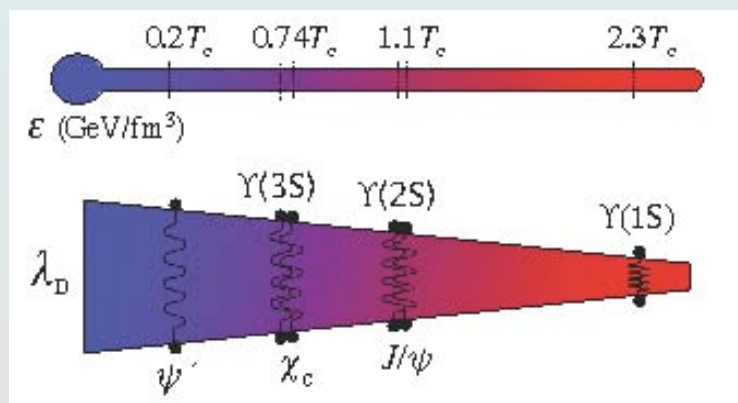
Is there a relevant screening length?

- **Plasma: interactions among charges of multiple particles spreads charge into characteristic (Debye) length, λ_D particles inside Debye sphere screen each other**
- **Strongly coupled plasmas: few ($\sim 1-2$) particles in Debye sphere Partial screening \rightarrow liquid-like properties sometimes even crystals!**
- **Test QGP screening with heavy quark bound states**

$c + \bar{c}$ and $b + \bar{b} : J/\psi$ & Y

Do they survive?

All? None? Some? Which size?



Color screening in quark gluon plasma

● q-qbar potential in vacuum: $V(r) = \sigma r - \frac{\alpha}{r}$

screened: $V(r) = -\frac{\alpha}{r} \times e^{(-r/\lambda_D)}$

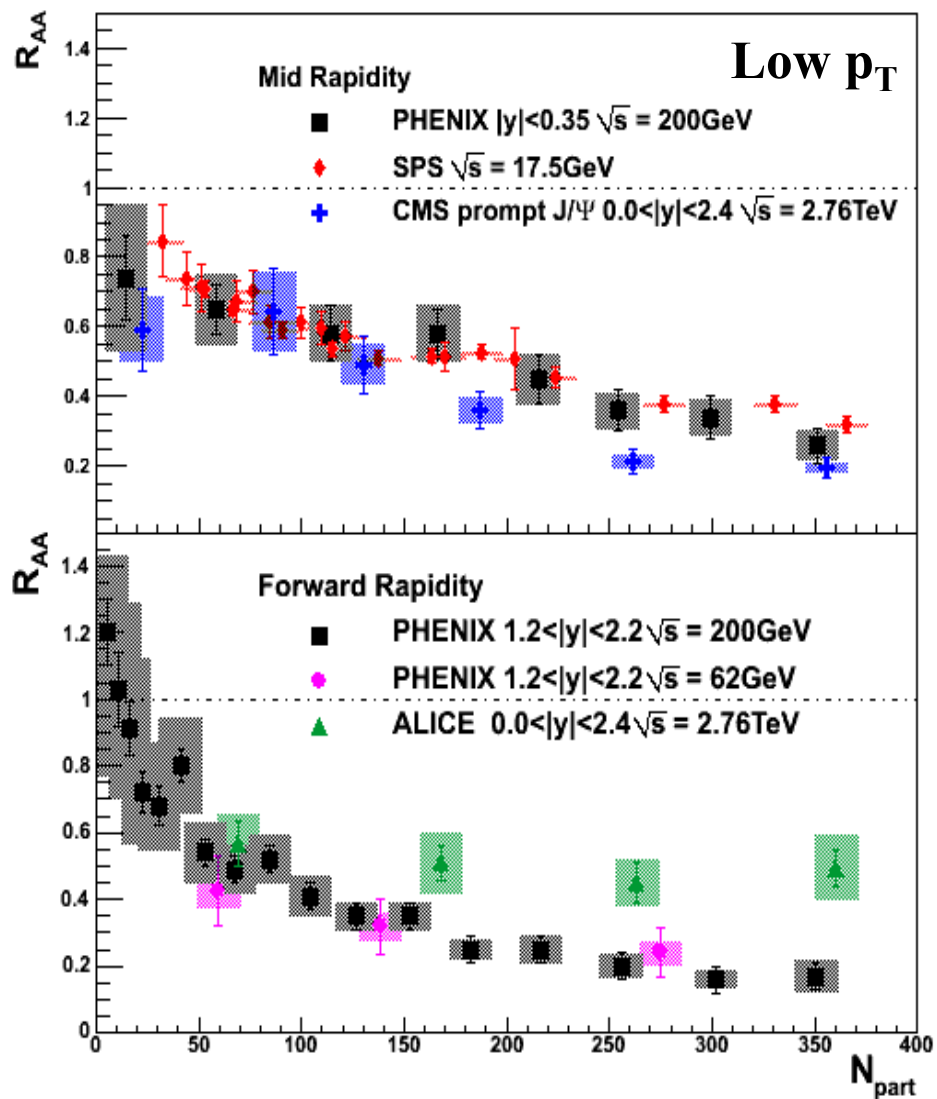
● Screening length in EM plasma: $\lambda_D = \sqrt{\frac{T}{8\pi\alpha\rho}}$

in (ideal) QGP, particle density $\rho \propto T^3$ so: $\lambda_D \sim \frac{1}{\sqrt{8\pi\alpha T}}$.

● For $\alpha_s \sim 0.2$, find:

Bound state	χ_c	ψ'	J/ ψ	$\Upsilon(2S)$	χ_b	$\Upsilon(1S)$
T_d	$\lesssim T_c$	$\lesssim T_c$	$\sim 1.2T_c$	$\sim 1.2T_c$	$\sim 1.3T_c$	$\sim 2.0T_c$

J/ψ vs. system size, √s



No clear suppression pattern
with \sqrt{s} , T!

Why more suppression at $y=2$?

Late break-up?

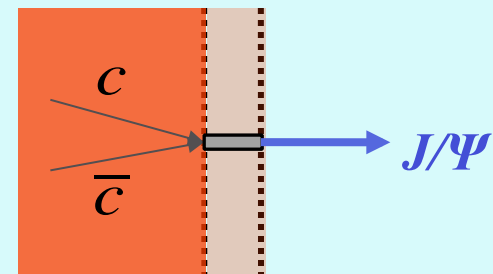
Final state coalescence of $q\bar{q}$?

J/ψ R_{AA} ~same from 17.5-200 GeV!

2.76 TeV direct J/ψ lower at mid-y, inclusive above at forward y

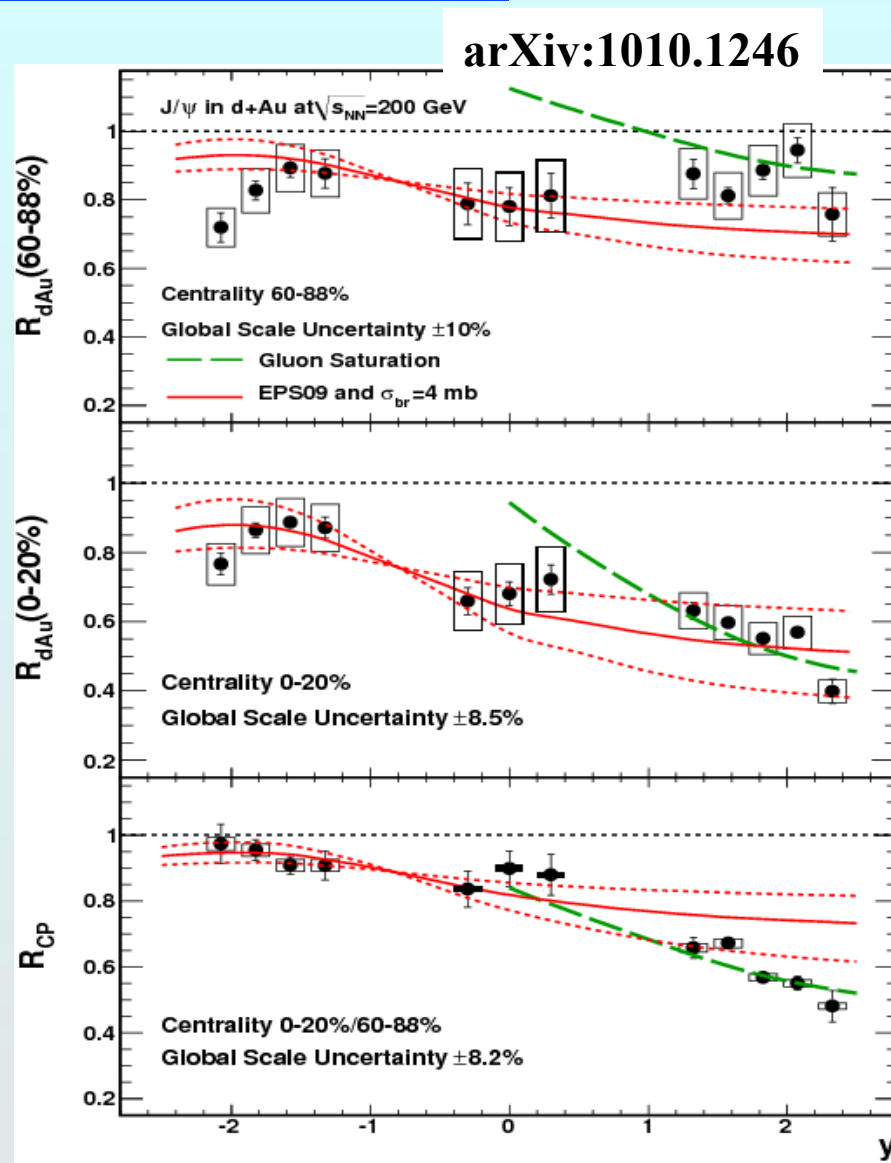
Lose some and gain some

- QGP screens primordial c - \bar{c} and b - \bar{b}
Melting depends on binding energy (i.e. size of bound state)
Look at J/ψ vs. ψ' vs. Υ states
- quarks find one another when system cools to $T \sim 150$ MeV
coalesce to form the hadrons we observe
- Occasionally a c can find a \bar{c}
“regeneration” or “coalescence”
Probability increases with c - \bar{c} pairs
Increases with beam energy
Confuses measurements of melting
- Measure energy & quark mass dependence to sort out
 - b 's are rare, so Υ is a good screening probe
- Initial state affects c & b production – so must measure $p+A$!

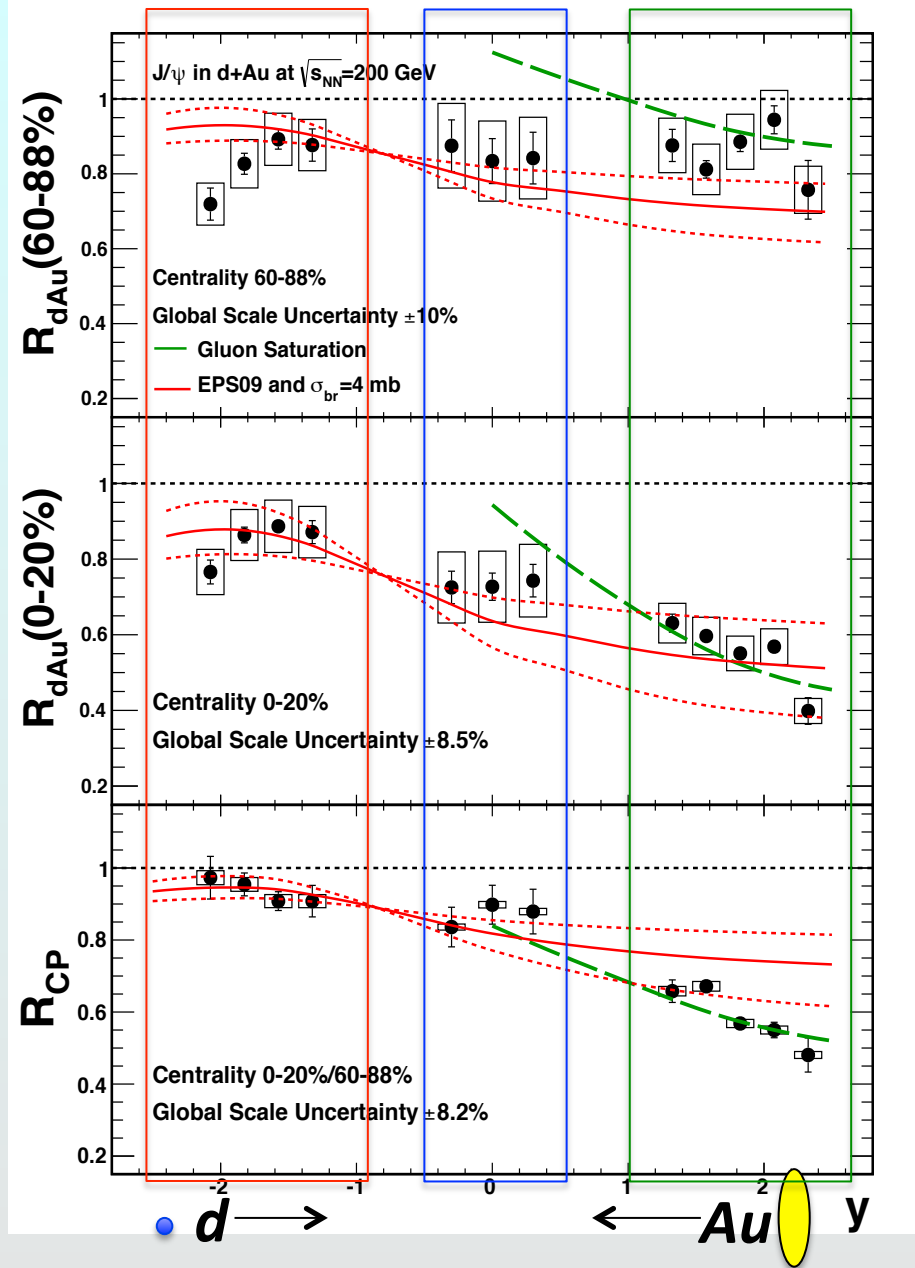
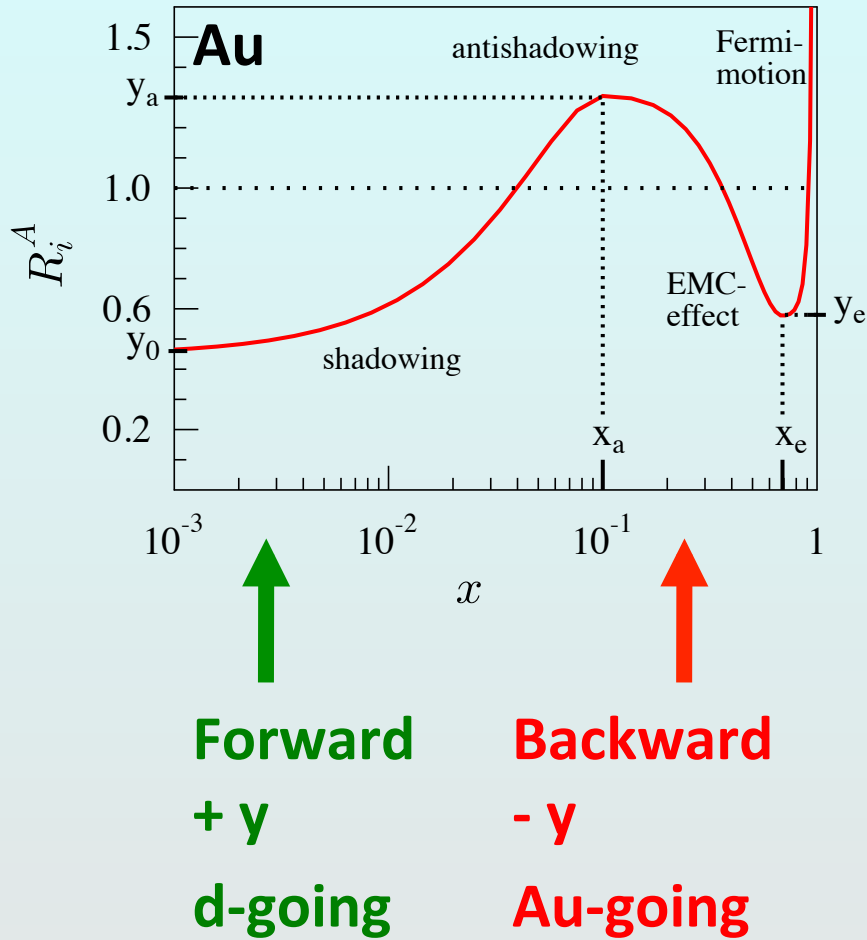


Suppression pattern ingredients

- Color screening
- Final state coalescence
- Initial state effects
 - Shadowing or saturation of incoming gluon distribution
 - Initial state energy loss (calibrate with p+A or e+A)
- Final state effects
 - Breakup of quarkonia due to co-moving hadrons (calibrate with A & centrality dependence)

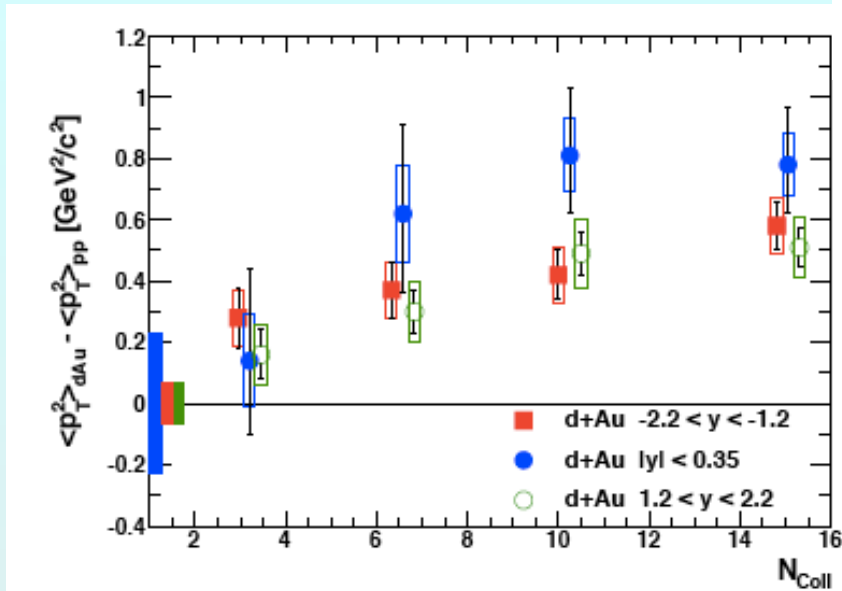


d+Au -> J/ψ from PHENIX

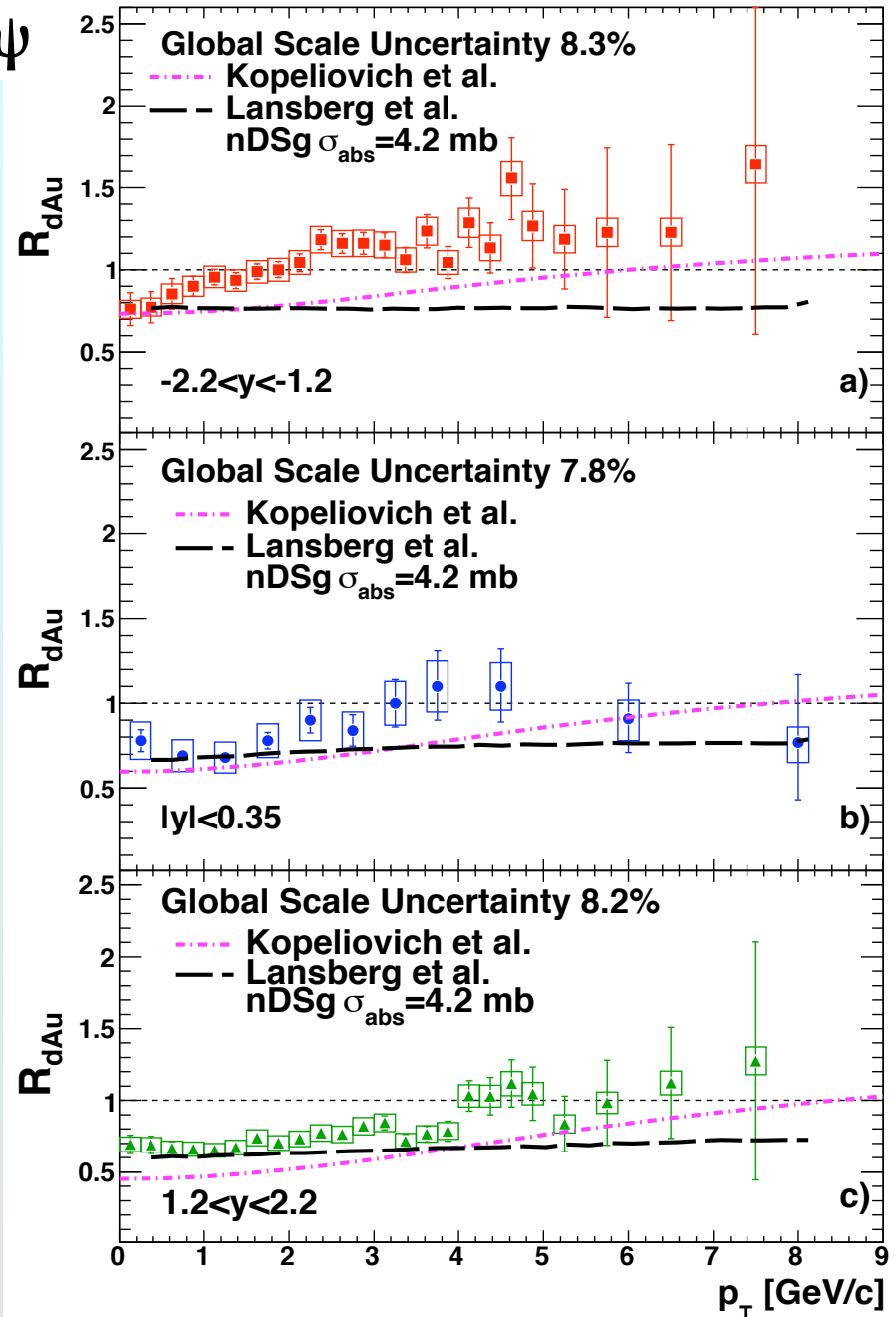


Shadowing, breakup & Cronin effect

PRC87, 034911 (2013)



J/ψ



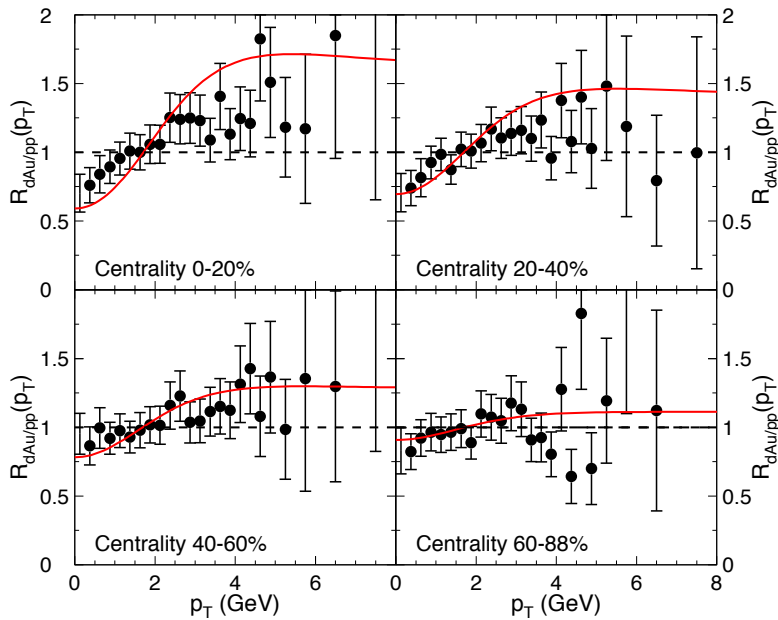
- ✦ p_T broadens (multiple scattering) w/ N_{coll} ; effect stronger at $y=0$
- ✦ J/ψ suppressed to higher p_T @ mid & forward y (lower x in Au);
- ✦ $R_{dA} > 1$ at high p_T backward (Cronin effect in Au nucleus)
- ✦ $p_T, y, \text{centrality}$ dependence was not reproduced by the models

d+Au -> J/ψ

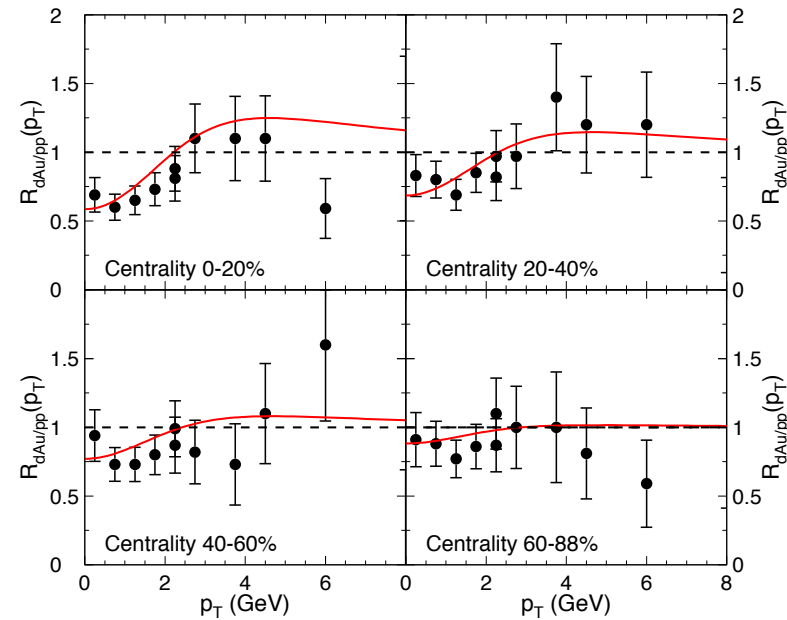
but

Arleo, et al 1304.090

$y = [-2.2 ; -1.2]$



$y = [-0.35 ; 0.35]$

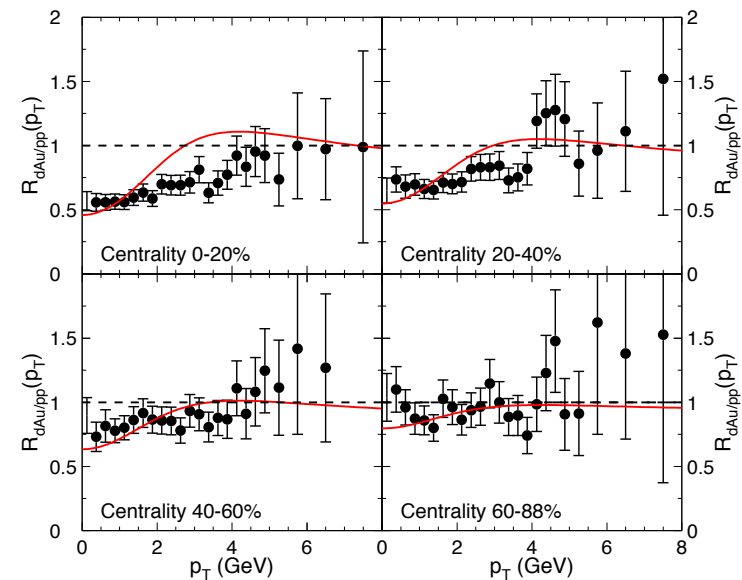


coherent parton energy loss and p_T broadening from multiple scattering in the nucleus is consistent with data!

$$\hat{q}_0 = 0.075 \text{ GeV}^2/\text{fm}$$

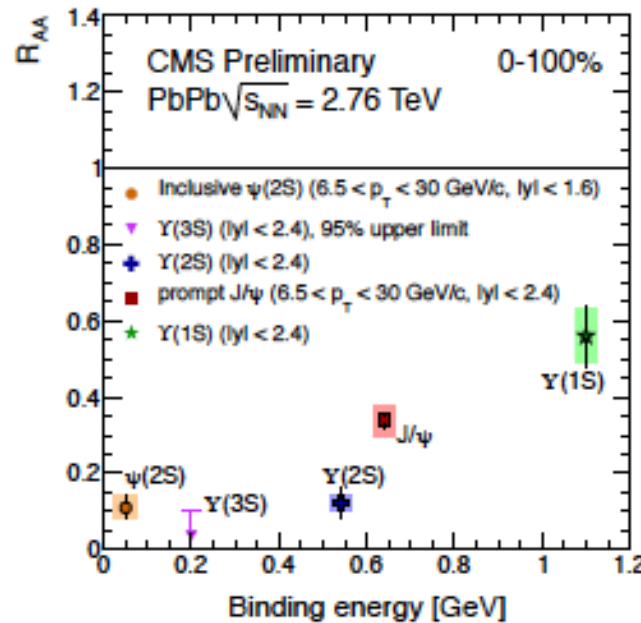
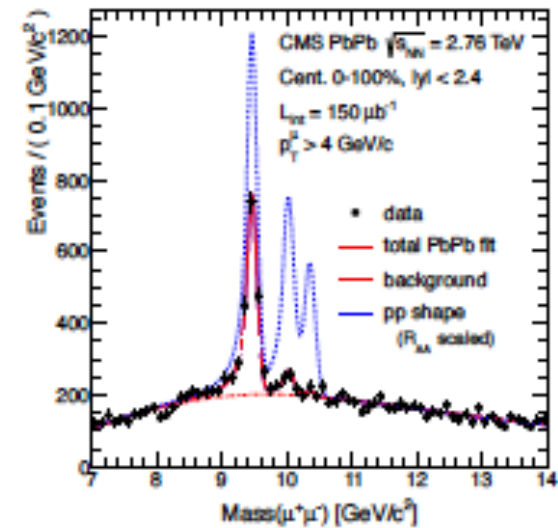
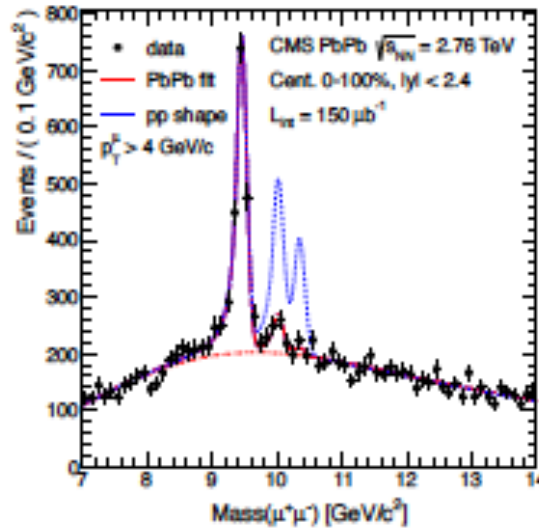
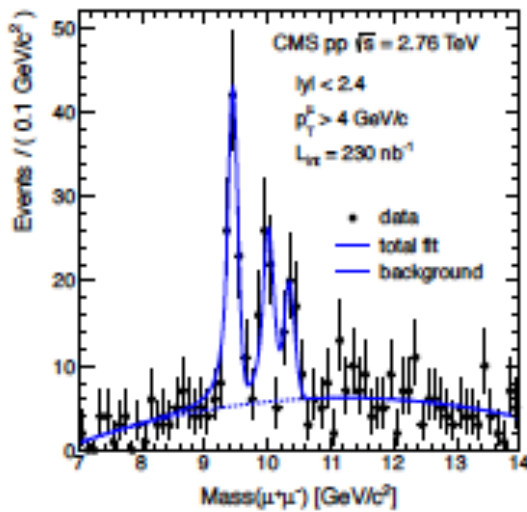
Dynamics of the probe & structure of the medium mix!!

$y = [1.2 ; 2.2]$



b-bar bound states at the LHC

- Many c-cbar at LHC: coalescence important. Use b-bar probe

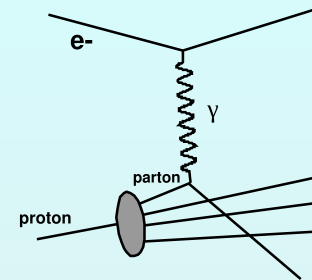


Y (2S,3S)
suppressed

Probe nucleons & nuclei with electrons

- How many gluons are there?

Measure $e+p, e+A \rightarrow e$



- How are they distributed?

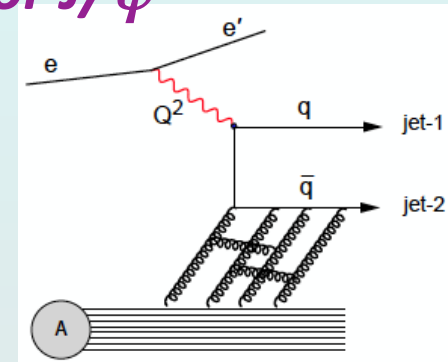
In space? In momentum?

Measure $e+p, e+A \rightarrow e + \text{hadrons}, e + \gamma$ or J/ψ

- How are they correlated inside nucleus?

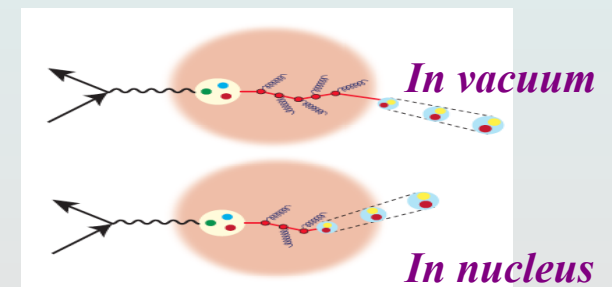
Measure $e+p, e+A \rightarrow$

$e + \geq 2 \text{ hadrons}$



- What's gluon range inside a nucleus?

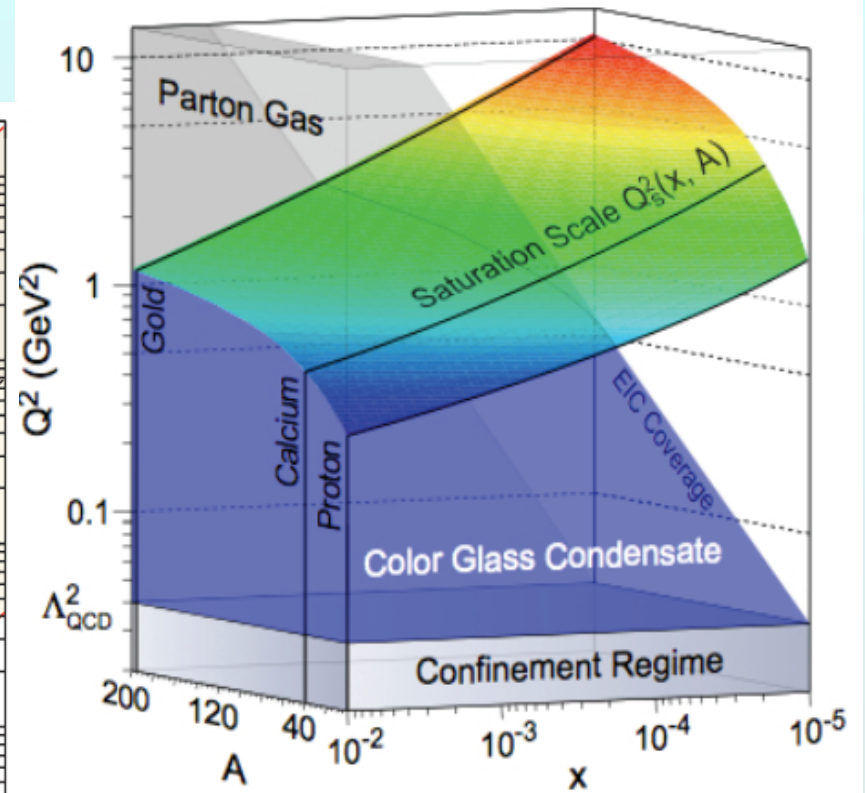
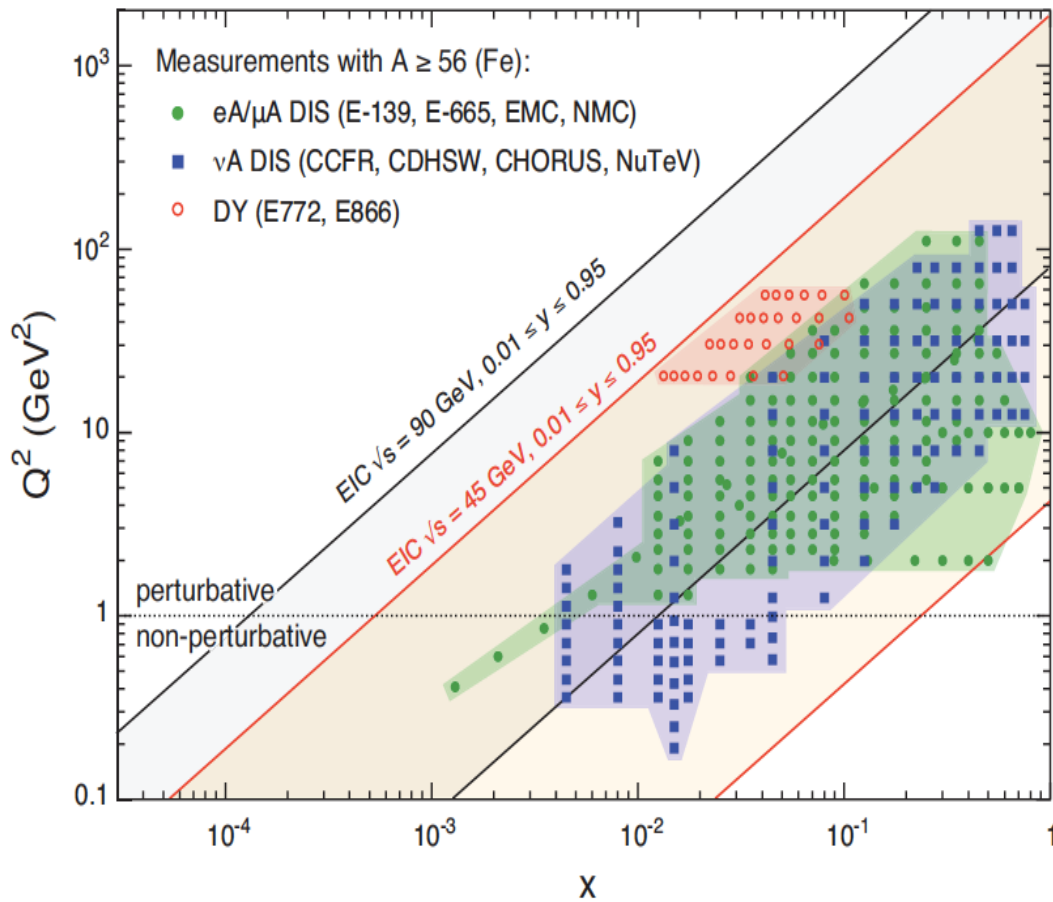
Measure hadron production



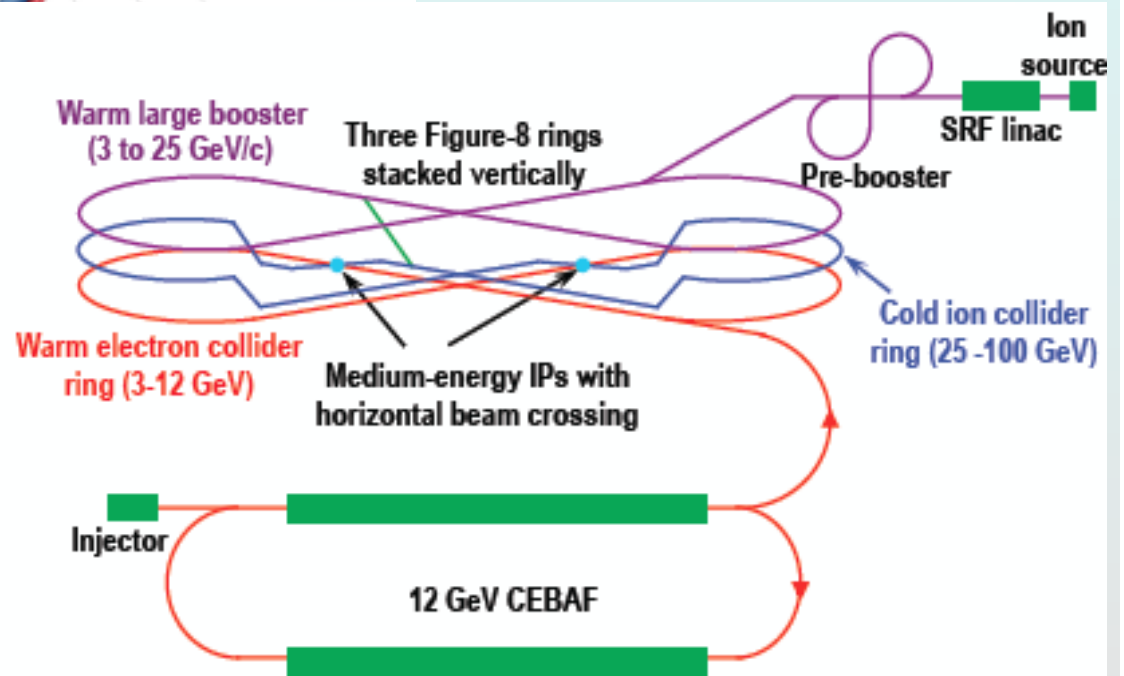
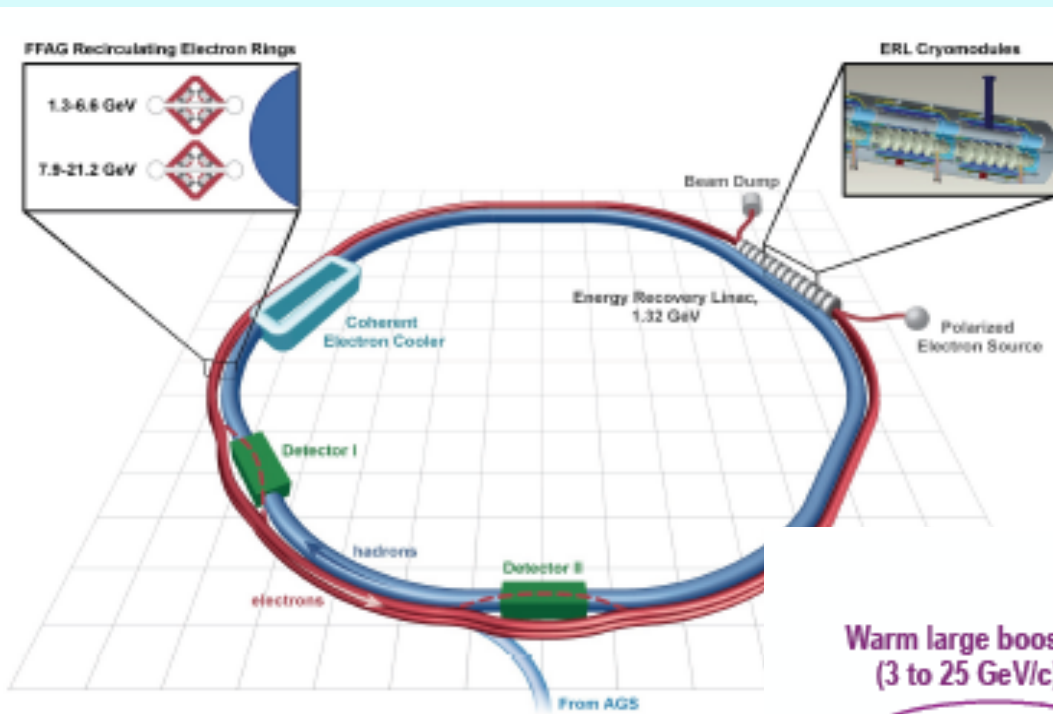
EIC kinematic reach

An electron-ion collider
probe structure directly

Enhance Q_s with A , not energy



Electron-ion collider



● **backup slides**

Questions

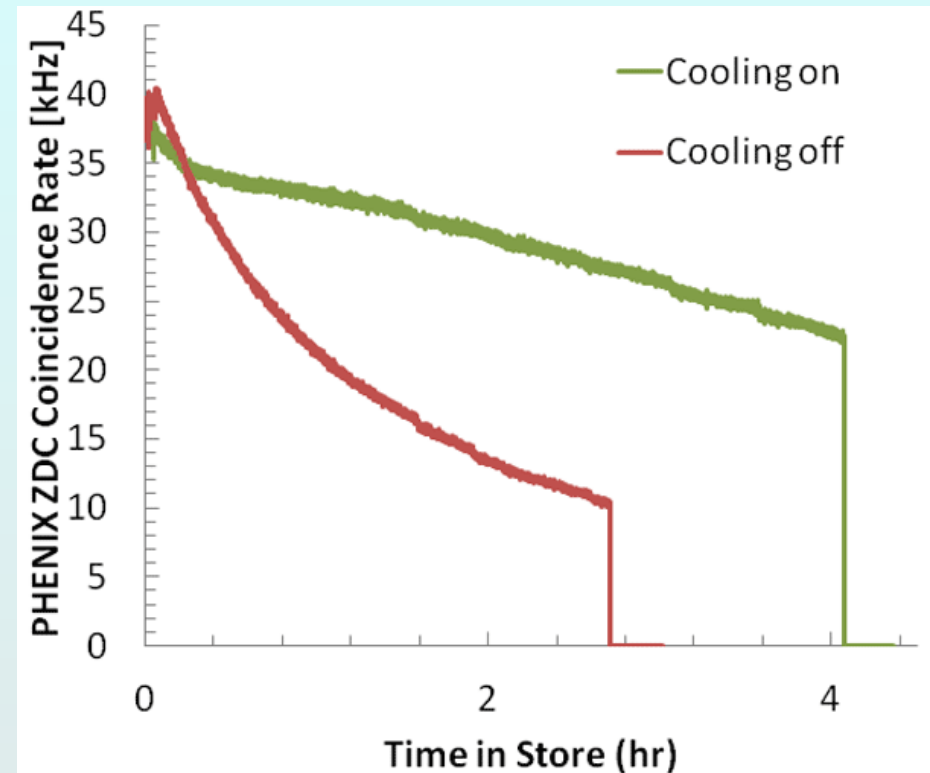
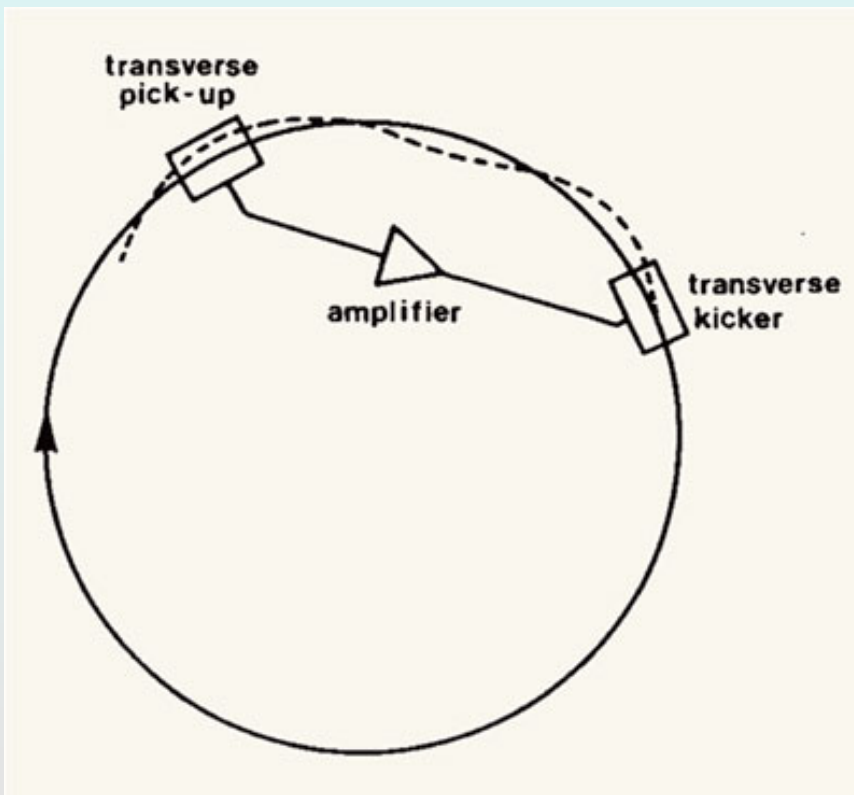
- Can we create quark gluon plasma in the lab?
- thermodynamic properties (equilibrium)
T, P, ρ
Equation Of State (relation btwn T, P, V, energy density)
 v_{sound} , static screening length
- transport properties (non-equilibrium)*
particle number, energy, momentum, charge
diffusion *sound* *viscosity* *conductivity*

**In plasma: interactions among charges of multiple particles
charge is spread, screened in characteristic (Debye) length, λ_D
*also the case for strong, rather than EM force***

***measuring these is new for nuclear/particle physics!**

Stochastic cooling keeps bunches tight

- ◆ *Coulomb interactions among highly charged ions blow up the beam bunches*
- ◆ *Measure and send correcting kick to outliers (at each go-round)*

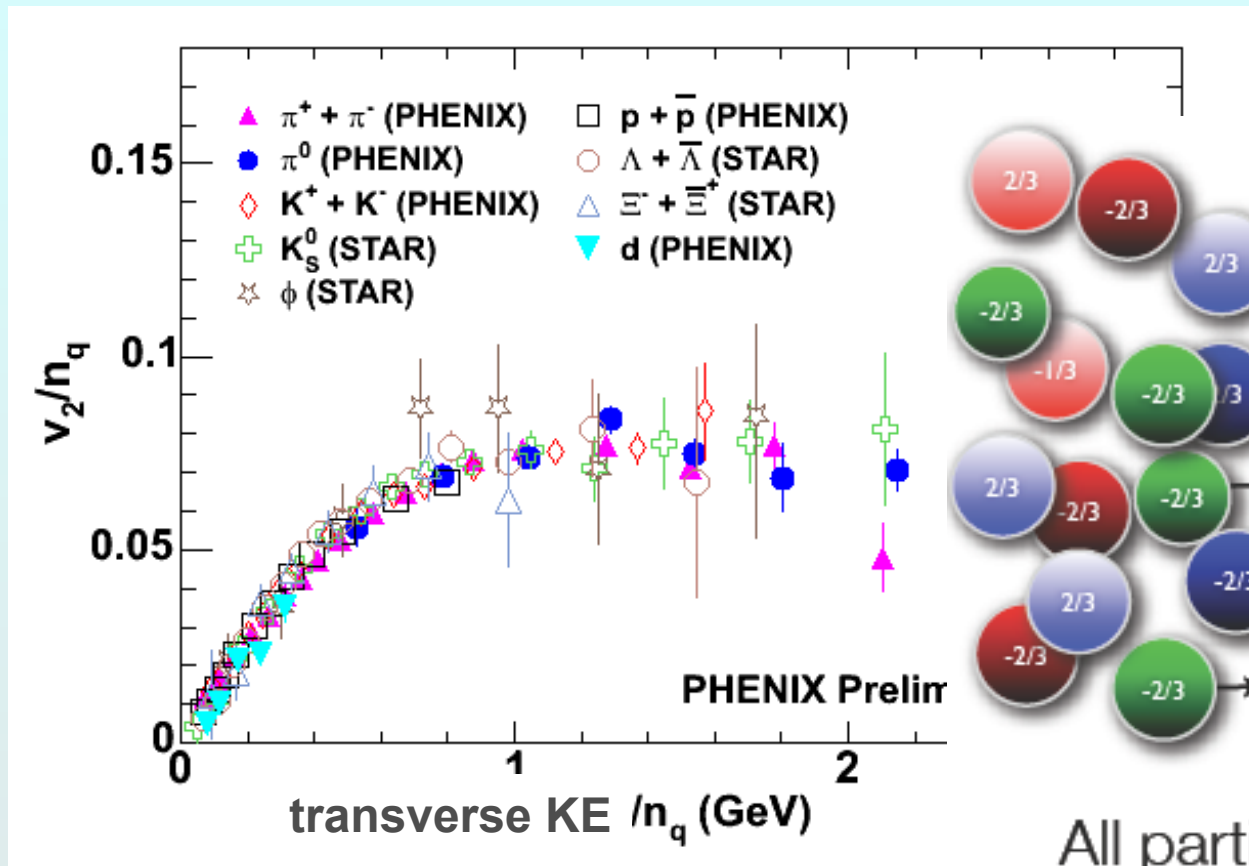


➤ *Huge increase in beam lifetime and luminosity!*

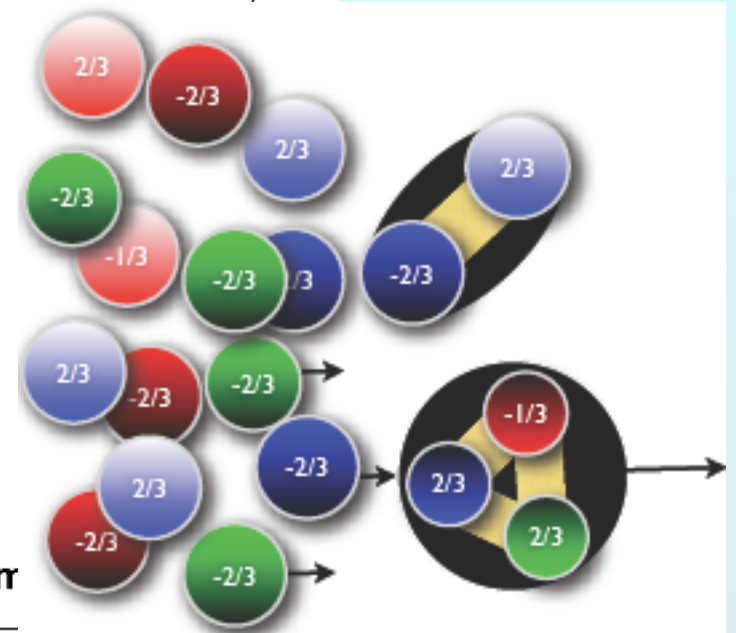
Mechanism for fast thermalization?

- **Must be thermalized in < 1 fm/c!**
Otherwise v_2 smaller than in the data
- **Can this be achieved with gg, qg, and qq binary scatterings?**
NO!
Making this picture yield sufficient v_2 , requires boosting the pQCD parton-parton cross sections by a factor of ~ 50 !
- **Mechanism not known yet. Perhaps start out with correlated gluonic matter at small x ?**

Elliptic flow scales with number of quarks



implication: valence quarks, not hadron pressure builds early, dressed quarks are

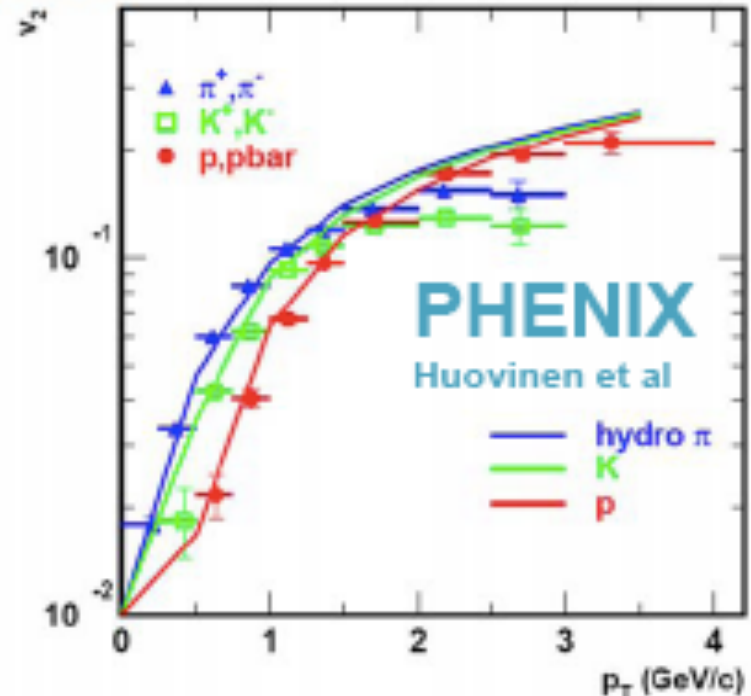
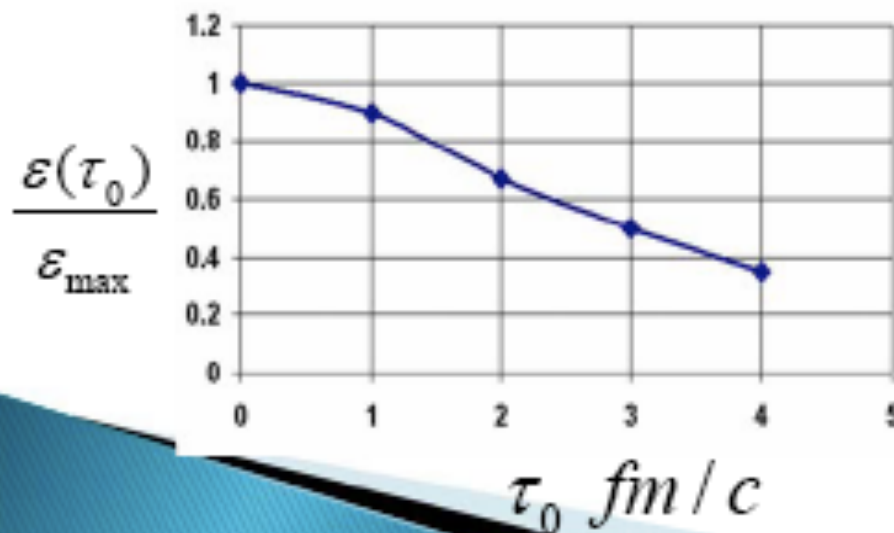


All particles flow as if frozen out from a flowing soup of constituent quarks

Strong flow Implies *early* thermalization

- ▶ If system free streams
 - spatial anisotropy is lost
 - v_2 is not developed

Eccentricity: $\varepsilon = \frac{\langle y^2 \rangle - \langle x^2 \rangle}{\langle y^2 \rangle + \langle x^2 \rangle}$



detailed hydro calculations
(QGP+mixed+RG, zero viscosity)

- $\tau_0 \sim 0.6 - 1.0$ fm/c
- $\varepsilon \sim 15 - 25$ GeV/fm³
- (ref: cold matter 0.16 GeV/fm³)

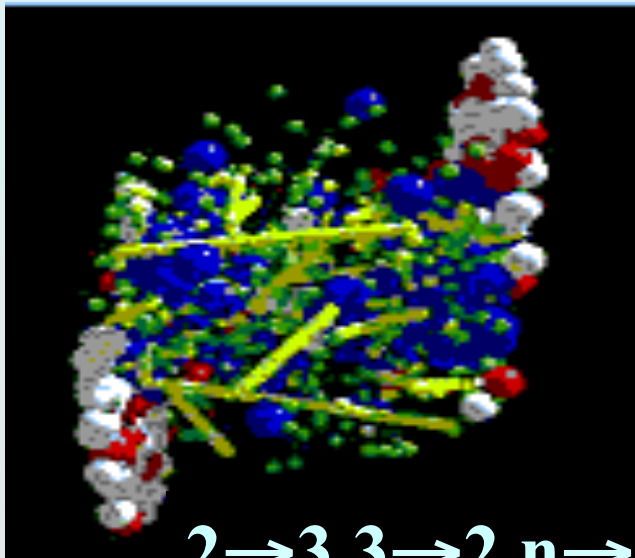
Calculating transport in QGP

weak coupling limit

perturbative QCD

kinetic theory, cascades

interaction of particles



$2 \rightarrow 3, 3 \rightarrow 2, n \rightarrow 2 \dots$

∞ strong coupling limit

not easy! Try a pure field...

gravity \leftrightarrow supersym 4-d

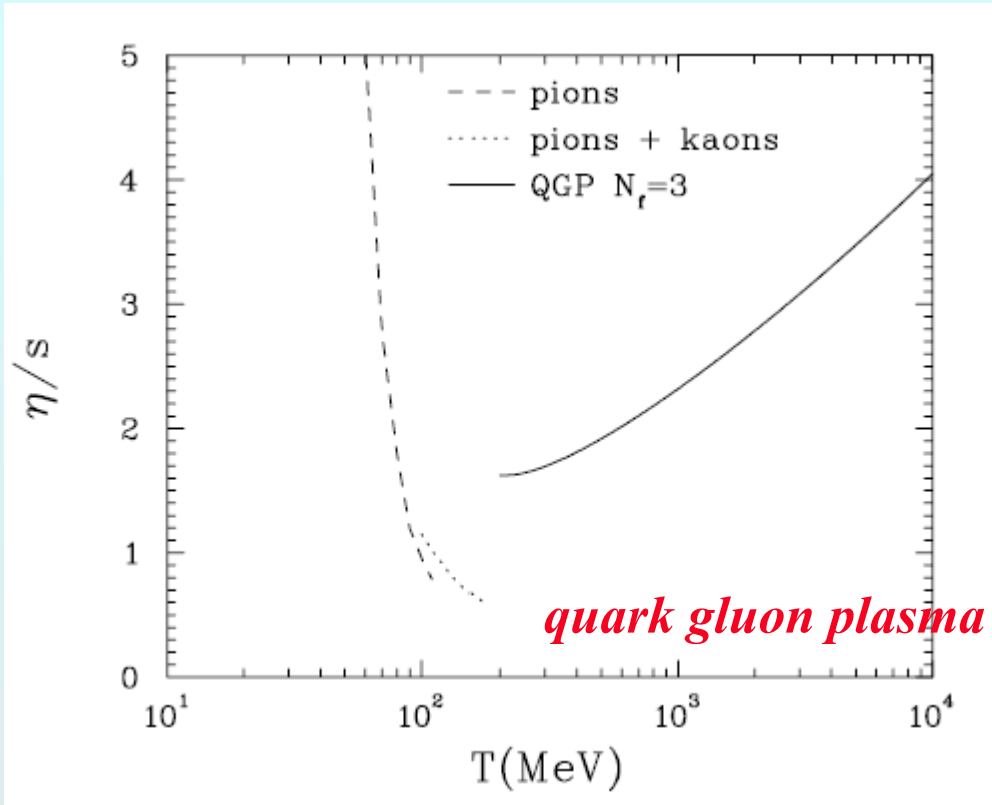
(AdS/CFT)



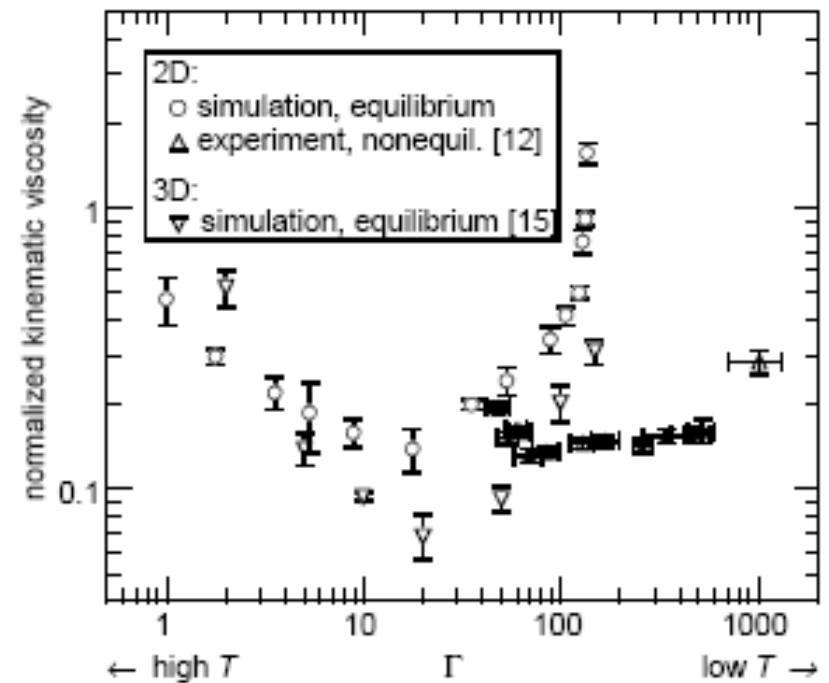
minimum η at phase boundary?

strongly coupled dusty plasma

B. Liu and J. Goree,

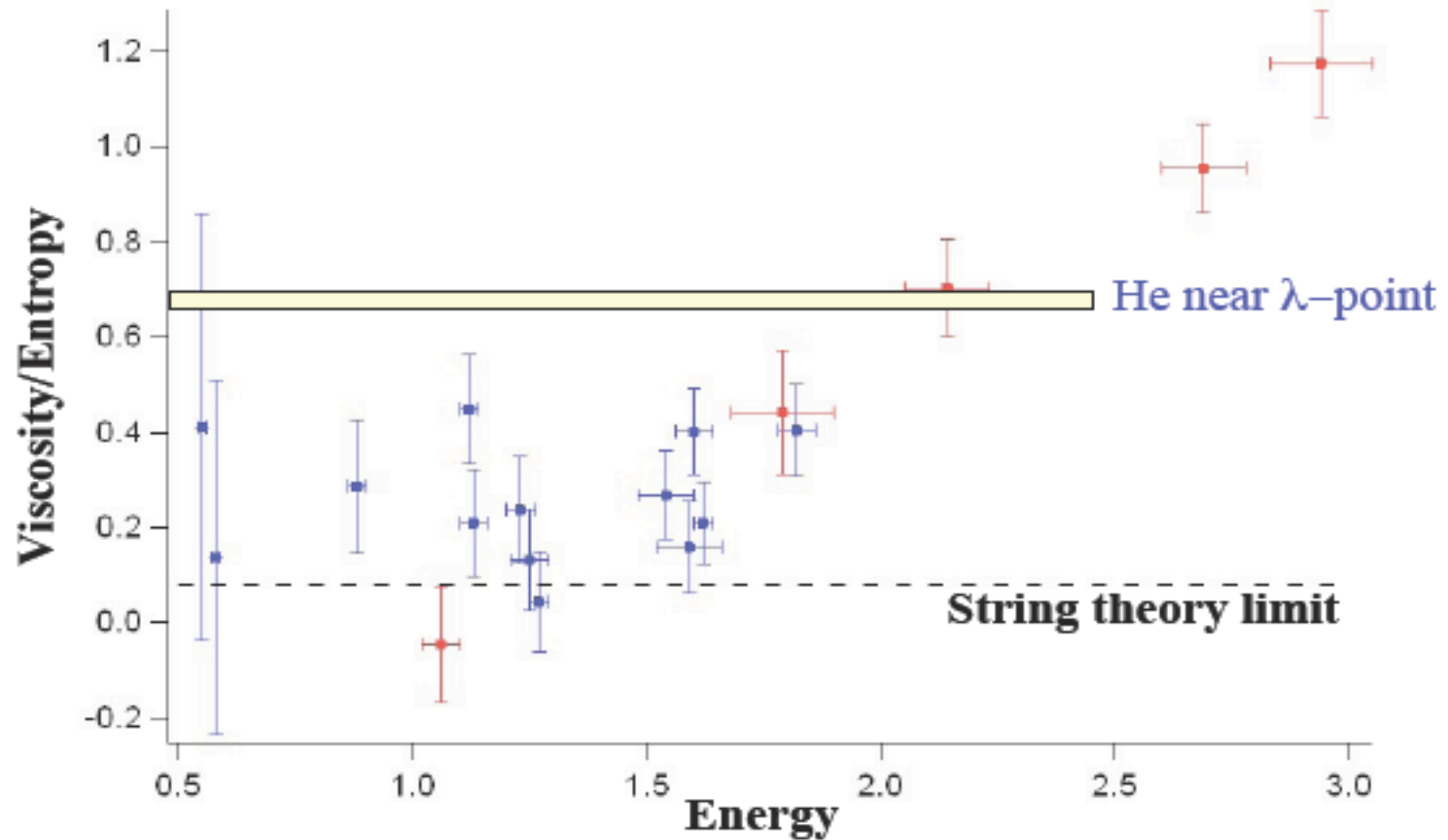


*Csernai, Kapusta & McLerran
PRL97, 152303 (2006)*



**minimum observed in other strongly coupled systems –
kinetic part of η decreases with Γ while potential part increases**

Viscosity/Entropy (natural units)

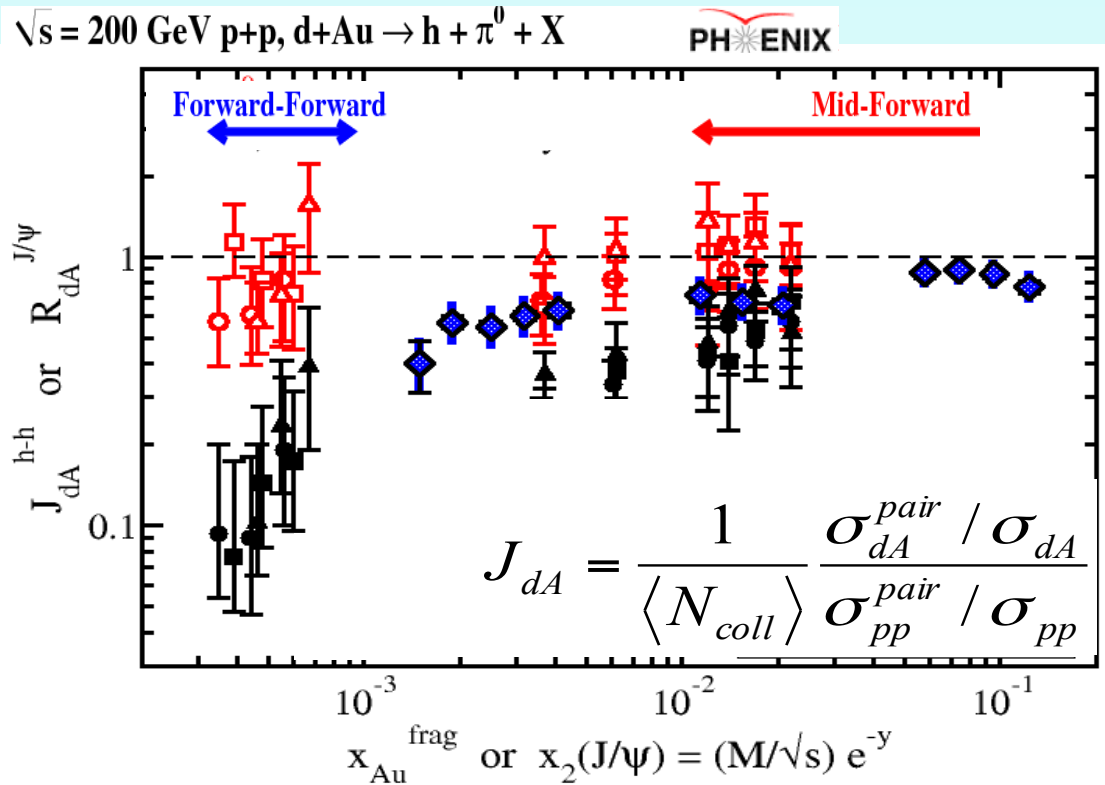
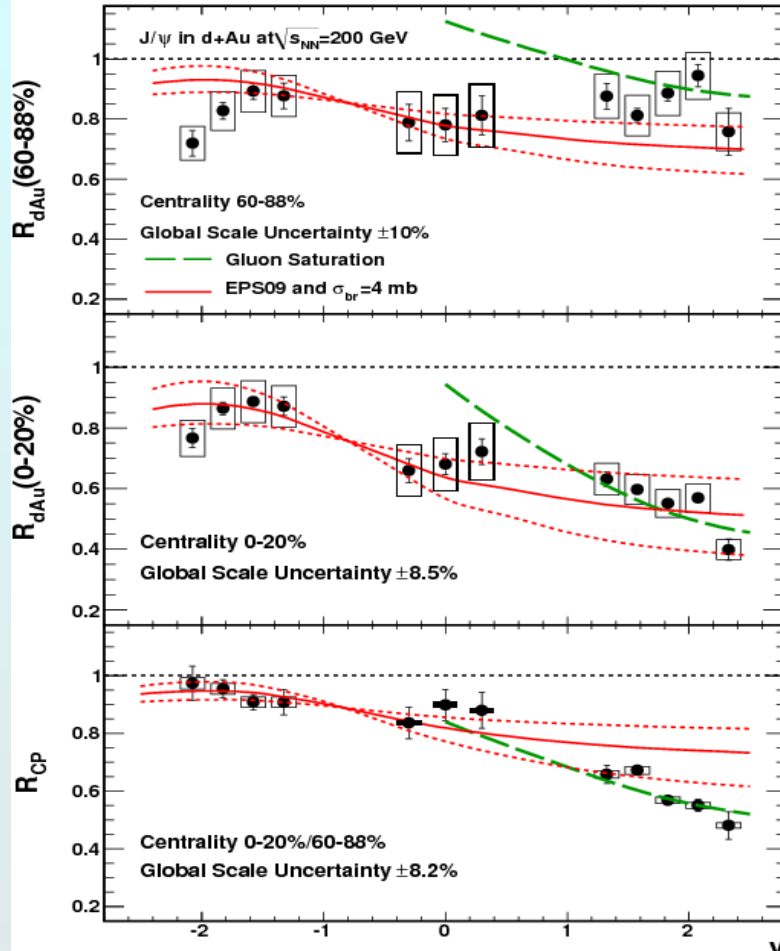


Dense gluonic matter (d+Au, forward γ):

large effects observed

arXiv:1010.1246

arXiv:1105.5112



Shadowing/absorption stronger than linear w/nuclear thickness

Di-hadron suppression at low x
pocket formula (for $2 \rightarrow 2$):

$$x_{Au}^{frag} = \frac{\langle p_{T1} \rangle e^{-\langle \eta_1 \rangle} + \langle p_{T2} \rangle e^{-\langle \eta_2 \rangle}}{\sqrt{s}}$$

trend as, e.g. in CGC ...

More jet probes = more insight

- **Hadrons**

Single high p_T hadrons (leading jet fragments)
 di-hadron correlations

- **Reconstructed jets** (*depend on reco algorithm, compare to theory*)

Single jets
 <di-jets> or jet-h correlations

- **Gamma-jet correlations** (*photon tags jet energy*)

γ -h correlations
 < γ -reconstructed jet>

- **Construct the variables: R_{AA} , I_{AA} , A_J , q -hat**

Nuclear modification:

$$R_{AA}(p_T) = \frac{d^2 N^{AA} / dp_T d\eta}{T_{AA} d^2 \sigma^{NN} / dp_T d\eta} \quad I_{AA} \equiv \frac{(1/N_{trig} dN/d\xi)_{AA}}{(1/N_{trig} dN/d\xi)_{pp}}$$

Jet asymmetry:

$$A_J = \frac{E_{T1} - E_{T2}}{E_{T1} + E_{T2}}, \Delta\phi > \frac{\pi}{2}$$

Jet transport coefficient: $q = \hat{\mu}^2 / \lambda$; $\mu = \langle p_T \text{ transfer} \rangle$ in 1 scattering

