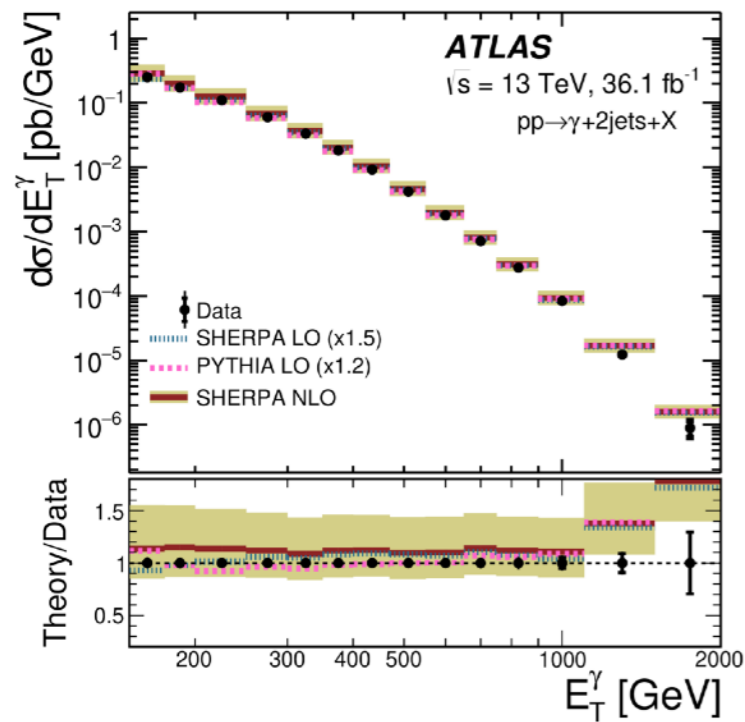


# Resummation for Photon Isolation

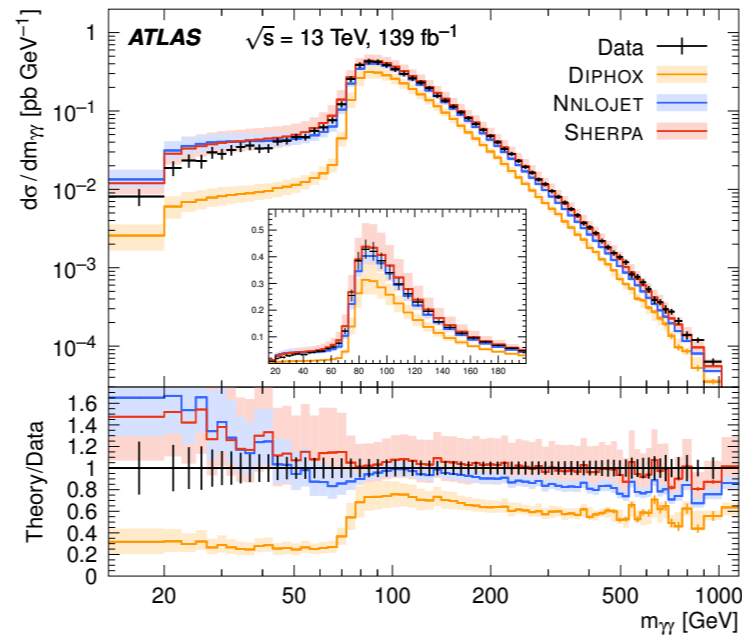
Thomas Becher  
University of Bern

*JHEP* 01 (2023) 005 with Samuel Favrod and Xiaofeng Xu

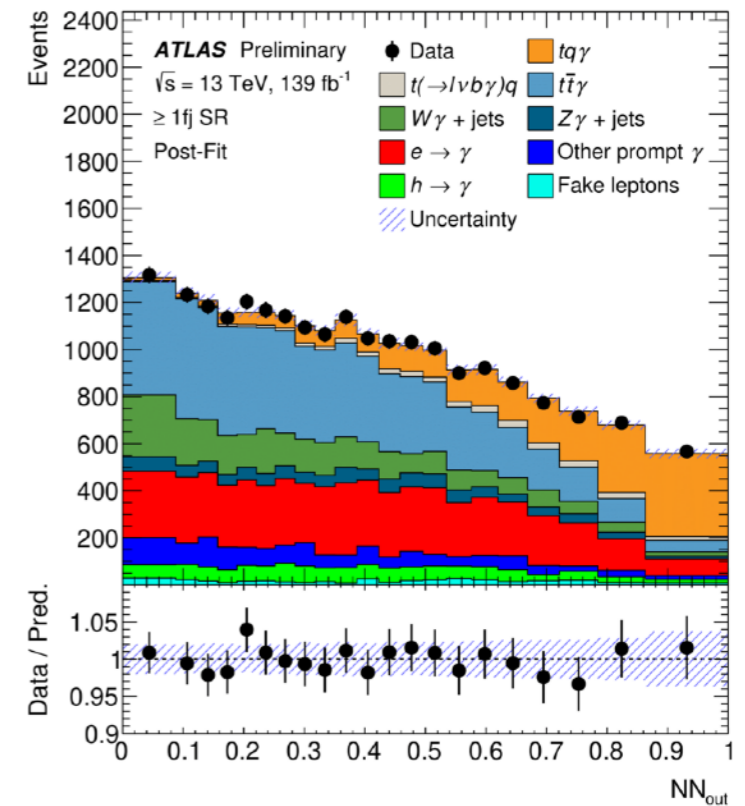
SCET 2023, Berkeley  
March 30, 2023



photon + 2 jets



diphotons



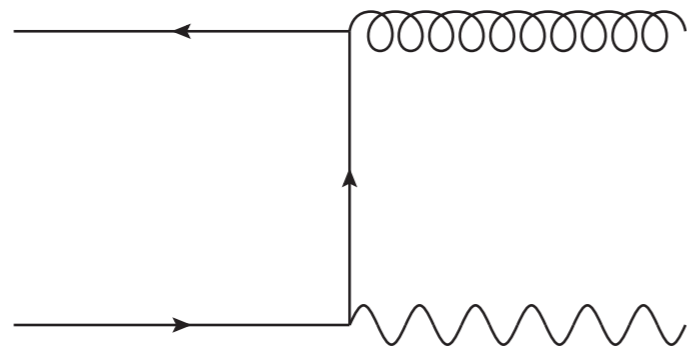
photon + top

Hard photons are a fundamental probe of short-distance physics, of interest both in the context of SM and BSM physics ...

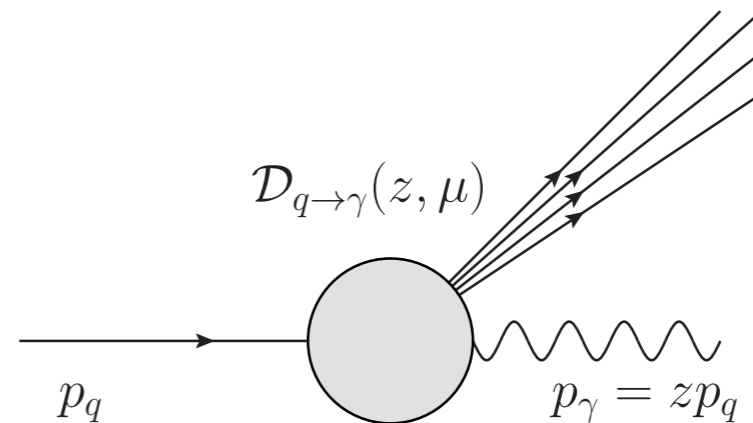
... but also pose experimental and theoretical challenges.

# Challenges

- **Fragmentation**



direct



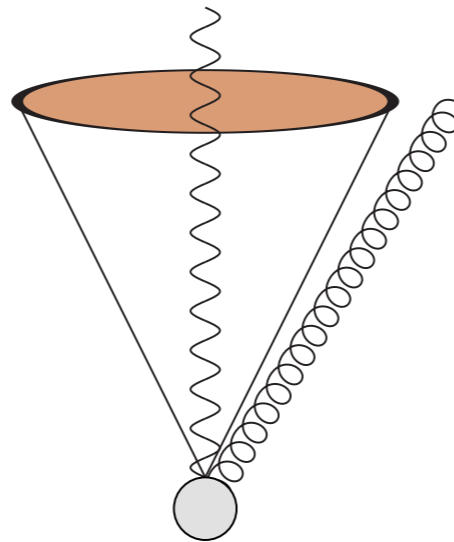
fragmentation

non-perturbative & poorly known

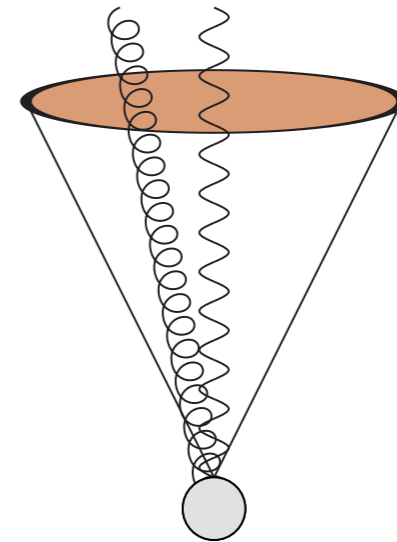
- Background from decays of energetic  $\pi^0$  and  $\eta$ , producing collimated photon pairs
- **Isolate photon** from hadronic radiation to suppress this background  $\rightarrow$  **large logs**

# Fixed-energy cone isolation

Isolation-cone  
of radius  $R$



Isolated

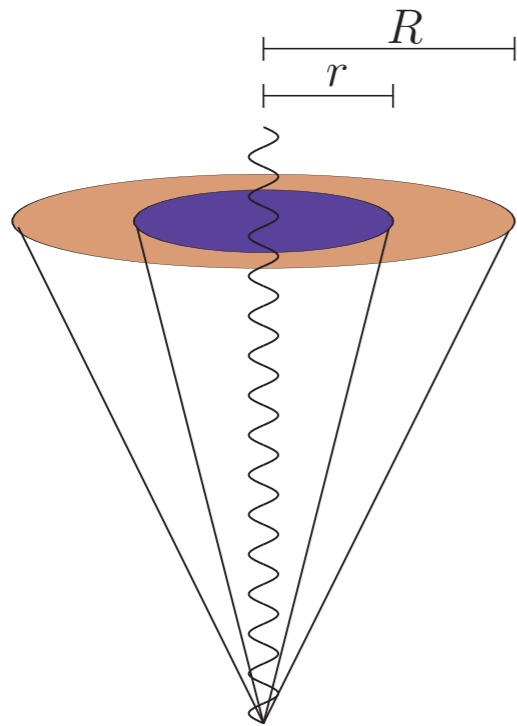


Isolated if  $E_g^T \leq E_{\text{iso}}^T$

- Traditionally  $E_{\text{cone}}^T(R) < E_0 = \epsilon_\gamma E_\gamma^T$
- ATLAS sets  $E_0 = \epsilon E_\gamma^T + E_{\text{th}}^T$  with  
 $\epsilon = 0.0042$  with  $E_{\text{th}}^T = 4.8 \text{ GeV}$
- Fragmentation contributes to cross section
- All LHC measurements use fixed cone.

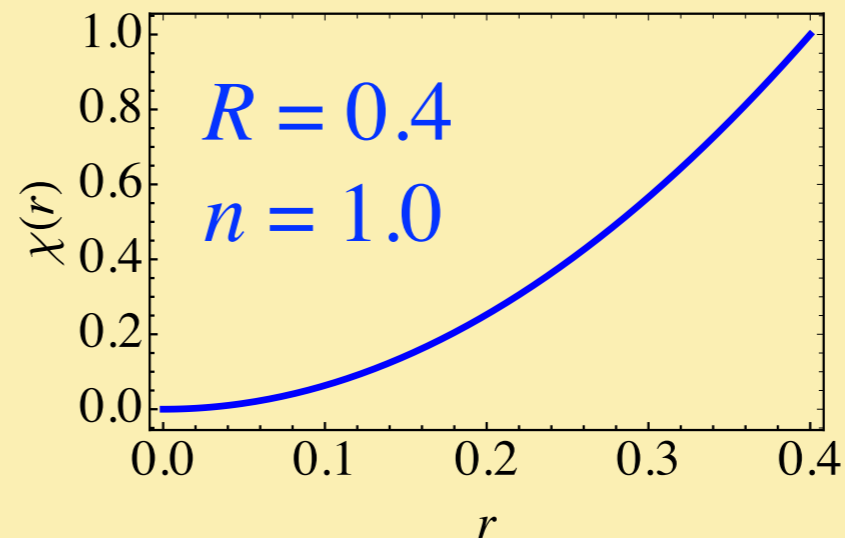
# Smooth-cone isolation

Frixione '98



$$E_{\text{cone}}^T(r) < E_0(r) = \epsilon_\gamma E_\gamma^T \chi(r) \text{ for all } r < R$$

$$\chi(r) = \left( \frac{1 - \cos r}{1 - \cos R} \right)^n$$



- No energetic collinear radiation  $\rightarrow$  no fragmentation
- big technical simplification for NNLO computations
- Experimentally not directly realizable. For a study of discretized version, see [hep-ph/1003.1241](https://arxiv.org/abs/hep-ph/1003.1241)

# NLO predictions

Publicly available fixed-cone NLO only from

- Jetphox (Catani et al. '99), Diphox (Binnoth et al. '99)
  - no longer actively maintained
- MCFM since 2011

Fragmentation functions (and related code) are 25 years old, based on simple models.

Other NLO codes such as MG5\_aMC@NLO restricted to smooth-cone isolation.

Have verified (thanks to Alex Huss!) that different codes produce compatible reference cross sections.

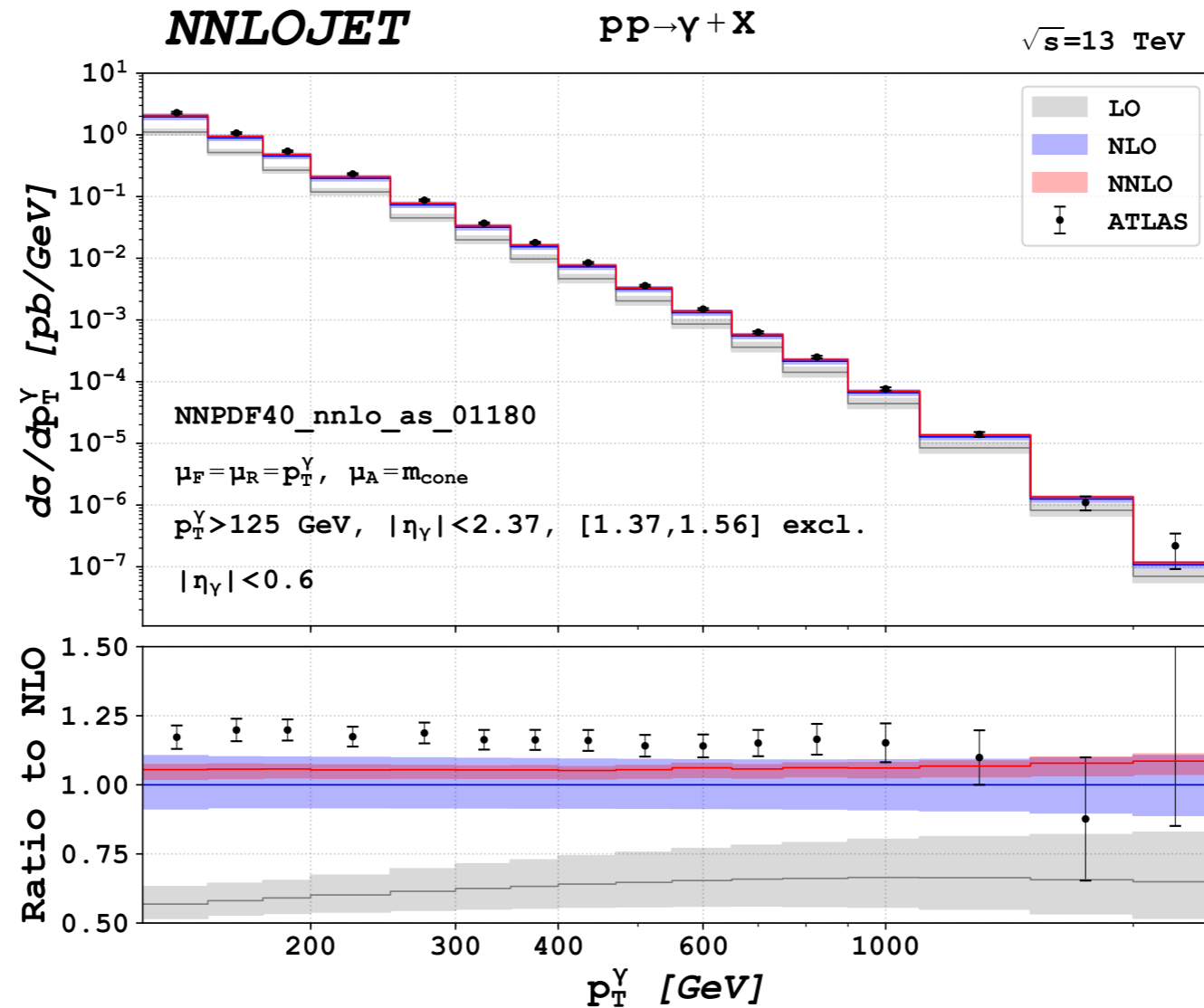
## NNLO predictions

- prompt photon [Campbell et al. '17](#), [Chen et al. 19](#)
- diphoton [Catani et al. 11](#), [Campbell et al. 16](#), [Gehrmann et al. '20](#)
- tri-photon [Chawdhry et al. '20](#), [Kallweit et al. '20](#)

but before last year **only with smooth-cone isolation**

- papers choose values of  $n$  and  $\varepsilon_\gamma$  that give similar NLO values as fixed-cone isolation
- unknown systematic uncertainty, unsatisfactory in view of the few % accuracy of measurements

Proposal to use fixed cone with smooth cone in the center “**hybrid cone**” but not completely satisfactory. [Gehrmann et al. '21](#)



**New:** First fixed-cone NNLO results [Chen, Gehrmann, Glover, Höfer, Huss, Schürmann '22](#) using antenna subtraction; extension to fragmentation: [Gehrmann, Schürmann '22](#).



# Outline

Isolation requirement induces small parameters into cross section

- higher-order corrections enhanced by powers of  $\ln(\epsilon_\gamma)$  and  $\ln(R)$
- will illustrate at NLO that this can lead to a breakdown of the fixed-order expansion

Factorization of isolation effects for small  $R$  using SCET yields

- simple **analytic understanding** of isolation effects
- **resummation** of  $\ln(\epsilon_\gamma)$  and  $\ln(R)$  using RG evolution
- **relation** between smooth- and fixed-cone in the limit of small  $\epsilon_\gamma$



Motivation: pathologies of NLO  
perturbation theory

For all cross section computations we will use

$$\begin{aligned} E_T^\gamma > E_T^{\min} &= 125 \text{ GeV} & |\eta_\gamma| < 2.37 \\ \alpha_s(M_Z) &= 0.119 & \alpha_{\text{EM}} = 1/132.507 \\ \sqrt{s} &= 13 \text{ TeV} & \text{NNPDF23\_nlo\_as\_0119\_qed\_mc} \end{aligned}$$

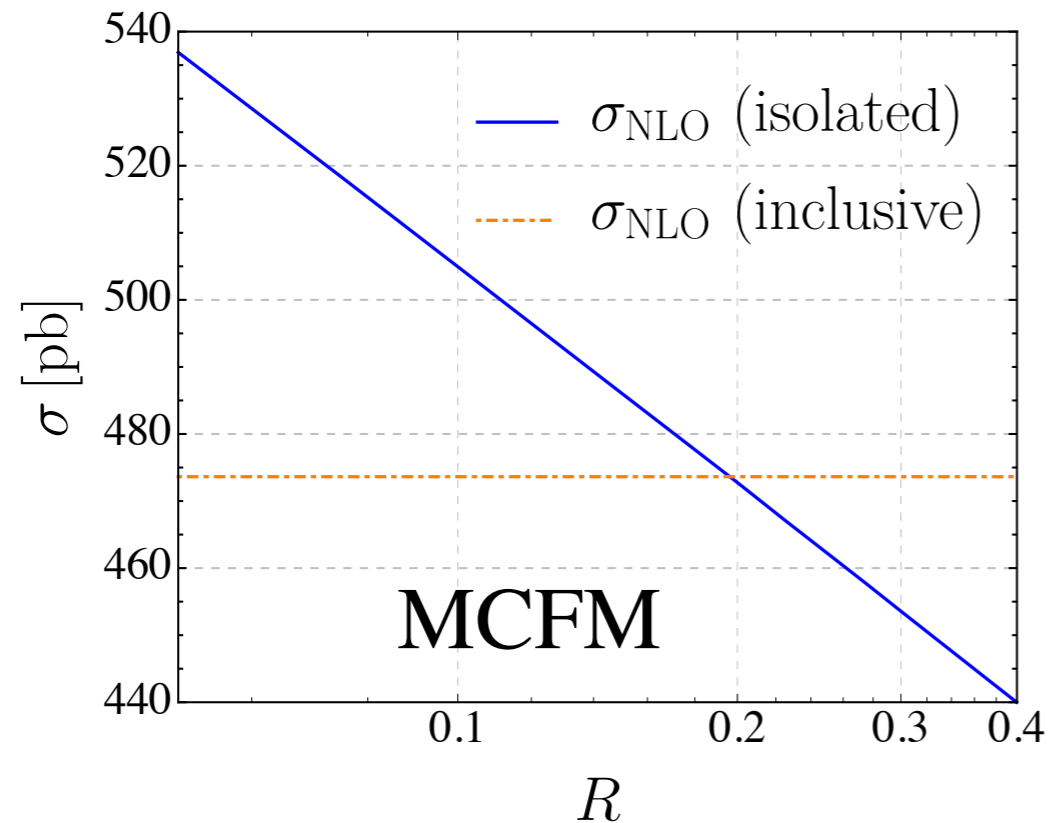
and for fixed-order results we set

$$\mu_f = \mu_r = 125 \text{ GeV}$$

Fixed-cone results involve fragmentation functions and associated scale. For fixed-order, we set

$$\mu_a = 125 \text{ GeV}$$

# Fixed-Order Pathologies (I)



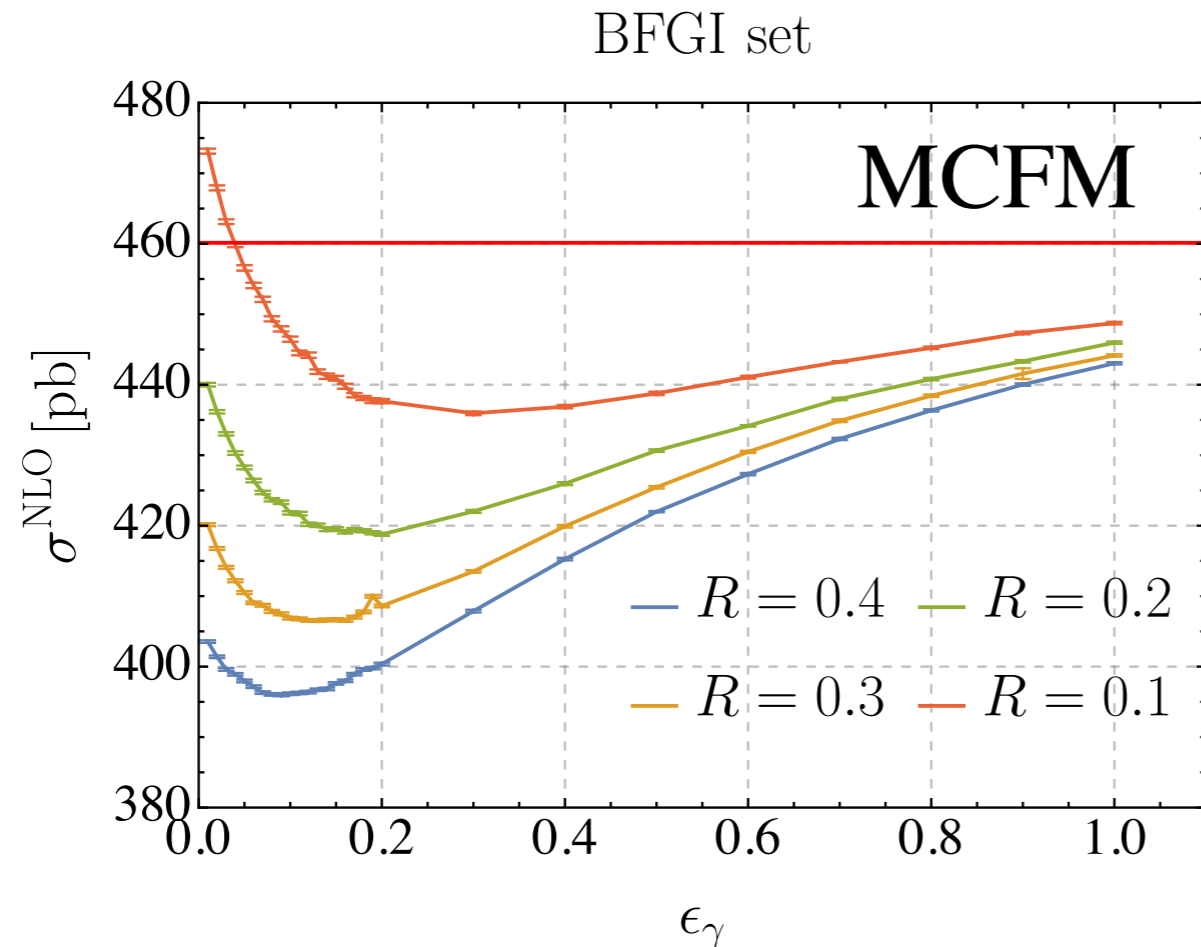
$\sigma(\text{isolated})$  with smooth-cone,  
 $n = 1, \varepsilon_\gamma = 1$

$\sigma(\text{inclusive})$  with [Gehrmann de Ridder, Glover '98](#)

fragmentation functions

- **Should** have:  $\sigma(\text{isolated}) < \sigma(\text{inclusive})$  but at NLO, the isolation dependent part of cross section is proportional to  $\ln(R)$
- Breakdown of FOPT for  $R \approx 0.2$ !  $R = 0.2$  is the default value for ATLAS diphoton analyses
- Same breakdown arises for fixed-cone isolation [Catani, Fontannaz, Guillet and Pilon in JHEP 05, 028 \(2002\)](#)

# Fixed-Order Pathologies (II)

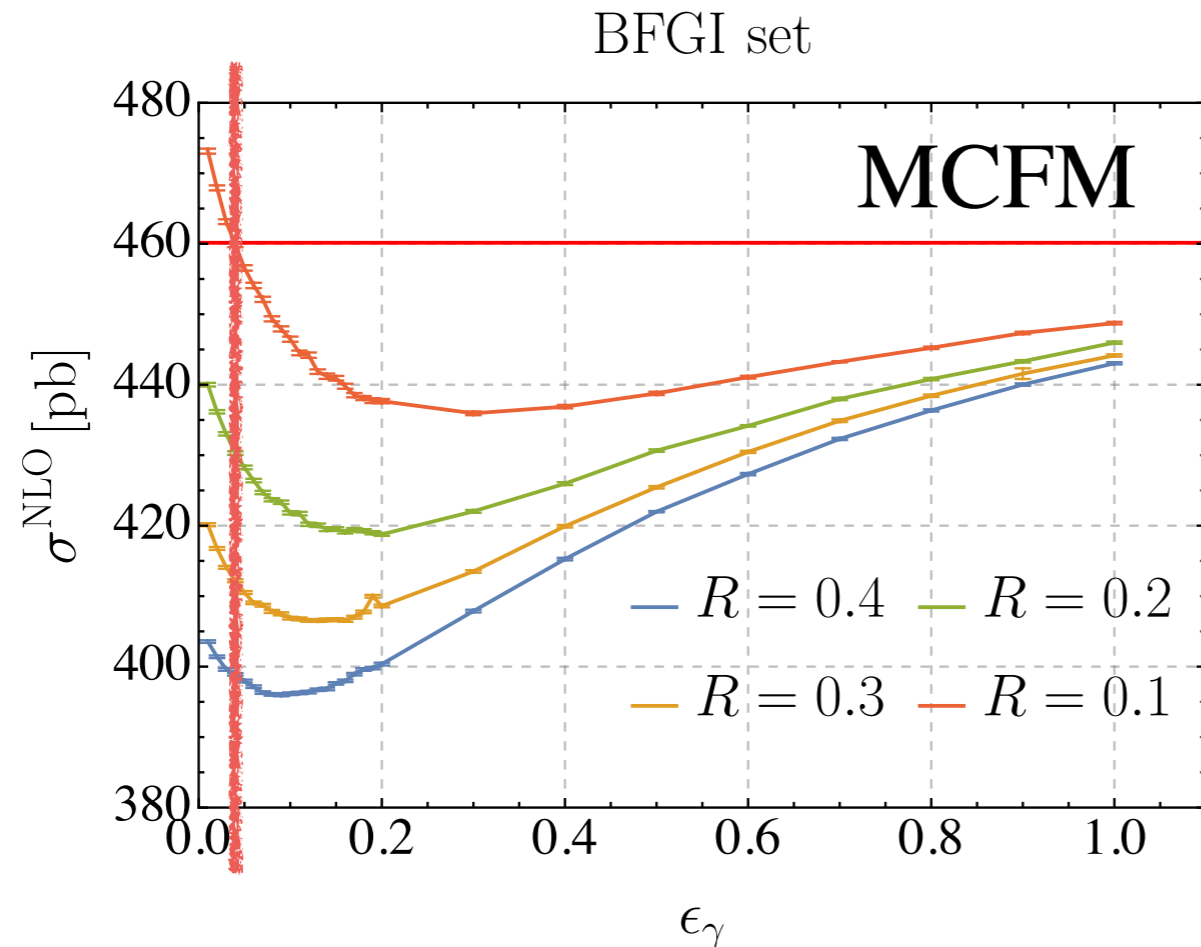


$\sigma(\text{isolated})$  with fixed-cone isolation.

BFG (Bourhis, Fontannaz and Guillet, '98) fragmentation functions

- $\sigma(\text{isolated})$  *should* monotonically decrease as  $\epsilon_\gamma$  is lowered
- NLO isolation effects are linear in  $\epsilon_\gamma$  for small  $\epsilon_\gamma$  (soft quark...)
- coefficient enhanced by  $\ln(R)$ , unphysical for small  $R$
- ATLAS isolation corresponds to  $\epsilon_\gamma = 0.04$  for  $E_T^\gamma = 125$  GeV

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Factorization and resummation  
for small cone radius  $R$

# Factorization

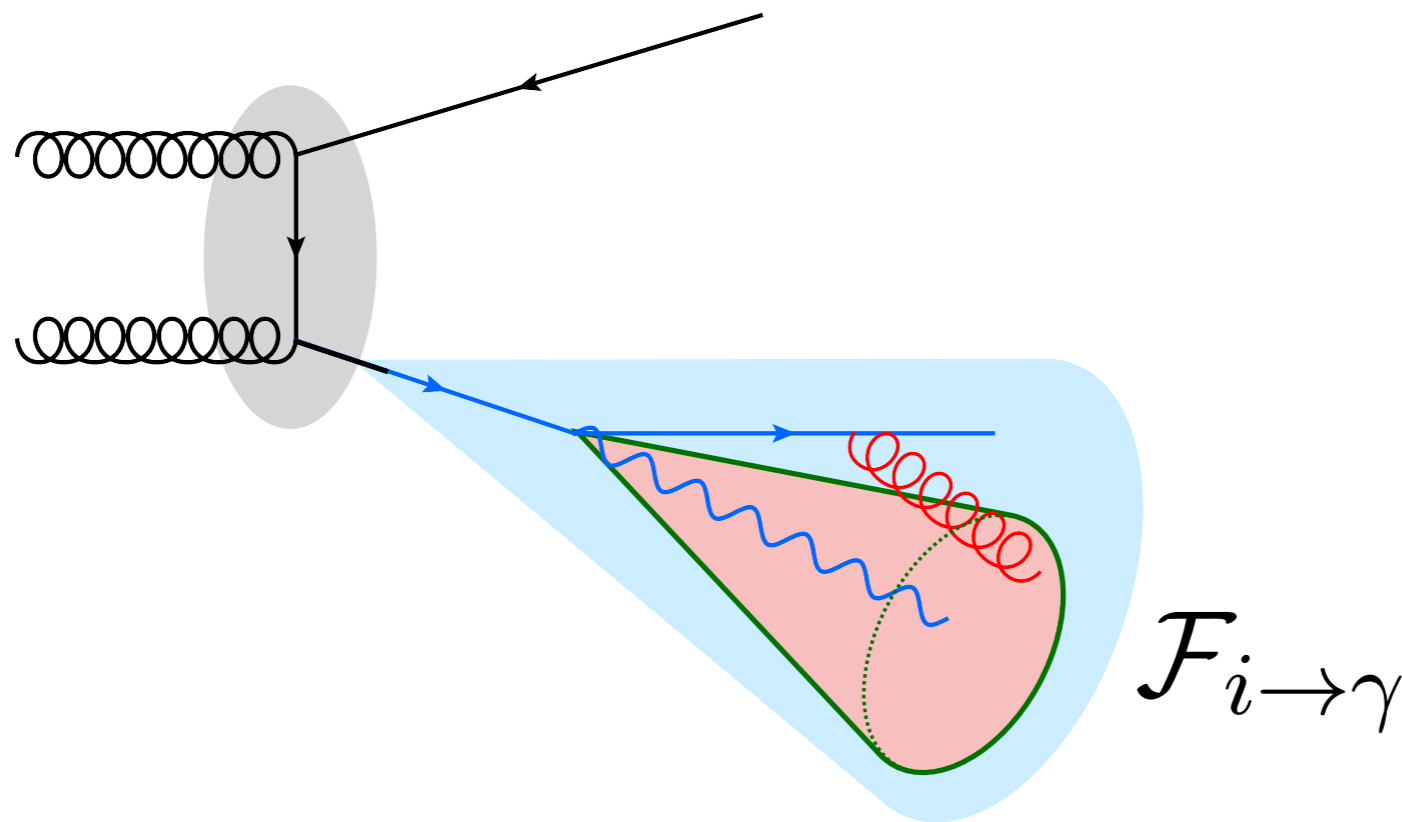
Becher, Favrod, Xu, 2208.01554

For small  $R$  all isolation effects can be factorized into a **cone fragmentation function**  $\mathcal{F}_{i \rightarrow \gamma}$

$$\frac{d\sigma(E_0, R)}{dE_\gamma} = \frac{d\sigma_{\gamma+X}^{\text{dir}}}{dE_\gamma} + \sum_{i=q, \bar{q}, g} \int dz \frac{d\sigma_{i+X}}{dE_i} \mathcal{F}_{i \rightarrow \gamma}(z, E_\gamma, E_0, R) + \mathcal{O}(R)$$

Analogous to factorization of non-perturbative effects, but  $\mathcal{F}_{i \rightarrow \gamma}$  includes perturbative part associated with isolation.





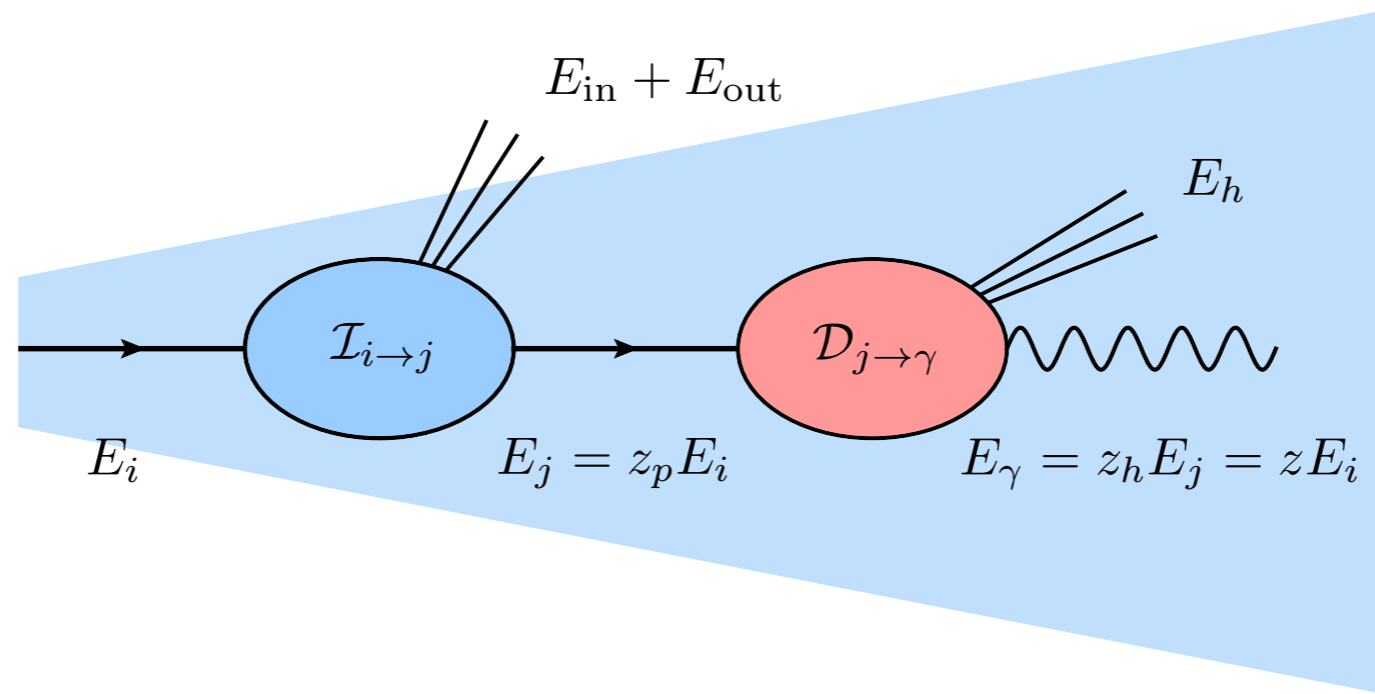
Cone fragmentation function  $\mathcal{F}_{i \rightarrow \gamma}$  contains all particles collinear to photon

$$\frac{d\sigma(E_0, R)}{dE_\gamma} = \frac{d\sigma_{\gamma+X}^{\text{dir}}}{dE_\gamma} + \sum_{i=q, \bar{q}, g} \int dz \frac{d\sigma_{i+X}}{dE_i} \mathcal{F}_{i \rightarrow \gamma}(z, E_\gamma, E_0, R) + \mathcal{O}(R)$$

inclusive, direct

all isolation effects + NP frag.

# Cone fragmentation function $\mathcal{F}_{i \rightarrow \gamma}$



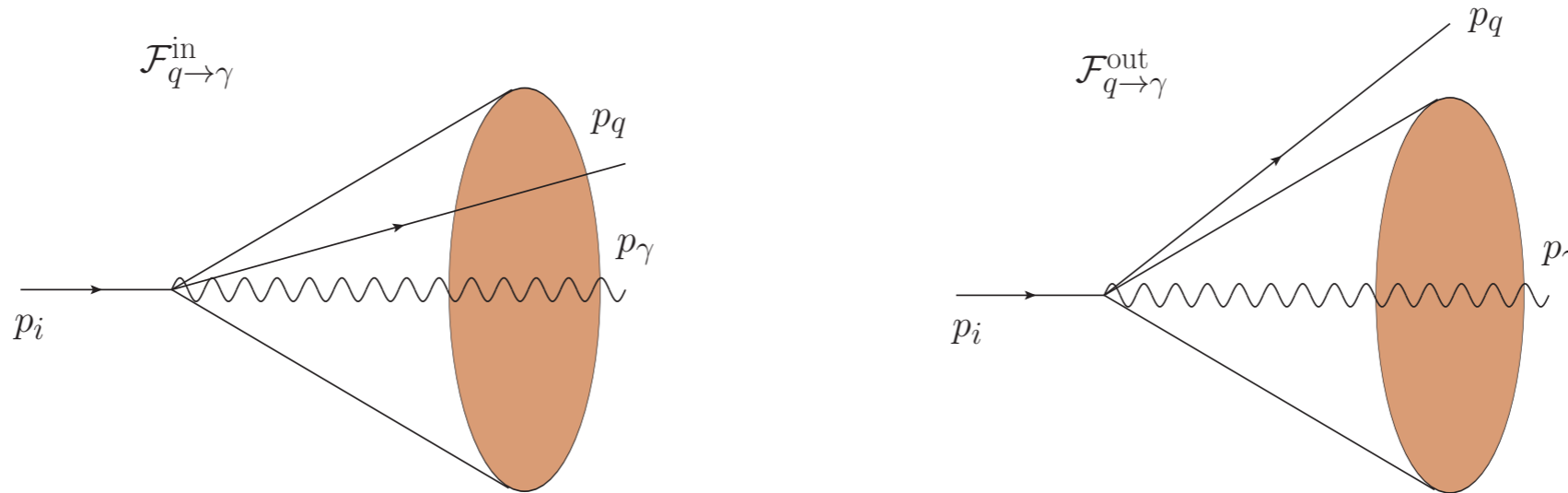
$$\mathcal{F}_{i \rightarrow \gamma}(z, \mu) = \sum_{j=\gamma, q, \bar{q}, g} \mathcal{I}_{i \rightarrow j}(E_\gamma R, E_0 R, \mu) \otimes \mathcal{D}_{j \rightarrow \gamma}(\mu)$$

non-pert. fragmentation

perturbative, scales  $E_\gamma R$  and  $E_0 R$

# NLO cone fragmentation functions

$$\mathcal{F}_{q \rightarrow \gamma}(z, E_\gamma, E_0, R, \mu) = \mathcal{F}_{q \rightarrow \gamma}^{\text{in}}(z, E_\gamma, E_0, R, \mu) + \mathcal{F}_{q \rightarrow \gamma}^{\text{out}}(z, R E_\gamma, \mu)$$



outside part is independent of isolation

$$\mathcal{F}_{q \rightarrow \gamma}^{\text{out}}(z, R E_\gamma) = \frac{\alpha_{\text{EM}} Q_q^2}{2\pi} \left\{ -P(z) \ln \left( \frac{R^2 (2E_T^\gamma)^2}{\mu^2} (1-z)^2 \right) - z \right\}$$

quark to photon splitting function  $P(z) = \frac{1 + (1-z)^2}{z}$

# Inside part

## Smooth-cone isolation

$$\mathcal{F}_{q \rightarrow \gamma}^{\text{in}}(z, E_\gamma, E_0, R, \mu) = \frac{\alpha_{\text{EM}} Q_q^2}{2\pi} P(z) \frac{1}{n} \ln \left( \frac{z \epsilon_\gamma}{1-z} \right) \theta \left( z - \frac{1}{1+\epsilon_\gamma} \right)$$

Note  $R$  independence!

## Fixed-cone isolation

$$\mathcal{F}_{i \rightarrow \gamma}^{\text{in}}(z, R, E_\gamma, E_0, \mu) = \left[ \mathcal{D}_{i \rightarrow \gamma}(z, \mu) + \sum_{k=q, \bar{q}} \delta_{ik} \mathcal{I}_{k \rightarrow \gamma}^{\text{in}}(z, R, E_\gamma, \mu) \right] \theta \left( z - \frac{1}{1+\epsilon_\gamma} \right)$$

$$\mathcal{I}_{q \rightarrow \gamma}^{\text{in}}(z, R, E_\gamma, \mu) = \frac{\alpha_{\text{EM}} Q_q^2}{2\pi} \left\{ P(z) \ln \left( \frac{R^2 (2E_T^\gamma)^2}{\mu^2} (1-z)^2 \right) + z \right\}$$

# Isolation parameter dependence

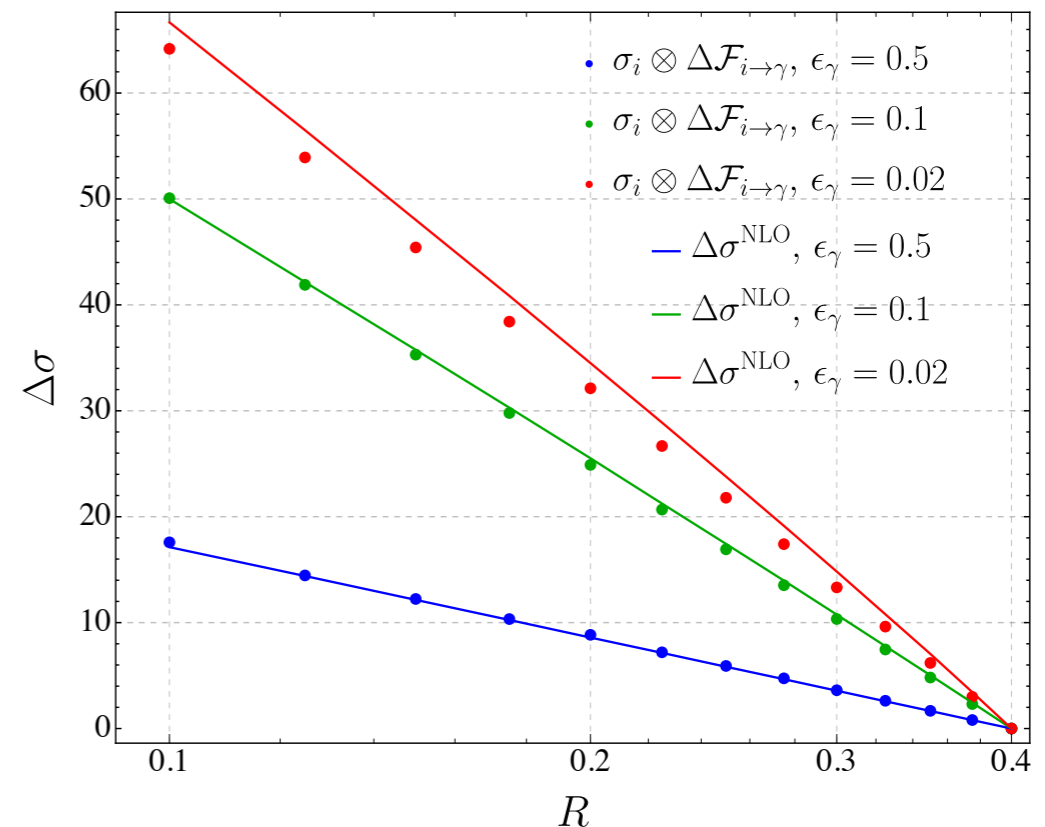
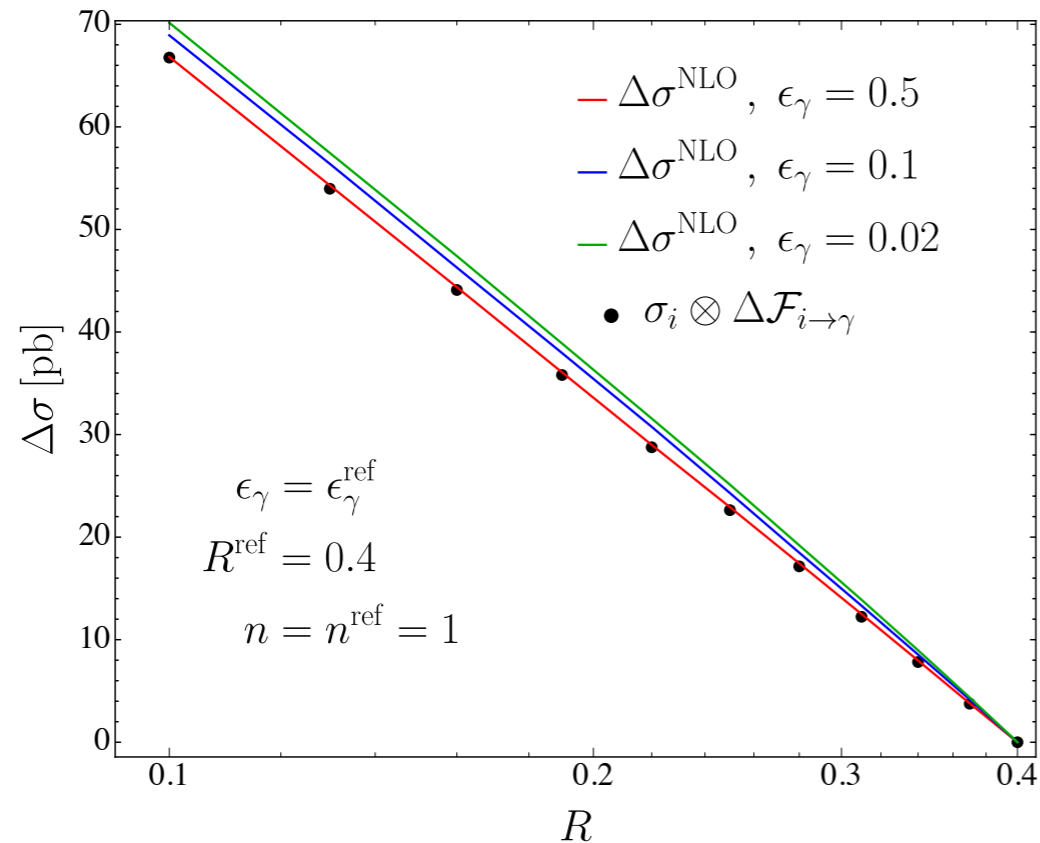
Interesting to look at difference to reference cross section

$$\Delta\sigma = \sigma(\epsilon_\gamma, n, R) - \sigma(\epsilon_\gamma^{\text{ref}}, n^{\text{ref}}, R^{\text{ref}})$$

since direct part drops out:

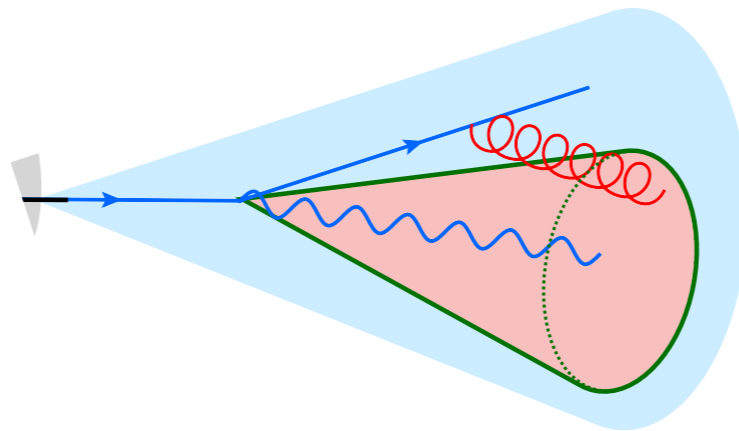
$$\Delta\sigma = \sum_{i=q,\bar{q}} \int_{E_T^{\text{min}}}^{\infty} dE_i \int_{z_{\text{min}}}^1 dz \frac{d\sigma_{i+X}}{dE_i} \Delta\mathcal{F}_{i\rightarrow\gamma}$$

# Smooth- vs fixed-cone isolation



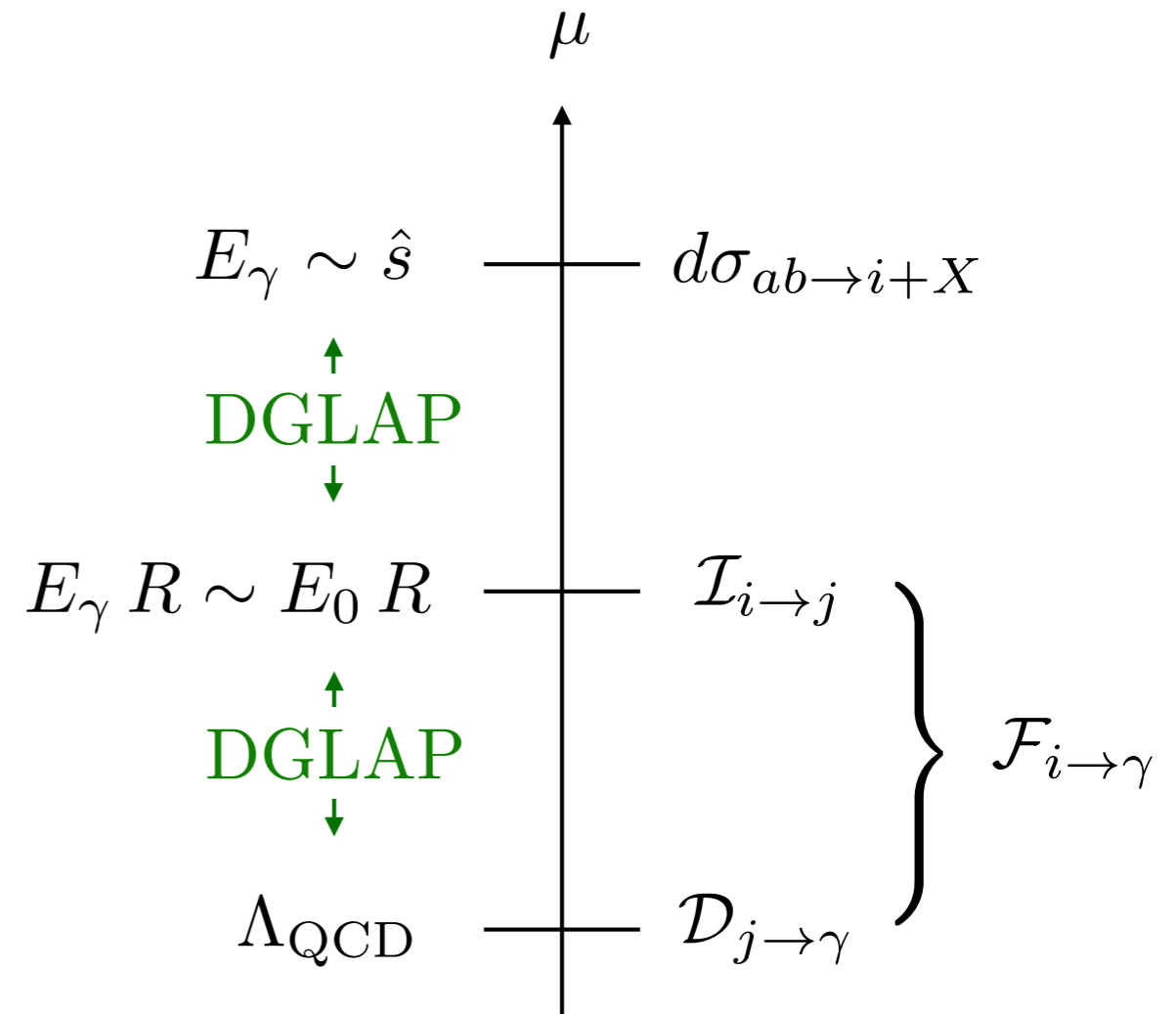
- For fixed cone also inside part of  $\mathcal{F}_{i \rightarrow \gamma}$  has  $\ln(R)$  contribution, which is  $\epsilon_\gamma$  dependent.
- For  $\epsilon_\gamma \rightarrow 0$  inside part vanishes and one recovers smooth-cone  $R$ -dep!

- More generally: for small  $\varepsilon_\gamma$  the inside part at NLO becomes small
- Non-perturbative fragmentation suppressed by  $\varepsilon_\gamma$
- and at NLO the following properties hold
  - $\ln(R)$  dependence only from outside part
  - **All isolation prescriptions become identical!**
- but at NNLO differences from out-in terms!



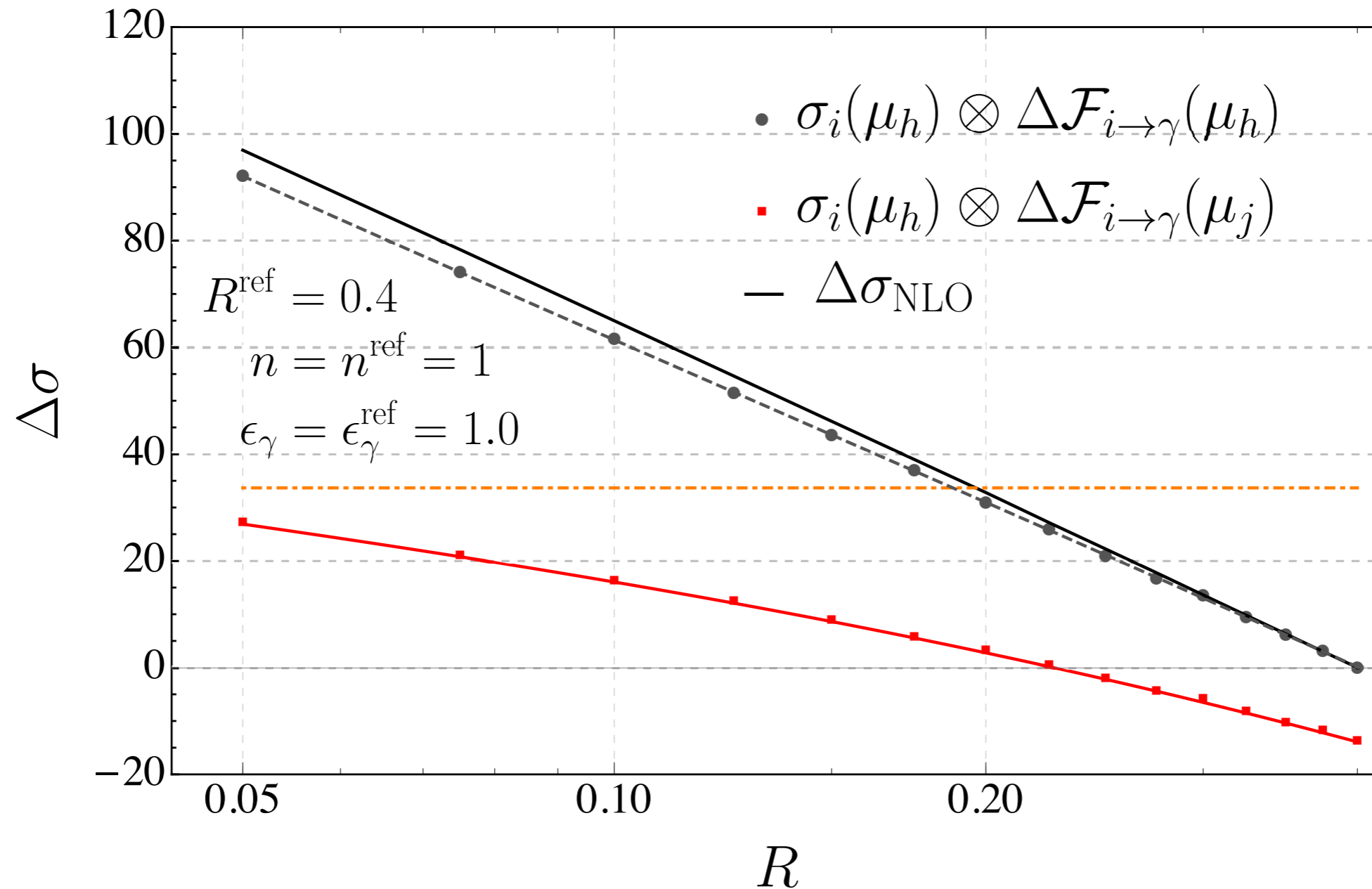
# Resummation of $\ln(R)$ terms

- $\mathcal{F}_{i \rightarrow \gamma}$  fulfills same DGLAP evolution equation as standard fragmentation function
- Solve DGLAP equation numerically to resum  $\ln(R)$  enhanced higher-order contributions
- Implemented evolution in moment space, interface to tree-level generator

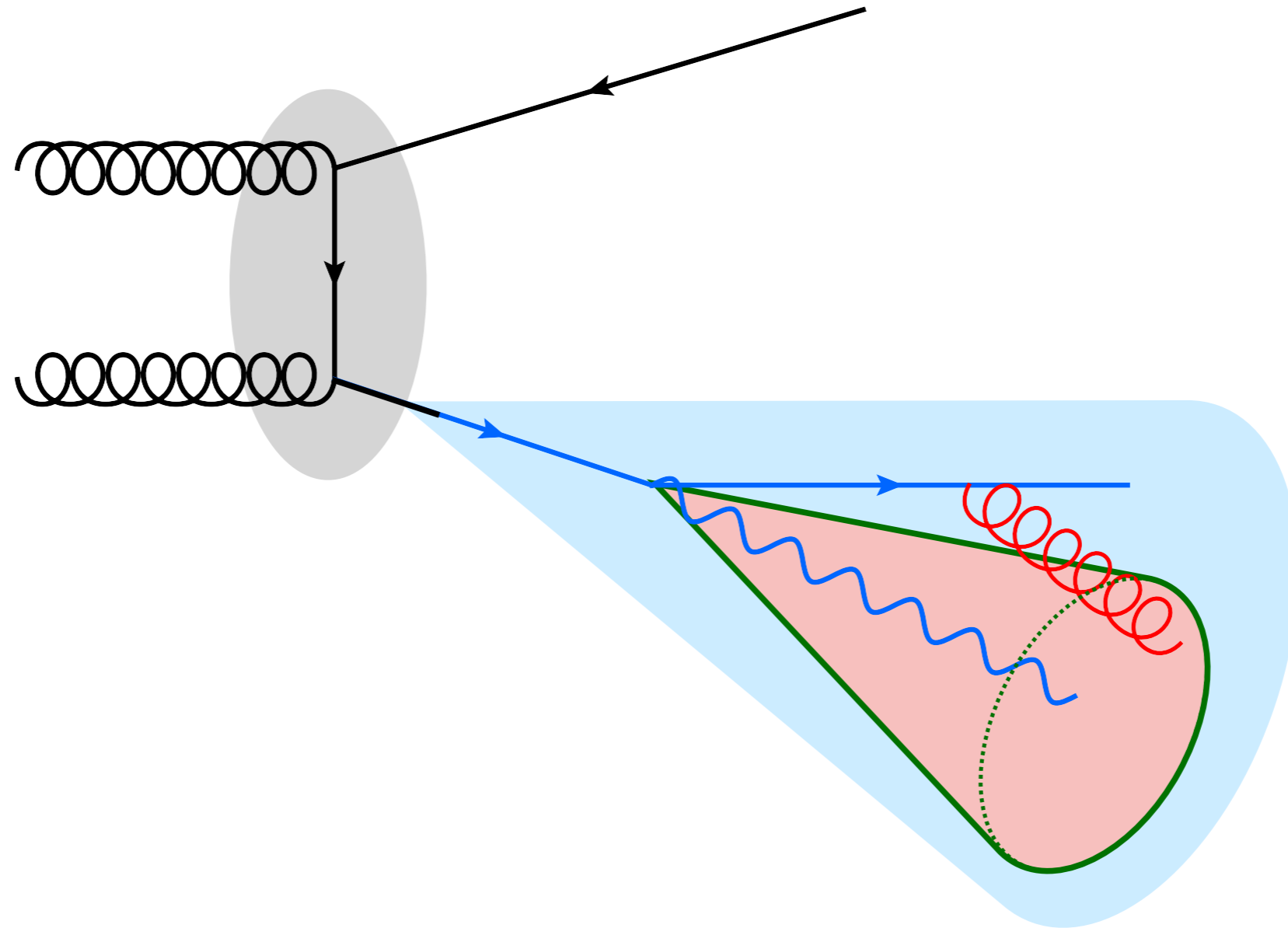




# $\ln(R)$ resummation



- Plot shows difference to reference NLO cross section
- Resummation cures pathological fixed-order behavior!



Factorization and resummation  
for small isolation energy  $E_0$

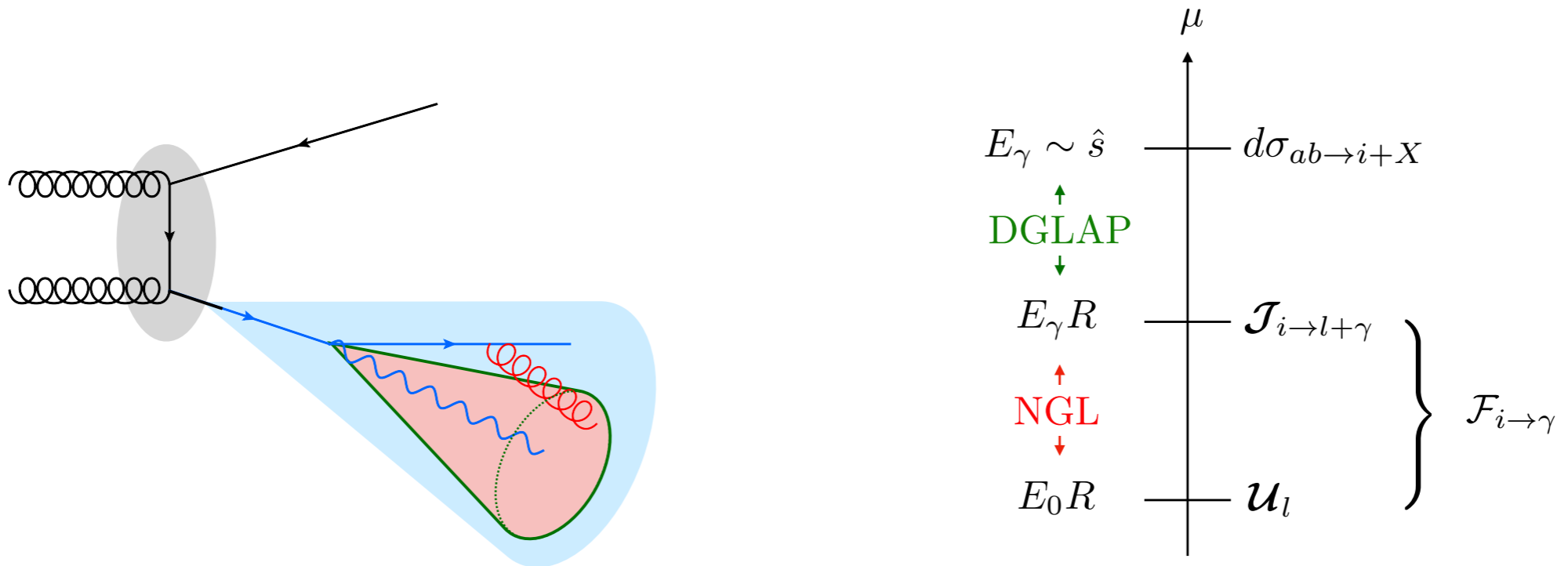
For fixed-cone isolation, the energy inside the cone is always much smaller than the photon energy, e.g.

$$E_{\text{cone}}^T(R) < E_0 = \epsilon_\gamma E_\gamma^T$$

For ATLAS  $E_0 \gtrsim 5$  GeV

- Only **soft radiation** inside cone
- **large non-global logarithms**, associated with the energy ratio  $\epsilon_\gamma$
- perturbation theory at **low scale  $RE_0$**
- **fragmentation is suppressed**

# Factorization of $\mathcal{F}_{i \rightarrow \gamma}$



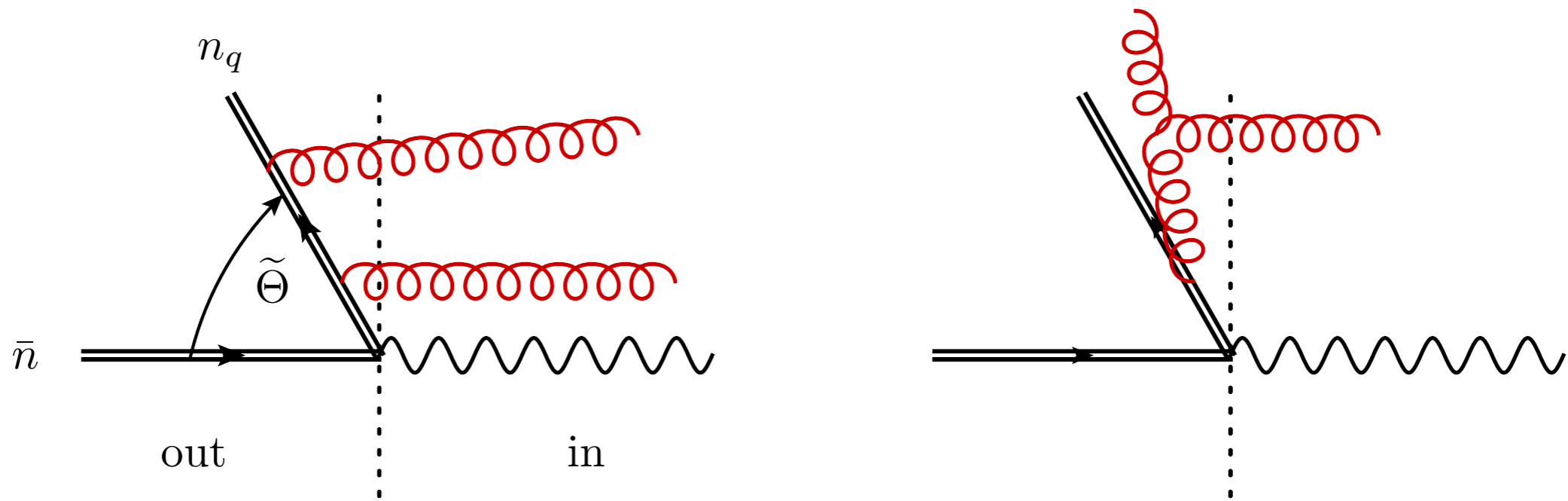
$$\mathcal{F}_{i \rightarrow \gamma}(z, R E_\gamma, R E_0, \mu) = \sum_{l=1}^{\infty} \langle \mathcal{J}_{i \rightarrow \gamma+l}(\{\underline{n}\}, R E_\gamma, z, \mu) \otimes \mathcal{U}_l(\{\underline{n}\}, R E_0, \mu) \rangle$$

energetic collinear  
outside cone

soft radiation  
inside cone

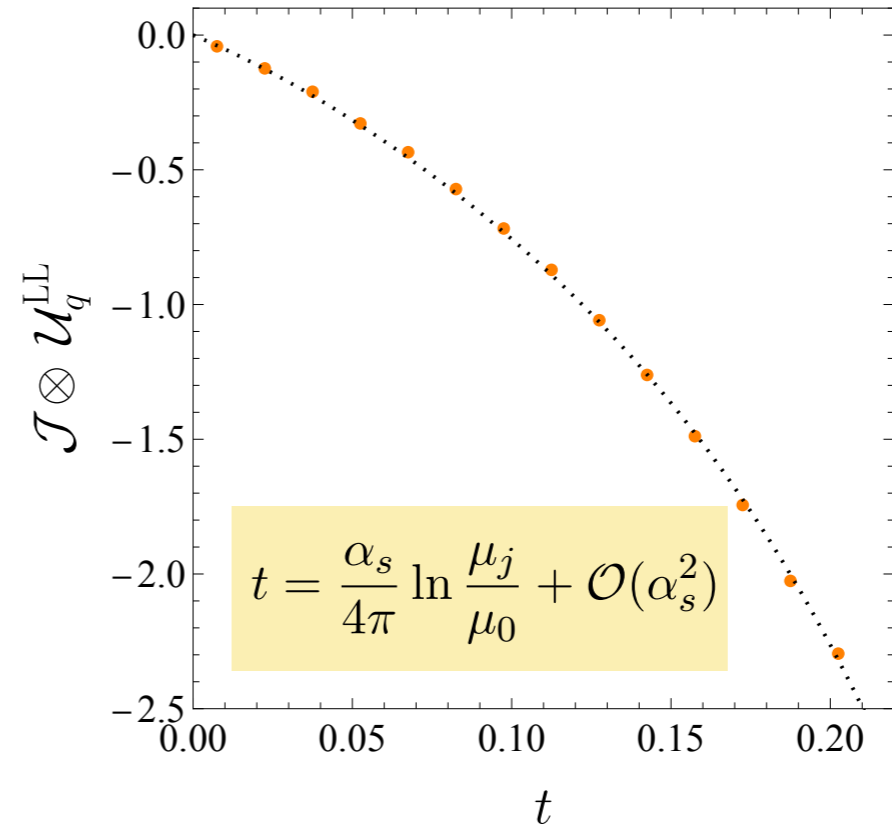
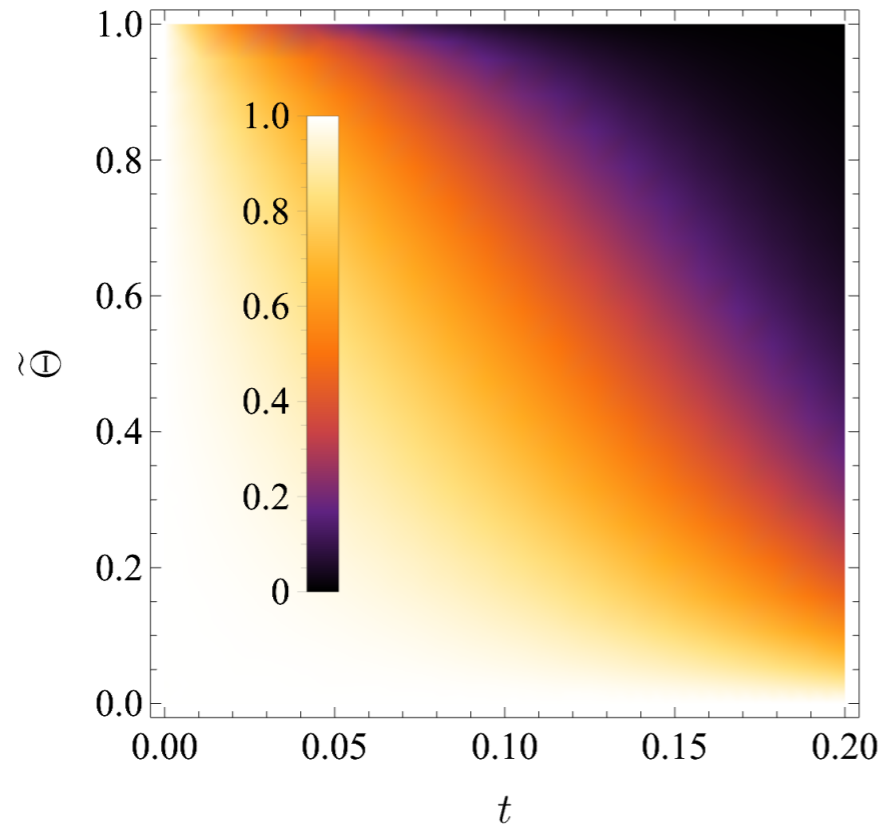
- Resum both  $\ln(\epsilon_\gamma)$  and  $\ln(R)$ .
- Lowest scale is  $R E_0 \gtrsim 1 \text{ GeV}$  for ATLAS !

# Non-global logarithms (NGLs)



- Isolation cone is prime example of **non-global observable**
- Complicated pattern of higher-order terms, not captured by standard resummation methods. Even leading NGLs  $(\alpha_s L)^n$  **do not** simply **exponentiate!** [Dasgupta, Salam '02](#)
  - Use **ngl-resum** to resum leading NGLs [Balsiger, Becher, Ferroglia '20](#) after boosting to frame where cone is hemisphere
  - see **Nicolas Schalch's talk** for progress on resummation of subleading NGLs.

quark-angle

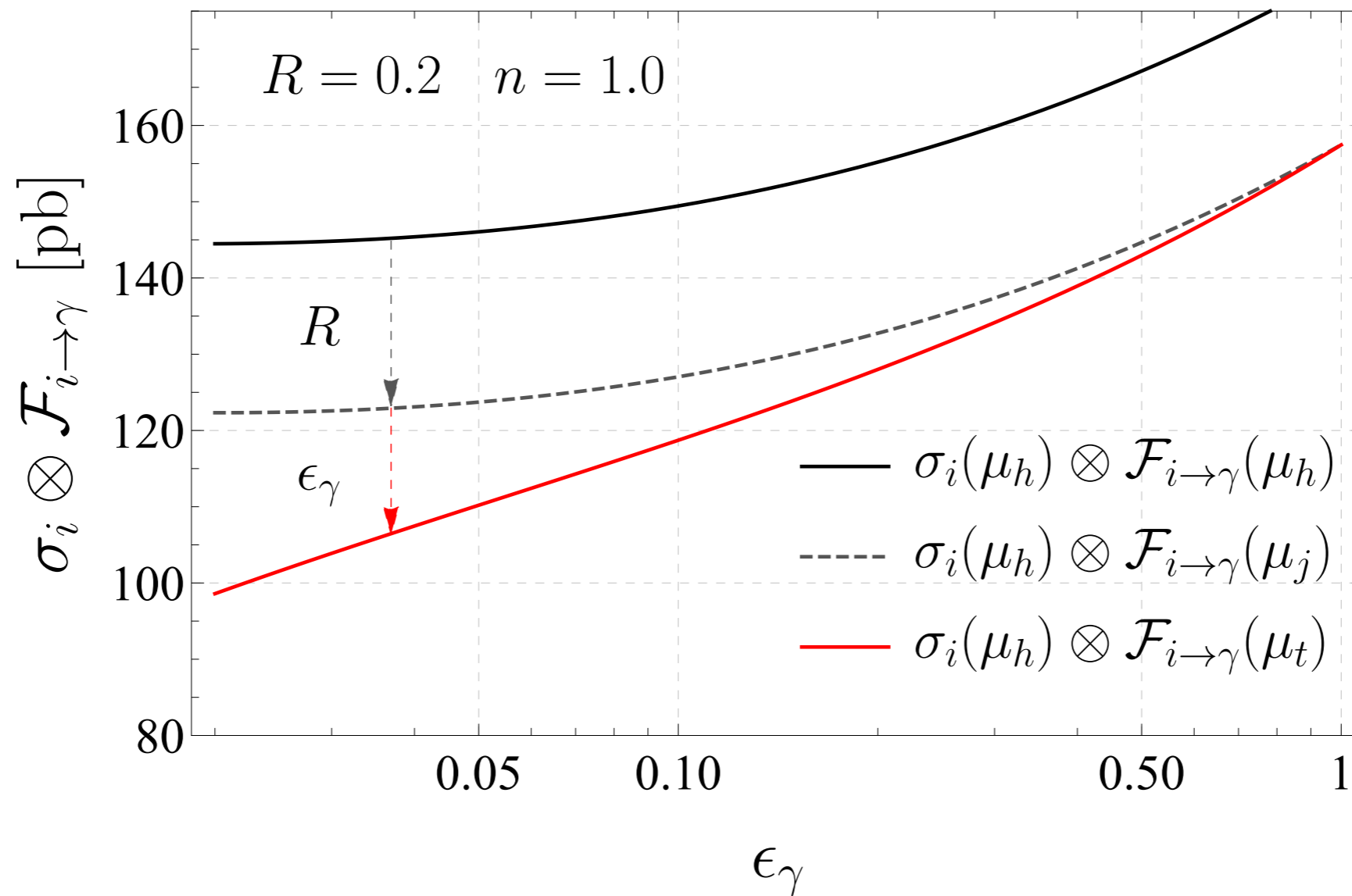


$$\int_0^1 d\tilde{\Theta} \left\langle \mathcal{J}_{q \rightarrow \gamma+q}(\tilde{\Theta}, R E_\gamma, z, \mu_j) \mathcal{U}_q^{\text{LL}}(\tilde{\Theta}, t) \right\rangle$$

$$= \frac{Q_i^2 \alpha_{\text{EM}}}{2\pi} \left[ -P(z) \ln \left( \frac{\delta^2 Q^2}{\mu_j^2} (z-1)^2 z^2 \right) - z + 2P(z) \int_0^1 d\tilde{\Theta} \left[ \frac{1}{\tilde{\Theta}} \right]_+ \mathcal{U}_q^{\text{LL}}(\tilde{\Theta}, t) \right]$$

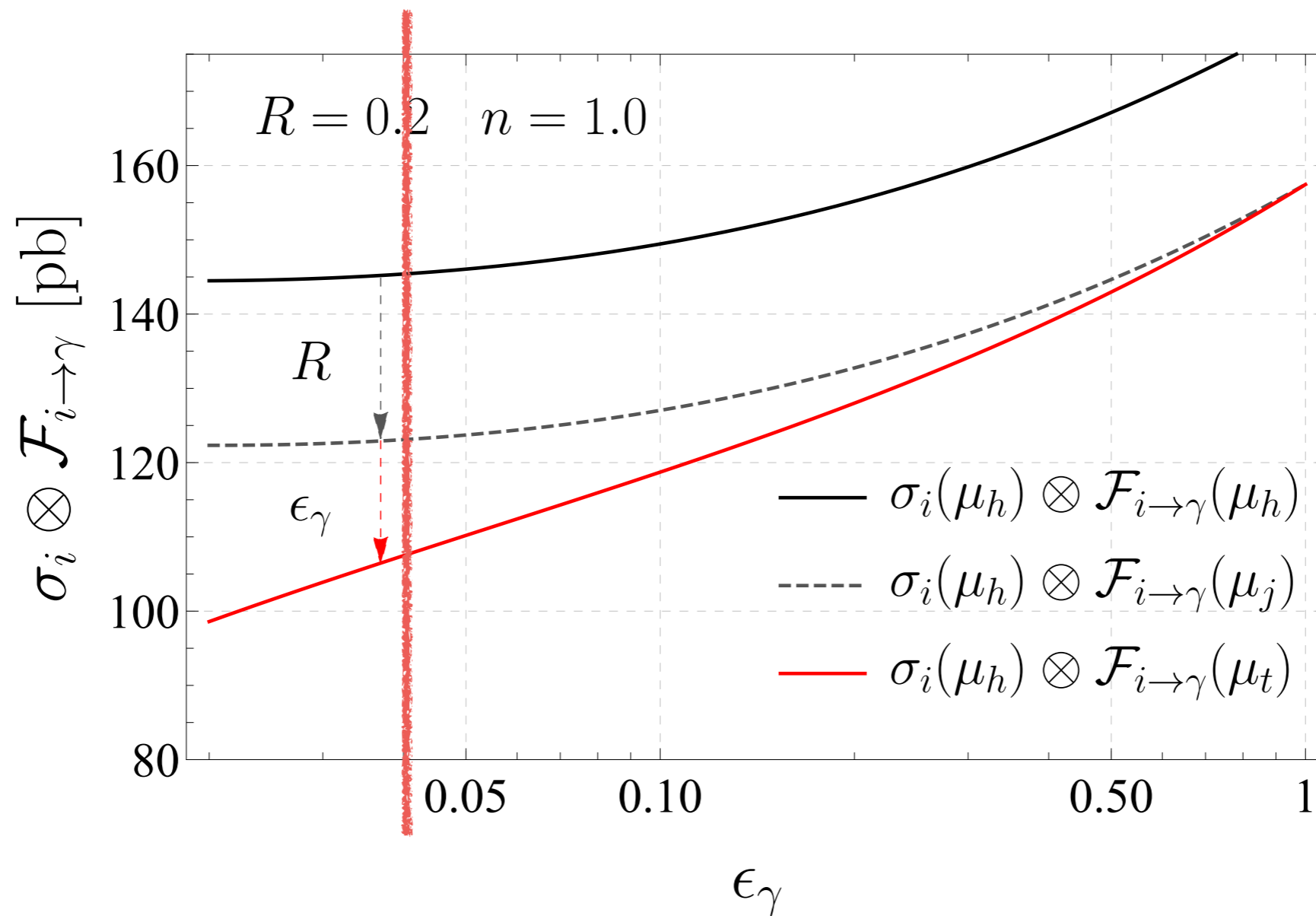
- NGLs much larger than global logs. Two-loop coefficient  
 $-31.5 = -43.7$  (“non-global”) +  $12.2$  (“global”)

# Resummation of $\ln(R)$ and $\ln(\epsilon_\gamma)$



- For the full cross section, add direct part  $\sigma^{\text{dir}} \approx 290$  pb
- Note: both resummations lower the cross section!

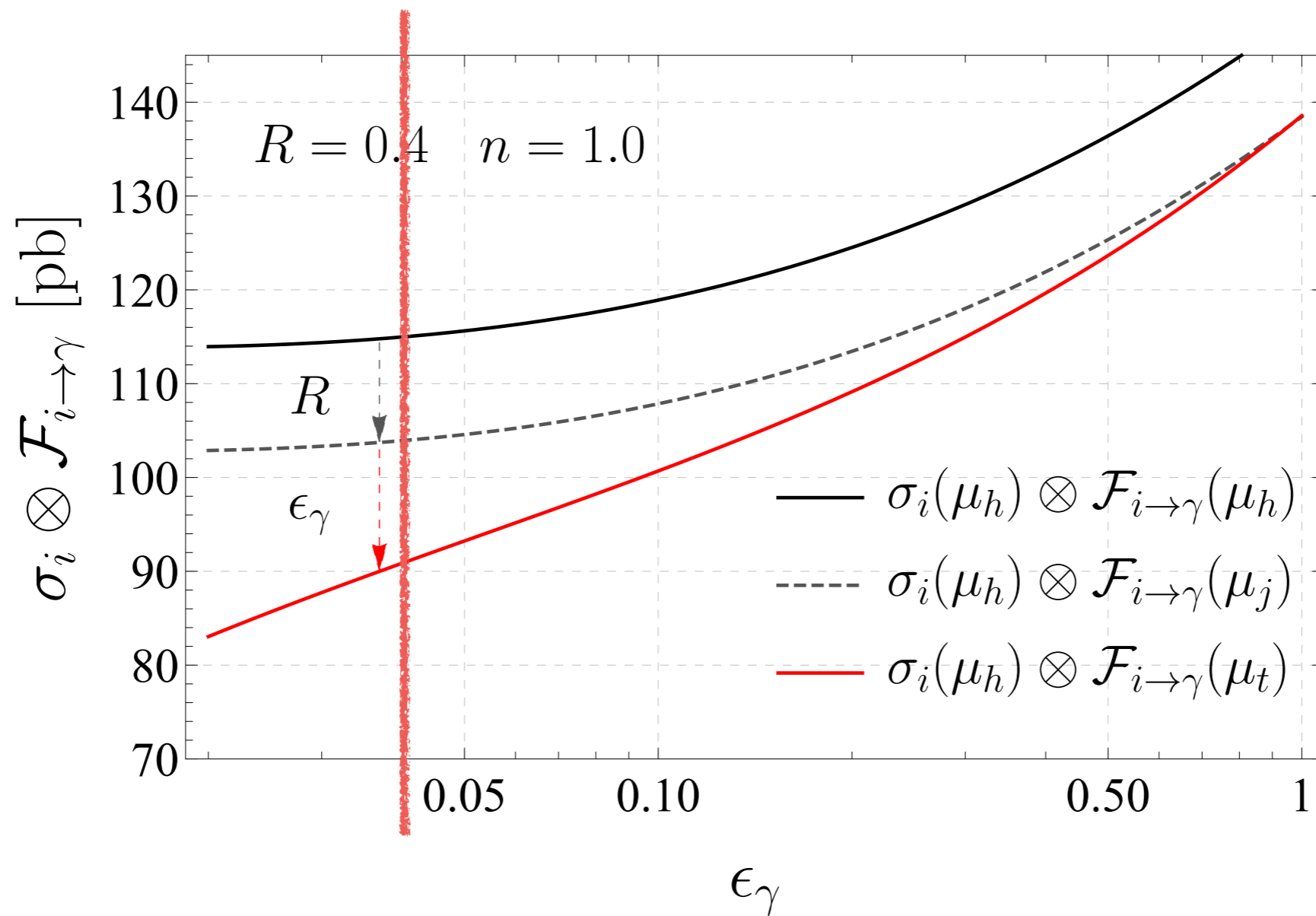
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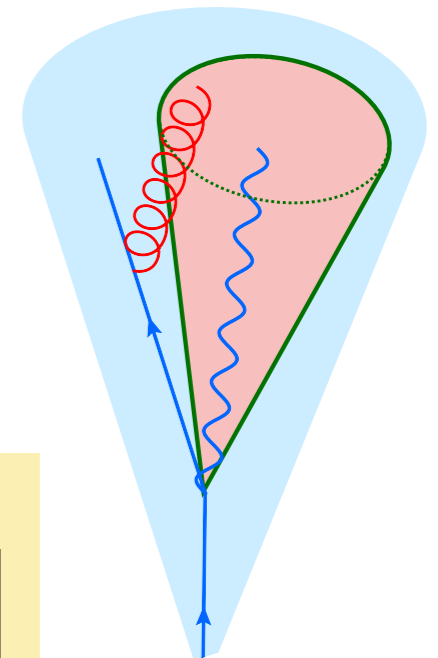
# A simple relation

From factorization theorem, we can derive a **relation between smooth- and fixed-cone isolation**

$$\Delta\sigma = \sigma_{\text{fixedcone}}(R, \epsilon_\gamma) - \sigma_{\text{smoothcone}}(R, \epsilon_\gamma^{\text{ref}}, n)$$

in the limit  $R \rightarrow 0$  and  $\epsilon_\gamma \rightarrow 0$  :

$$\Delta\sigma = \sum_{i=q,\bar{q}} \int_{E_T^{\min}}^{\infty} dE_i \int_{z_{\min}}^1 dz \frac{d\sigma_{i+X}}{dE_i} \frac{Q_q^2 \alpha_{\text{EM}}}{\pi} \frac{C_F \alpha_s}{4\pi} P(z) \left[ \frac{\pi^2}{3} \ln \frac{\epsilon_\gamma}{\epsilon_\gamma^{\text{ref}}} + 2n \zeta_3 \right]$$

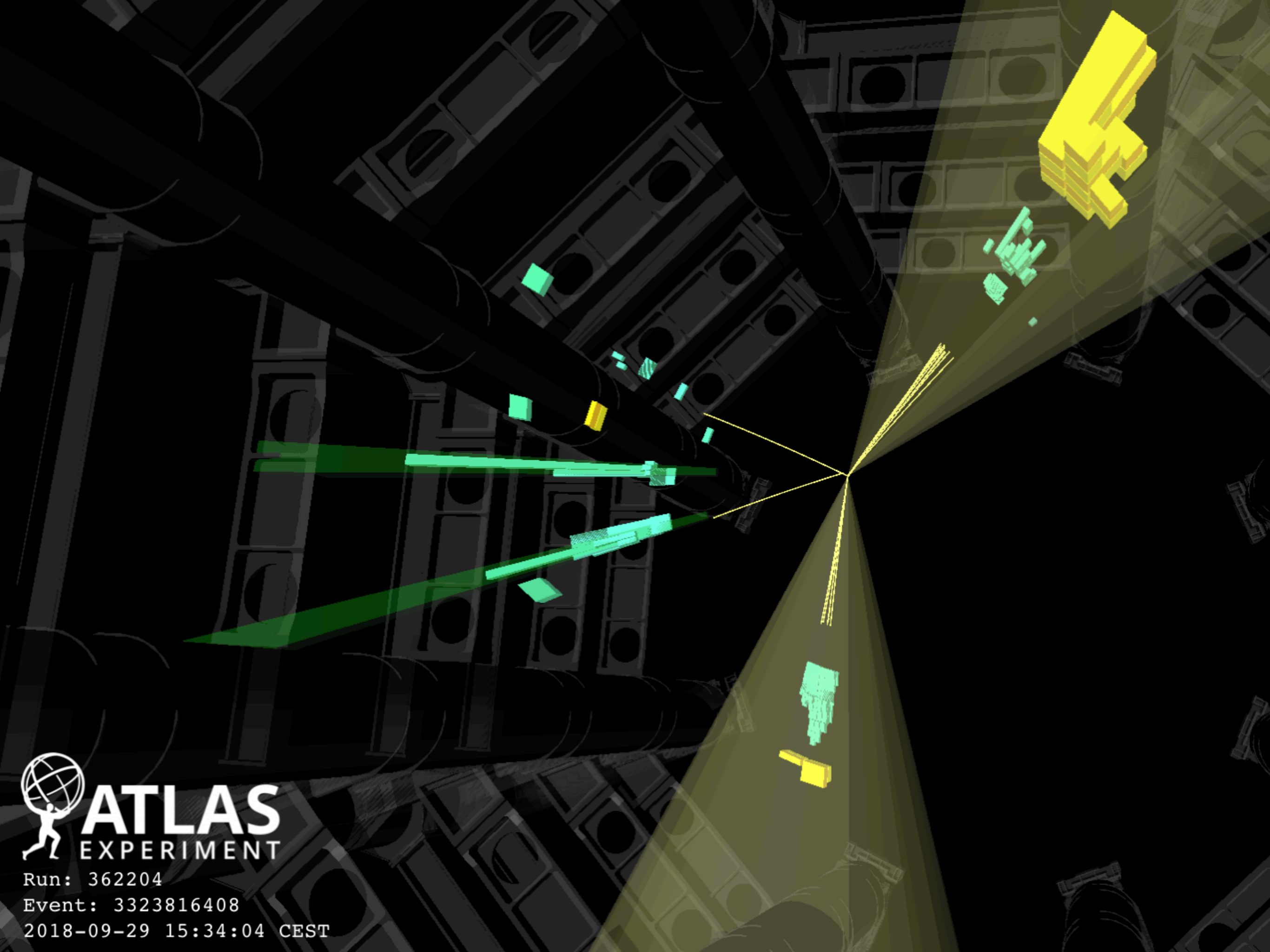


Can be used to convert NNLO smooth-cone into fixed-cone results. For standard setup and  $\epsilon_\gamma = \epsilon_\gamma^{\text{ref}}$

$$\Delta\sigma = -1.3 \text{ pb}$$

# Conclusions & Outlook

- Have performed a detailed analysis of QCD effects associated with photon isolation
  - Factorization of isolation effects for small  $R$ 
    - Lowest scale is  $R E_0 \gtrsim 1 \text{ GeV}$
  - First resummation of both  $\ln(\epsilon_\gamma)$  and  $\ln(R)$  effects
    - numerically relevant for ATLAS isolation, crucial for  $R = 0.2$
    - With some effort, could extend accuracy of resummations by one order and match to NNLO
- Formalism provides analytical understanding of isolation
  - Study parameter dependence, convert between isolation schemes, ...



 **ATLAS**  
EXPERIMENT

Run: 362204

Event: 3323816408

2018-09-29 15:34:04 CEST