

Quantum entanglement and PDFs in 1+1D

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What problem do we want to solve ?

- Can we compute structure functions such as the PDF by appealing to some simple principle that emerges from strong interactions among partons?

$$f_q(x) = \int_{-\infty}^{\infty} \frac{dt}{2\pi} e^{-itx(n \cdot P)} \langle \mathcal{H} | \mathcal{O}_2 | \mathcal{H} \rangle,$$

- A description of the hadron state.
- Renormalization of the operator .
- At a more basic level, understand how collective long distance d.o.f. emerge from the microscopic ones connecting different Effective descriptions of gauge theories.
- For QCD, no systematic expansion parameter available, neither is there a dual theory.
- Several simple models in lower dimensions that exhibit QCD like behavior and are solvable in various limits.
- Can we extract some universal information that does not depend on the special symmetries or specifics of these theories and apply them to higher dimensions?

Gauge Theories in 1+1 D

Abelian : Schwinger model 1+1 D QED

Non-Abelian : 't Hooft model 1+1 D QCD

- Gauge coupling is dimensionful $[g] = M^1 \rightarrow \Lambda_{QCD}$
- The strength of the interaction is controlled by $\bar{g} = g/m_q$
- Running effects of the PDF operator suppressed by g/Q
- No spin
- Gauge field is not a propagating degree of freedom but gives rise to a confining potential.
- There is no renormalization

Problem reduces to describing the hadron bound state

Gauge theories in 1+1D

Hamiltonian in light cone gauge $A^+ = 0$, x^+ is the time variable.

Schwinger model : 1+1 QED

$$P^- = \frac{1}{2} \int dx^- \bar{\psi}_+ \gamma^+ \frac{m^2}{i\partial^+} \psi_+ + \frac{e^2}{2} \int dx^- J^+ \frac{1}{(i\partial^+)^2} J^+$$

Kinetic energy

In the regime of strong coupling $e/m \rightarrow \infty$, the spectrum can be solved by looking at the dual theory \rightarrow A theory of free bosons with mass $e/\sqrt{\pi}$

Instantaneous
Linear potential in
position space

confinement even at
small coupling

't hooft model : 1+1 D QCD Hamiltonian

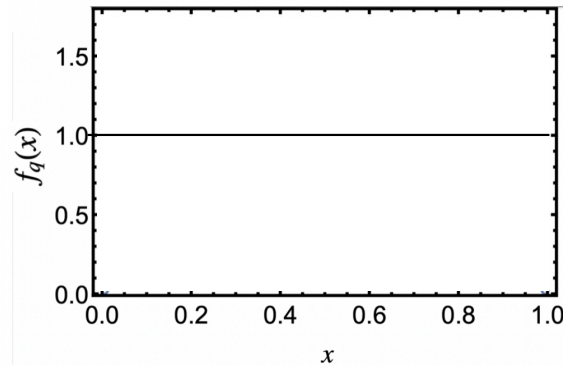
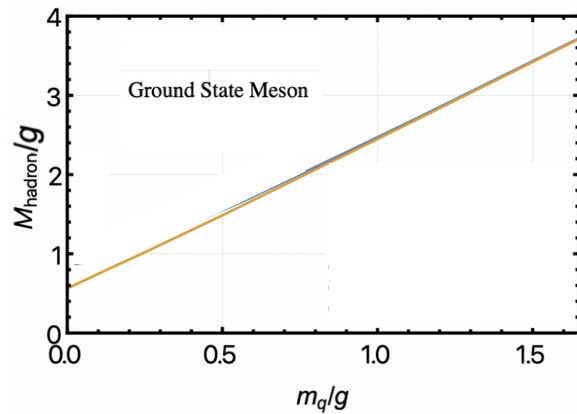
$$P^- = \frac{1}{2} \int dx^- \bar{\psi}_+ \gamma^+ \frac{m^2}{i\partial^+} \psi_+ + \frac{g^2}{2} \int dx^- J^{+,a} \frac{1}{(i\partial^+)^2} J^{+,a}$$

- Large N limit meson spectrum expressed as solution of Bethe Salpeter eqn \rightarrow 2 body ($q\bar{q}$) state schrodinger equation in a confining potential
- Finite N meson, Baryon spectrum solved by numerical methods

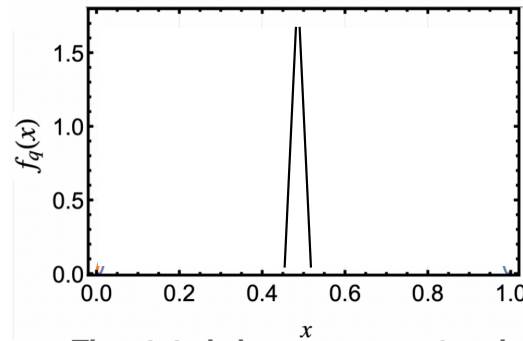
- Ground state meson is a $q\bar{q}$ state; contributions from higher fock states is suppressed
- So there is no contribution from sea partons \rightarrow Same is true for ground state Baryon

Spectrum and PDF in 1+1 D

Schwinger model

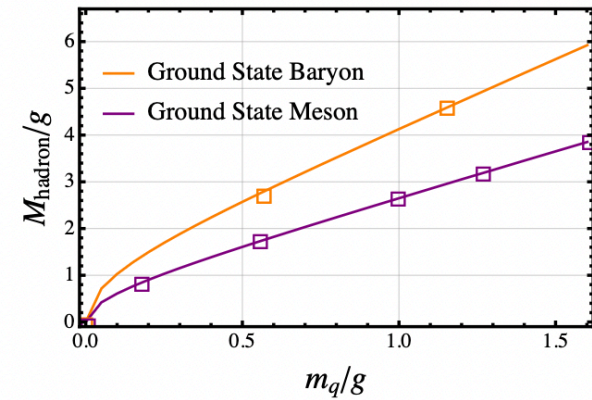


The state is maximally entangled in momentum space at $g/m \rightarrow \infty$



The state becomes unentangled in momentum space at $g/m \rightarrow 0$

't hooft model



Conclusion: We have universality of behavior for the PDF in the two limits across theories .

Perhaps in both cases the system is organizing itself based on some coupling strength dependent measure of Quantum entanglement.

A guess at the answer

- The simplest guess is the Von Neumann entropy S as a measure of entanglement.
- For $g/m \rightarrow \infty$, we will extremize $S \rightarrow$ Maximum entropy principle.
- To describe the system as it moves away from $g/m \rightarrow \infty$, we can draw an analogy with thermodynamics,
- Conjecture a free energy $F = E - TS$ Expectation value of the **Kinetic Energy**

$$E = \langle P^- \rangle_0 \propto m^2$$

The Meson wave function

- $q\bar{q}$ state

i (\bar{j}) \rightarrow fraction of the longitudinal momentum P^+ carried by quark (antiquark)

$$|\psi\rangle = \sum_{i,\bar{j}} p_{i,\bar{j}} \delta_{i+\bar{j},1} |i,\bar{j}\rangle.$$

- Reduced density matrix is diagonal

$$\rho_1 = \sum_i |p_{i,1-i}|^2 |i\rangle\langle i|.$$

- Entanglement entropy

$$S = - \sum_i |p_{i,1-i}|^2 \ln |p_{i,1-i}|^2,$$

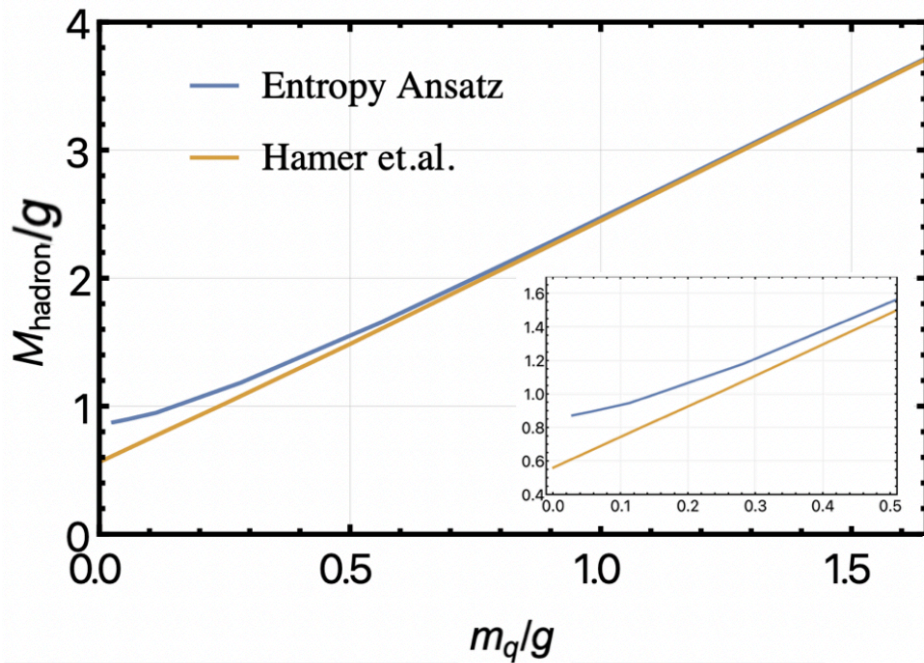
- Kinetic energy

$$\langle E \rangle \equiv \langle P^- \rangle_{\text{kinetic}} = \frac{m^2}{P^+} \sum_i |p_i|^2 \left(\frac{1}{i} + \frac{1}{1-i} \right)$$

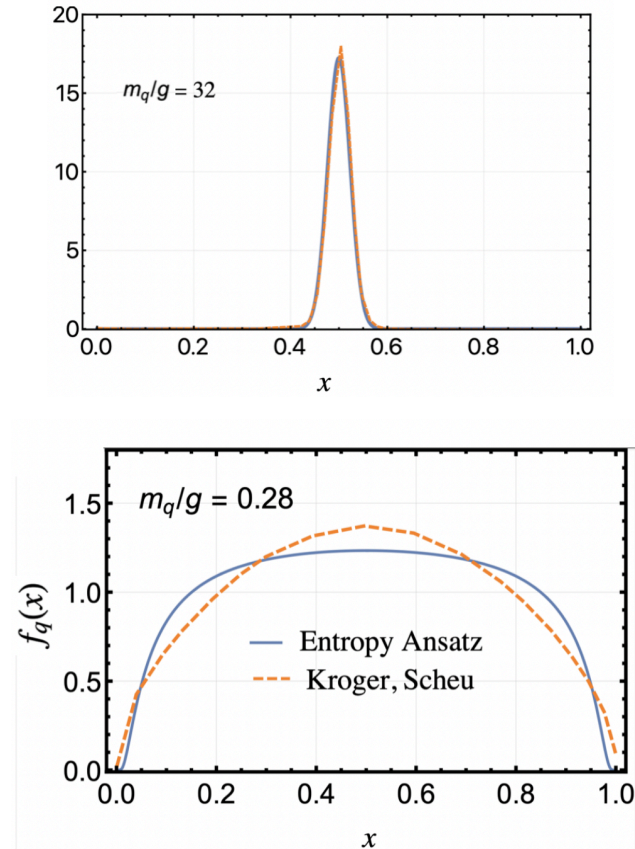
- Minimizing $F = \langle E \rangle - TS$ yields the thermal ansatz. $|p_{i,1-i}|^2 = \mathcal{N} e^{-\frac{m^2}{TP^+} \left(\frac{1}{i} + \frac{1}{1-i} \right)} \equiv e^{-\frac{\langle E \rangle_0}{T}}$

Minimizing $P^+ \langle P^- \rangle = M^2$ allows us to calculate the invariant bound state mass and T.

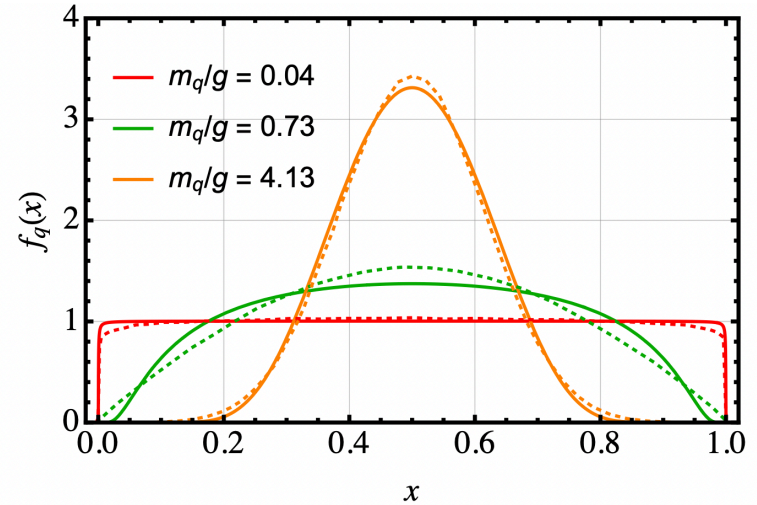
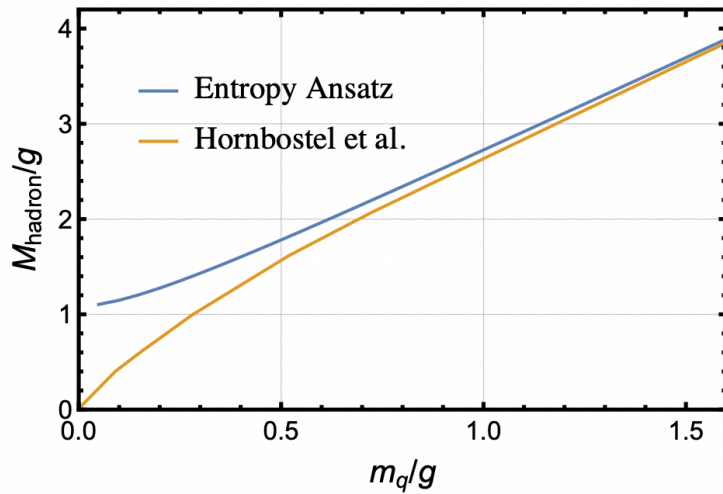
How accurate is a thermal description?



Schwinger model



How good is a thermal description?



't hooft model

- A good description for moderate to weak couplings .
- A Thermal description does not work in the non- perturbative regime.

An Alternative Entanglement measure

- Can we do better than this?
- Perhaps we can use a different measure of entanglement that reduces to the Von Neumann entropy in the weak limit?

- Renyi Entropy.
$$S_\alpha(\rho) = \frac{1}{1-\alpha} \ln \left(\text{Tr}[\rho^\alpha] \right) \quad \alpha \rightarrow 1 \rightarrow \text{Tr}[\rho \ln \rho]$$

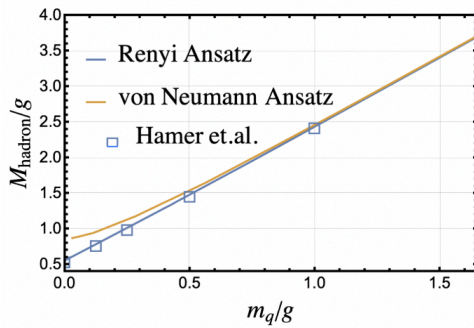
$$F_\alpha = E - TS_\alpha;$$

$$|\psi(x, 1-x)|^2 = \mathcal{N} e^{-\frac{m^2}{TP^+} \left(\frac{1}{x} + \frac{1}{1-x} \right)} \rightarrow \mathcal{N} [x(1-x)]^{\frac{1}{1-\alpha}}$$

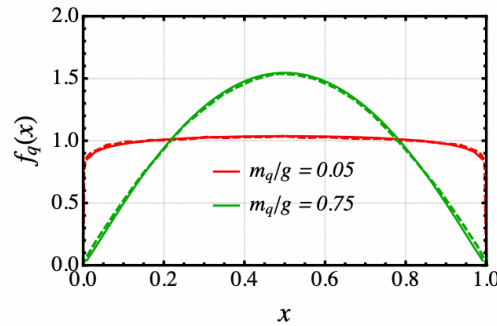
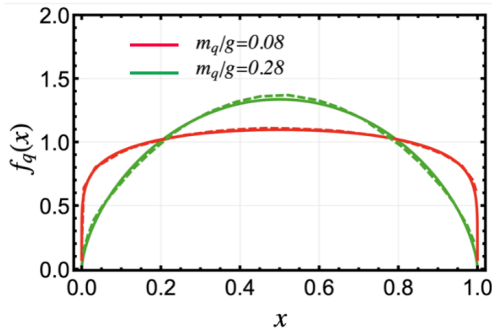
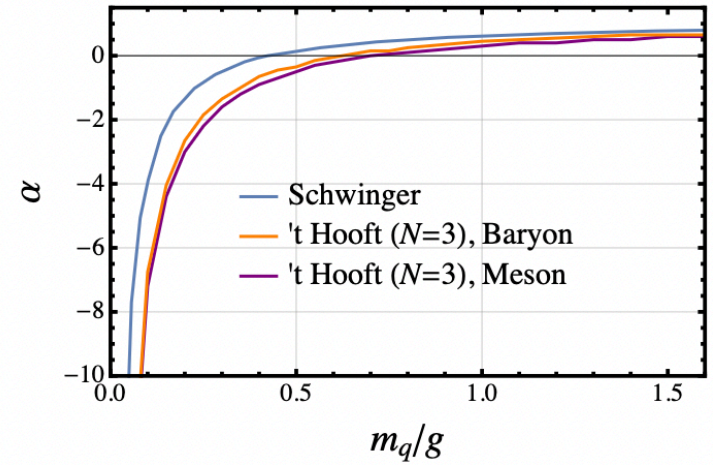
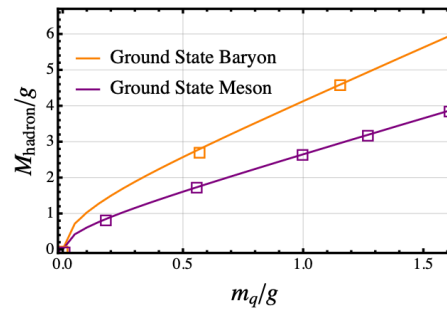
- We still have a single variational parameter α , instead of T.

The Meson mass and PDF

Schwinger model



't hooft model



- The order α becomes negative as we enter the non-perturbative regime.
- The system appears to switch from maximizing high probability configurations to minimizing low probability ones.

The Baryon mass and PDF

- qqq color singlet state .

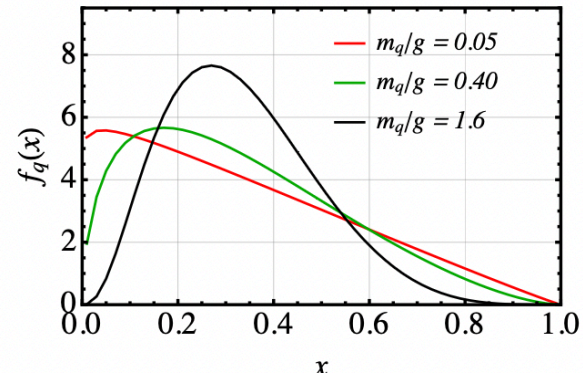
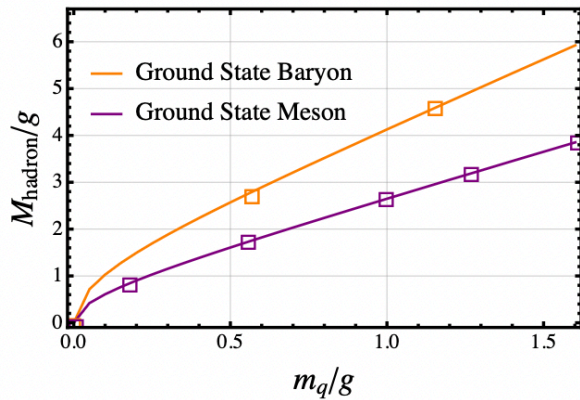
$$|\psi\rangle = \sum_{ijk=0}^1 p_{i,j,k} \delta_{i+j+k,1} |ijk\rangle.$$

- Reduced density matrix is a weighted sum over two-quark matrices with a total momentum fraction $1 - l$

$$\rho_2 = \sum_l \mathcal{N}_l |\psi_l\rangle \langle \psi_l| \text{ density}$$

- Minimize free energy for **each** bi-parton density matrix ,
- Dictated by the 2-2 interaction in the Hamiltonian

$$|p_{i,j,1-i-j}|^2 = \mathcal{N}[i \times j \times (1 - i - j)]^{\frac{1}{1-\alpha}}$$



Outlook

- Does a simple organizing principle persist in higher dimensions?
 - Incorporate renormalization of the theory.
 - Incorporate renormalization of the PDF operator.
 - Contribution from higher occupation Fock states.