

# NNLO jet and beam functions

Automation, distributions and jet clustering

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# Introduction

- Typical factorisation theorem in SCET

$$d\sigma \simeq H(\mu_F) \cdot \prod_i B_i(\mu_F) \otimes \prod_k J_k(\mu_F) \otimes S(\mu_F)$$

- Resummation requires knowledge of anomalous dimensions and matching corrections

$$\underbrace{\Gamma_{\text{Cusp}}, \gamma_H, c_H}_{\text{Observable independent}}$$

$$\underbrace{\gamma_B, \gamma_J, \gamma_S, c_B, c_J, c_S}_{\text{Observable dependent}}$$

- SoftSERVE: [Bell,Rahn,Talbert;19,20]
  - Automated framework to calculate NNLO soft functions
  - Applies to SCET-I and SCET-II observables
- We developed a similar framework for the beam and jet functions at NNLO
  - First application:  $p_T$ -veto for quark beam function [Bell,KB,Das,Wald;22]

# Status-Previous SCET Workshop

## Beam functions

- Framework
  - Focus on quark-to-quark matching kernel
  - SCET-I and SCET-II observables in Mellin-space
- Results
  - Beam Thrust
  - $p_T$ -resummation
  - $p_T$ -veto

## Jet functions

- Framework
  - Complete framework for both quark and gluon jet functions
  - Only considered SCET-I observables
- Results
  - Thrust
  - Angularities
  - Transverse Thrust

## Beam functions

- Framework
  - Focus on quark-to-quark matching kernel
  - SCET-I and SCET-II observables in Mellin-space
  - Complete framework for all beam function channels
  - SCET-I and SCET-II observables in  $x$ -space
- Results
  - Beam Thrust
  - $p_T$ -veto

## Jet functions

- Framework
  - Complete framework for both quark and gluon jet functions
  - Only considered SCET-I observables
  - Further generalisation of measurement function
  - Extended Framework to SCET-II observables
- Results
  - WTA-axis Broadening
  - WTA-axis  $p_T$ -resummation

# Definitions

- Quark beam function

$$\frac{1}{2} \left[ \frac{\not{n}}{2} \right]_{\beta\alpha} B_{q/h}(x, \tau, \mu) = \sum_X \delta \left( (1-x)P_- - \sum_i k_i^- \right) \langle h(P) | \bar{\chi}_\alpha | X \rangle \langle X | \chi_\beta | h(P) \rangle \mathcal{M}(\tau; \{k_i\})$$

- Quark jet function

$$\left[ \frac{\not{n}}{2} \right]_{\beta\alpha} J_q(\tau, \mu) = \frac{1}{\pi} \sum_X (2\pi)^d \delta \left( Q - \sum_i k_i^- \right) \delta^{(d-2)} \left( \sum_i k_i^\perp \right) \langle 0 | \chi_\beta | X \rangle \langle X | \bar{\chi}_\alpha | 0 \rangle \mathcal{M}(\tau; \{k_i\})$$

- Similar definitions for the gluon initiated channels
- SCET-II observables require additional analytical regulator [Becher, Bell; 11]

$$\prod_i \int \frac{d^d k_i}{(2\pi)^d} \left( \frac{\nu}{k_i^- + k_i^+} \right)^\alpha \delta(k_i^2) \theta(k_i^{(0)})$$

# Approach at NLO

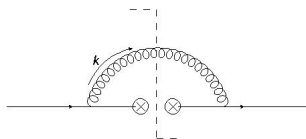
## Beam function

- Parametrisation:

$$\begin{aligned}k_- &= (1-x)P_-, \\ |\vec{k}_\perp| &= k_T, \\ \cos(\theta_k) &= 1 - 2t_k\end{aligned}$$

- Measurement function:

$$\mathcal{M}_1^B(\tau; k) = \exp\left[-\tau k_T \left(\frac{k_T}{(1-x)P_-}\right)^n f(t_k)\right]$$



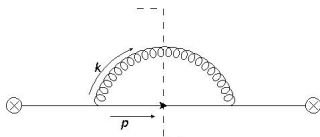
## Jet function

- Parametrisation:

$$\begin{aligned}k_- &= zQ, \quad p_- = \bar{z}Q, \\ |\vec{k}_\perp| &= |\vec{p}_\perp| = k_T, \\ \cos(\theta_k) &= 1 - 2t_k\end{aligned}$$

- Measurement function:

$$\mathcal{M}_1^J(\tau; k, p) = \exp\left[-\tau (k_T)^m \left(\frac{k_T}{Q}\right)^n (z\bar{z})^{-\frac{m+n}{1+n}} f(z, t_k)\right]$$

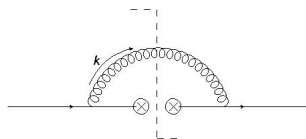


# Approach at NLO

## Beam function

- Master formula:

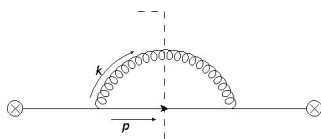
$$B_{i/j}^{(1)}(x, \tau) \simeq \frac{\Gamma\left(\frac{-2\epsilon}{1+n}\right)}{1+n} (1-x)^{-1-\frac{2n\epsilon}{1+n}-\alpha} W_{i/j}^{B_1}(x) \\ \times \int_0^1 dt_k (4t_k \bar{t}_k)^{-1/2-\epsilon} f(t_k) \frac{2\epsilon}{1+n}$$



## Jet function

- Master formula:

$$J_i^{(1)}(\tau) \simeq \frac{\Gamma\left(\frac{-2\epsilon}{m+n}\right)}{m+n} \int_0^1 dz (z\bar{z})^{-1-\frac{2n\epsilon}{1+n}-\alpha} W_i^J(z) \\ \times \int_0^1 dt_k (4t_k \bar{t}_k)^{-1/2-\epsilon} f(z, t_k) \frac{2\epsilon}{m+n}$$



# Approach at NNLO

- At NNLO we need to consider two different contributions

## 1) Real-Virtual contribution

- Matrix element: NLO collinear splitting functions
- Same parametrisation as NLO

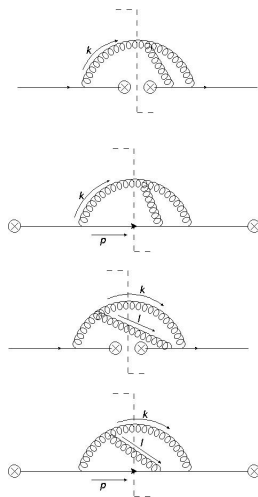
## 2) Real-Real contribution

- Matrix element: LO triple collinear splitting functions
- Parametrisation:

$$q_T = \sqrt{(k_- + l_-)(k_+ + l_+)}, \quad b = \frac{k_T}{l_T}$$

$$a = \frac{k_- l_T}{k_T l_-}, \quad t_{kl} = \frac{1 - \cos \theta_{kl}}{2}$$

- Beam function:  $x = \frac{k_- + l_-}{P_-}$
- Jet function:  $z = \frac{k_- + l_-}{Q}$





# Approach at NNLO

- Measurement function at NNLO

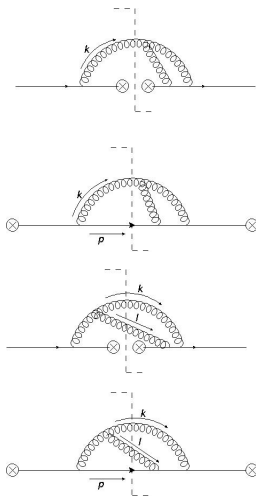
$$\mathcal{M}_2^B(\tau; k, l) = \exp \left[ -\tau q_T \left( \frac{q_T}{(1-x)P_-} \right)^n F(a, b, t_{kl}, t_k, t_l) \right]$$

$$\mathcal{M}_2^J(\tau; k, l, p) = \exp \left[ -\tau (q_T)^m \left( \frac{q_T}{(z\bar{z})^{\frac{m+n}{1+n}} Q} \right)^n F(z, a, b, t_{kl}, t_k, t_l) \right]$$

- Many overlapping singularities remain

- Sector decomposition
- Selector functions
- Non-linear transformations

→ All singularities factorised

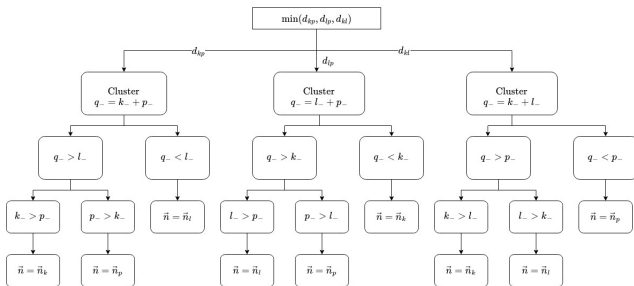


# Winner-take-all recombination scheme

- Select the two closest particles and merge [Bertolini, Chan, Thaler; 13]

$$E_{(ij)} = E_i + E_j, \quad \vec{p}_{(ij)} = E_{(ij)} \left[ \frac{\vec{p}_i}{|\vec{p}_i|} \theta(E_i - E_j) + \frac{\vec{p}_j}{|\vec{p}_j|} \theta(E_j - E_i) \right]$$

- Massless pseudoparticle that points in the direction of the most energetic particle
- This direction is the "winner-take-all" axis
- At leading power the soft function is the same as for the standard jet axis



# SCET-II renormalisation for the jet function

- Additional scale  $\nu \rightarrow$  Additional large logarithms
  - Resum them via collinear anomaly approach [Becher, Neubert; 13]

$$[J_i(\tau, \mu, \nu) \bar{J}_i(\tau, \mu, \nu) S(\tau, \mu, \nu)] \Big|_{\alpha=0} \equiv (Q^2 \bar{\tau}^2)^{-F_i(\tau, \mu)} W_i(\tau, \mu)$$

- Anomaly exponent  $F_i$  and remainder function  $W_i$  fulfill RGEs

$$\frac{dF_i(\tau, \mu)}{d \ln \mu} = 2\Gamma_{\text{Cusp}}$$

$$\frac{dW_i(\tau, \mu)}{d \ln \mu} = 2[\Gamma_{\text{Cusp}} L - \gamma_H] W_i(\tau, \mu)$$

$\rightarrow$  Extraction of  $\gamma_H$ , anomaly coefficient  $d_i$  and  $c_i$  at NNLO

- Considered both quark- and gluon jet function
- Distance measure:  $k_T$ -algorithm
- Measurement: Scalar sum of the transverse momentum projected onto the WTA axis
- Anomaly exponent [Bell,Rahn,Talbert;19,20]

$d_2$	This work	SoftSERVE
$C_F T_F$	-18.737(1)	-18.737
$C_F^2$	-0.116(230)	0
$C_F C_A$	15.950(154)	15.981(1)

$d_2$	This work	SoftSERVE
$C_A T_F$	-18.737(7)	-18.737
$C_A^2$	15.976(15)	15.981(2)

- Remainder function

$c_2$	This work
$C_F T_F$	62.958(3)
$C_F^2$	511.094(298)
$C_F C_A$	-240.436(274)

$c_2$	This work
$T_F^2$	17.435(1)
$C_F T_F$	-85.222(58)
$C_A T_F$	135.965(21)
$C_A^2$	293.955(71)

- Considered only quark jet function
- Distance measure:  $k_T$ -algorithm, anti- $k_T$ -algorithm and Cambridge-Aachen-algorithm
- Measurement: Vector sum of the transverse momentum projected onto the WTA axis
- Anomaly exponent agrees with literature [Bell,Rahn,Talbert;19,20]
- Remainder function comparison with Event2 GSWZ:[Gutierrez-Reyes et al.;20]

$c_2$	This work	GSWZ
$C_F T_F$	12.72(5)	-13.0(3)
$C_F^2$	-13.84(37)	12.2(11)
$C_F C_A$	14.40(80)	-9.3(2)
$k_T$ -algorithm		

$c_2$	This work	GSWZ
$C_F T_F$	12.27(31)	-12.5(3)
$C_F^2$	-25.90(59)	25.3(6)
$C_F C_A$	3.72(131)	-6.3(2)
anti- $k_T$ -algorithm		

$c_2$	This work	GSWZ
$C_F T_F$	12.29(31)	-12.5(3)
$C_F^2$	-25.76(60)	24.5(6)
$C_F C_A$	7.41(132)	-6.7(2)
Cambridge-Aachen-algorithm		

# Beam function approaches

## 1) Our previous approach

- Resolve all distributions in Mellin space

$$\hat{B}_{i/j}(N, \tau) = \int_0^1 dx x^{N-1} B_{i/j}(x, \tau)$$

- Perform all integrals numerically  $\rightarrow$  Similar to the jet function

## 2) Our new approach

- Direct calculation in  $x$ -space
- Expand in terms of distributions

$$\begin{aligned}(1-x)^{-1-2\alpha} f(x) &= \frac{1}{-2\alpha} f(x) \delta(1-x) + \left[ \frac{1}{1-x} \right]_+ f(x) + \dots \\ &= \frac{1}{-2\alpha} f(1) \delta(1-x) + \left[ \frac{1}{1-x} \right]_+ f(1) + \frac{f(x) - f(1)}{1-x} + \dots\end{aligned}$$

# SCET-II renormalisation for the beam function

- Matching on parton distribution functions

$$B_{i/h}(x, \tau, \mu) = \sum_k \int_x^1 \frac{dz}{z} I_{i \leftarrow k}(x/z, \tau, \mu) f_{k/h}(z, \mu)$$

- Renormalisation follows as in the jet function case

$$\begin{aligned} [I_{i \leftarrow k}(x_1, \tau, \mu, \nu) I_{j \leftarrow l}(x_2, \tau, \mu, \nu) S_{ij}(\tau, \mu, \nu)] \Big|_{\alpha=0} &\equiv \\ &(q^2 \bar{\tau}^2)^{-F_{ij}(\tau, \mu)} I_{i \leftarrow k}(x_1, \tau, \mu) I_{j \leftarrow l}(x_2, \tau, \mu) \end{aligned}$$

- Slightly different RGE for the matching kernel  $I_{i \leftarrow k}(x, \tau, \mu)$

$$\frac{dI_{i \leftarrow k}(x, \tau, \mu)}{d \ln \mu} = 2 [\Gamma_{\text{Cusp}} L - \gamma_H] I_{i \leftarrow k}(x, \tau, \mu) - 2 \sum_j I_{i \leftarrow j}(x, \tau, \mu) \otimes P_{j \leftarrow k}(x, \mu)$$

→ Extraction of  $\gamma_H$ , anomaly coefficient  $d_i$  and  $I_{i \leftarrow k}(x)$  at NNLO

- Measurement:  $\omega_2(k, l) = k_+ + l_+$
- Matching kernel can be written as

$$I_{i \leftarrow k}(x) = a_{-1}^{ik} \delta(1-x) + a_0^{ik} \left[ \frac{1}{1-x} \right]_+ + a_1^{ik} \left[ \frac{\log(1-x)}{1-x} \right]_+ \\ + a_2^{ik} \left[ \frac{\log^2(1-x)}{1-x} \right]_+ + a_3^{ik} \left[ \frac{\log^3(1-x)}{1-x} \right]_+ + I_{i \leftarrow k}^{\text{Grid}}(x)$$

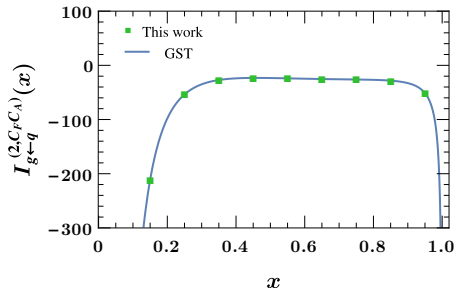
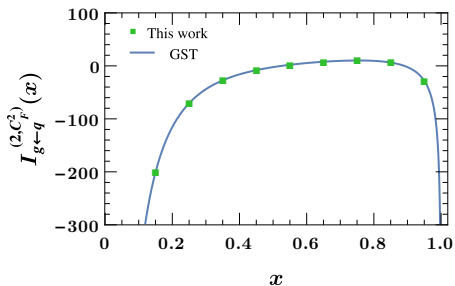
- Focus on the gluon-to-gluon and gluon-to-quark channels

$C_A T_F n_F$	This work	GST
$a_{-1}^{gg}$	-1.238(5)	-1.238
$a_0^{gg}$	3.910(3)	3.910
$a_1^{gg}$	-8.889(1)	-8.889
$a_2^{gg}$	2.667(1)	2.667
$a_3^{gg}$	0	0

$C_A^2$	This work	GST
$a_{-1}^{gg}$	-14.896(29)	-14.897
$a_0^{gg}$	35.027(17)	35.027
$a_1^{gg}$	-9.701(9)	-9.701
$a_2^{gg}$	-7.333(2)	-7.333
$a_3^{gg}$	8.000(1)	8

GST: [Gaunt,Stahlhofen,Tackmann;14]

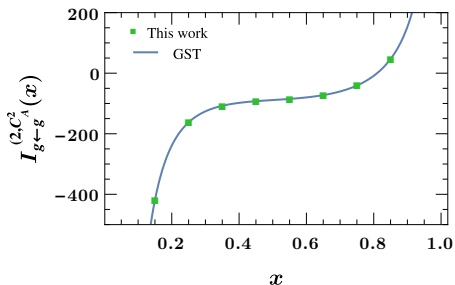
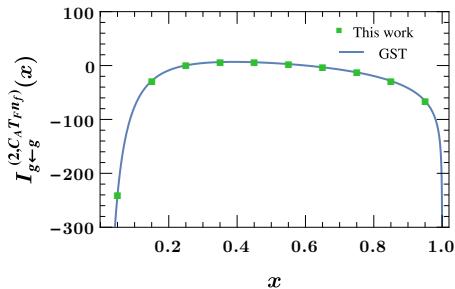
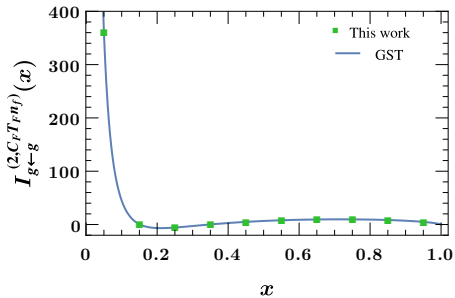




GST: [Gaunt,Stahlhofen,Tackmann;14]

# Beam Thrust- $I_{g \leftarrow g}^{\text{Grid}}(x)$

[Preliminary]



GST: [Gaunt,Stahlhofen,Tackmann;14]

- Measurement:  $\omega_2(k, l) = \theta(\Delta - R) \max(|\vec{k}^\perp|, |\vec{l}^\perp|) + \theta(R - \Delta) |\vec{k}^\perp + \vec{l}^\perp|$   
with  $\Delta = \sqrt{\frac{1}{4} \ln^2 \frac{k^- l^+}{k^+ l^-} + \theta_{kl}^2}$
- Matching kernel can be written as

$$I_{i \leftarrow k}(x) = a_{-1}^{ik} \delta(1-x) + a_0^{ik} \left[ \frac{1}{1-x} \right]_+ + I_{i \leftarrow k}^{\text{Grid}}(x)$$

- Independent calculation by Dingyu Shao
  - Based on computing the difference wrt to a known reference observables
  - Asymmetric phase-space regulator
    - Method is similar to the one from Abreu et al.[2207.07037]

- Matching kernel can be written as

$$I_{i \leftarrow k}(x) = a_{-1}^{ik} \delta(1-x) + a_0^{ik} \left[ \frac{1}{1-x} \right]_+ + I_{i \leftarrow k}^{\text{Grid}}(x)$$

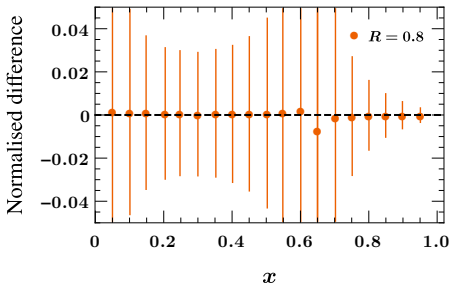
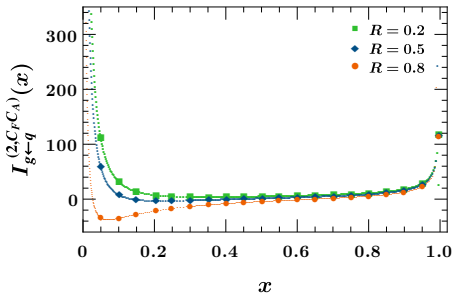
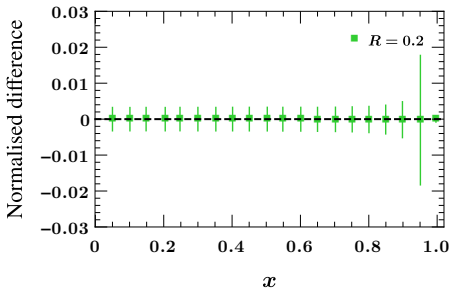
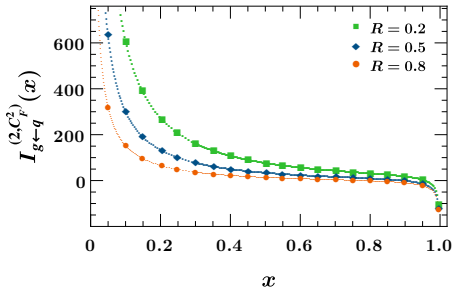
$R = 0.2$	This work	AGMRS
$a_{-1, C_{ATF}}^{gg}$	0.31(1)	0.32
$a_{-1, C_A^2}^{gg}$	-23.32(9)	-23.32
$a_{0, C_{ATF}}^{gg}$	16.44(1)	16.44
$a_{0, C_A^2}^{gg}$	79.02(6)	79.02

$R = 0.5$	This work	AGMRS
$a_{-1, C_{ATF}}^{gg}$	1.38(1)	1.38
$a_{-1, C_A^2}^{gg}$	-16.98(9)	-16.98
$a_{0, C_{ATF}}^{gg}$	11.40(1)	11.40
$a_{0, C_A^2}^{gg}$	43.20(6)	43.20

$R = 0.8$	This work	AGMRS
$a_{-1, C_{ATF}}^{gg}$	1.94(1)	1.94
$a_{-1, C_A^2}^{gg}$	-12.29(9)	-12.29
$a_{0, C_{ATF}}^{gg}$	9.01(1)	9.01
$a_{0, C_A^2}^{gg}$	20.19(6)	20.19

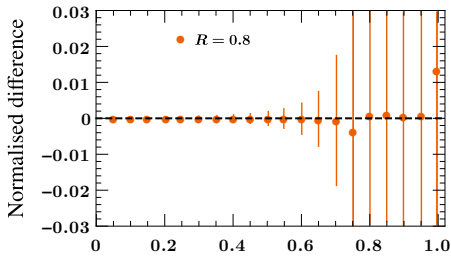
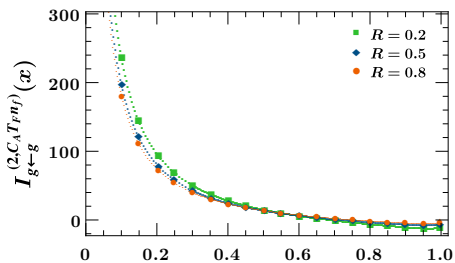
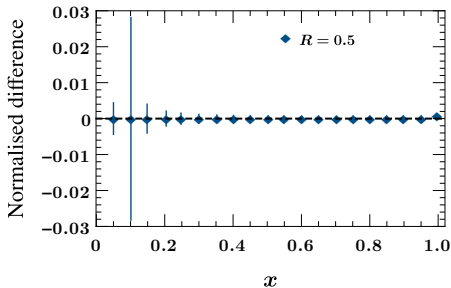
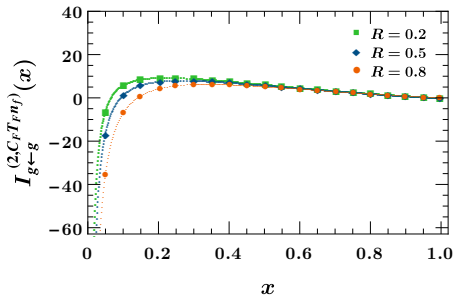
# Comparison $p_T$ -veto $I_{g \leftarrow q}^{\text{Grid}}(x)$

[Preliminary]



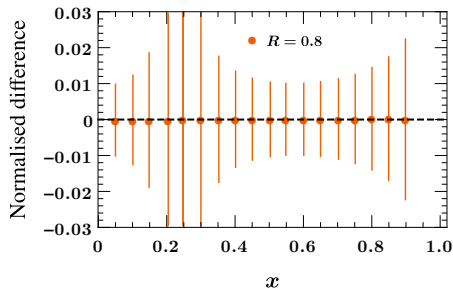
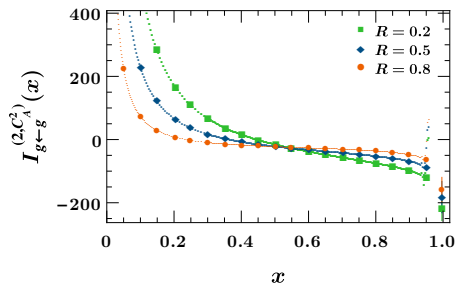
# Comparison $p_T$ -veto $I_{g \leftarrow g}^{\text{Grid}}(x)$

[Preliminary]



# Comparison $p_T$ -veto $I_{g \leftarrow g}^{\text{Grid}}(x)$

[Preliminary]



# Conclusion

- We have reported on progress of our automated setup for the calculation of beam and jet functions at NNLO

## Beam functions

- Setup extended to all matching kernels
- Direct calculation in  $x$ -space
- Beam thrust (SCET-I),  $p_T$ -veto (SCET-II)
- Comparison to Abreu et al.

## Jet functions

- Setup extended to SCET-II observables
- Measurement functions based on jet algorithms
- WTA-axis Broadening, WTA-axis  $p_T$ -resummation

→ We are ready to compute other jet and beam functions!