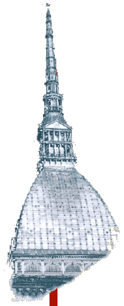
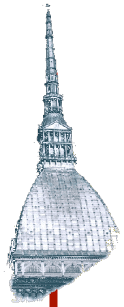


Ultra-Fast Silicon Detector

- The “4D” challenge
- A parameterization of time resolution
- The “Low Gain Avalanche Detectors” project
- Laboratory measurements
- UFSD: LGAD optimized for timing measurements
- WeightField2: a simulation program to optimize UFSD
- First measurements
- Future directions

Nicolo Cartiglia

With
INFN Gruppo V, LGAD group of RD50, FBK and Trento University, Micro-Electronics Turin group
Rome2 - INFN.



Acknowledgement

This research was carried out with the contribution of the Ministero degli Affari Esteri, "Direzione Generale per la Promozione del Sistema Paese" of Italy.

Ministero degli Affari Esteri
e della Cooperazione Internazionale

DIREZIONE GENERALE
PER LA PROMOZIONE DEL SISTEMA PAESE
Unità per la cooperazione scientifica
e tecnologica bilaterale e multilaterale

This work is currently supported by INFN Gruppo V, UFSD project (Torino, Trento Univ., Roma2, Bologna, FBK).

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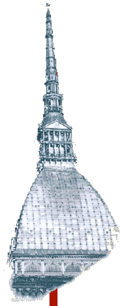
The work at SCIPP was partially supported by the United States Department of Energy, grant DE-FG02-04ER41286.

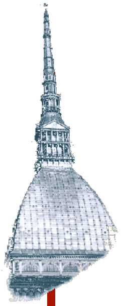


The **4D** challenge

Is it possible to build a detector with concurrent excellent time and position resolution?

Can we provide in the same detector and readout chain:

- **Ultra-fast timing resolution [~ 10 ps]**
- **Precision location information [10's of μm]**

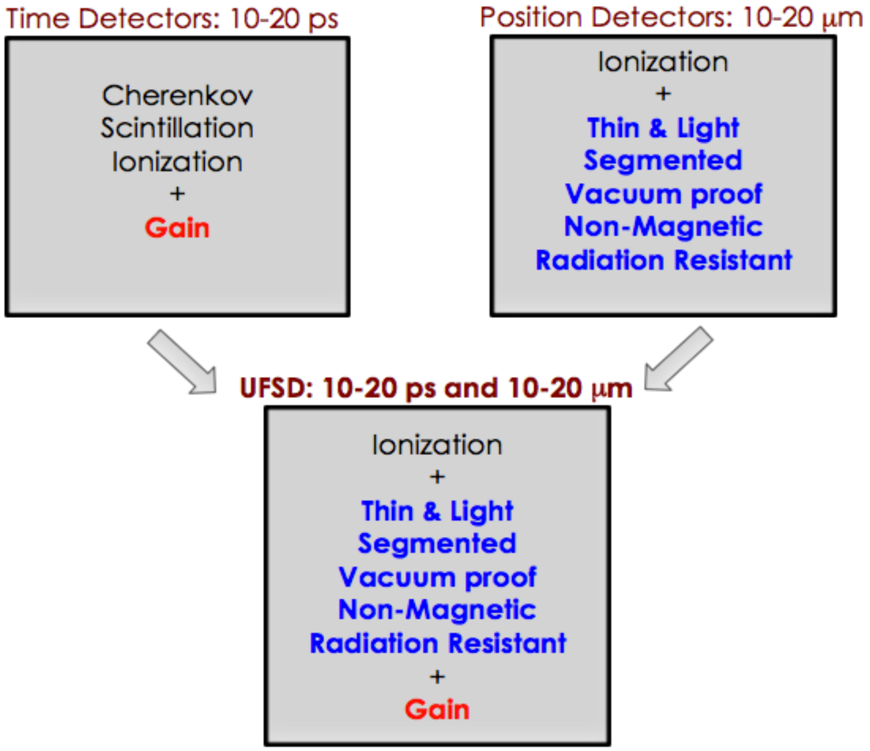


Our path: Ultra-fast Silicon Detectors

Is it possible to build **a silicon detector** with concurrent excellent timing and position resolutions?

Why silicon?

- It already has excellent position resolution
- Very well supported in the community
- Finely segmented
- Thin
- Light
- A-magnetic
- Small
- Radiation resistant



Time Detectors: 10-20 ps

Cherenkov
Scintillation
Ionization
+
Gain

Position Detectors: 10-20 μm

Ionization
+
Thin & Light
Segmented
Vacuum proof
Non-Magnetic
Radiation Resistant

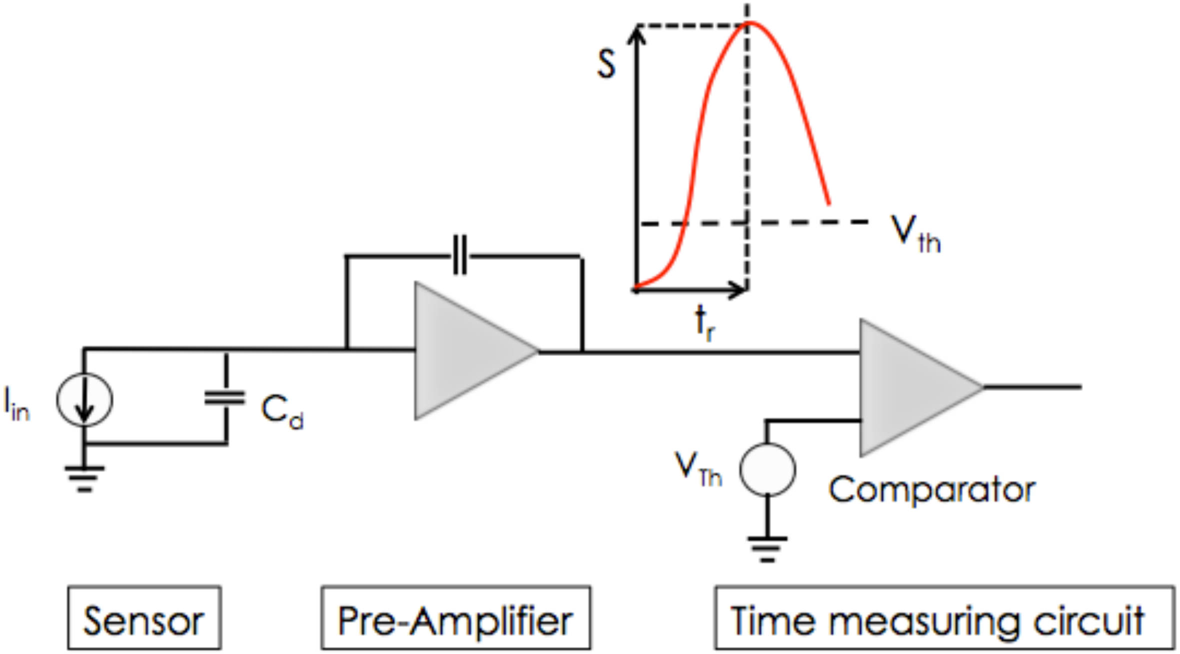
UFSD: 10-20 ps and 10-20 μm

Ionization
+
Thin & Light
Segmented
Vacuum proof
Non-Magnetic
Radiation Resistant
+
Gain

But can it be precise enough?

A time-tagging detector

(a simplified view)



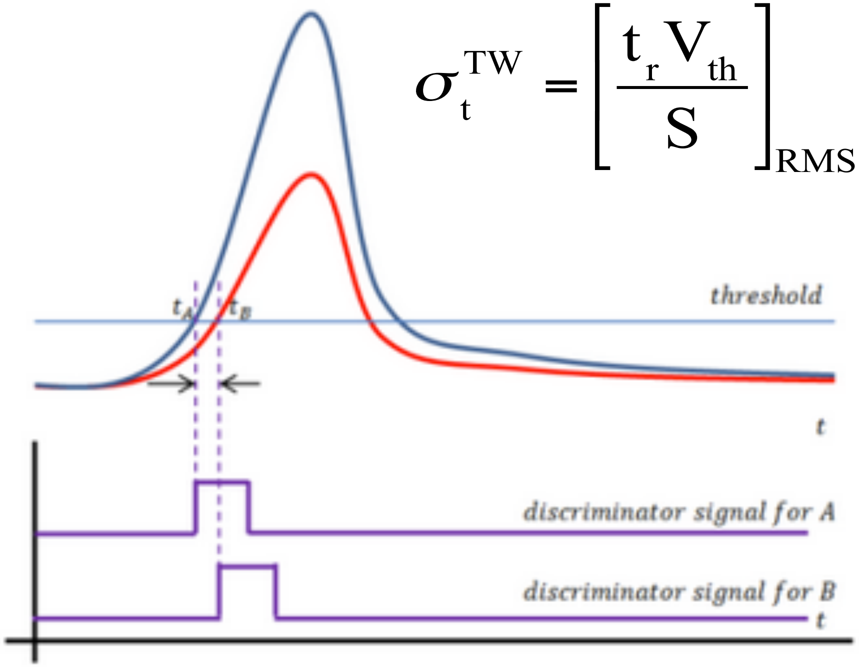
Time is set when the signal crosses the comparator threshold

The timing capabilities are determined by the characteristics of the signal at the output of the pre-Amplifier and by the TDC binning.

Noise source: Time walk and Time jitter

Time walk: the voltage value V_{th} is reached at different times by signals of different amplitude

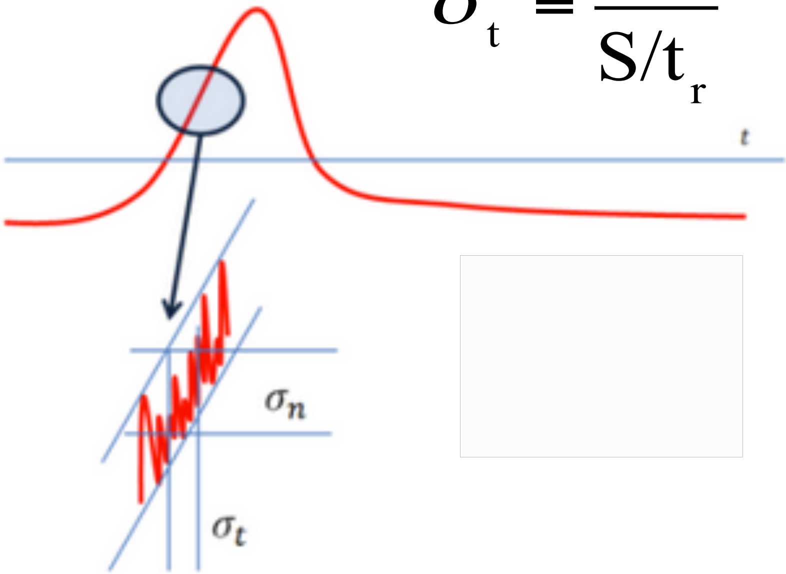
$$\sigma_t^{TW} = \left[\frac{t_r V_{th}}{S}\right]_{RMS}$$



Due to the physics of signal formation

Jitter: the noise is summed to the signal, causing amplitude variations

$$\sigma_t^{J} = \frac{N}{S/t_r}$$



Mostly due to electronic noise

$$\sigma_{Total}^2 = \sigma_{Jitter}^2 + \sigma_{Time\ Walk}^2 + \sigma_{TDC}^2$$

Time Resolution and slew rate

Using the expressions in the previous page, we can write

$$\sigma_t^2 = \left(\left[\frac{V_{th}}{S/t_r}\right]_{RMS}\right)^2 + \left(\frac{N}{S/t_r}\right)^2 + \left(\frac{TDC_{bin}}{\sqrt{12}}\right)^2$$

where:

- S/t_r = dV/dt = slew rate
- N = system noise
- V_{th} = 10 N

Assuming constant noise, to minimize time resolution
we need to maximize the S/t_r term
(i.e. the slew rate dV/dt of the signal)

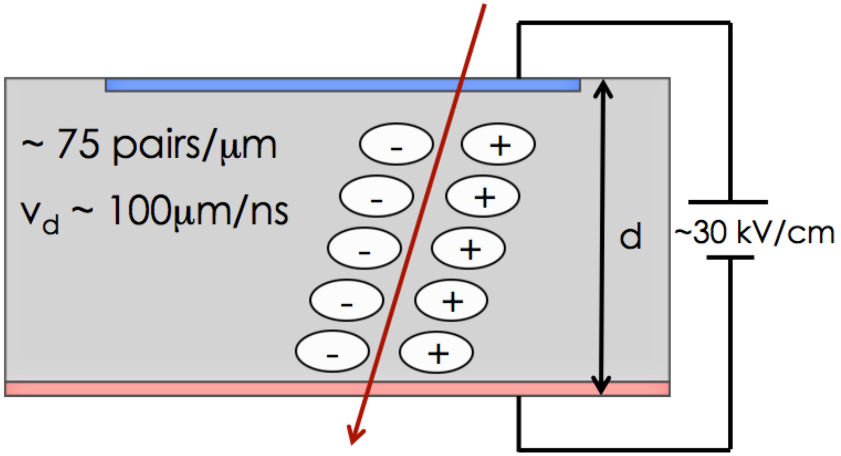
➔ We need large and short signals ➔

Signal formation in silicon detectors

We know we need a large signal, but **how is the signal formed?**

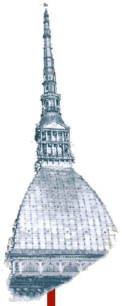
What is controlling the slew rate?

$$\frac{dV}{dt} \propto ?$$



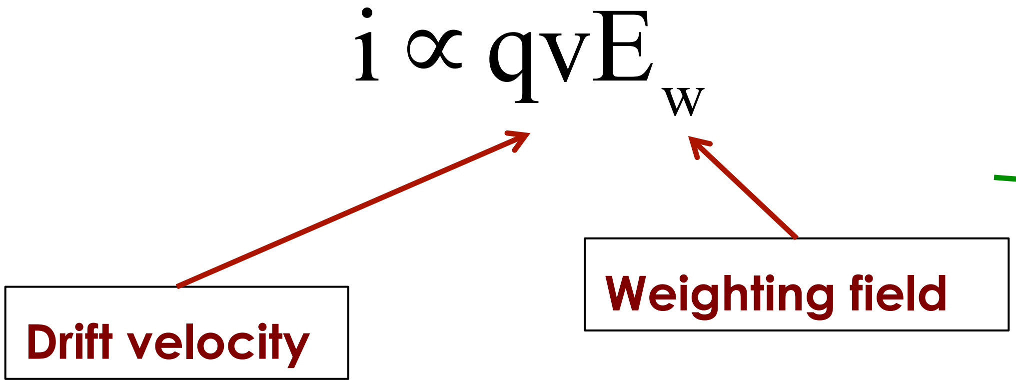
A particle creates charges, then:

- The charges start moving under the influence of an external field
- The motion of the charges induces a current on the electrodes
- The signal ends when the charges reach the electrodes



How to make a **good** signal

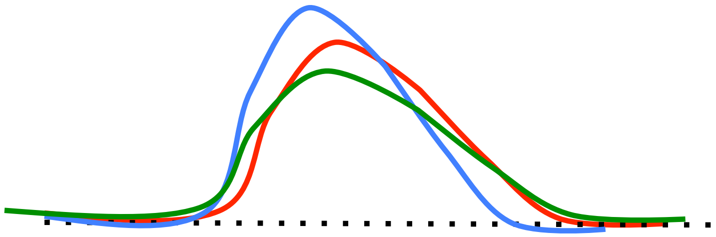
Signal shape is determined by Ramo's Theorem:



$$i \propto qvE_w$$

Drift velocity

Weighting field



A key to good timing is the uniformity of signals:

Drift velocity and **Weighting field** need to be **as uniform as possible**

Drift Velocity

$$i \propto q \boxed{v} E_w$$

- ➔ Highest possible E field to saturate velocity
- ➔ Highest possible resistivity for velocity uniformity

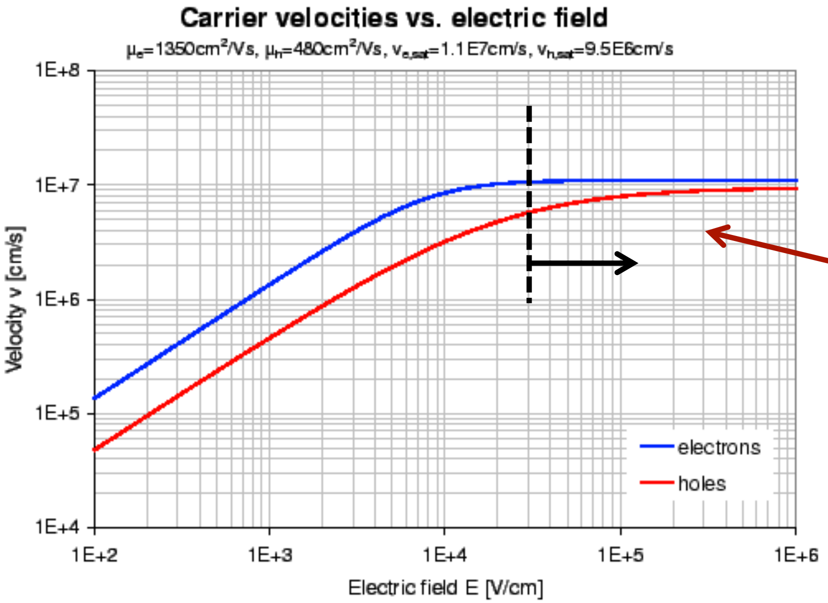


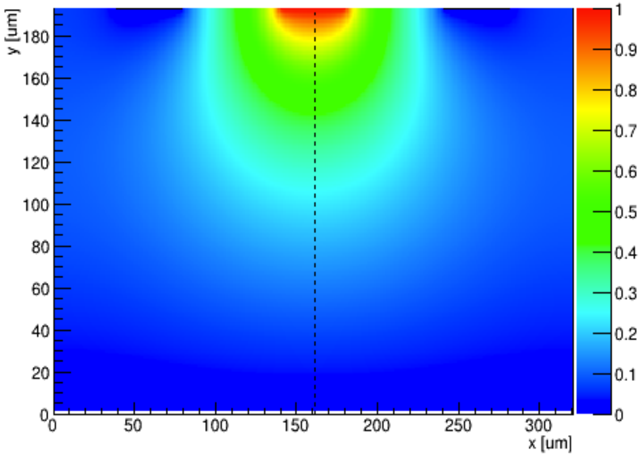
Figure: Electron and hole velocities vs. the electric field strength in silicon.

We want to operate in this regime

Weighting Field: coupling the charge to the electrode

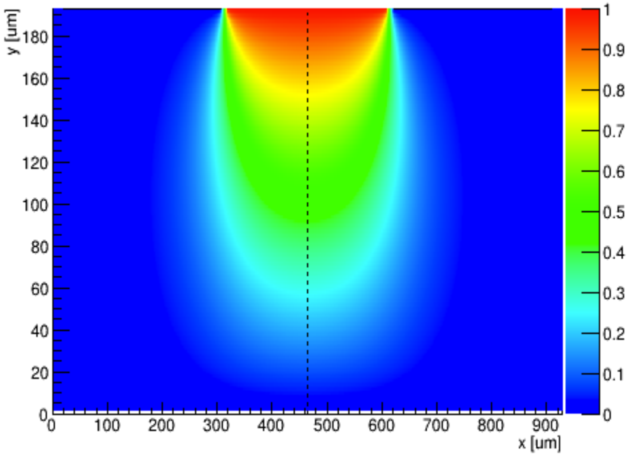
$$i \propto qv\boxed{E_w}$$

Strip: 100 μm pitch, 40 μm width



Bad: almost no coupling away from the electrode

Pixel: 300 μm pitch, 290 μm width



Good: strong coupling almost all the way to the backplane

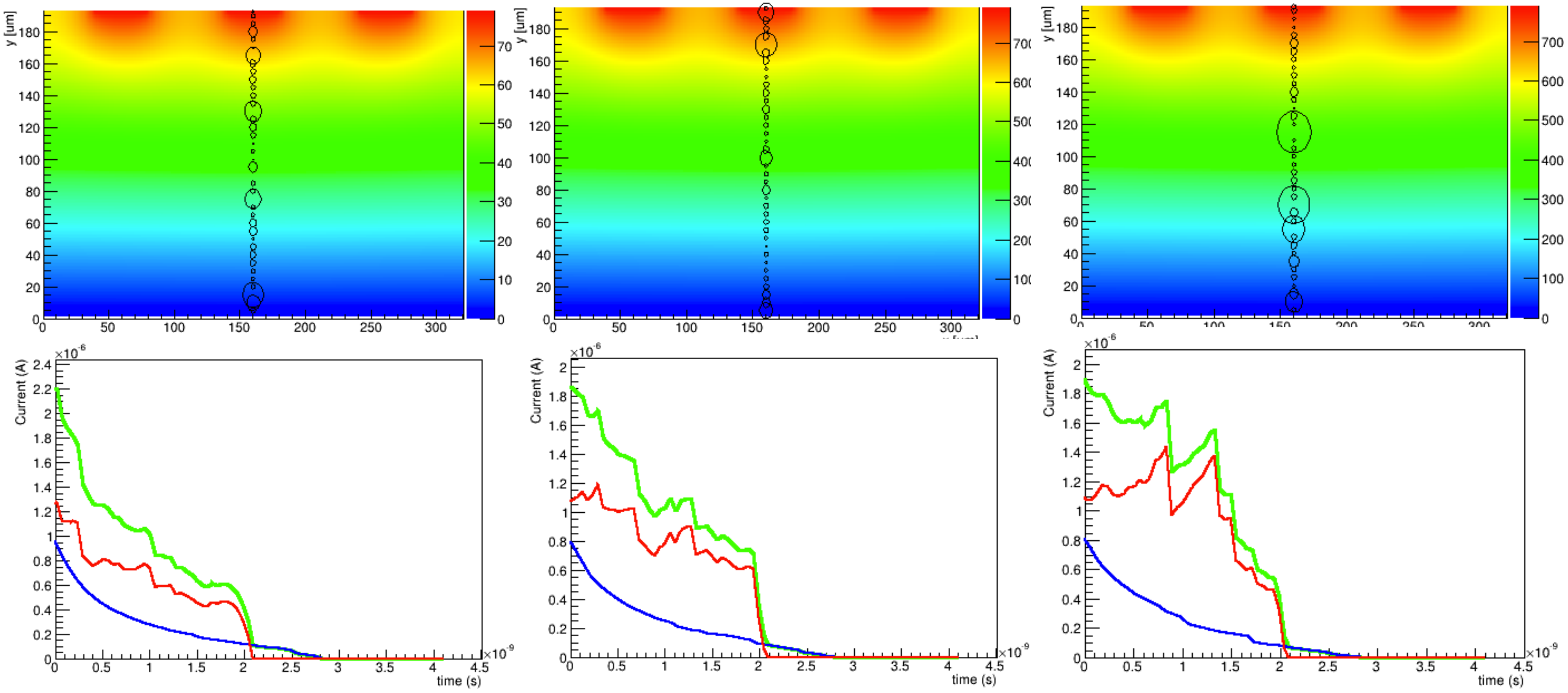
The weighting field needs to be as uniform as possible, so that the coupling is always the same, regardless of the position of the charge

Non-Uniform Energy deposition

Landau Fluctuations cause two major effects:

- Amplitude variations, that can be corrected with time walk compensation
- For a given amplitude, the charge deposition is non uniform.

These are 3 examples of this effect:



What is the signal of one e/h pair?

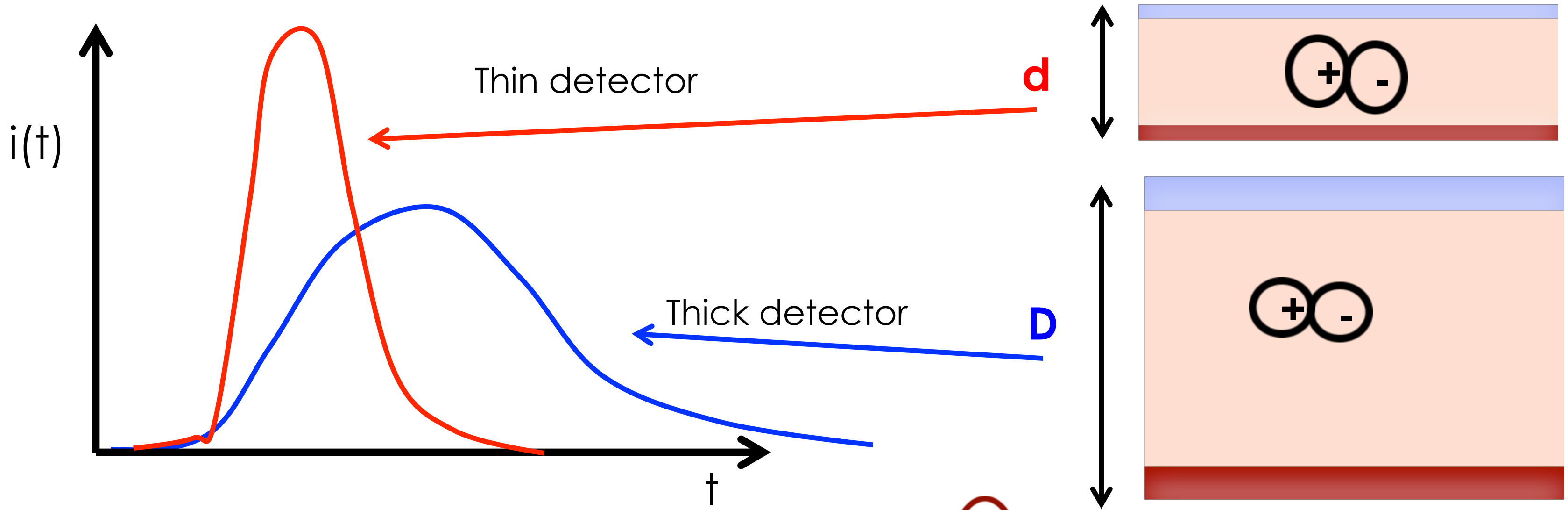
(Simplified model for pad detectors)

Let's consider **one single electron-hole pair.**

The integral of their currents is equal to the electric charge, q:

$$\int [i_{el}(t)+i_h(t)]dt = q$$

However **the shape of the signal depends on the thickness** d:
thinner detectors have higher slew rate



→ One e/h pair generates higher current in thin detectors

$$i \propto qv\boxed{\frac{1}{d}}$$

Weighting field

Large signals from thick detectors?

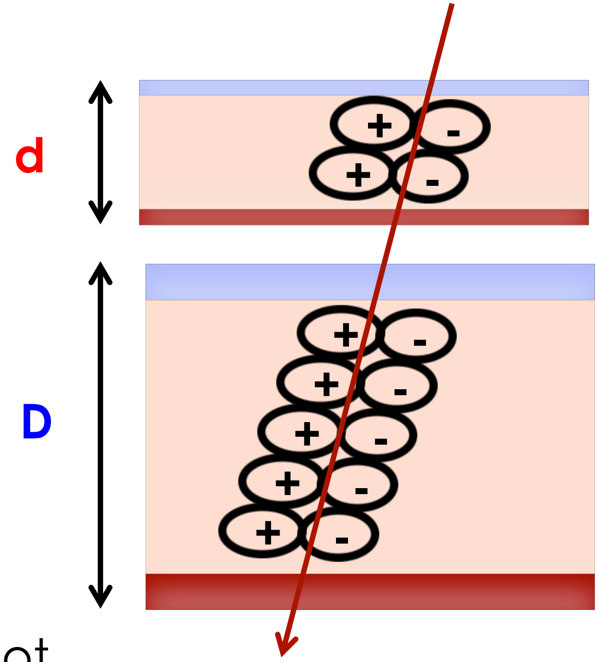
(Simplified model for pad detectors)

Thick detectors have higher number of charges:

$$Q_{tot} \sim 75\ q*d$$

However each charge contributes to the initial current as:

$$i \propto qv\frac{1}{d}$$



The initial current for a silicon detector does not depend on how thick (d) the sensor is:

$$i = Nq\frac{k}{d}v = (75\cancel{d}q)\frac{k}{\cancel{d}}v = 75kqv \sim 1-2*10^{-6}A$$

Number of e/h = 75/micron

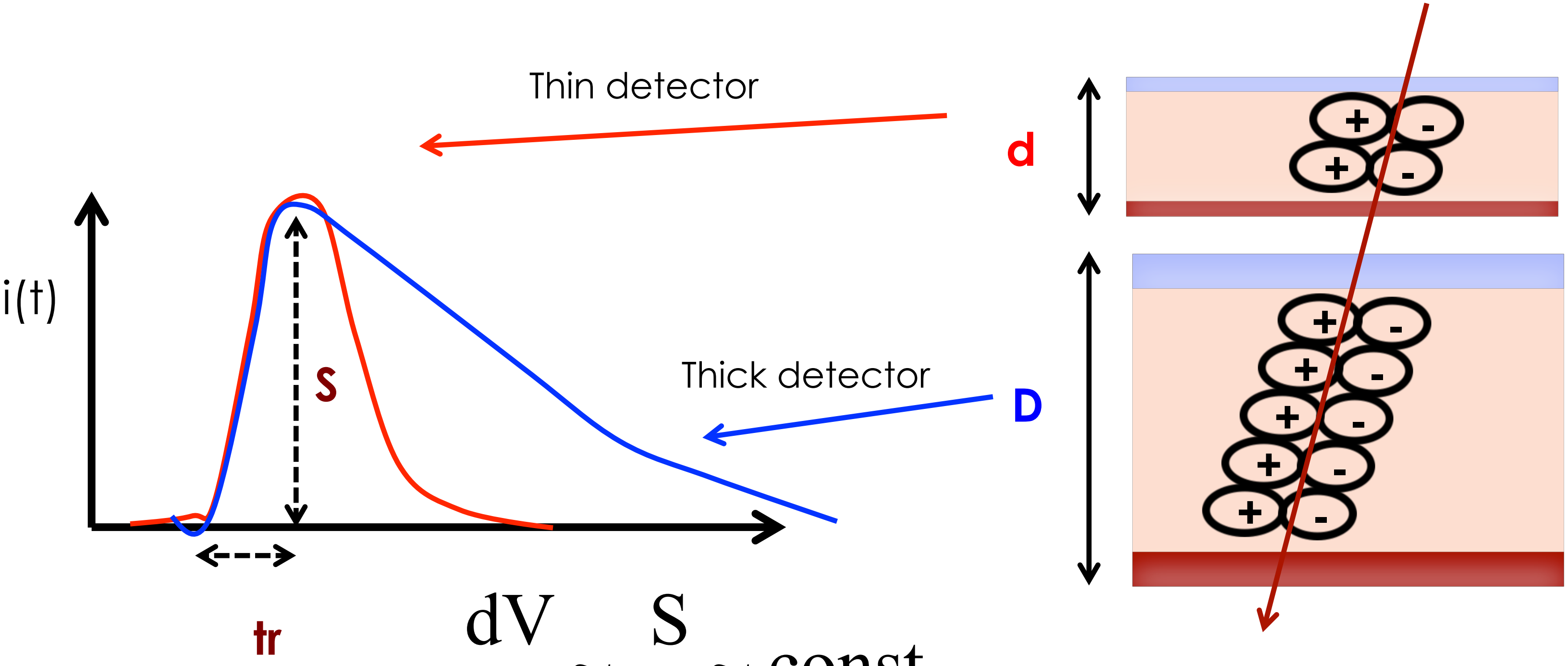
Weighting field

velocity

➔ Initial current = constant

Thin vs Thick detectors

(Simplified model for pad detectors)



$$\frac{dV}{dt} \sim \frac{S}{t_r} \sim \text{const}$$

Thick detectors have longer signals, not higher signals

Best result : NA62, 150 ps on a 300 x 300 micron pixels

To do better, we need to add gain

The "Low-Gain Avalanche Detector" project

Is it possible to manufacture a silicon detector that looks like a normal pixel or strip sensor, but with a much larger signal (RD50)?

- 750 e/h pair per micron instead of 75 e/h?
- Finely Segmented
- Radiation hard
- No dead time
- Very low noise (low shot noise)
- No cross talk
- Insensitive to single, low-energy photon

Many applications:

- Low material budget (30 micron == 300 micron)
- Excellent immunity to charge trapping (larger signal, shorter drift path)
- Very good S/N: 5-10 times better than current detectors
- Good timing capability (large signal, short drift time)

Gain in Silicon detectors

Gain in silicon detectors is commonly achieved in several types of sensors. It's based on the avalanche mechanism that starts in high electric fields: **E ~ 300 kV/cm**

Charge multiplication

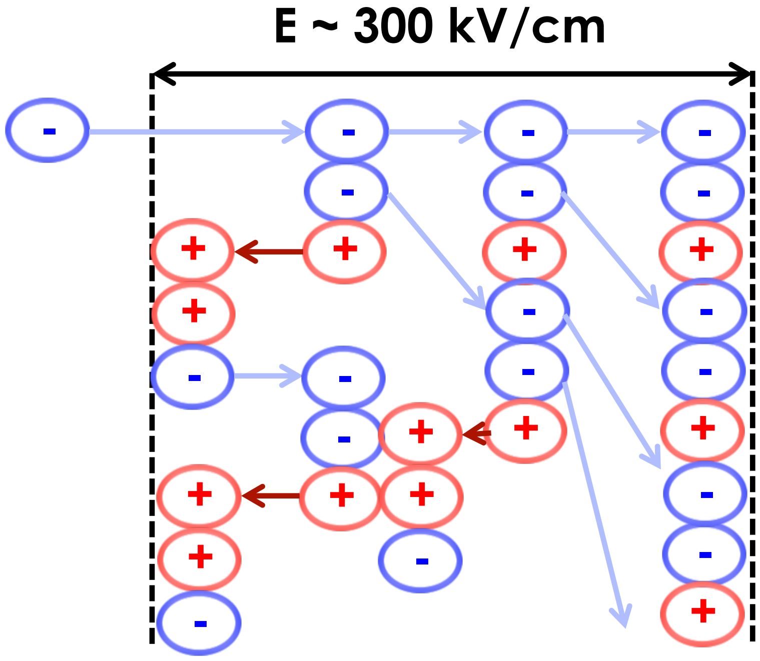
$$N(l) = N_0 \cdot e^{\alpha \cdot l}$$

Gain:

$$G = e^{\alpha\, l}$$

$$\alpha_{e,h}(E) = \alpha_{e,h}(\infty) * \exp\left(-\frac{b_{e,h}}{|E|}\right)$$

α = strong E dependance
α ~ 0.7 pair/μm for electrons,
α ~ 0.1 for holes



Concurrent multiplication of electrons and holes generate very high gain

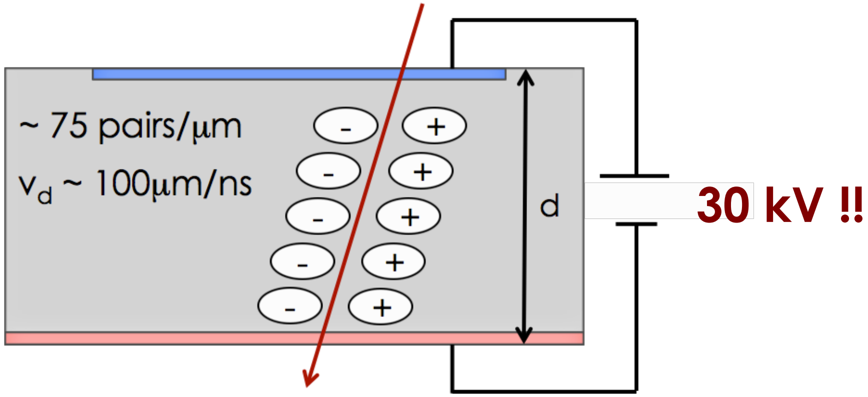
Silicon devices with gain:

- **APD: gain 50-500**
- **SiPM: gain ~ 10^4**

How can we achieve E ~ 300kV/cm?

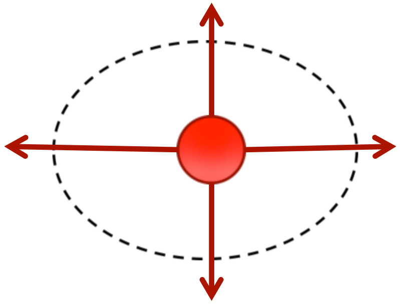
1) Use external bias: assuming a 300 micron silicon detector, we need $\mathbf{V_{bias} = 30\ kV}$

Not possible



2) Use Gauss Theorem:

$$\sum q = 2\pi r * E$$



E = 300 kV/cm ➔ q ~ 10^{16} /cm^3

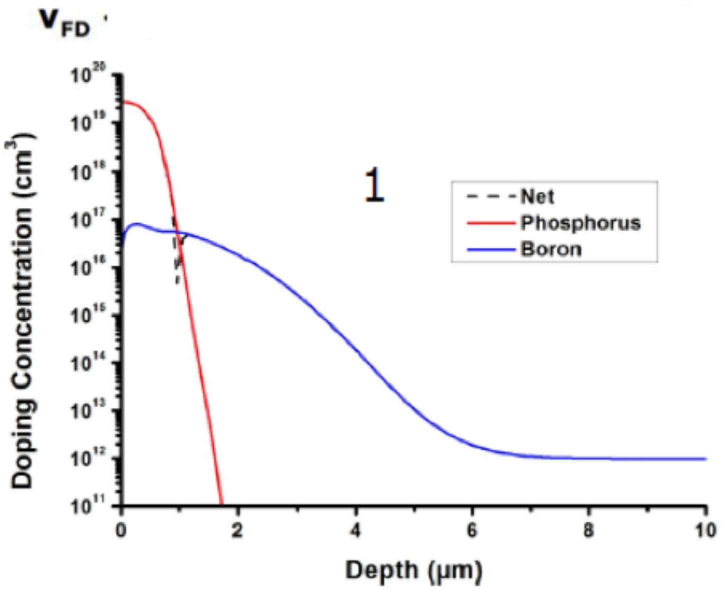
Low Gain Avalanche Detectors (LGADs)

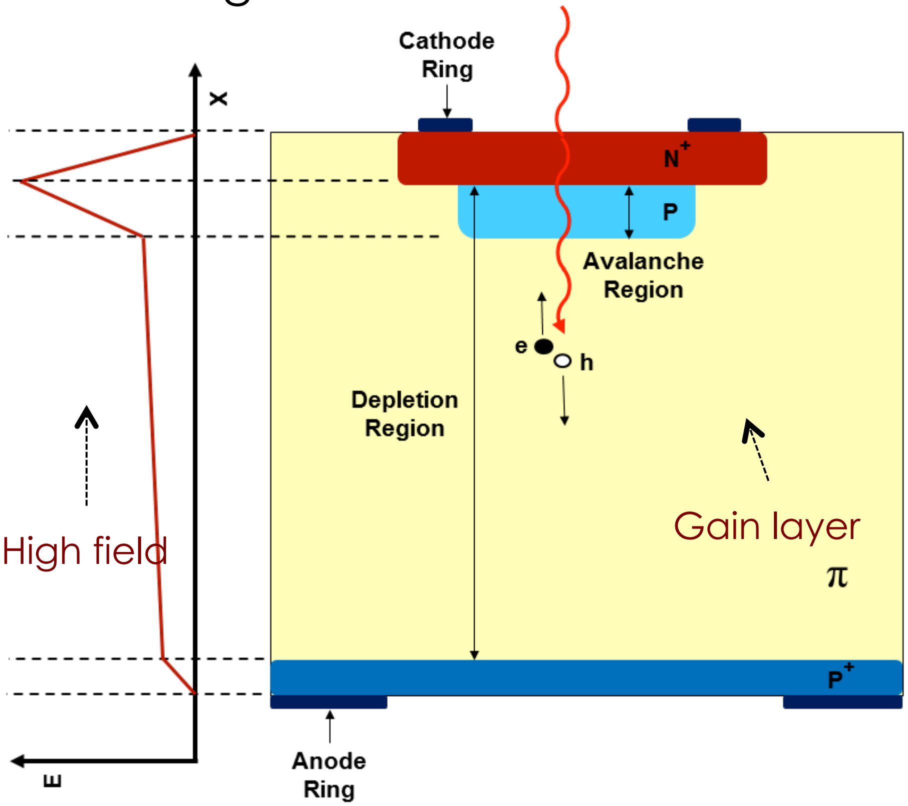
The LGAD sensors, as proposed and manufactured by CNM

(National Center for Micro-electronics, Barcelona):

High field obtained by adding an extra doping layer

E ~ 300 kV/cm, closed to breakdown voltage





Why low gain? Can we use APD or SiPM instead?

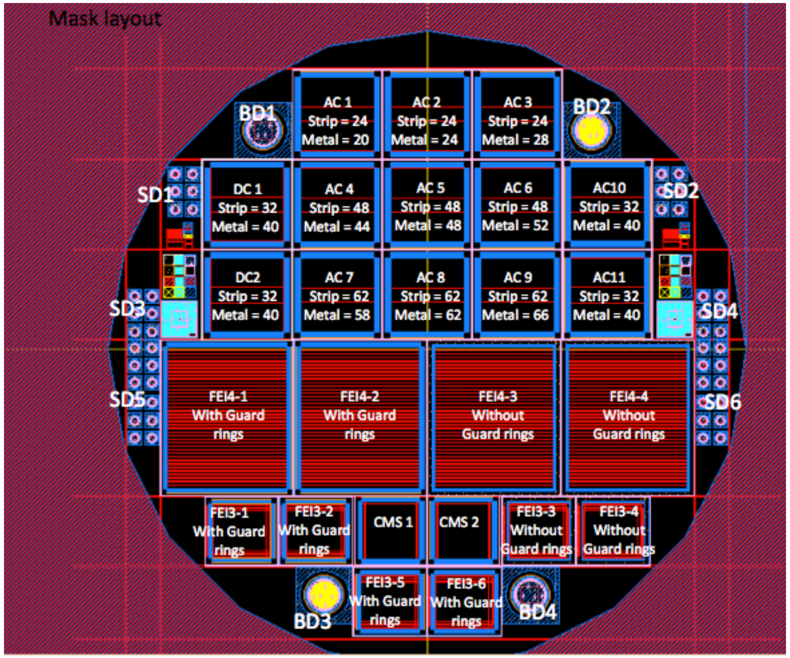
My personal conclusion: I think it's possible to obtain very good timing: APDs, SiPMs have very high gain, so they are excellent in "single shot" timing.

However, we are seeking to obtain something more powerful: a very low noise, finely pixelated device, able to provide excellent timing in any geometry, and also able to work in the presence of many low energy photons without giving fake hits.

These requirements make the use of high gain devices challenging

CNM LGADs mask

CNM, within the RD50 project, manufactured several runs of LGAD, trying a large variety of geometries and designs



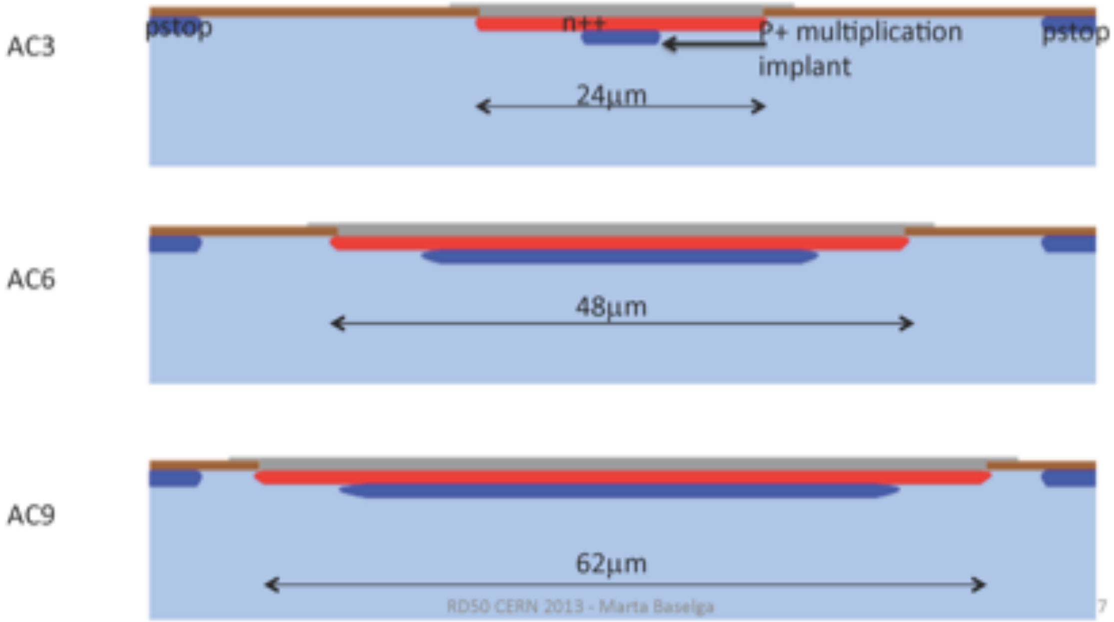
This implant controls the value of the gain

Wafer Number	P-layer Implant (E = 100 keV)	Substrate features	Expected Gain
1-2	1.6 × 10^{13} cm^{-2}	HRP 300 (FZ; ρ >10 KΩ ·cm; <100>; T = 300±10 µm)	2 – 3
3-4	2.0 × 10^{13} cm^{-2}	HRP 300 (FZ; ρ >10 KΩ ·cm; <100>; T = 300±10 µm)	8 – 10
5-6	2.2 × 10^{13} cm^{-2}	HRP 300 (FZ; ρ >10 KΩ ·cm; <100>; T = 300±10 µm)	15
7	(---) PiN Wafer	HRP 300 (FZ; ρ >10 KΩ ·cm; <100>; T = 300±10 µm)	No Gain

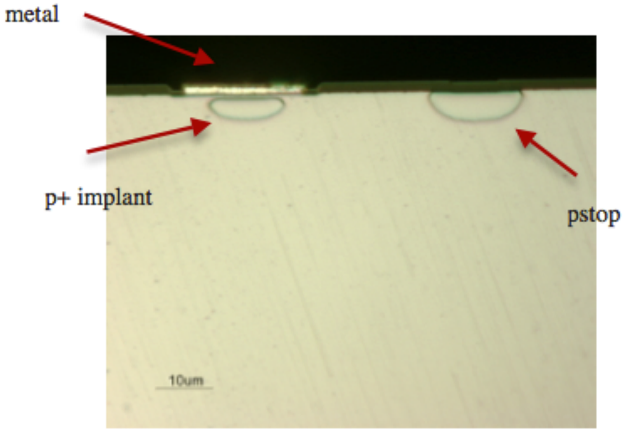
LGADs Pads, Pixels and Strips

The LGAD approach can be extended to any silicon structure, not just pads.

This is an example of LGAD strips



RD50 CERN 2013 - Marta Baselga



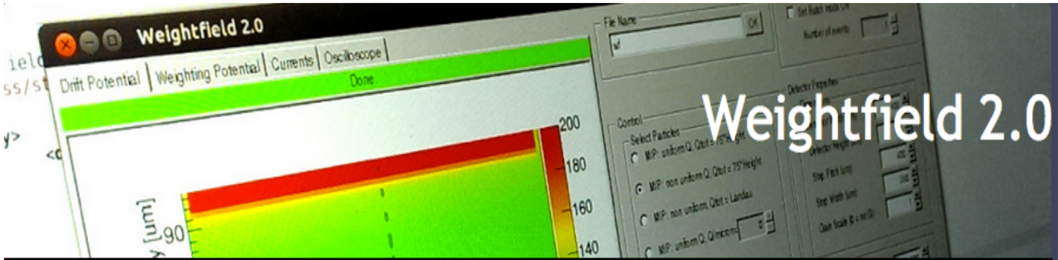
Sensor: Simulation

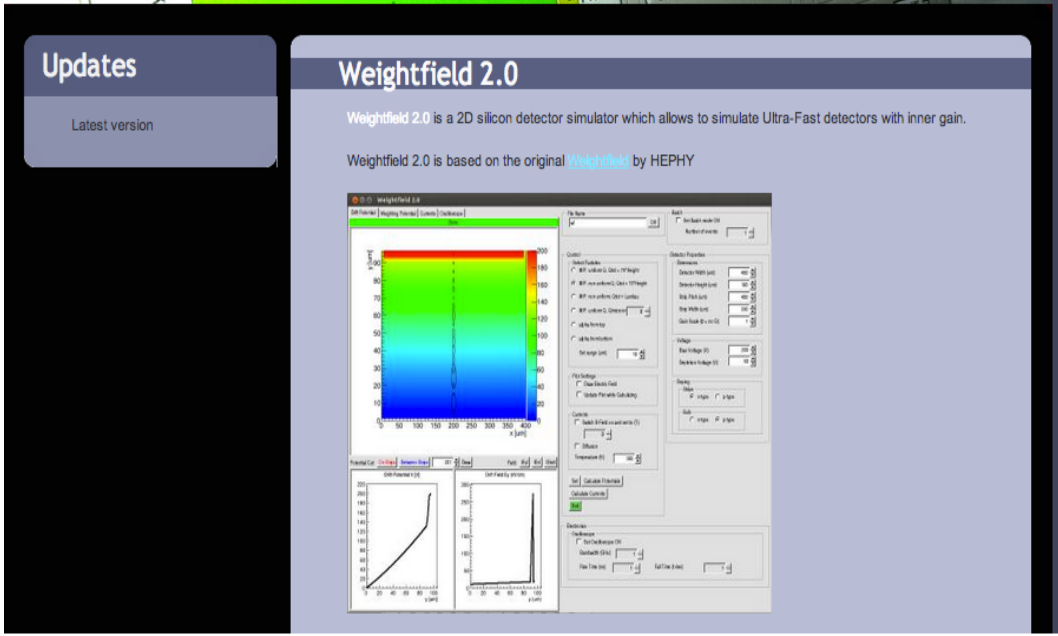
We developed a full sensor simulation to optimize the sensor design

WeightField2, F. Cenna, N. Cartiglia 9th Trento workshop, Genova 2014
Available at http://personalpages.to.infn.it/~cartigli/weightfield2

It includes:

- Custom Geometry
- Calculation of drift field and weighting field
- Currents signal via Ramo's Theorem
- Gain
- Diffusion
- Temperature effect
- Non-uniformdeposition
- Electronics

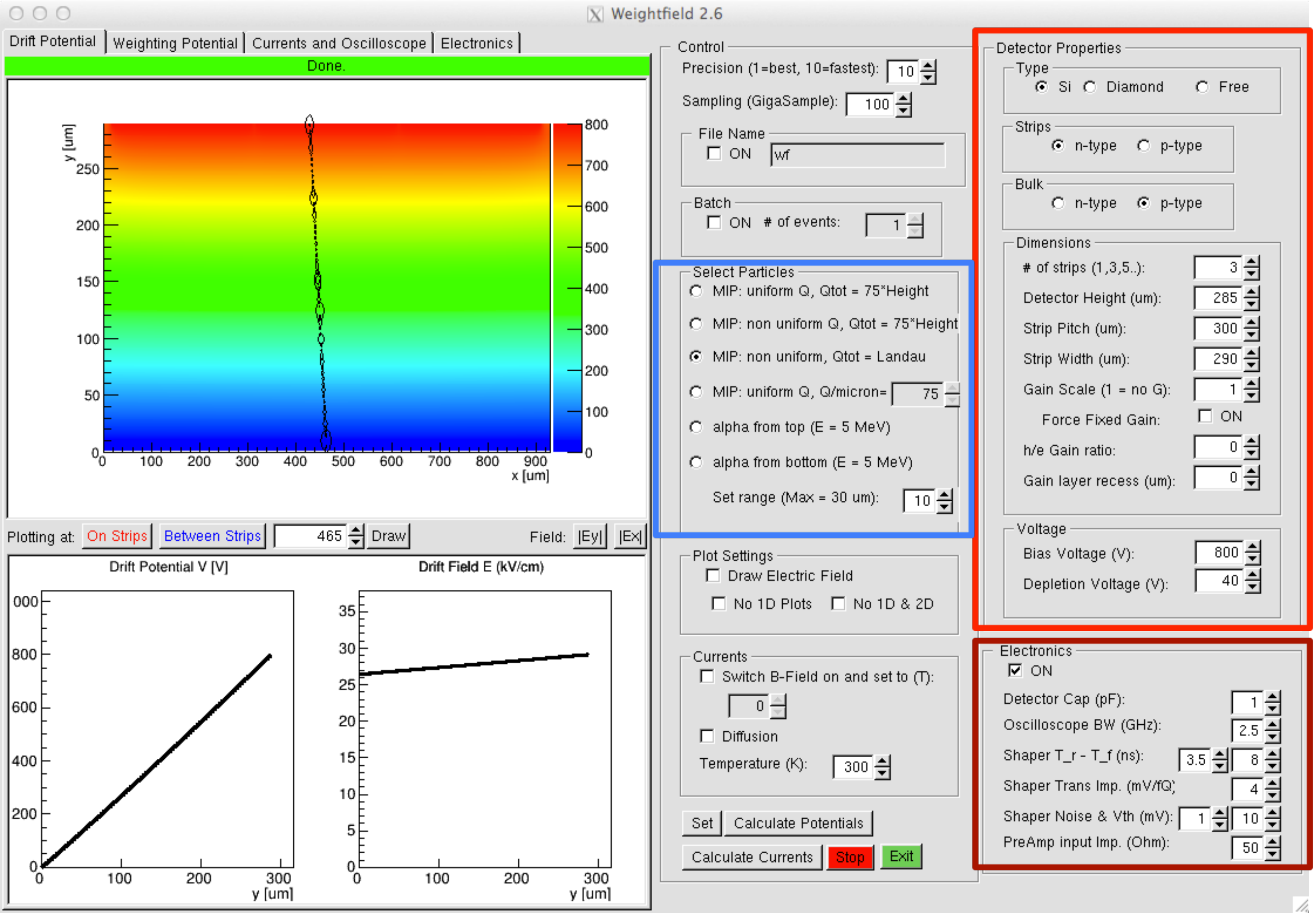




Weightfield 2.0 is a 2D silicon detector simulator which allows to simulate Ultra-Fast detectors with inner gain.

Weightfield 2.0 is based on the original Weightfield by HEPHY

WeightField2: a program to simulate silicon detectors



Plotting at: 465

Control

Precision (1=best, 10=fastest): 10

Sampling (GigaSample): 100

File Name
☐ ON wf

Batch
☐ ON # of events: 1

Select Particles

- ○ MIP: uniform Q, Qtot = 75*Height
- ○ MIP: non uniform Q, Qtot = 75*Height
- ● MIP: non uniform, Qtot = Landau
- ○ MIP: uniform Q, Q/micron= 75
- ○ alpha from top (E = 5 MeV)
- ○ alpha from bottom (E = 5 MeV)

Set range (Max = 30 um): 10

Plot Settings

☐ Draw Electric Field

☐ No 1D Plots ☐ No 1D & 2D

Currents

☐ Switch B-Field on and set to (T): 0

☐ Diffusion

Temperature (K): 300

Detector Properties

Type
● Si ○ Diamond ○ Free

Strips
● n-type ○ p-type

Bulk
○ n-type ● p-type

Dimensions

# of strips (1,3,5..): 3

Detector Height (um): 285

Strip Pitch (um): 300

Strip Width (um): 290

Gain Scale (1 = no G): 1

Force Fixed Gain: ☐ ON

h/e Gain ratio: 0

Gain layer recess (um): 0

Voltage

Bias Voltage (V): 800

Depletion Voltage (V): 40

Electronics

☑ ON

Detector Cap (pF): 1

Oscilloscope BW (GHz): 2.5

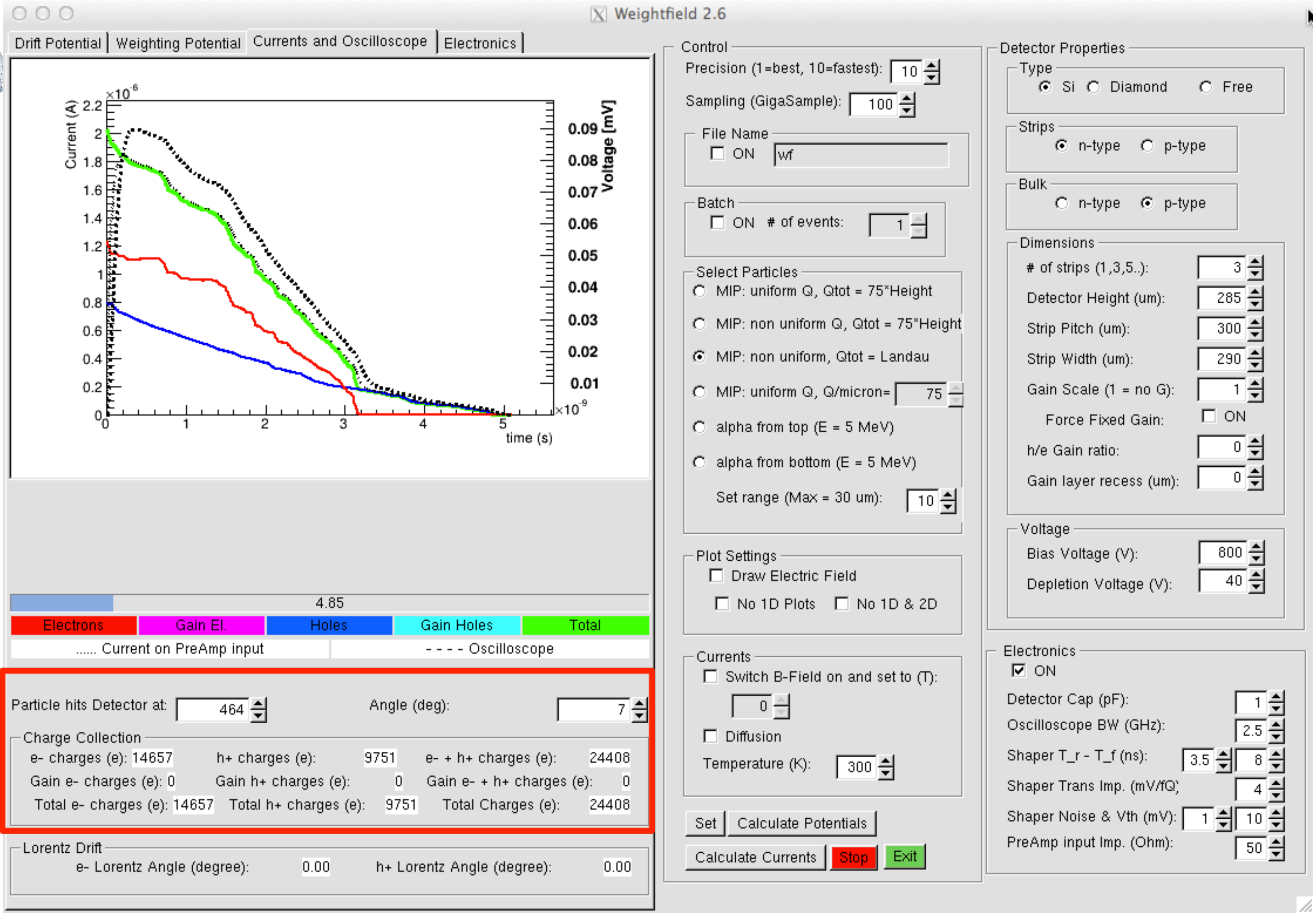
Shaper T_r - T_f (ns): 3.5 8

Shaper Trans Imp. (mV/fQ) 4

Shaper Noise & Vth (mV): 1 10

PreAmp input Imp. (Ohm): 50

WeightField2: output currents



Particle hits Detector at: 464
Angle (deg): 7

Charge Collection

e- charges (e):	14657	h+ charges (e):	9751	e- + h+ charges (e):	24408
Gain e- charges (e):	0	Gain h+ charges (e):	0	Gain e- + h+ charges (e):	0
Total e- charges (e):	14657	Total h+ charges (e):	9751	Total Charges (e):	24408

Lorentz Drift
e- Lorentz Angle (degree): 0.00
h+ Lorentz Angle (degree): 0.00

Control
Precision (1=best, 10=fastest): 10
Sampling (GigaSample): 100

File Name
☐ ON wf

Batch
☐ ON # of events: 1

Select Particles
○ MIP: uniform Q, Qtot = 75*Height
○ MIP: non uniform Q, Qtot = 75*Height
● MIP: non uniform, Qtot = Landau
○ MIP: uniform Q, Q/micron= 75
○ alpha from top (E = 5 MeV)
○ alpha from bottom (E = 5 MeV)
Set range (Max = 30 um): 10

Plot Settings
☐ Draw Electric Field
☐ No 1D Plots ☐ No 1D & 2D

Currents
☐ Switch B-Field on and set to (T): 0
☐ Diffusion
Temperature (K): 300

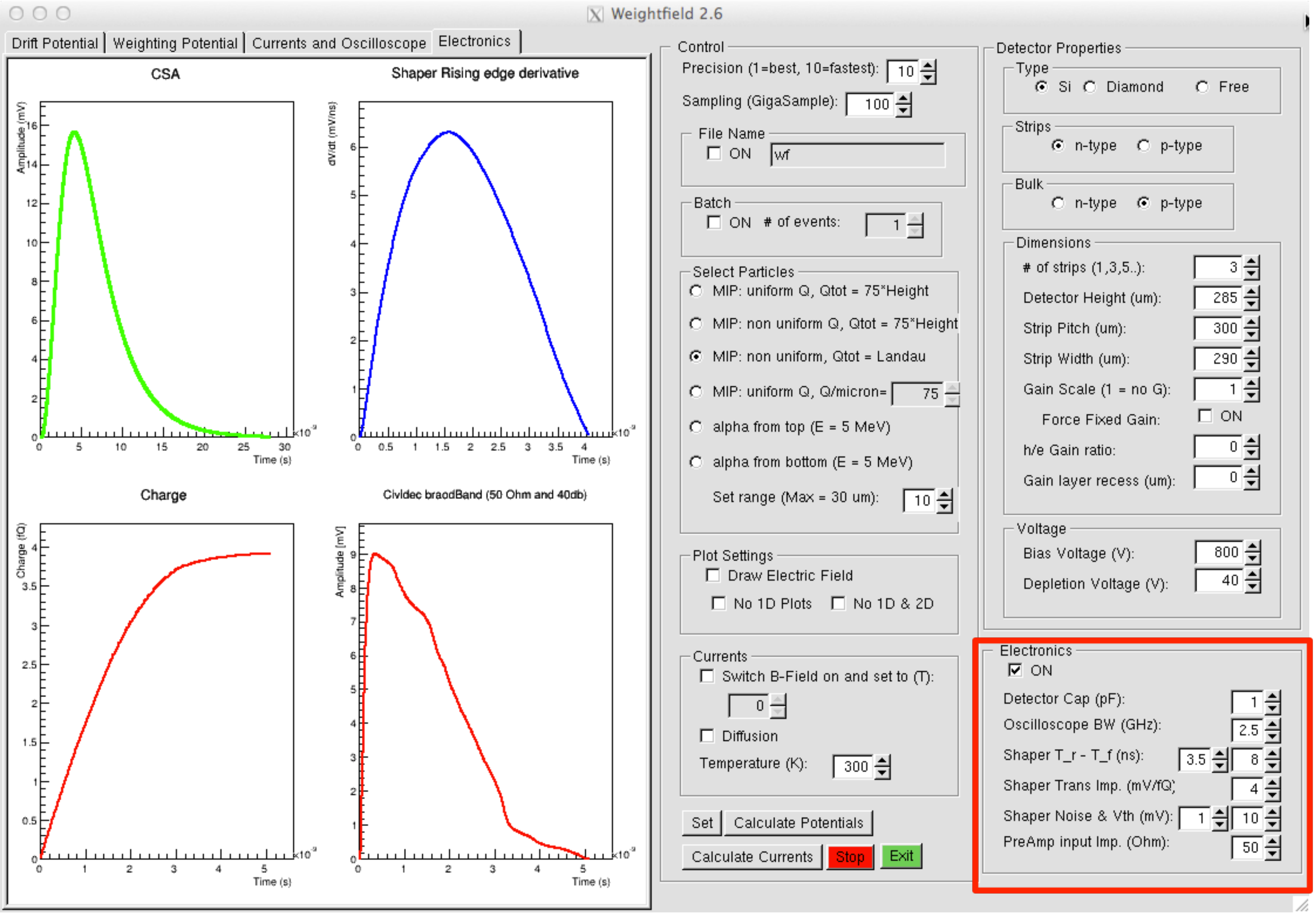
Detector Properties
Type
● Si ○ Diamond ○ Free
Strips
● n-type ○ p-type
Bulk
○ n-type ● p-type

Dimensions
# of strips (1,3,5..): 3
Detector Height (um): 285
Strip Pitch (um): 300
Strip Width (um): 290
Gain Scale (1 = no G): 1
Force Fixed Gain: ☐ ON
h/e Gain ratio: 0
Gain layer recess (um): 0

Voltage
Bias Voltage (V): 800
Depletion Voltage (V): 40

Electronics
☑ ON
Detector Cap (pF): 1
Oscilloscope BW (GHz): 2.5
Shaper T_r - T_f (ns): 3.5 8
Shaper Trans Imp. (mV/fQ) 4
Shaper Noise & Vth (mV): 1 10
PreAmp input Imp. (Ohm): 50

WeightField2: response of the read-out electronics



Control

Precision (1=best, 10=fastest): 10

Sampling (GigaSample): 100

File Name

☐ ON wf

Batch

☐ ON # of events: 1

Select Particles

○ MIP: uniform Q, Qtot = 75*Height

○ MIP: non uniform Q, Qtot = 75*Height

● MIP: non uniform, Qtot = Landau

○ MIP: uniform Q, Q/micron= 75

○ alpha from top (E = 5 MeV)

○ alpha from bottom (E = 5 MeV)

Set range (Max = 30 um): 10

Plot Settings

☐ Draw Electric Field

☐ No 1D Plots ☐ No 1D & 2D

Currents

☐ Switch B-Field on and set to (T): 0

☐ Diffusion

Temperature (K): 300

Detector Properties

Type

● Si ○ Diamond ○ Free

Strips

● n-type ○ p-type

Bulk

○ n-type ● p-type

Dimensions

# of strips (1,3,5..): 3

Detector Height (um): 285

Strip Pitch (um): 300

Strip Width (um): 290

Gain Scale (1 = no G): 1

Force Fixed Gain: ☐ ON

h/e Gain ratio: 0

Gain layer recess (um): 0

Voltage

Bias Voltage (V): 800

Depletion Voltage (V): 40

Electronics

☑ ON

Detector Cap (pF): 1

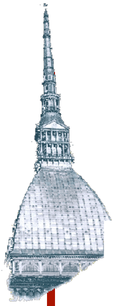
Oscilloscope BW (GHz): 2.5

Shaper T_r - T_f (ns): 3.5 8

Shaper Trans Imp. (mV/fQ) 4

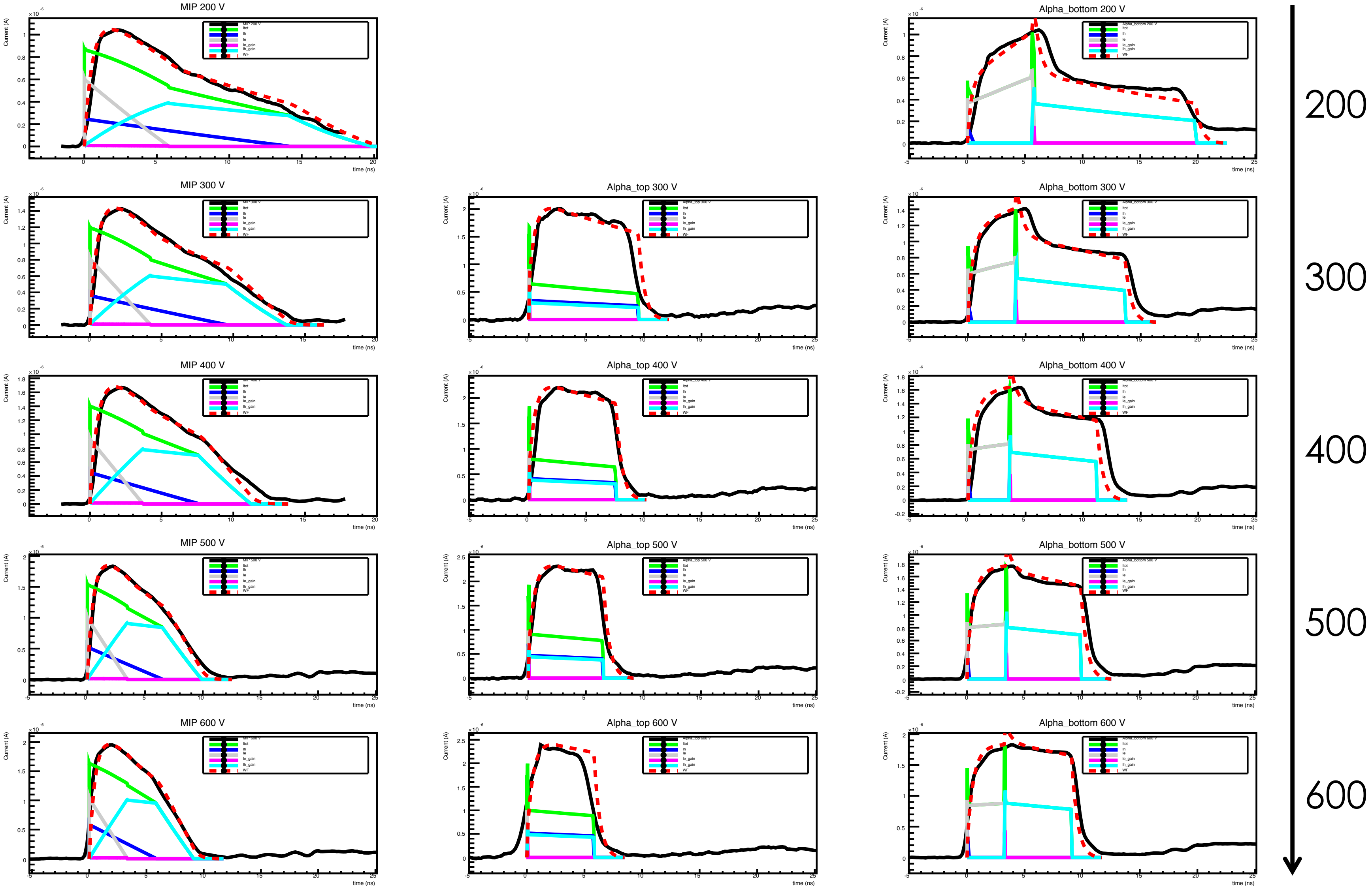
Shaper Noise & Vth (mV): 1 10

PreAmp input Imp. (Ohm): 50

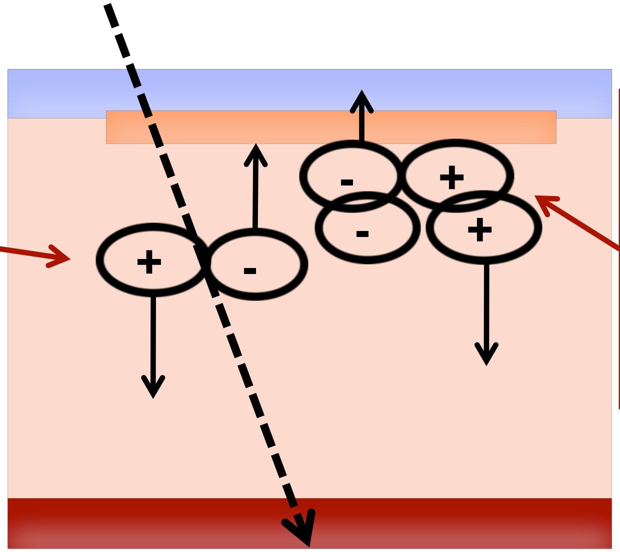


Comparison Data Simulation

MIP **Alpha from Top** **Alpha from bottom** **V bias**



How gain shapes the signal



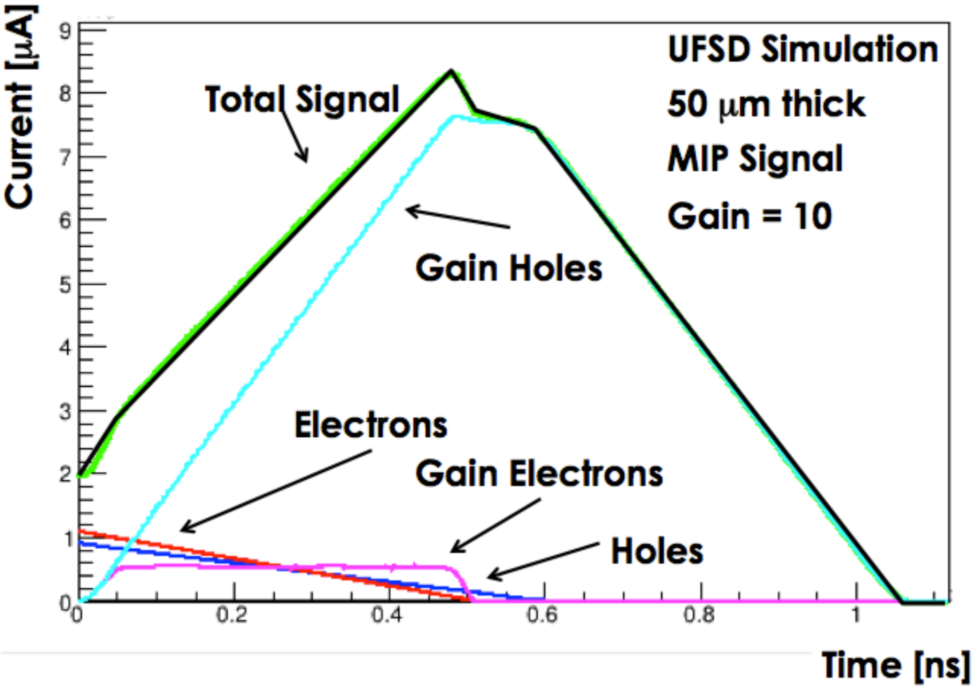
Initial electron, holes

Gain electron:

absorbed immediately

Gain holes:

long drift home



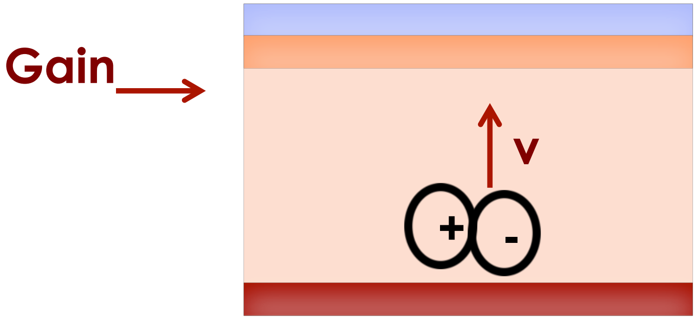
Electrons multiply and produce additional electrons and holes.

- **Gain electrons have almost no effect**
- **Gain holes dominate the signal**

➔ **No holes multiplications**

Interplay of gain and detector thickness

The rate of particles produced by the gain does not depend on *d*
(assuming saturated velocity v_{sat})



Particles per micron

Gain

$$dN_{Gain} \propto 75(v_{sat}dt)G$$

➔ **Constant rate of production**

However the initial value of the **gain current depends on d**
(via the weighing field)

$$di_{gain} \propto dN_{Gain} q v_{sat} \left(\frac{k}{d}\right)$$

➔ **Gain current ~ 1/d**

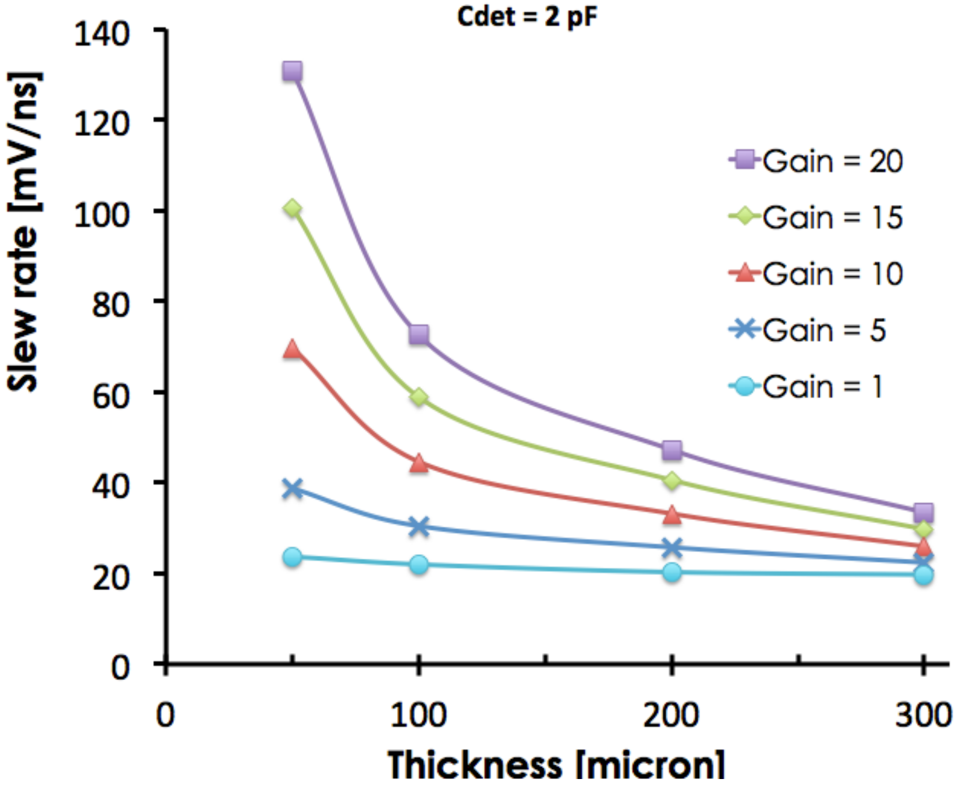
A given value of gain has much more effect on thin detectors

Gain current vs Initial current

$$\frac{di_{gain}}{i} \propto \frac{dN_{Gain}\, q v_{sat} \frac{k}{d}}{k q v_{sat}} = \frac{75(v_{sat} dt)\, G q v_{sat} \frac{k}{d}}{k q v_{sat}} \propto \boxed{\frac{G}{d}} dt$$

!!!

➔ Go thin!!



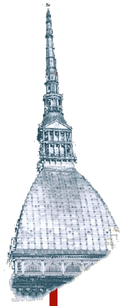
(Real life is a bit more complicated, but the conclusions are the same)

Full simulation

(assuming 2 pF detector capacitance)

300 micron:
~ 2-3 improvement
with gain = 20

Significant improvements in time resolution require thin detectors



Ultra Fast Silicon Detectors

UFSD are LGAD detectors optimized to achieve the best possible time resolution

Specifically:

1. Thin to maximize the slew rate (dV/dt)
2. Parallel plate – like geometries (pixels..) for most uniform weighting field
3. High electric field to maximize the drift velocity
4. Highest possible resistivity to have uniform E field
5. Small size to keep the capacitance low
6. Small volumes to keep the leakage current low (shot noise)

First Measurements and future plans

LGAD laboratory measurements

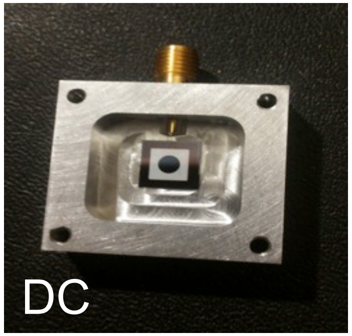
- Doping concentration
- Gain
- Time resolution measured with laser signals

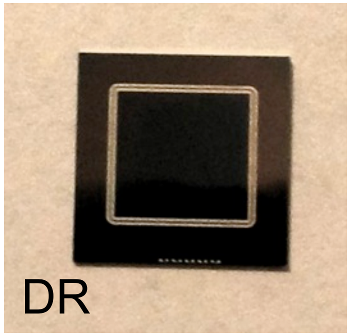
LGAD Testbeam measurements

- Landau shape at different gains
- Time resolution measured with MIPs

LGAD Sensors in Torino

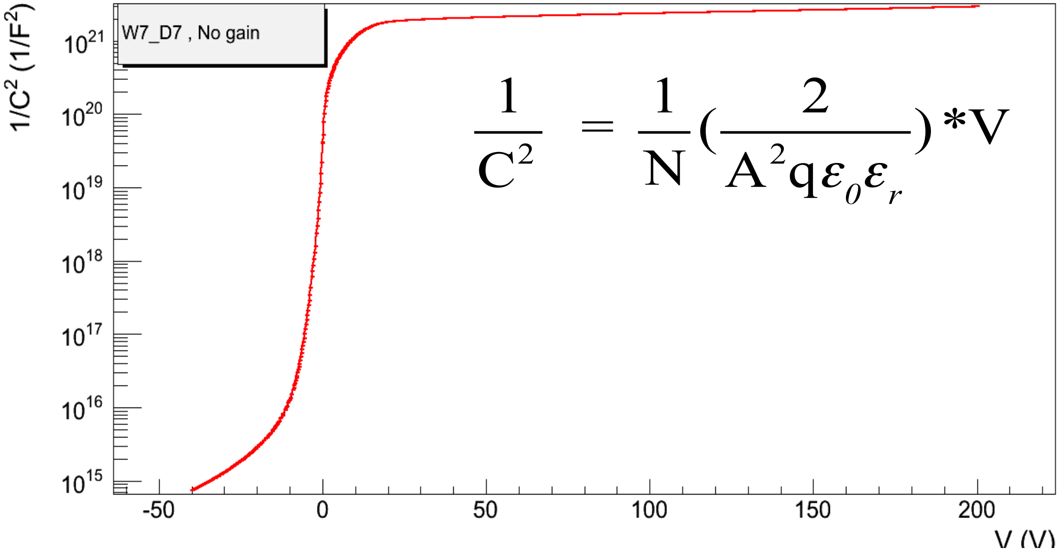
Thickness:
300 µm





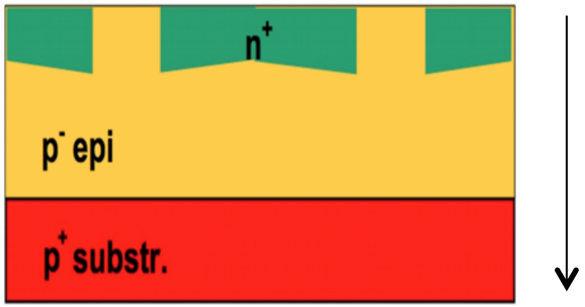
Run	**Sensor**	**P-Layer Implant (E=100 KeV)**	**Gain**	**V_{break}**	**Metal Layer**
6474	W8_B4	?	~ 10	> 500 V	DR
6474	W8_C6	?	~ 10	> 500 V	DC
6474	W9_B6	No implant	No Gain	> 500 V	DR
7062	W1_F3	1.6×10^{13} cm^{-2}	~ 1-2	> 500 V	DR
7062	W3_H5	2.0×10^{13} cm^{-2}	~ 10	> 500 V	DR
7062	W7_D7	No implant	No Gain	> 500 V	DR

Doping profile from CV measurement - I

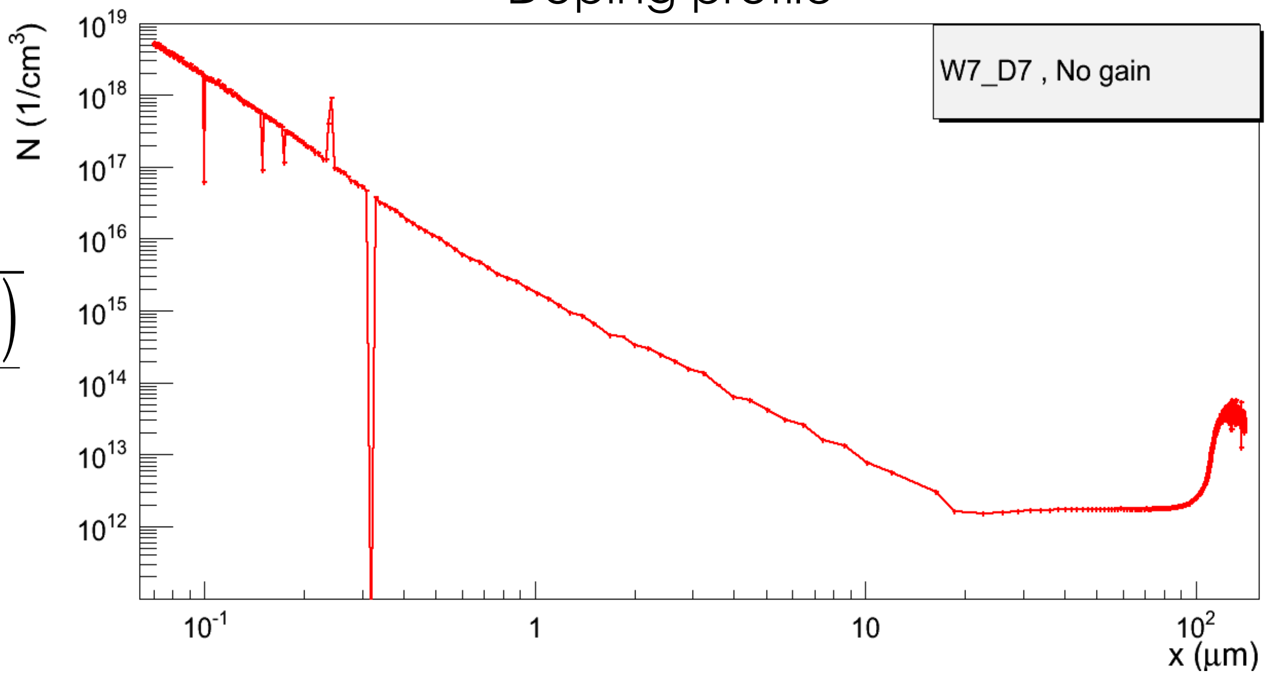


$$\frac{1}{C^2} = \frac{1}{N}\left(\frac{2}{A^2 q\varepsilon_0\varepsilon_r}\right)*V$$

No-gain sensor



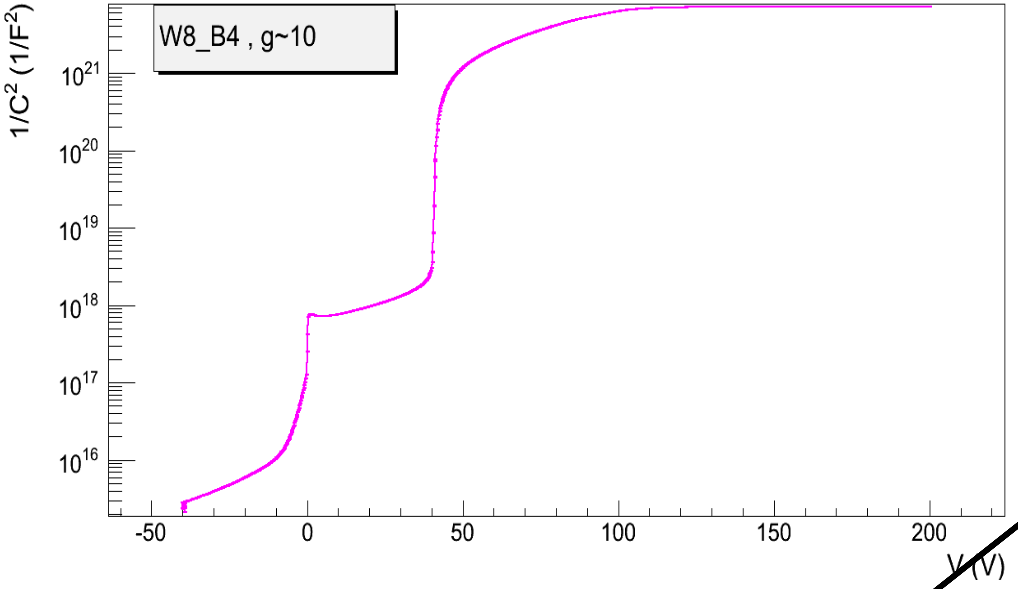
Doping profile



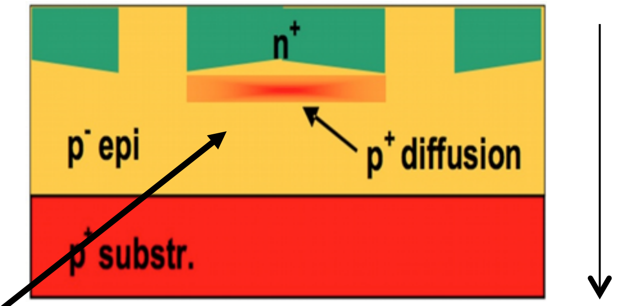
$$N = \frac{2}{q\varepsilon_0\varepsilon_r A^2 \frac{d\left(1/C^2\right)}{dV}}$$

Doping

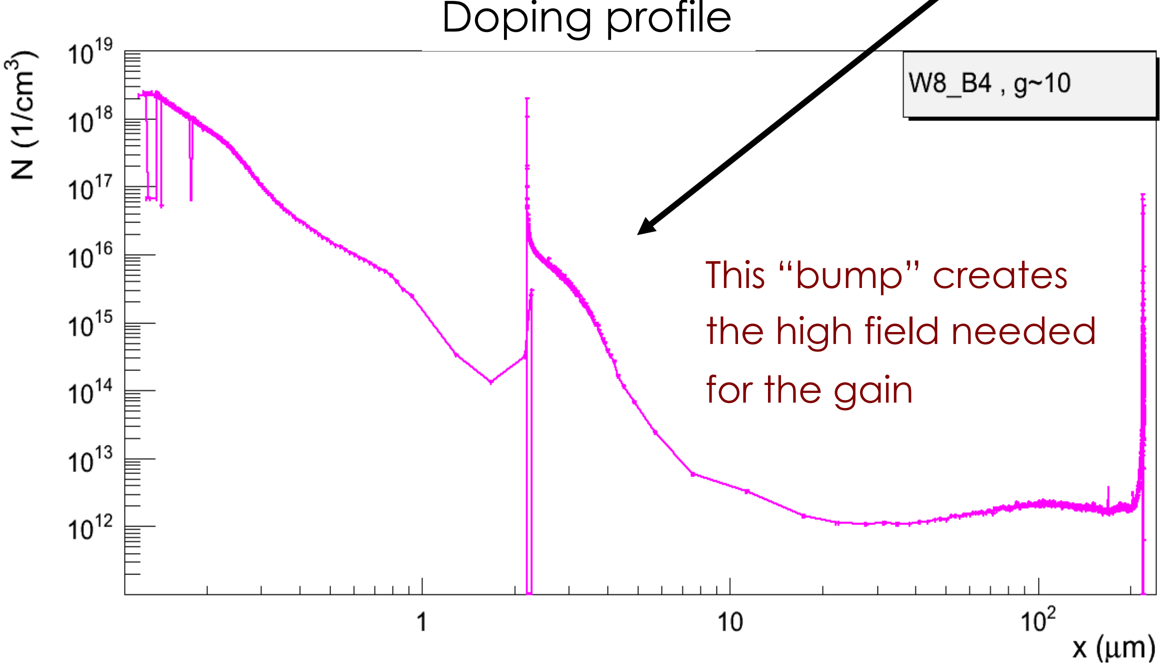
Doping profile from CV measurement - II



Gain sensor

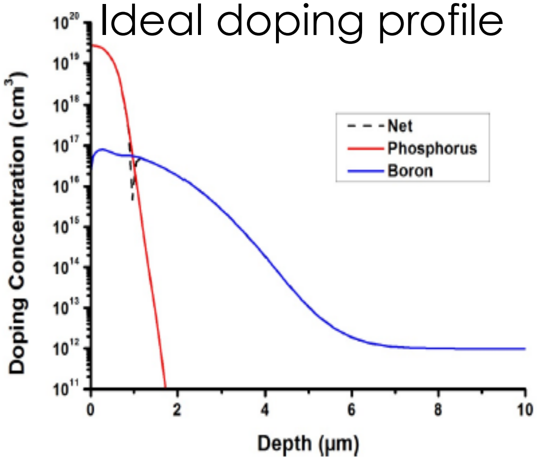


Doping profile



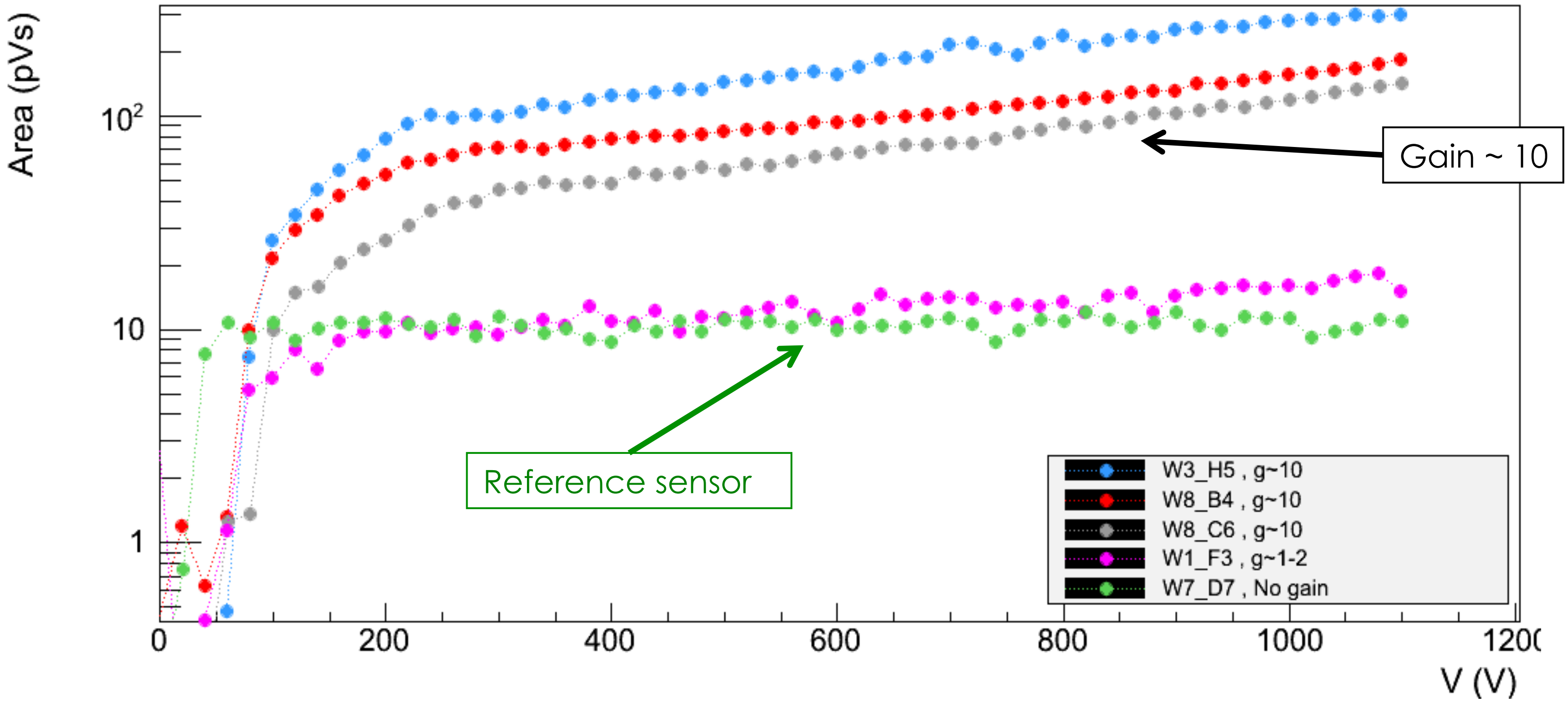
This “bump” creates the high field needed for the gain

Ideal doping profile



Signal amplitude

Using laser signals we are able to measure the different responses of LGAD and traditional sensors



Gain ~ 10

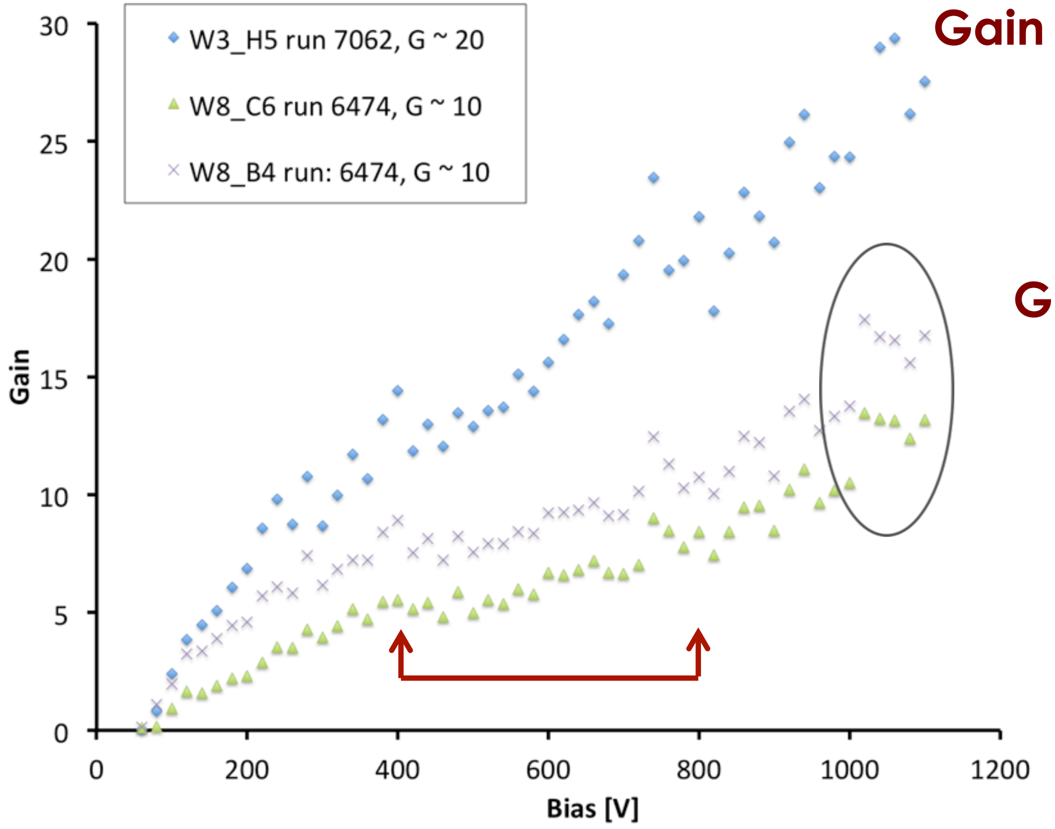
Reference sensor

Gain

The gain is estimated as the ratio of the output signals of LGAD detectors to that of traditional one

The gain increases linearly with Vbias (not exponentially!)

$$\frac{\text{Gain @ 800V}}{\text{Gain @ 400V}} \sim 2$$



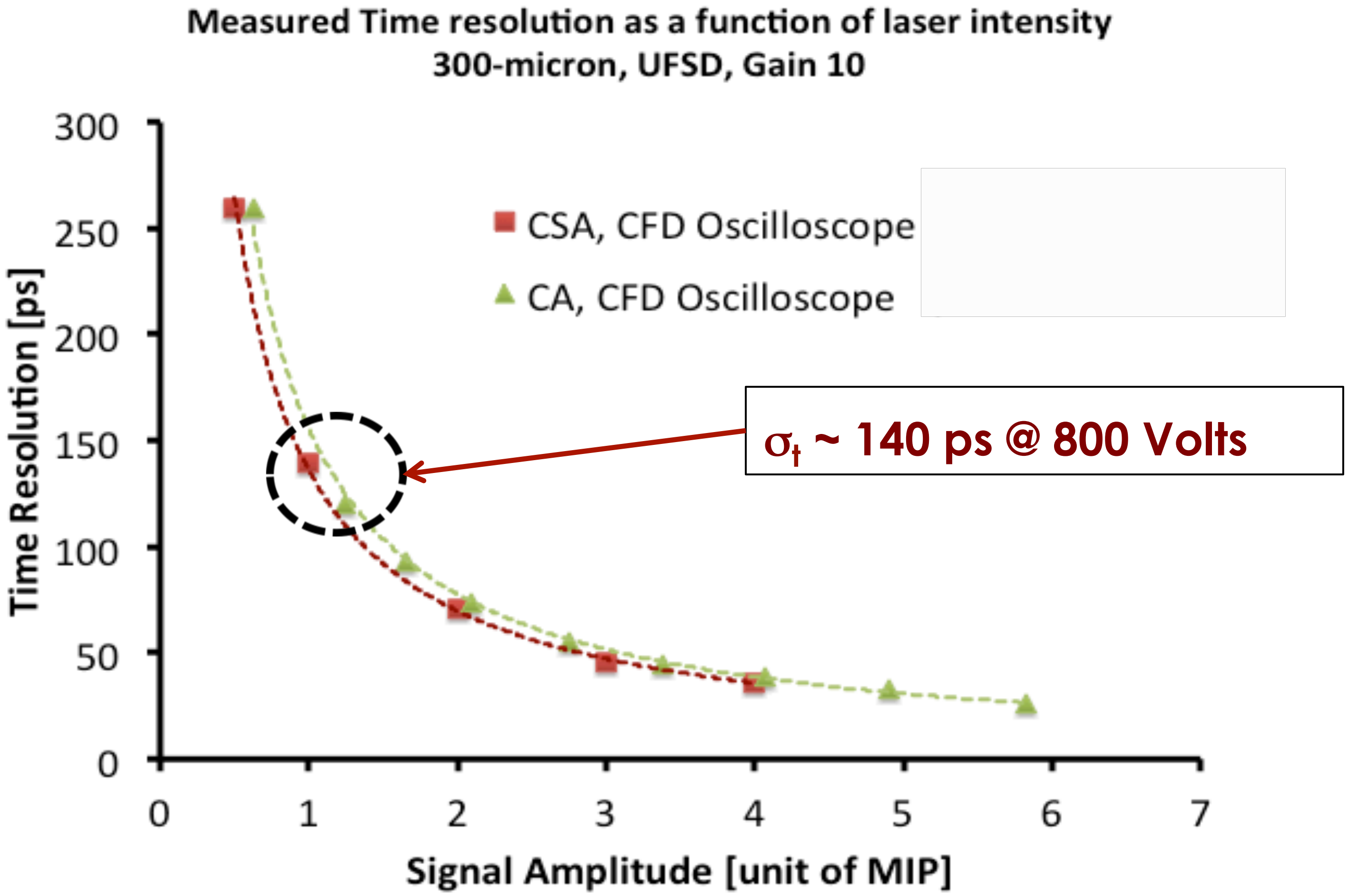
Gain ~ 20

Gain ~ 10

Laser Measurements on CNM LGAD

We use a 1064 nm picosecond laser to emulate the signal of a MIP particle (without Landau Fluctuations)

The signal output is read out by either a Charge sensitive amplifier or a Current Amplifier (Cividec)

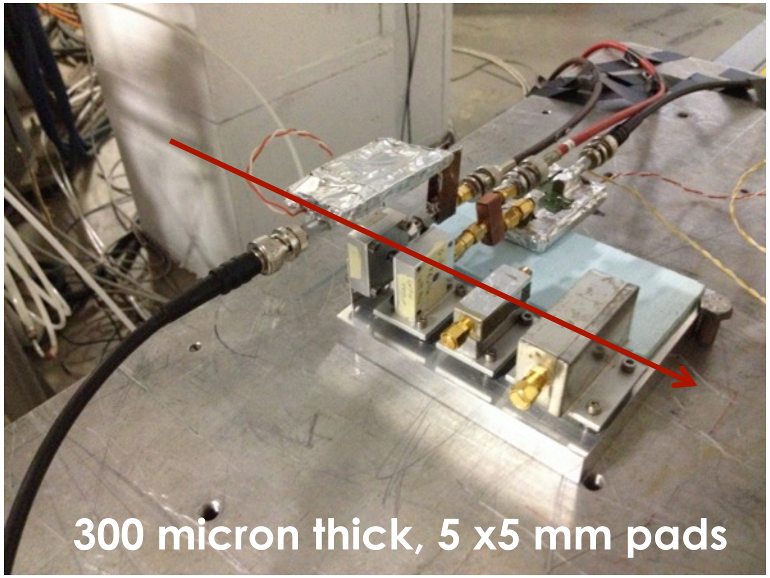


Measured Time resolution as a function of laser intensity
300-micron, UFSD, Gain 10

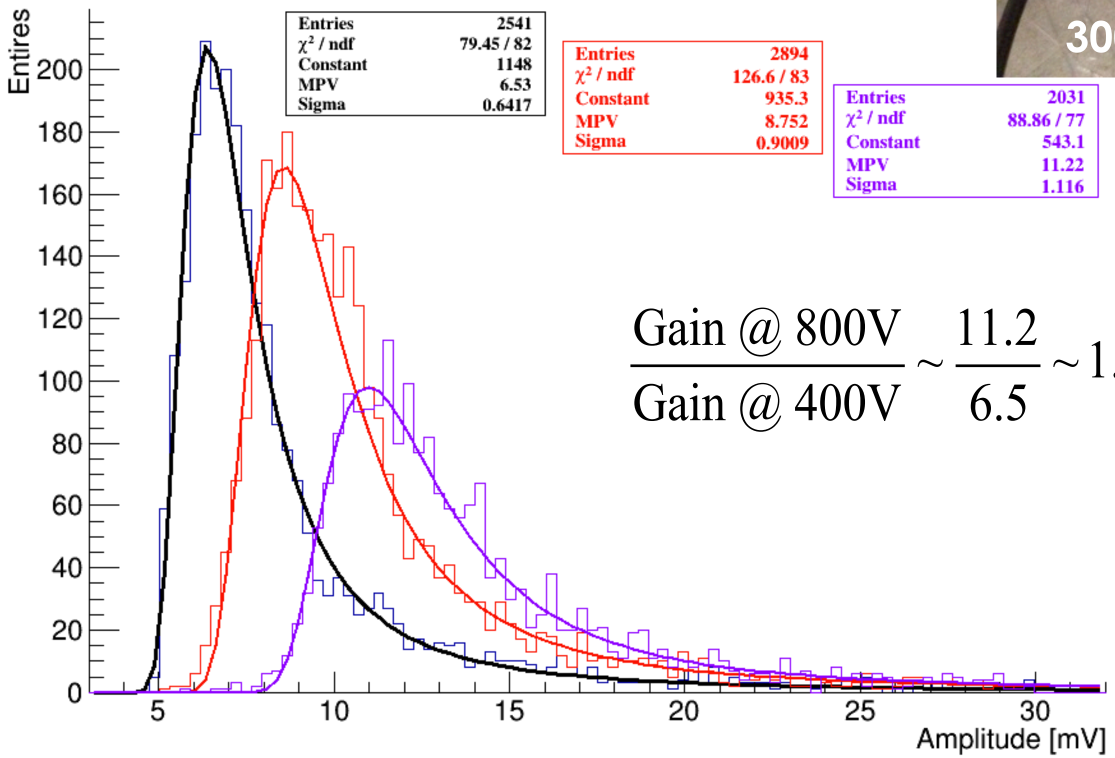
σ_t ~ 140 ps @ 800 Volts

Testbeam Measurements on CNM LGAD

In collaboration with Roma2, we went to Frascati for a testbeam using 500 MeV electrons



As measured in the lab, the gain ~ doubles going from 400 -> 800 Volt.



$$\frac{\text{Gain @ 800V}}{\text{Gain @ 400V}} \sim \frac{11.2}{6.5} \sim 1.7$$

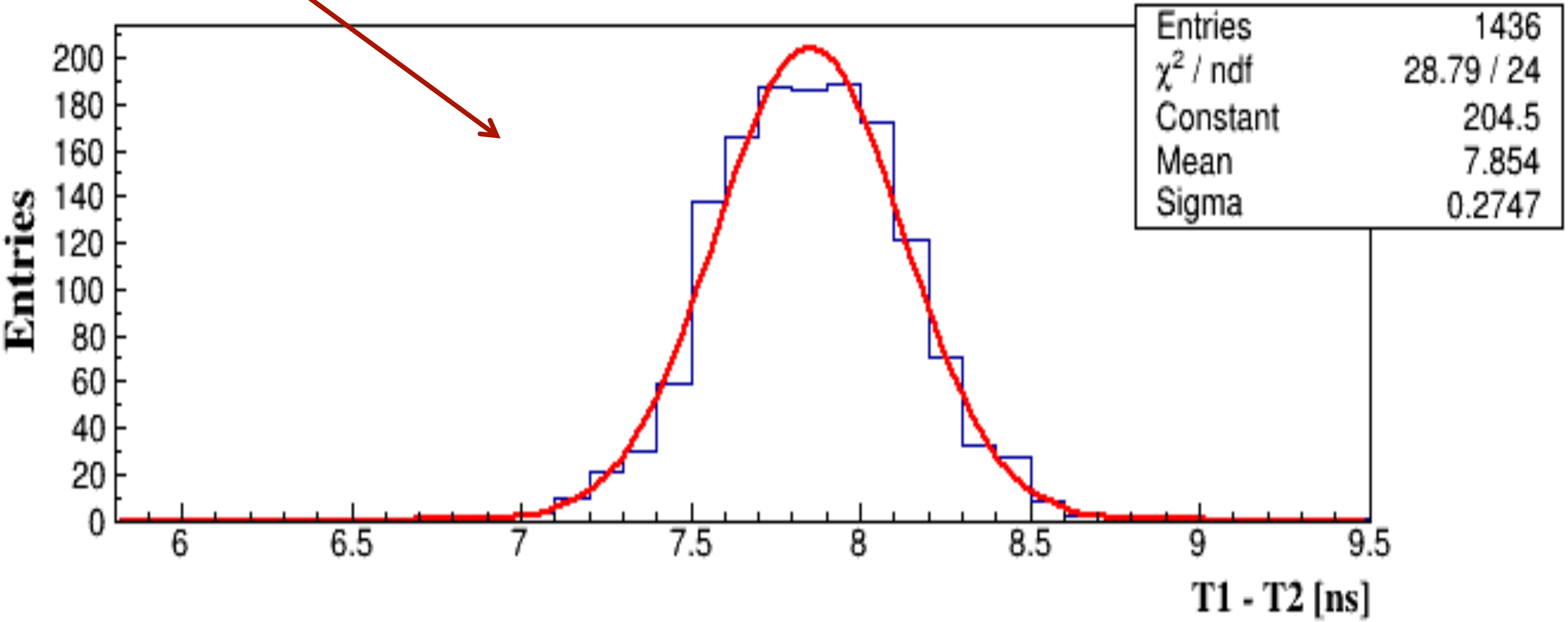
The gain mechanism preserves the Landau amplitude distribution of the output signals

Testbeam Measurements on CNM LGAD

Time difference between two LGAD detectors crossed by a MIP

Tested different types of electronics (Rome2 SiGe, Cividec),
Not yet optimized for these detectors

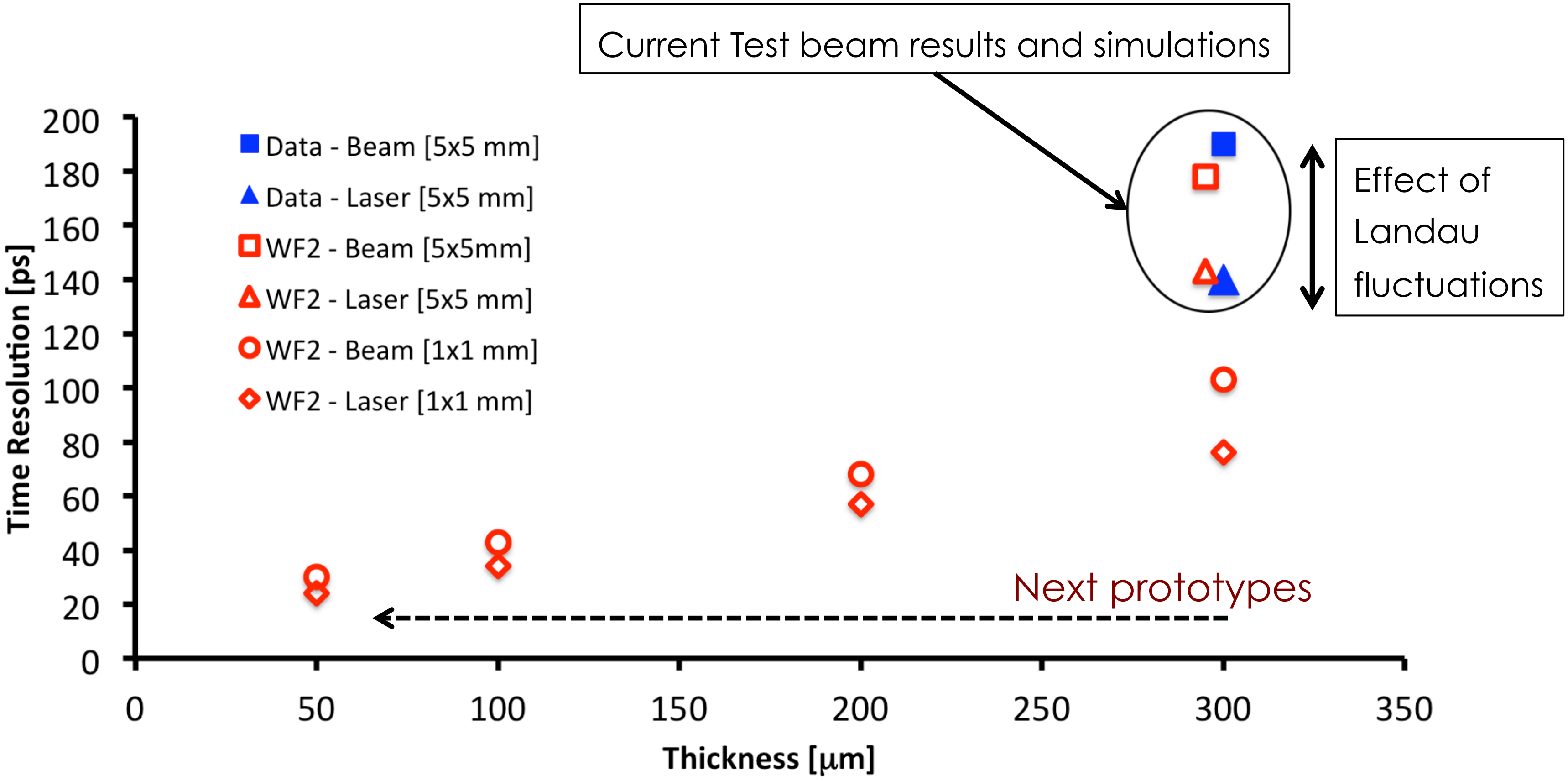
σ_t ~ 190 ps @ 800 Volts



Present results and future productions

With WF2, we can reproduce very well the laser and testbeam results.

Assuming the same electronics, and 1 mm^2 LGAD pad with gain 10, we can predict the timing capabilities of the next sets of sensors.


Current Test beam results and simulations

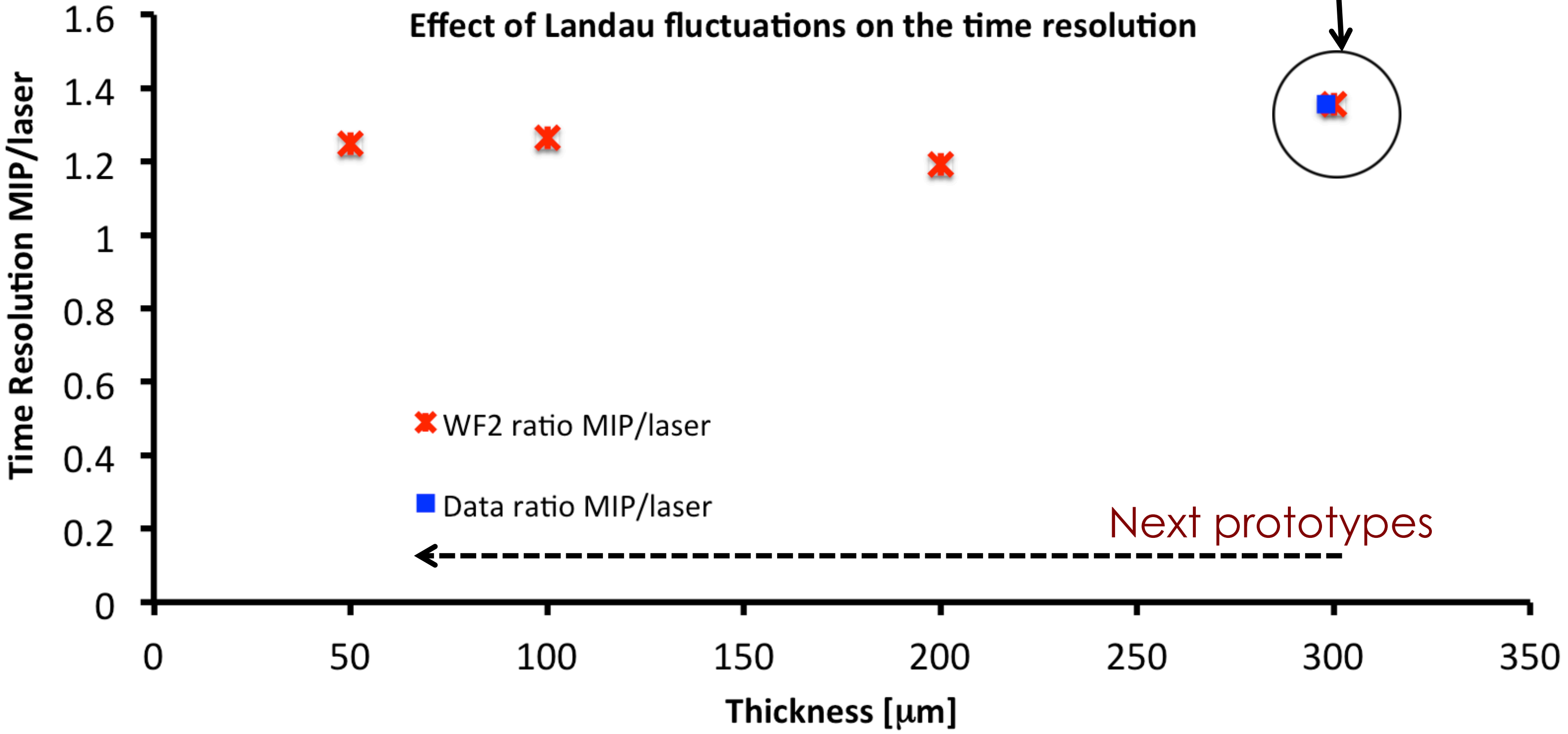
Effect of Landau fluctuations

Next prototypes

Effect of Landau Fluctuations on the time resolution

The effect of Landau fluctuations in a MIP signal are degrading the time resolution by roughly 30 % with respect of a laser signal

Current Test beam results and simulations

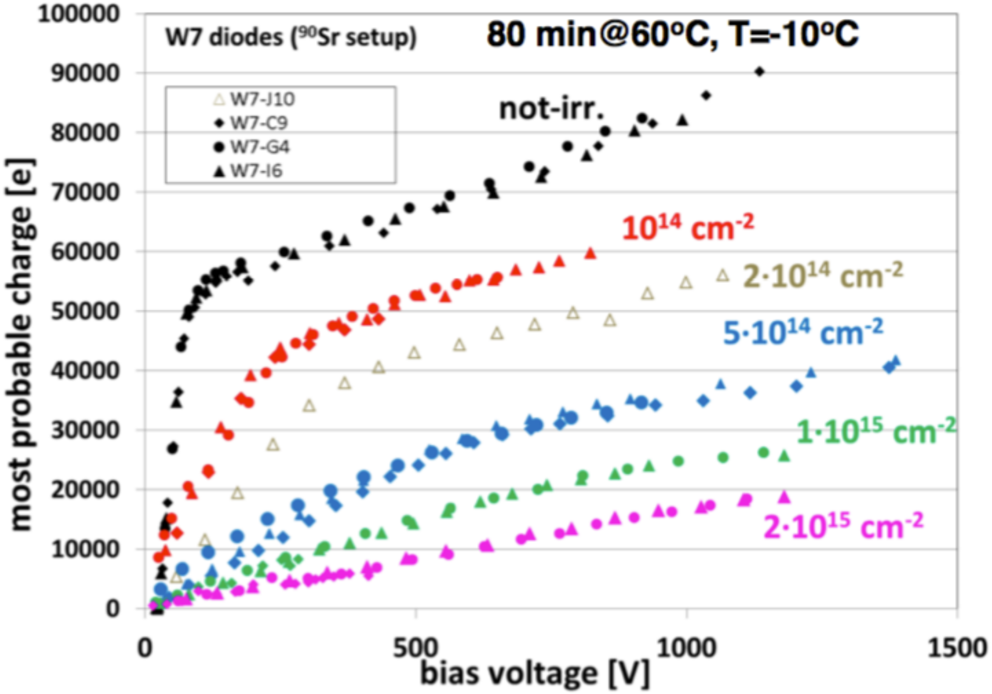


Next prototypes

Irradiation tests

The gain decreases with irradiations:
at 10^{14} n/cm^2 is 20% lower

➔ Due to boron disappearance



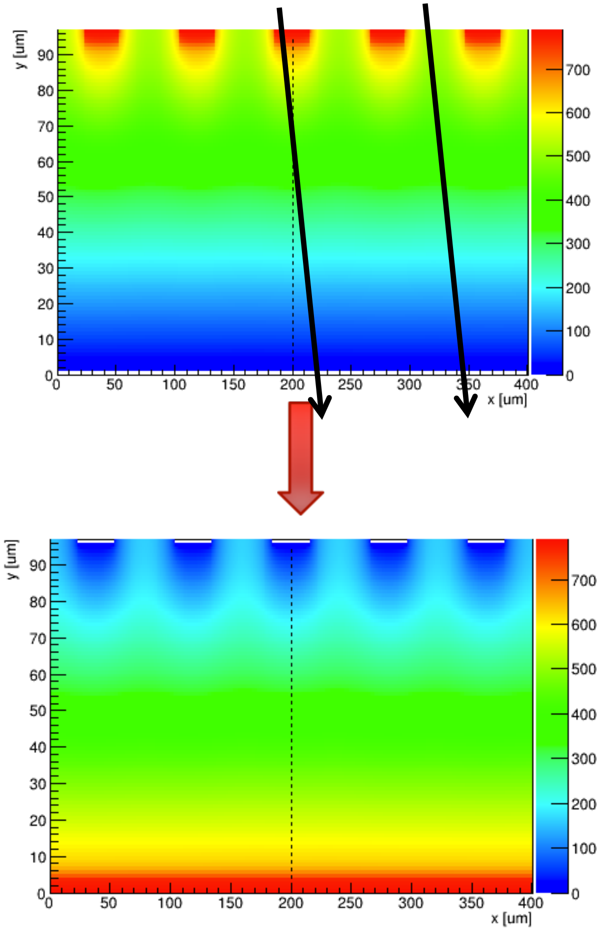
What-to-do next:

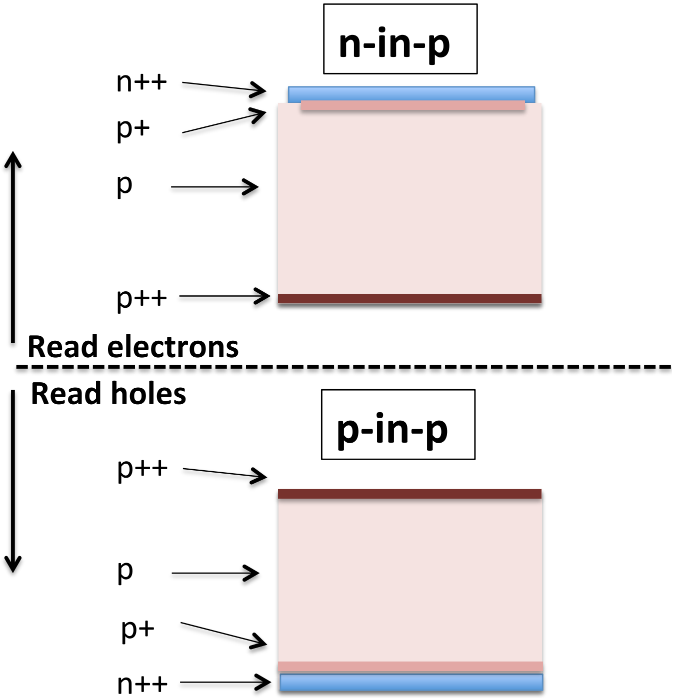
Planned new irradiation runs (neutrons, protons), new sensor geometries

Use Gallium instead of Boron for gain layer (in production now)

Gain in finely segmented sensors

Segmentation makes the effect of gain more difficult to predict, and most likely very dependent on the hit position




n-in-p

n++

p+

p

p++

Read electrons

Read holes

p-in-p

p++

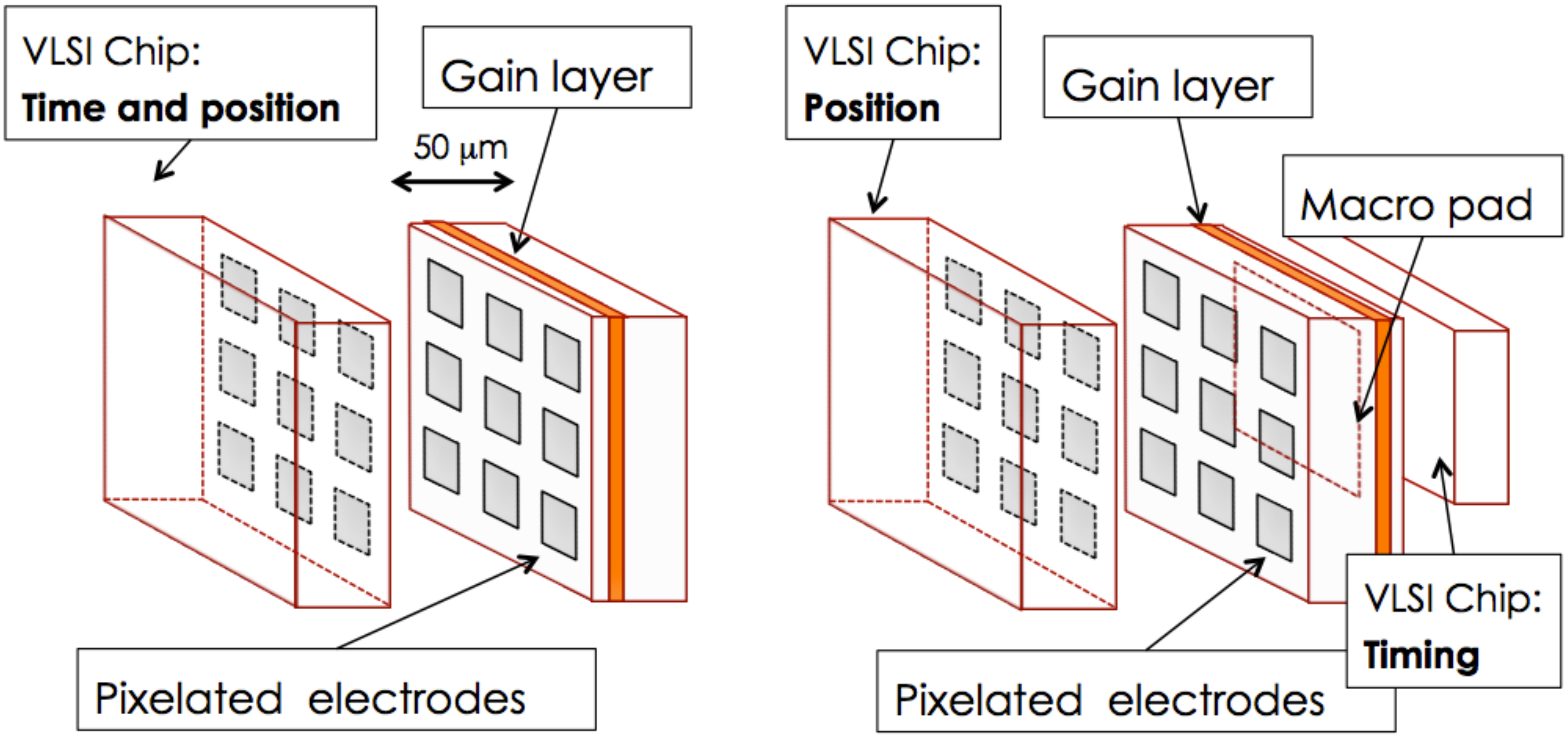
p

p+

n++

Moving the junction on the deep side allows having a very uniform multiplication, regardless of the electrode segmentation

Splitting gain and position measurements

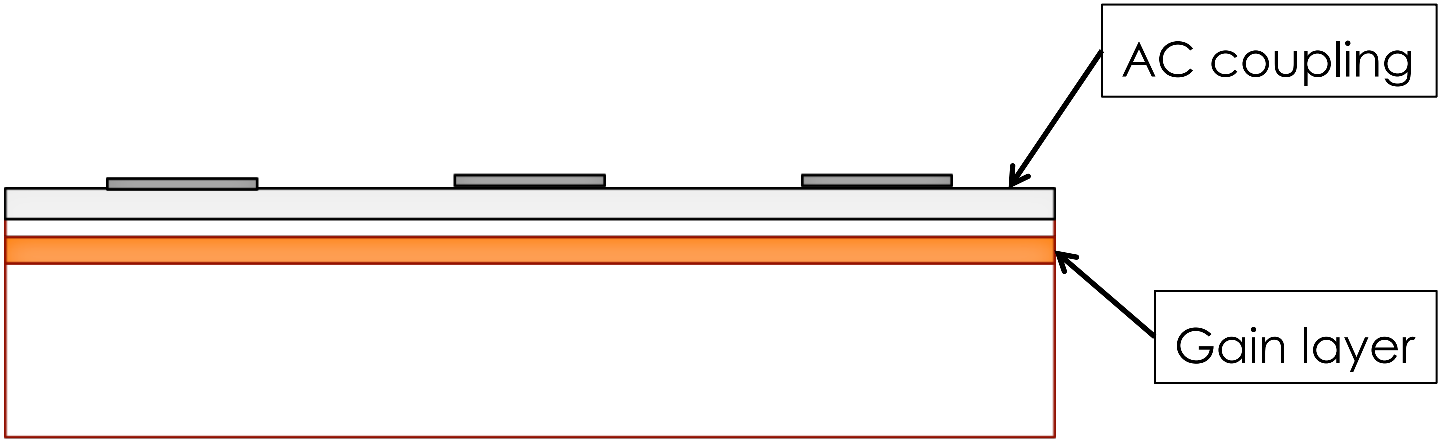


The ultimate time resolution will be obtained with a custom ASIC.

However we might split the position and the time measurements

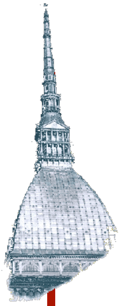
Using AC coupling to achieve segmentation

Standard n-in-p LGAD, with AC read-out



Very uniform field due to large pads,

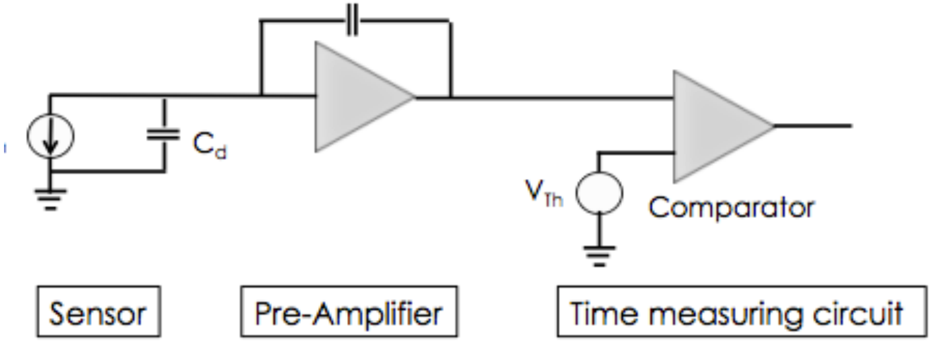
Segmentation due to AC coupling pick-up

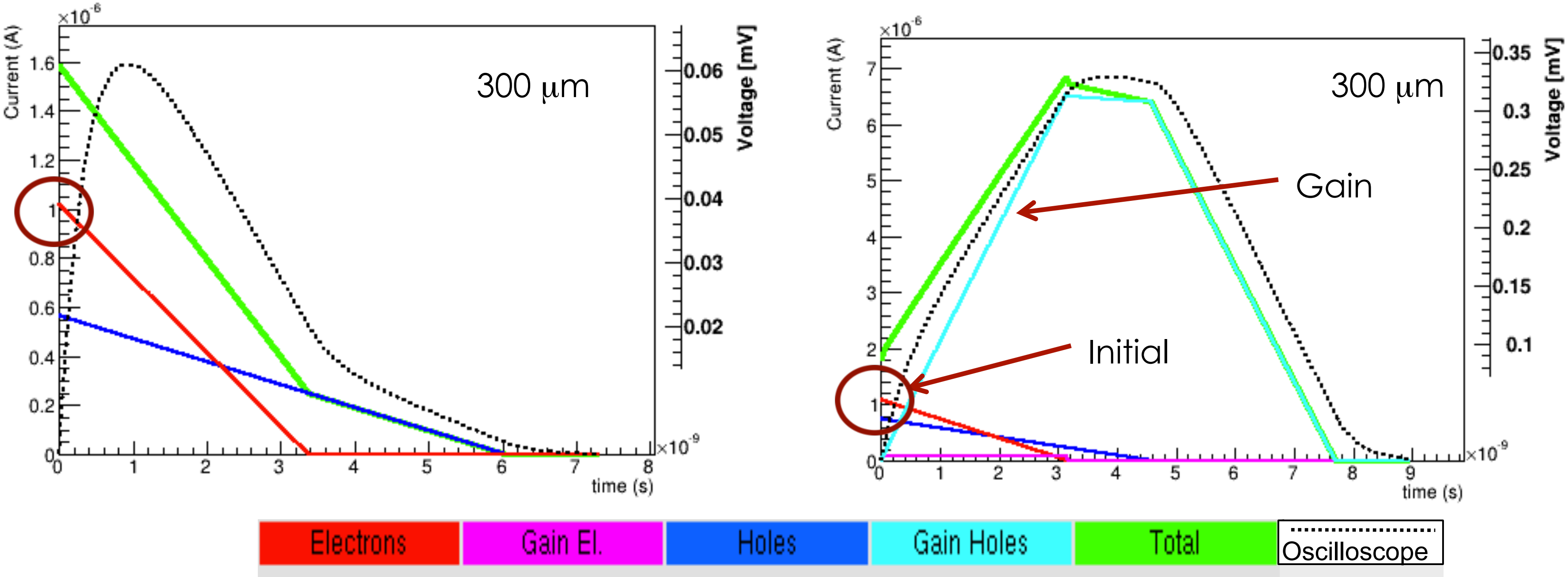


Electronics

To fully exploit UFSDs, dedicated electronics needs to be designed.

The signal from UFSDs is different from that of traditional sensors





Simulated Weightfield2

Pads with no gain

Charges generated uniquely by the incident particle

Pads with gain

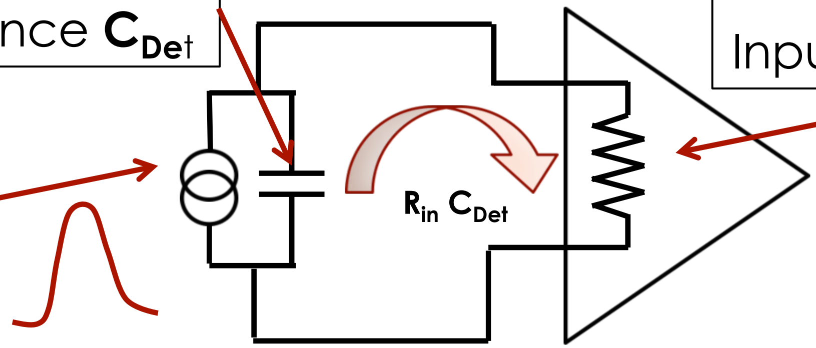
Current due to gain holes creates a longer and higher signal

Interplay of T_{Col} and $\tau = R_{in} C_{Det}$

Detector Capacitance C_{Det}

Input impedance R_{in}

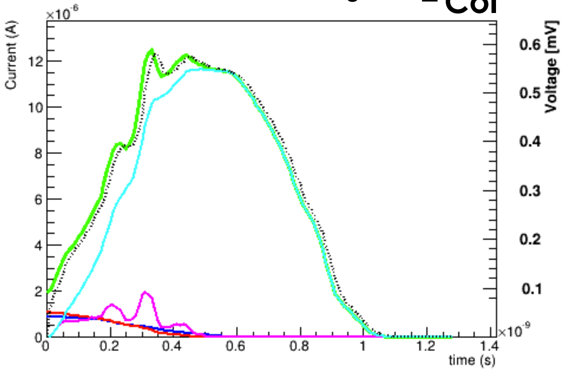
Collection time T_{Col}



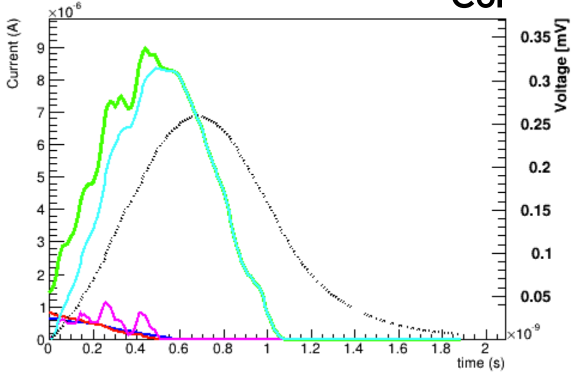
There are two time constants at play:

- T_{Col} : the signal collection time (or equivalently the rise time)
- $\tau = R_{in} C_{Det}$: the time needed for the charge to move to the electronics

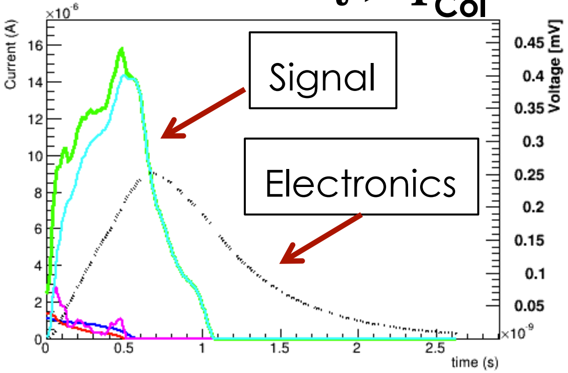
$\tau < T_{Col}$



$\tau \sim T_{Col}$



$\tau > T_{Col}$



Signal

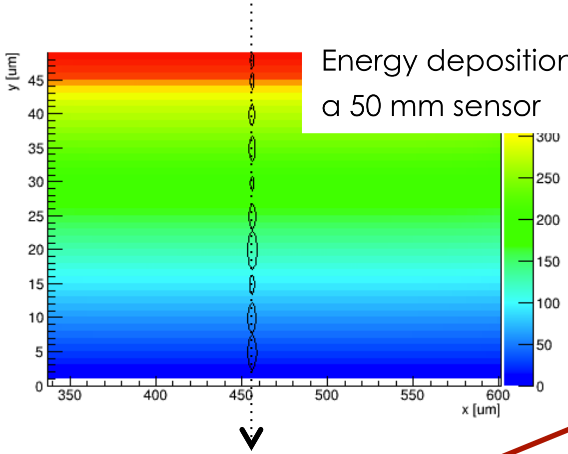
Electronics

τ / T_{Col} increases ➔ dV/dt decreases

➔ Smoother current

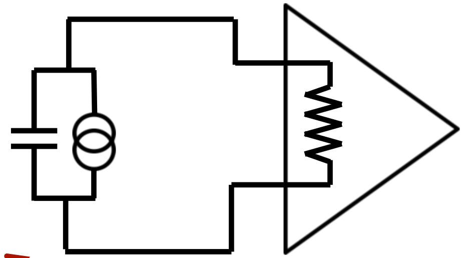
Need to find the optimum balance

Electronics: What is the best pre-amp choice?

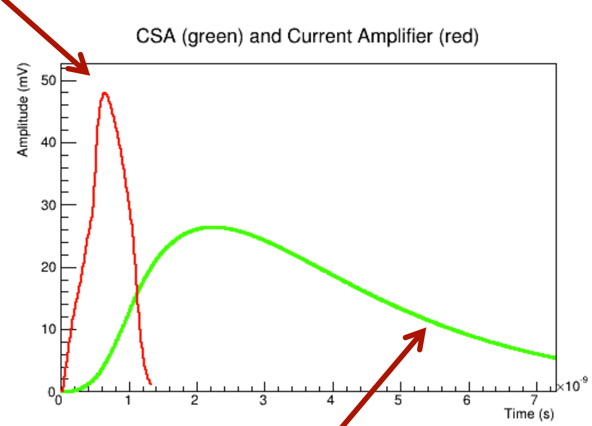


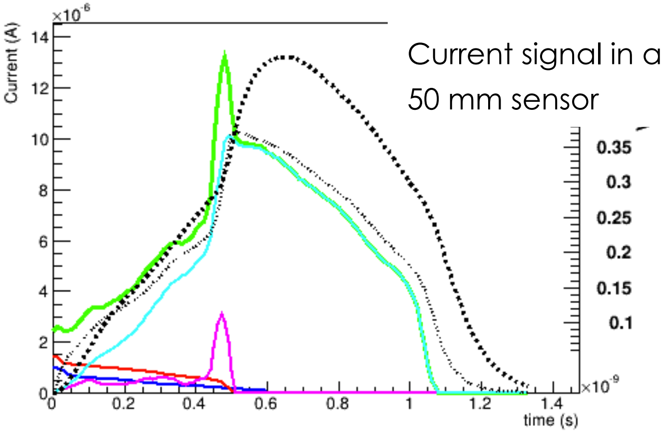
Energy deposition in a 50 mm sensor

Current Amplifier



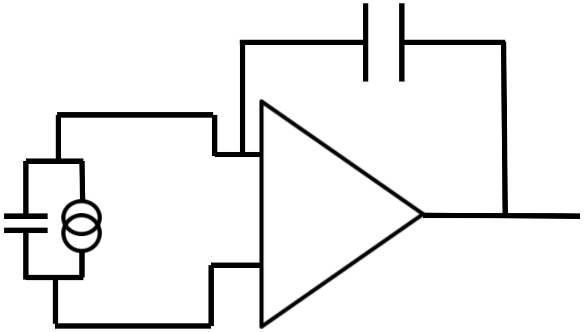
- Fast slew rate
- Higher noise
- Sensitive to Landau bumps





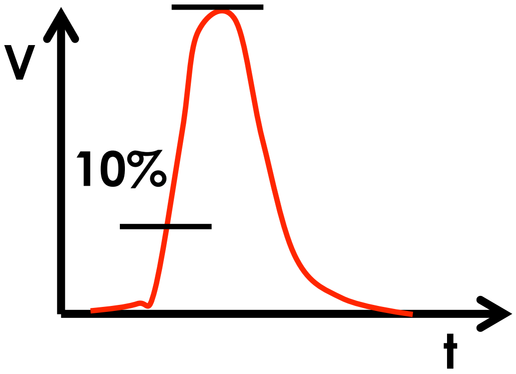
Current signal in a 50 mm sensor

Integrating Amplifier



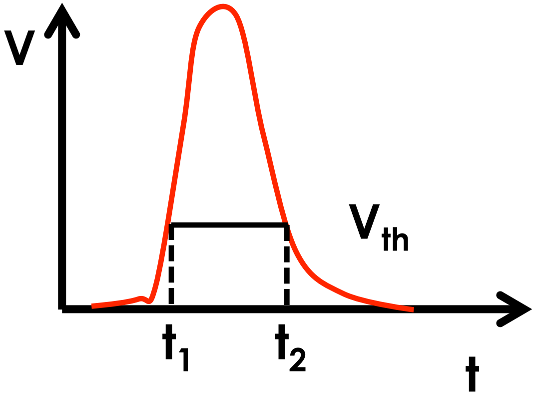
- Slower slew rate
- Quieter
- Integration helps the signal smoothing

What is the best "time measuring" circuit?



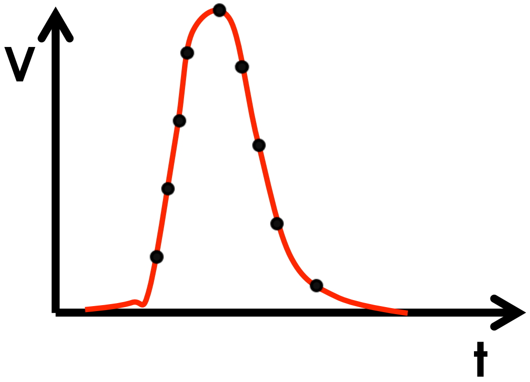
Constant Fraction Discriminator

The time is set when a fixed fraction of the amplitude is reached



Time over Threshold

The amount of time over the threshold is used to correct for time walk

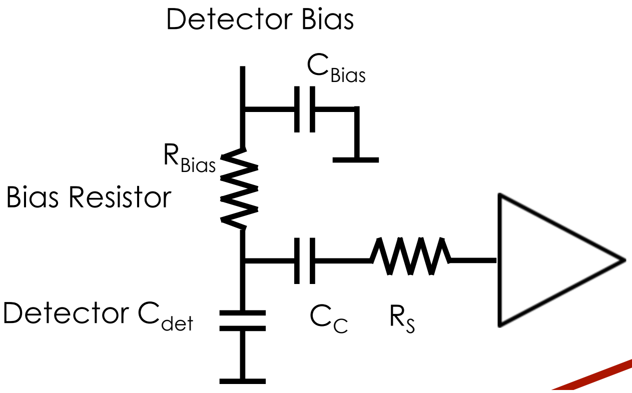


Multiple sampling

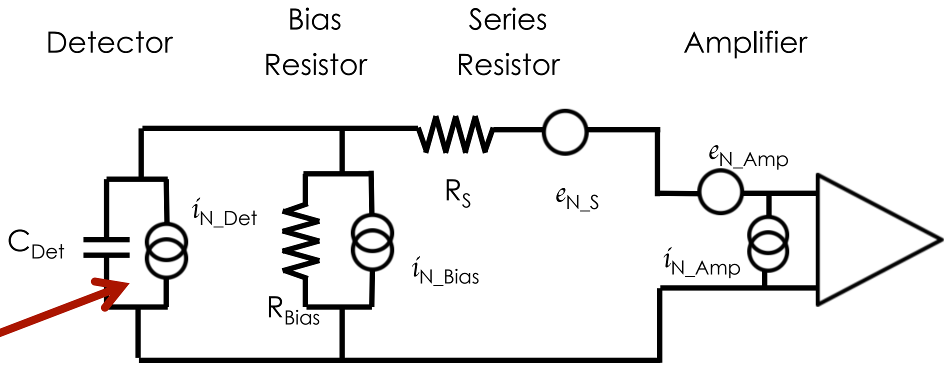
Most accurate method, needs a lot of computing power

Noise - I

Real life



Noise Model



This term, the detector current shot noise, depends on the gain

$$Q_n^2 = (2eI_{Det} + \frac{4kT}{R_{Bias}} + i^2_{N\_Amp})F_i T_s + \underbrace{(4kTR_s + e^2_{N\_Amp})F_v \frac{C^2_{Det}}{T_S}} + F_{vf} A_f C^2_{Det}$$

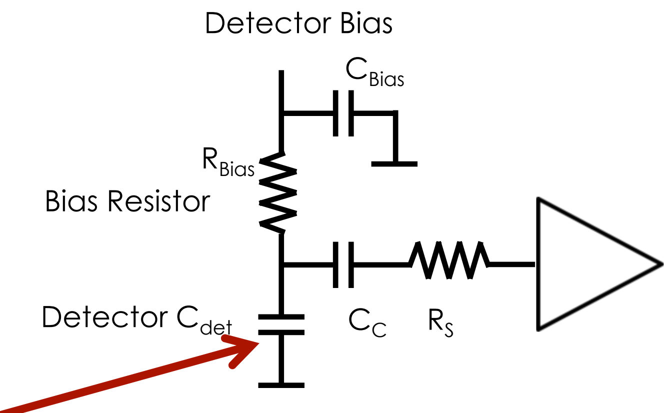
$$\boxed{2eI_{Det} * \text{Gain}}$$

low gain!

This term dominates for short shaping time

Noise - II

Real life



$$ENF = kG + (2 - \frac{1}{G})(1 - k)$$

k = ratio h/e gain

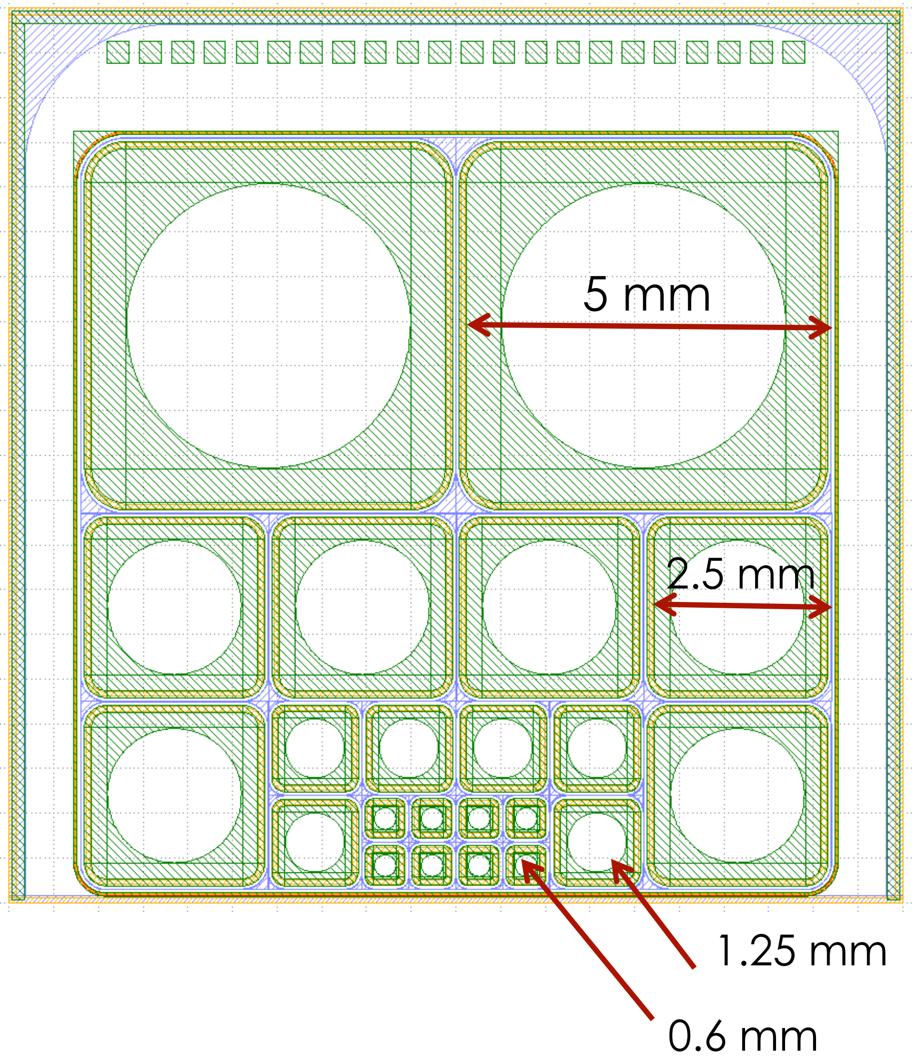
NOISE DUE TO GAIN:

Excess noise factor:

low gain, very small k

Low leakage current and low gain (~ 10) together with short shaping time are necessary to keep the noise down.

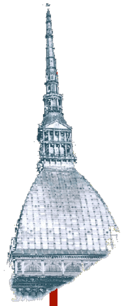
Next CNM productions



These new productions will allow a detailed exploration of the UFSD timing capabilities, including border effects between pads, and distance from the sensor edge.

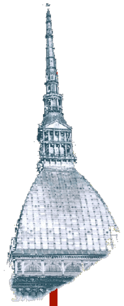
Timescale:

- **Spring 2015**: 200 micron
- **Summer 2015**: 100 micron
- **Summer 2015**: 50 micron



Next Steps

1. Wafer Production
   200 micron thick sensors by **Spring-2015**
   100 and 50 micron thick sensors by **Summer 2015.**
2. Production of UFSD doped with Gallium instead of Boron.
3. Study of reversed-UFSD started for the production of pixelated UFSD sensors (FBK, Trento).
4. UFSD are included in the CMS TDR CT-PPS as a solution for forward proton tagging
5. Use of UFSD in beam monitoring for hadron beam. INFN patent and work on-going
6. Interest in UFSD for 4D tracking at high luminosity
7. Testbeam analyses just started. Results coming soon…



UFSD – Summary

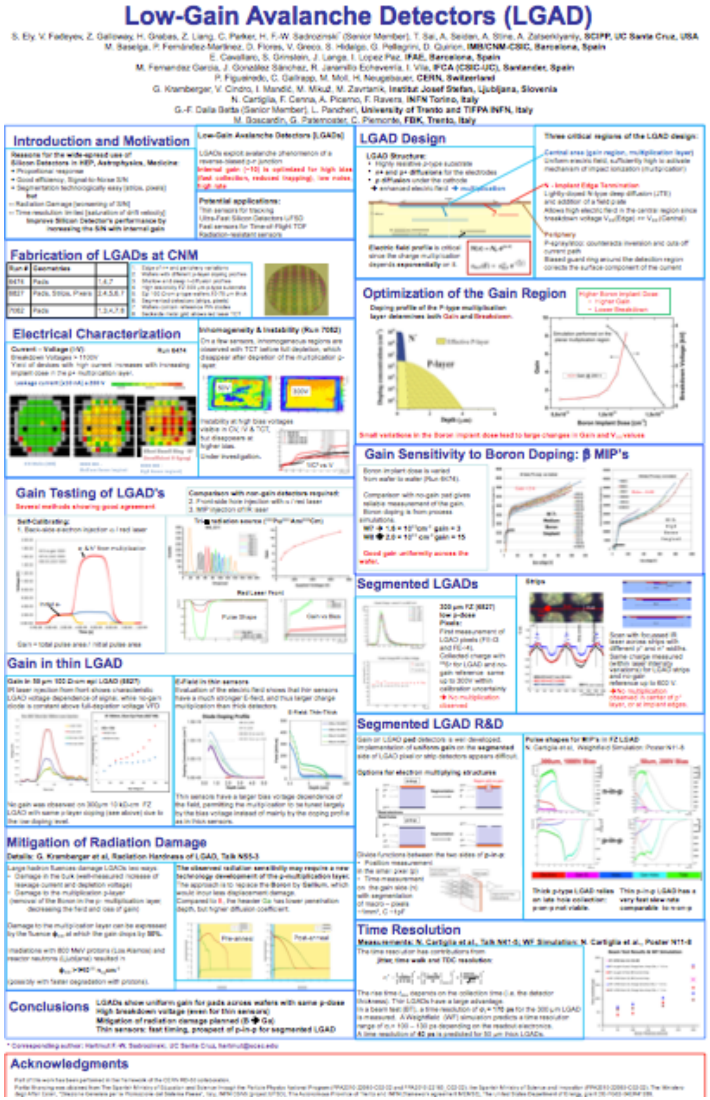
We are just starting to understand the timing capability of UFSD

- Low-gain avalanche diodes (LGAD) offer silicon sensors with an enhanced signal amplitude: UFSD are LGAD detectors optimized for timing resolution.
- Several options under studies to obtain concurrently excellent space and time resolutions.

- We developed a program, **Weightfield2** to simulate the behaviors of LGAD and optimized them for fast timing (available at http://personalpages.to.infn.it/~cartigli/Weightfield2.0/)

Timescale: 1 year to asses UFSD timing capabilities

Presented at IEEE, oral and posters, presentations



Low-Gain Avalanche Detectors (LGAD)

Introduction and Motivation

LGAD Design

Fabrication of LGADs at CNM

Optimization of the Gain Region

Electrical Characterization

Gain Sensitivity to Boron Doping: β MIP's

Gain Testing of LGAD's

Segmented LGADs

Gain in thin LGAD

Segmented LGAD R&D

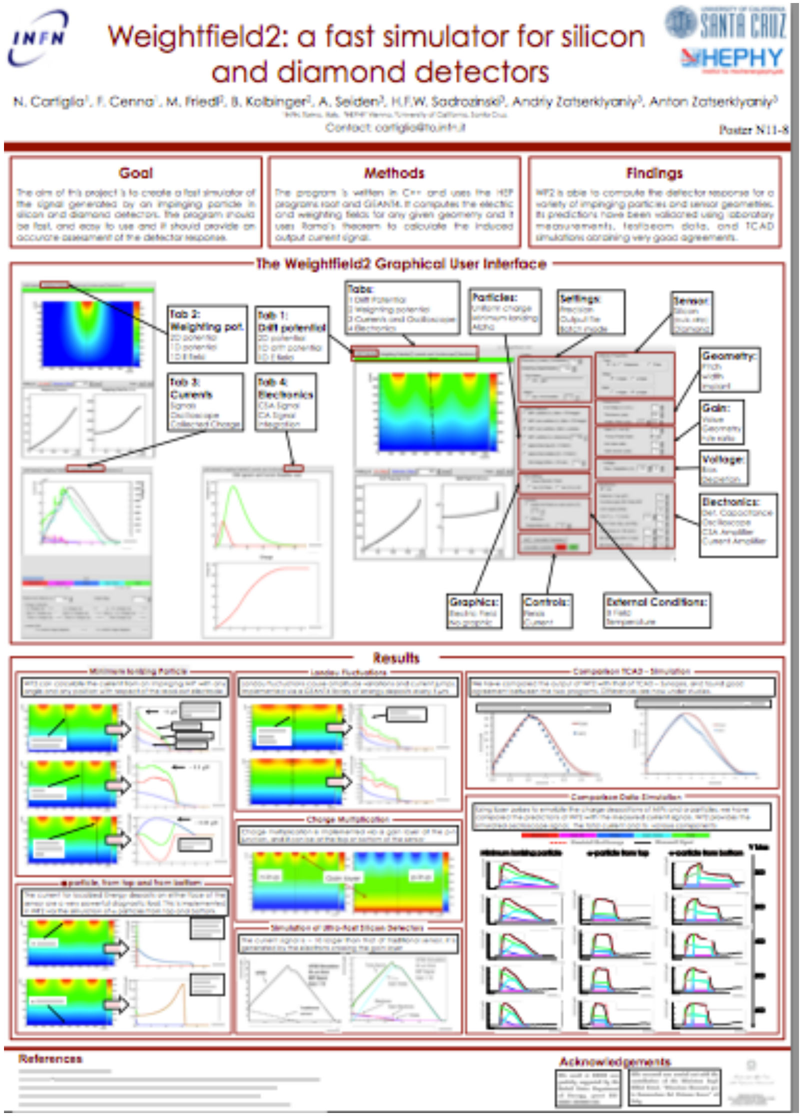
Mitigation of Radiation Damage

Time Resolution

Conclusions

Acknowledgments

Poster Session IEEE N26-13



Weightfield2: a fast simulator for silicon and diamond detectors

N. Cartiglia, F. Cenna, M. Friedl, B. Kolbinger, A. Seiden, H.F.W. Sadrozinski, Andriy Zatserklyaniy, Anton Zatserklyaniy

Poster N11-8

Goal

Methods

Findings

The Weightfield2 Graphical User Interface

Results

References

Acknowledgements

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Additional references

Several talks at the 22nd, 23rd and 24th RD50 Workshops:

23rd RD50: https://indico.cern.ch/event/265941/other-view?view=standard
22nd RD50: http://panda.unm.edu/RD50_Workshop/

9Th Trento Workshop, Genova, Feb 2014.

F. Cenna “**Simulation of Ultra-Fast Silicon Detectors**”

N. Cartiglia “**Timing capabilities of Ultra-Fast Silicon Detector**”

Papers:

[1] N. Cartiglia, Ultra-Fast Silicon Detector, 13th Topical Seminar on Innovative Particle and Radiation Detectors (IPRD13), 2014 JINST 9 C02001, http://arxiv.org/abs/1312.1080

[2] H.F.-W. Sadrozinski, N. Cartiglia et al., Sensors for ultra-fast silicon detectors, Proceedings "Hiroshima" Symposium HSTD9, DOI: 10.1016/j.nima.2014.05.006 (2014).

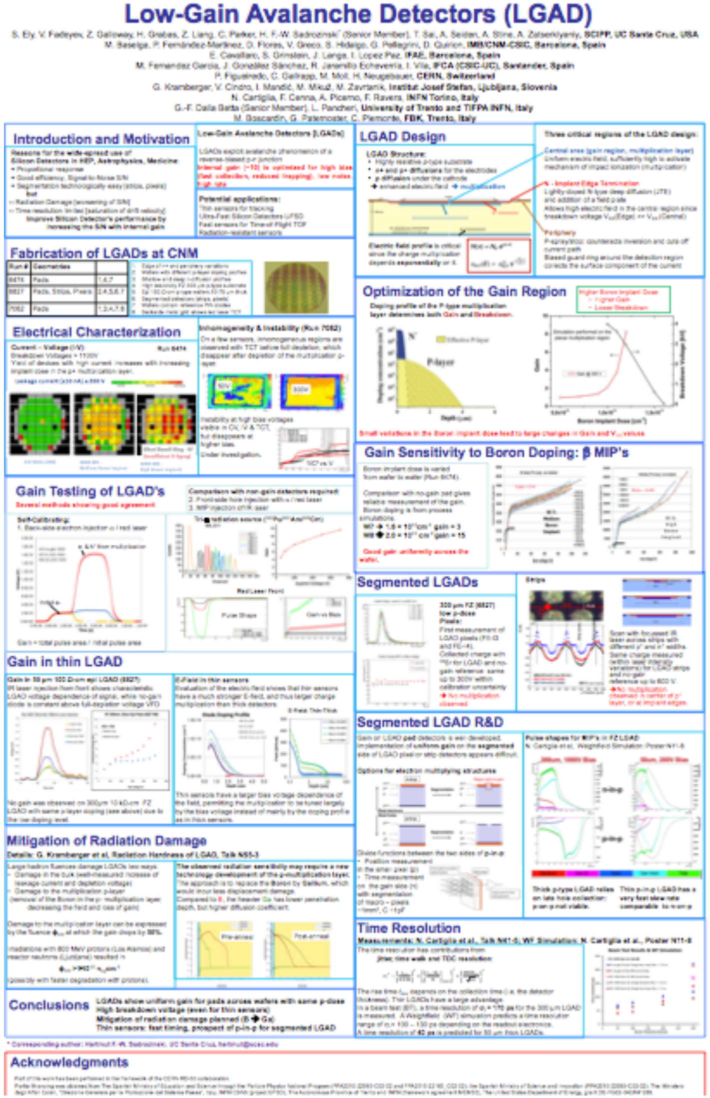
Backup

The “Low-Gain Avalanche Detector” project

Is it possible to manufacture a silicon detector that looks like a normal pixel or strip sensor, but

with a much larger signal (RD50)?

- 730 e/h pair per micron instead of 73 e/h
- Finely segmented
- Radiation hard
- No dead time
- Very low noise (low shot noise)
- No cross talk



Low-Gain Avalanche Detectors (LGAD)

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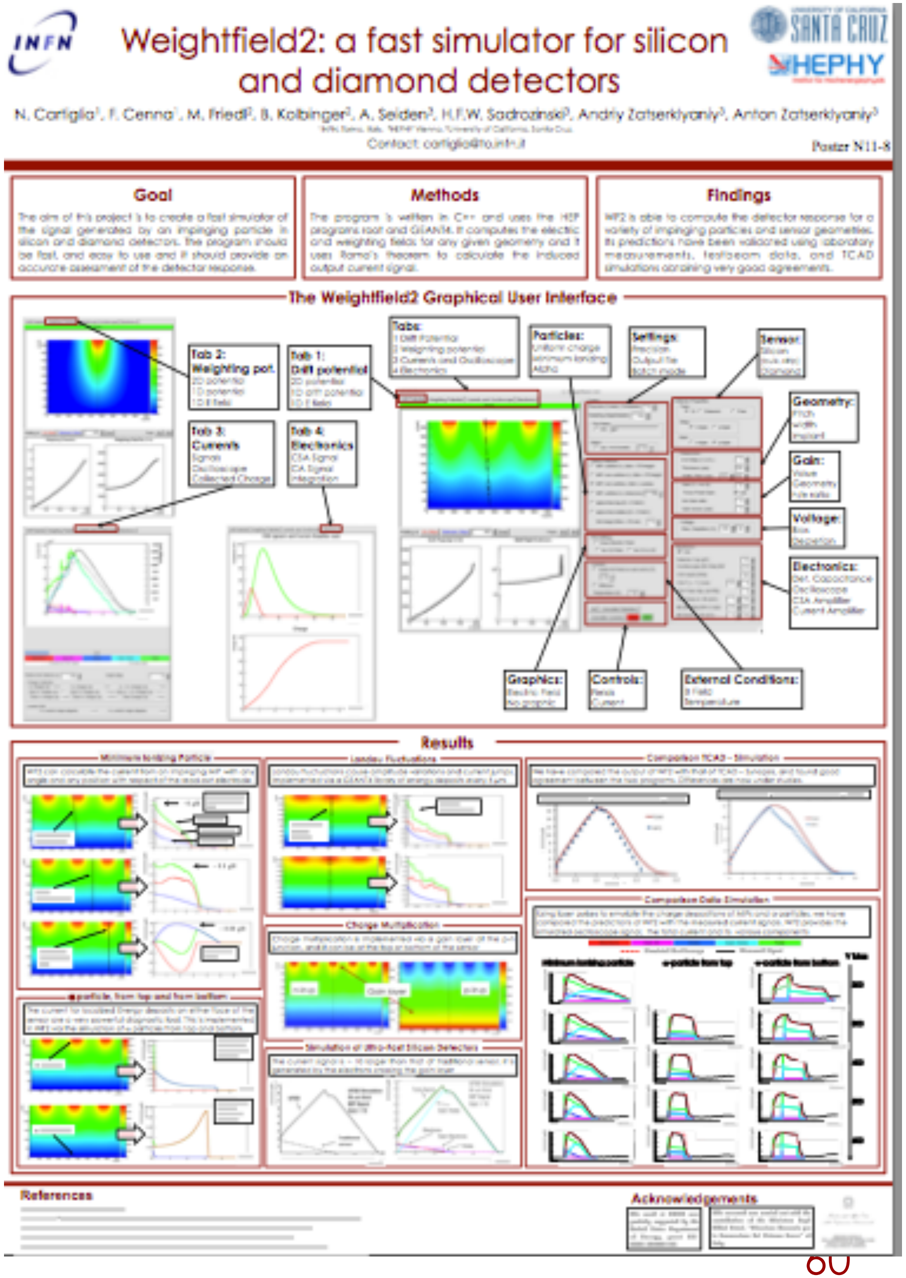
How can we progress? Need simulation

We developed a full simulation program to optimize the sensor design, WeightField2, (http://cern.ch/weightfield2)

It includes:

- Custom Geometry
- Calculation of drift field and weighting field
- Currents signal via Ramo's Theorem
- Gain
- Diffusion
- Temperature effect
- Non-uniform deposition
- Electronics



Weightfield2: a fast simulator for silicon and diamond detectors

N. Cartiglia[1], F. Cenna[1], M. Friedl[2], B. Kolbinger[2], A. Seiden[3], H.F.W. Sadrozinski[3], Andriy Zatserklyaniy[3], Anton Zatserklyaniy[3]

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Poster N11-8

Goal

The aim of this project is to create a fast simulator of the signal generated by an impinging particle in silicon and diamond detectors. The program should be fast, and easy to use and it should provide an accurate assessment of the detector response.

Methods

The program is written in C++ and uses the HEP programs root and GEANT4. It computes the electric and weighting fields for any given geometry and it uses Ramo's theorem to calculate the induced output current signal.

Findings

WF2 is able to compute the detector response for a variety of impinging particles and sensor geometries. Its predictions have been validated using laboratory measurements, testbeam data, and TCAD simulations obtaining very good agreements.

The Weightfield2 Graphical User Interface

Results

References

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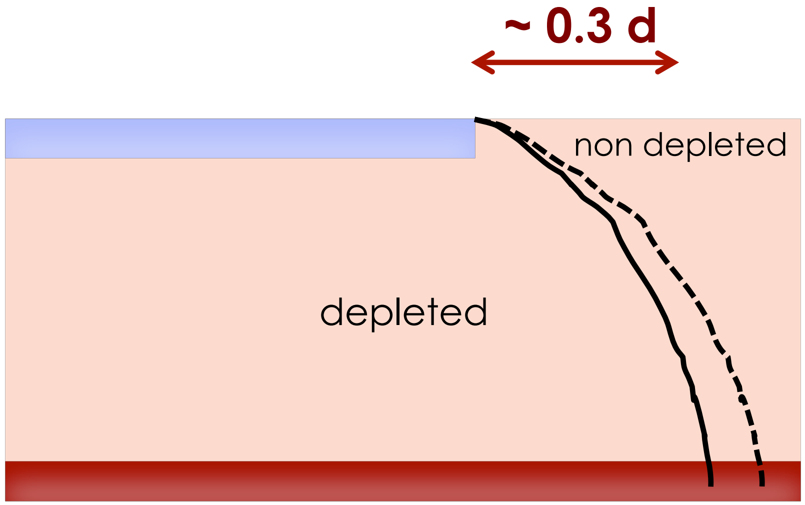
Sensor thickness and slim edge

Rule: when the depletion volume reaches the edge, you have electrical breakdown.

It's customary to assume that the field extends on the side by ~ 1/3 of the thickness.



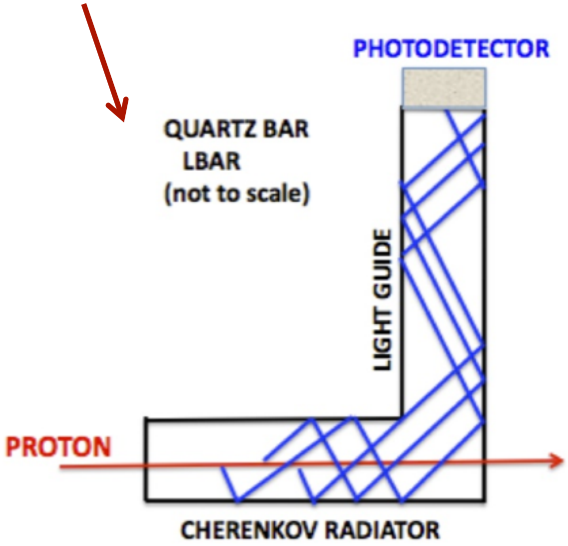
edge = k* thickness

- k = 1 very safe
- k = 0.5 quite safe
- K = 0.3 limit

By construction, thin detectors (~ 100 micron) might have therefore slim edge

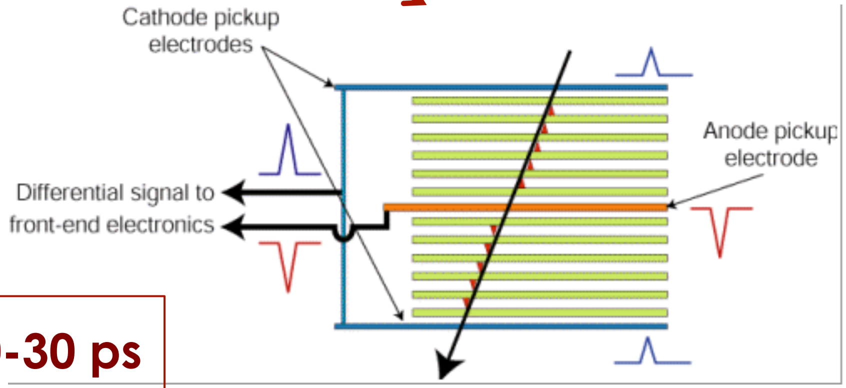
State-of-the-art Timing Detectors

Timing detectors exploit very fast physics processes such as **Cherenkov light** emission or **electronic avalanches** to create prompt signals



CMS/ATLAS

σ_t **~ 20-30 ps**
σ_x **~ 1-2 mm**



ALICE

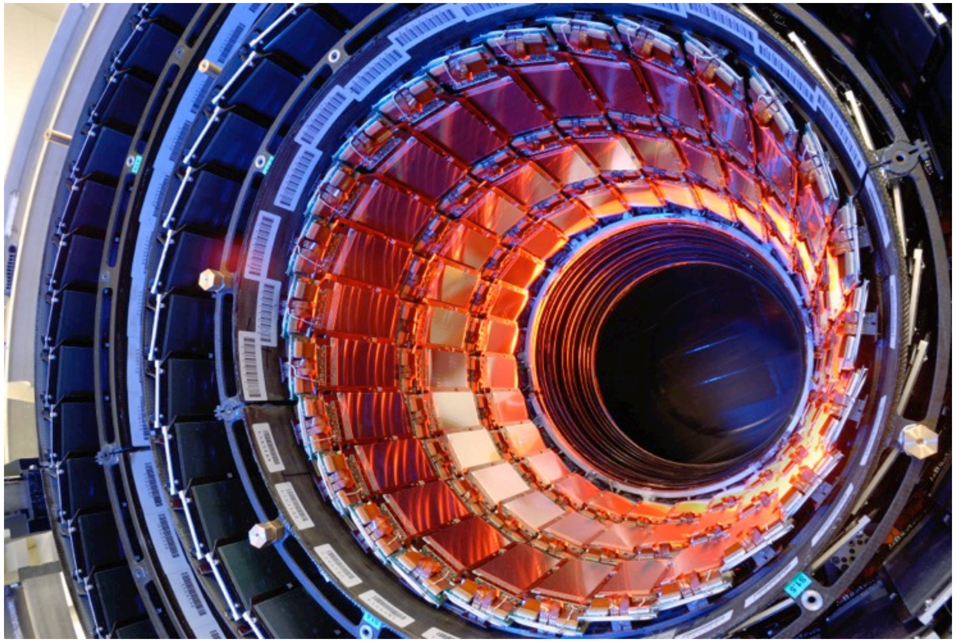
- These detectors measure time very accurately but locate particles with the precision of ~ 1 mm
- Good timing is obtain by using a gain mechanism, either in the detector or in the electronics

State-of-the-art Position Detectors

Extremely good position detectors are currently in use in every major high energy physics experiment:

- Millions of channels
- Very reliable
- Very radiation hard

The timing capability is however limited to ~ 100-150 ps (NA62 @CERN)



σ_t ~ 100-150 ps

σ_x ~ 20-30 μm