

Generalized Parton Distributions (GPDs)

Jennet Dickinson

Physics 290e

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Outline

- A review of electron-proton scattering
 - At different values of Q^2
- What are GPDs?
- How do we measure GPDs?
 - Deeply virtual Compton scattering (DVCS)
- Getting back what we started with

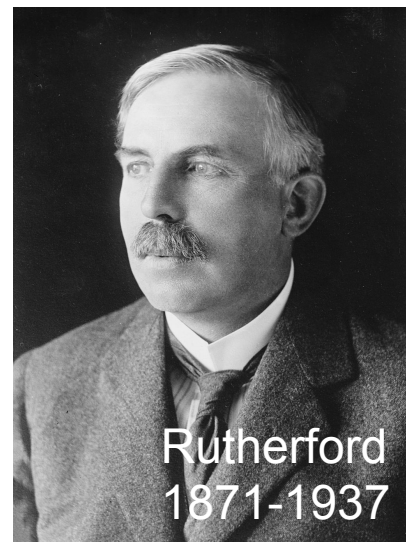
Electron-proton scattering

- Q^2 (virtuality of exchanged photon) ↓
- $Q^2 \ll 1/r_p$
 - Electron recoils from point-like spinless object
 - $Q^2 \sim 1/r_p$
 - Electron recoils from extended charged object with spin 1/2
 - $Q^2 > 1/r_p$
 - Electron can resolve proton structure

Rutherford Scattering

$$Q^2 \ll 1/r_p$$

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{16(p_e^2/2m_e)^2 \sin^4(\theta/2)}$$



- Scattering of charged point particles via Coulomb interaction
- Assume:
 - The electron is non-relativistic
 - The proton does not recoil and we can ignore proton spin
 - The proton is point-like

Mott Scattering

$$Q^2 \sim 1/r_p$$

proton recoil

spin-spin interactions

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E_1^2 \sin^4(\theta/2)} \frac{E_3}{E_1} \left(\cos^2 \frac{\theta}{2} - \frac{q^2}{2M_p^2} \sin^2 \frac{\theta}{2} \right)$$

Rutherford scattering
with relativistic electron energy

Taking electron spin
states into account

- Assume:

- ~~The electron is non-relativistic~~
- ~~The proton does not recoil and we can ignore proton spin~~
- The proton is point-like

Rosenbluth Formula

$$Q^2 \sim 1/r_p$$

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E_1 \sin^4(\theta/2)} \frac{E_3}{E_1}$$

Mott scattering + terms
describing proton's structure

$$\left\{ \left(F_1^2 - \frac{\kappa_p^2 q^2}{4M_p^2} F_2^2 \right) \cos^2 \frac{\theta}{2} - \left(F_1 + \kappa_p F_2 \right) \frac{q^2}{2M_p^2} \sin^2 \frac{\theta}{2} \right\}$$

- Assume:

- ~~The electron is non-relativistic~~
- ~~The proton does not recoil and we can ignore proton spin~~
- ~~The proton is point-like~~

Elastic Form Factors

- All information about the proton's structure is contained in **form factors** F_1 and F_2

$$\left(F_1^2 - \frac{\kappa_p^2 q^2}{4M_p^2} F_2^2 \right)$$

$$(F_1 + \kappa_p F_2)$$

- The form factors are functions of Q^2
- The proton also has anomalous magnetic moment $\kappa_p = 1.79$

Deep Inelastic Scattering

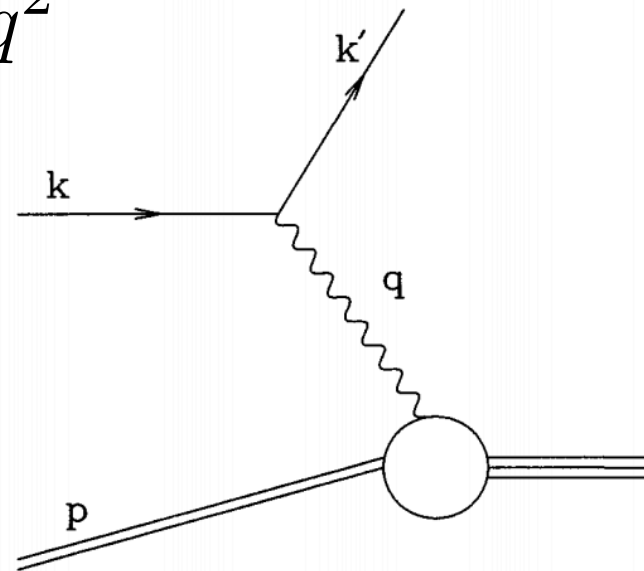
- Know p and k (from your beam/target)
- Measure k'
- This is enough to determine all of the following, with $Q^2 = -q^2$

$$M^2 = p^2$$

$$\nu = p \cdot q = M(E' - E)$$

$$x = \frac{Q^2}{2\nu} = \frac{Q^2}{2M(E - E')}$$

$$y = \frac{q \cdot p}{k \cdot p} = 1 - E'/E,$$



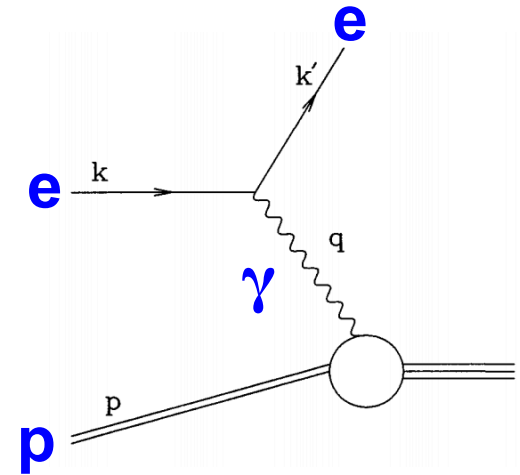
Bjorken x : $p_{\text{parton}} = x p_{\text{proton}}$

Deep Inelastic Scattering

- Charged lepton scattering

$$e + p \rightarrow e + X$$

$$\frac{d^2\sigma^{em}}{dxdy} = \frac{8\pi\alpha^2 ME}{Q^4} \left[\left(\frac{1 + (1-y)^2}{2} \right) 2xF_1^{em} + (1-y)(F_2^{em} - 2xF_1^{em}) - (M/2E)xyF_2^{em} \right]$$

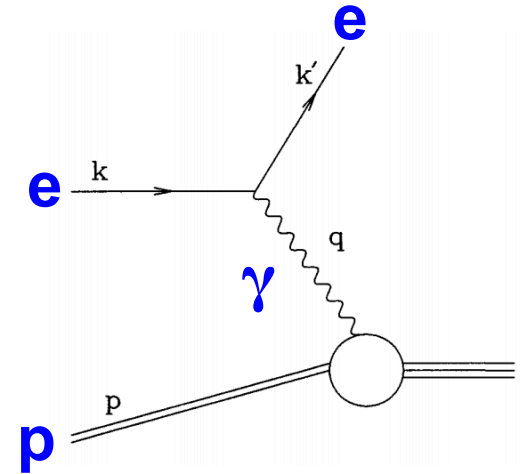


Deep Inelastic Scattering

- Charged lepton scattering

$$e + p \rightarrow e + X$$

$$\frac{d^2\sigma^{em}}{dxdy} = \frac{8\pi\alpha^2 ME}{Q^4} \left[\left(\frac{1 + (1-y)^2}{2} \right) 2x F_1^{em} + (1-y)(F_2^{em} - 2xF_1^{em}) - (M/2E)xyF_2^{em} \right]$$



- All information about the proton's structure is contained in **structure functions**

$$F_i(x, Q^2)$$

Bjorken limit

$$Q^2 \rightarrow \infty$$

- In this limit, the parton momentum is parallel to the proton momentum
 - Structure functions and PDFs are independent of Q^2
- The structure functions are sensitive to the quark PDFs by

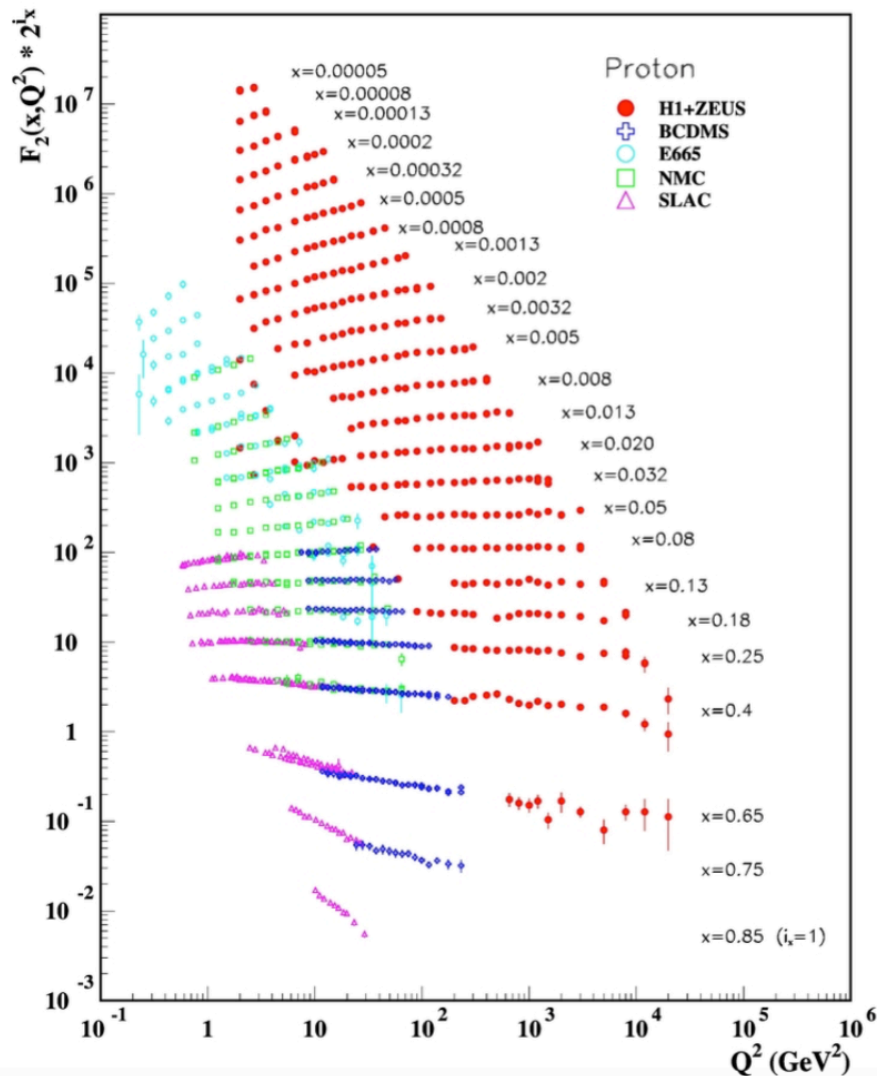
$$F_2^{em}(x) = 2xF_1^{em}(x) = \sum_{q, \bar{q}} e_q^2 x q(x)$$

Bjorken limit

$$Q^2 \rightarrow \infty$$

- No longer applies if we allow constituent quarks to emit a gluon
 - Gluon emission allows quarks to acquire momentum perpendicular to proton momentum
- Scaling violation: must consider dependence of structure functions (and PDFs) on Q^2
 - If we calculate the structure functions to \geq first order in $\alpha_S \sim g^2$, PDFs are $q(x, Q^2)$

Summary of DIS Experiments



- Can see the dependence of the structure function F_2 on x and Q^2
- PDFs are extracted from cross section measurements
 - e.g. H1 and ZEUS at the ep collider HERA

Cool, but...

isn't this talk about GPDs?

Form factors

$$F_1(Q^2) \text{ \& } F_2(Q^2)$$

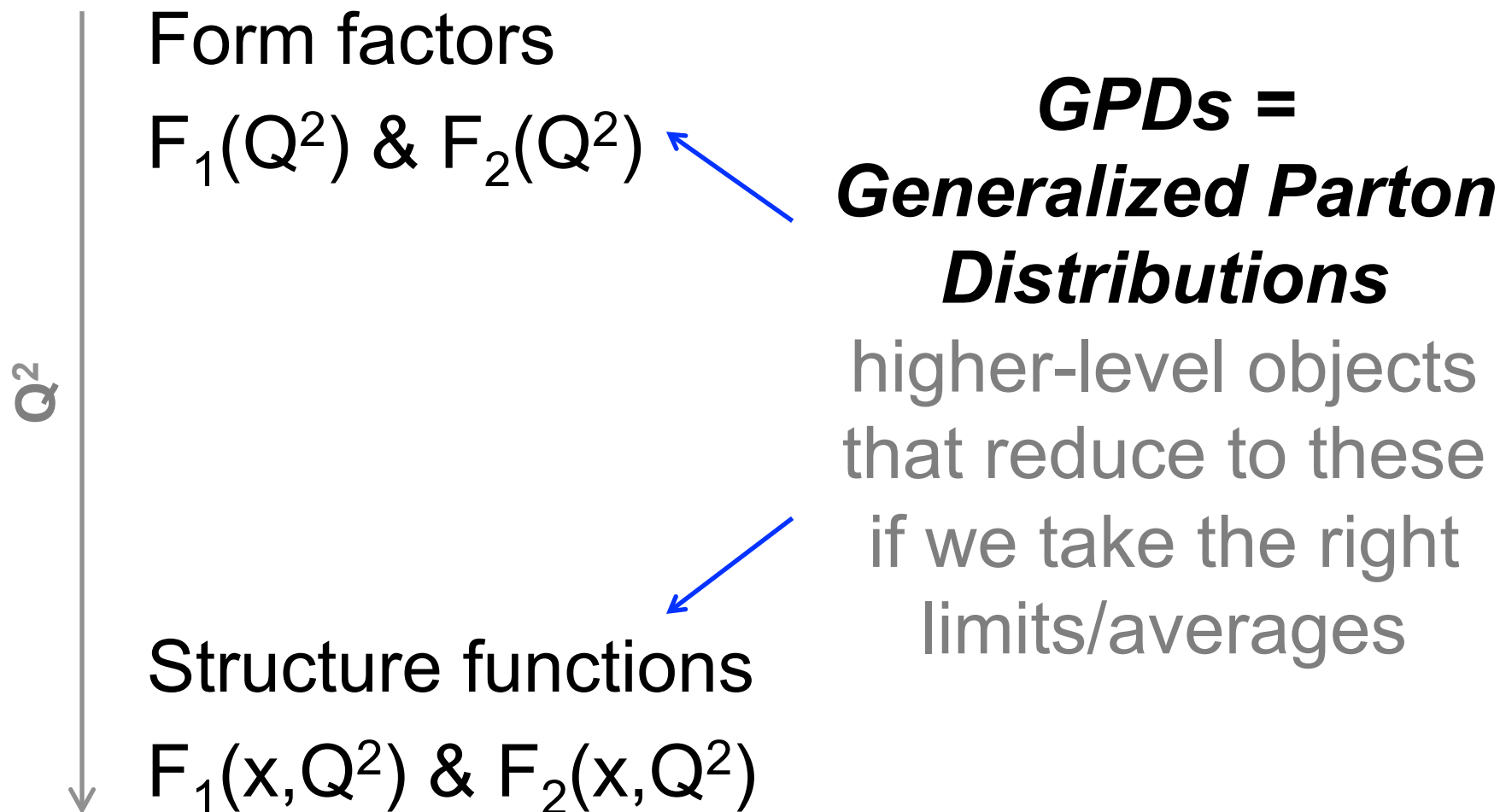
Q^2

Structure functions

$$F_1(x, Q^2) \text{ \& } F_2(x, Q^2)$$

Cool, but...

isn't this talk about GPDs?



Generalized Parton Distributions

- Each parton flavor has two GPDs
 - $H^q(x, \xi, t, Q^2)$: for when the proton helicity is unchanged
 - $E^q(x, \xi, t, Q^2)$: for when the proton helicity flips
- To understand the variables the GPDs depend on, let's look at the main process useful for probing them
 - Deeply Virtual Compton Scattering (DVCS)

Deeply Virtual Compton Scattering

$$e + p \rightarrow e + \gamma + p$$

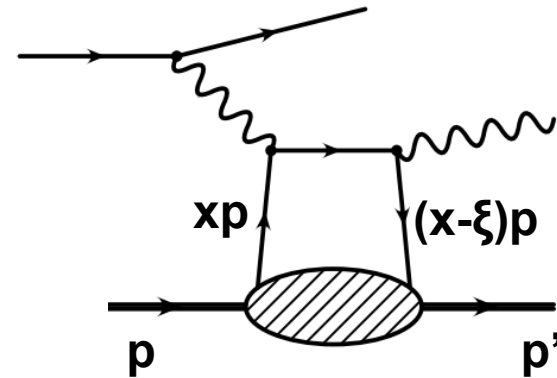
What variables do we use to describe the leading order DVCS diagram?

Q^2 = photon virtuality

Bjorken x

ξ tells you about the quark momentum carried away by γ

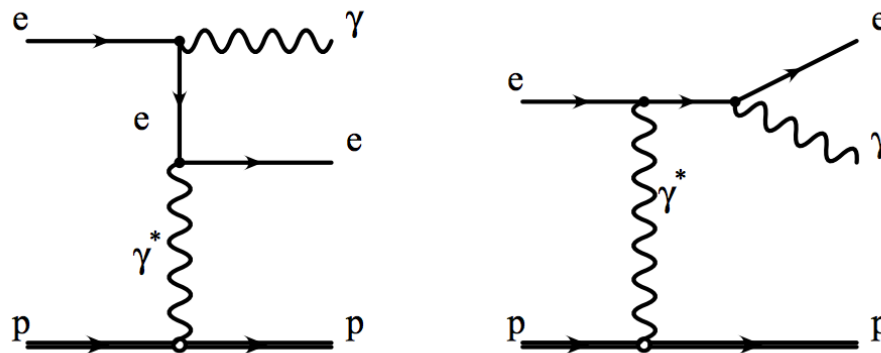
Mandelstam $t = (p - p')^2$



Background: Bethe-Heitler Process

Also $e + p \rightarrow e + \gamma + p$

- Here, a photon is emitted from the electron/positron line



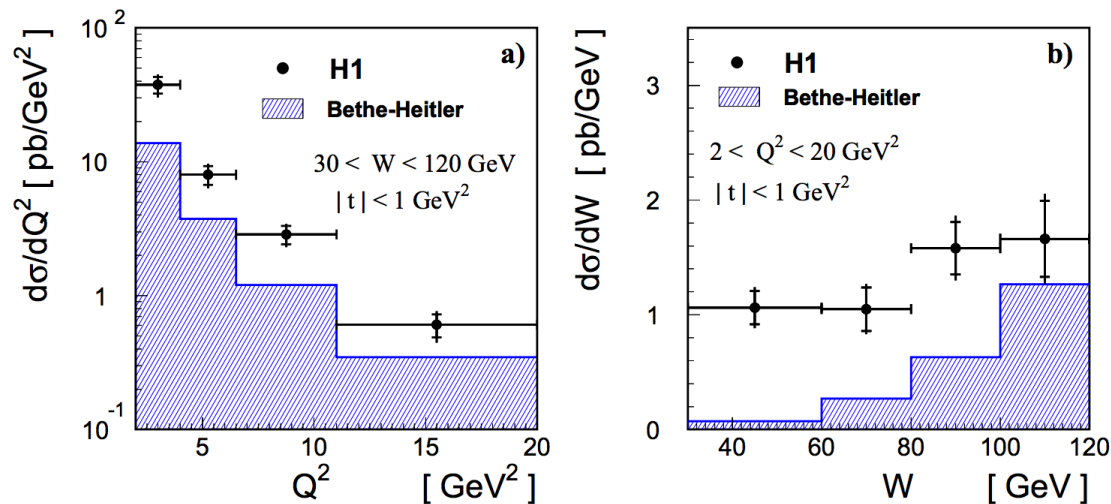
- BH contribution to the DVCS final state is known from QED and can be subtracted off
 - Interference vanishes when integrated over ϕ

DVCS results from H1 at HERA

- Positron-proton collisions
- For DVCS, require exactly two calorimeter clusters
 - Outgoing positron & photon, but no hadrons
- Small background from inelastic collisions
 - Proton remnants not detected
- Measuring cluster angles and energies gives information about x , ξ , and t

DVCS results from H1 at HERA

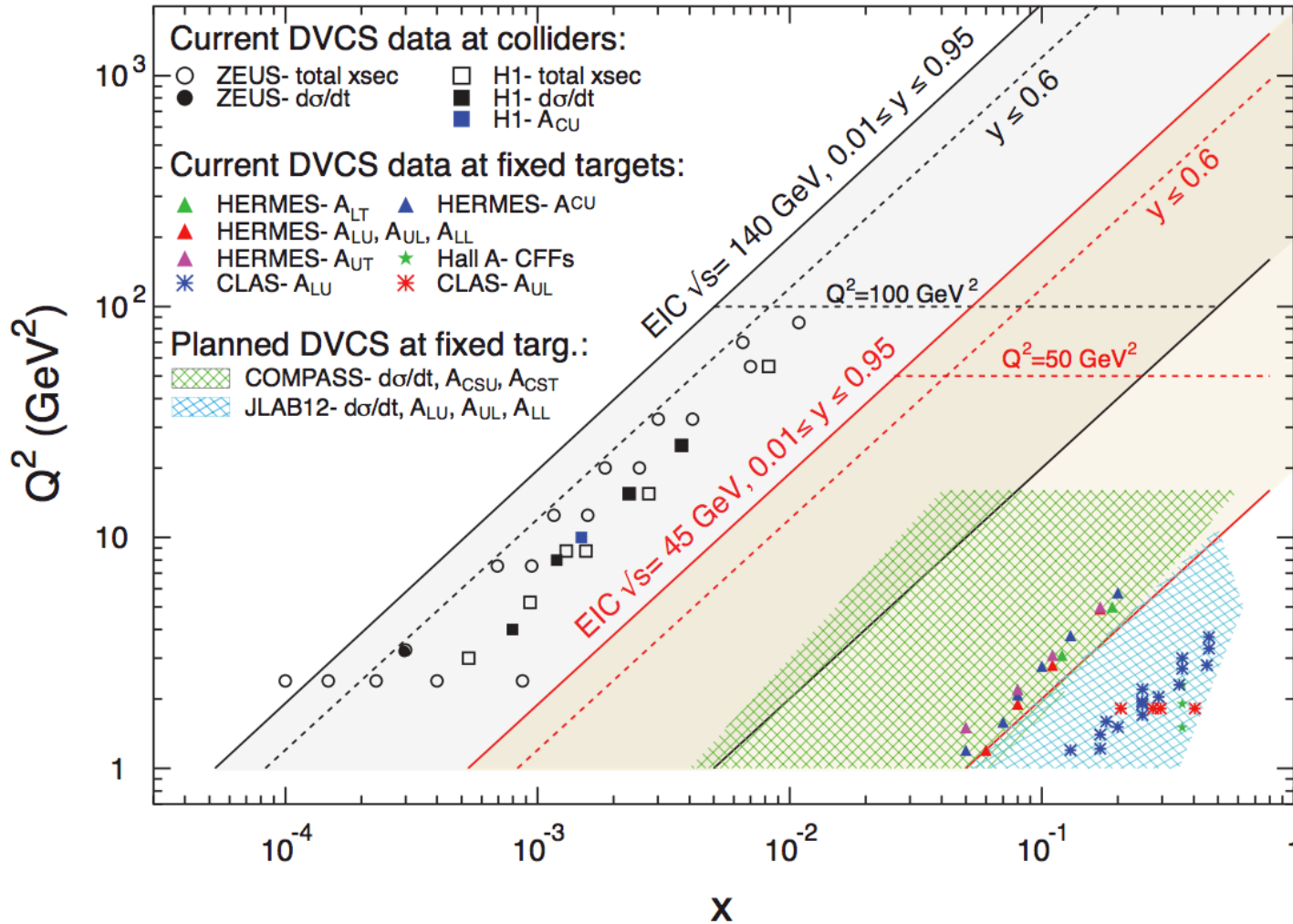
- Differential DVCS cross sections



- Q^2 and t are now familiar variables, but they introduce

$$W^2 = \frac{Q^2}{x} (1 - x)$$

Summary of DVCS data



Reducing the GPDs

To elastic form factors

- Take the limit $\xi = 0$, $Q^2 = t$.
 - This gets us back to $e + p \rightarrow e + p$
- Integrate over x , weighting each GPD by the charge of the corresponding quark
 - Since elastic scattering is not sensitive to the parton structure of the proton

$$\sum_q e_q \int dx H^q(x, 0, Q^2, Q^2) = F_1(Q^2)$$

$$\sum_q e_q \int dx E^q(x, 0, Q^2, Q^2) = F_2(Q^2)$$

Reducing the GPDs

To PDFs

- Throw away the GPDs E^q
 - It doesn't make sense to talk about proton helicity flip in DIS
- Take the limit $\xi = 0$
- Fourier transform t (transverse momentum info) to b (transverse position info)
- Integrate over b

$$\int db \tilde{H}^q(x, 0, b, Q^2) = q(x, Q^2)$$

Bigger scale: nuclear GPDs

When should we treat the nucleus as a bag of nucleons vs. as a bag of partons?

- GPDs and DVCS become very useful here
- One goal of a future EIC is to determine how nuclear GPDs are built up
 - Summing over nucleon GPDs? Convoluting nucleon GPDs with other functions?
Calculating nuclear PDFs?

Summary & Conclusions

- GPDs provide a high-level description of proton structure that simplifies to the form factors and structure functions
- GPDs are probed through Deeply Virtual Compton Scattering
 - At ep colliders and a future EIC
- Understanding proton GPDs is important for describing the structure of larger nuclei

Backup

PDFs at hadron-hadron colliders

- If we measure PDFs in ep and pp collisions, do we expect them to agree?
 - Do strong interactions between hadrons distort the PDFs?
- These interactions give corrections \sim powers of m^2/E_{CM}^2
 - Ok to neglect these at high energies
- So PDFs will be the same in ep and high energy pp experiments